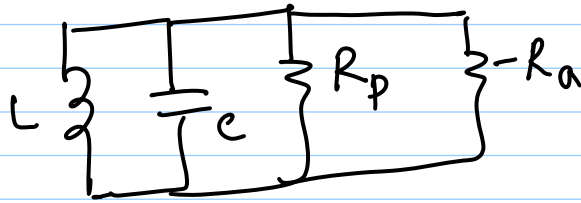


# Lecture 31 : Phase Noise (cont.)

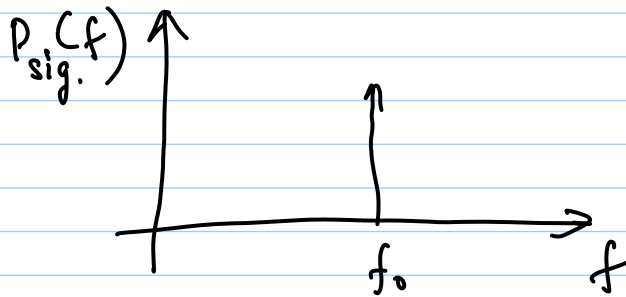
Note Title

10/13/2011

Last class : LTI analysis



Remember  $\Rightarrow |-Ra| = R_p$  only @ resonance  
 If  $-Ra$  were to cancel  $R_p$  @ all freq.,  
 $\Rightarrow$  ideal LC osc.!



at other freq.  $|-Ra| > R_p$ , but  $|-Ra| \sim R_p$

$$R_{eq.} = \frac{-Ra \cdot R_p}{R_p - Ra} = \frac{R_p Ra}{\underbrace{Ra - R_p}_{\text{very small}}} \gg R_p$$

$$Q_{eq.} = \omega_0 C \cdot R_{eq.} \gg Q_0$$

$$Y(\omega_0 + \Delta\omega) = G_{eq.} + j \frac{2\omega_0 \Delta\omega}{(\omega_0 + \Delta\omega)} \cdot C$$

$$\begin{aligned} \Rightarrow Z(\omega_0 + \Delta\omega) &= \frac{1}{G_{eq.} + j \frac{2\omega_0 \Delta\omega}{(\omega_0 + \Delta\omega)} \cdot C} \\ &= \frac{1}{G_{eq.}} \cdot \frac{1}{1 + j \frac{2\omega_0 \Delta\omega}{(\omega_0 + \Delta\omega)} R_{eq.} C} \end{aligned}$$

$$= R_{eq} \cdot \frac{1}{1 + 2Q_{eq} \cdot j \cdot \frac{\Delta \omega}{\omega_0}}$$

$$\approx \frac{\omega_0 R_{eq}}{j(2Q_{eq} \Delta \omega)} \quad \{ Q_{eq} \gg 1 \}$$

$$|Z(\omega_0 + \Delta \omega)| = \frac{\omega_0 R_{eq}}{2Q_{eq} \Delta \omega}$$

$$Q = \omega_0 R_p C$$

$$Q_{eq} = \omega_0 R_{eq} C$$

$$\Rightarrow \frac{\omega_0 R_{eq}}{Q_{eq}} = \frac{\omega_0 R_p}{Q} (= C)$$

$$\Rightarrow |Z(\omega_0 + \Delta \omega)| = \frac{\omega_0 R_p}{2Q \Delta \omega}$$

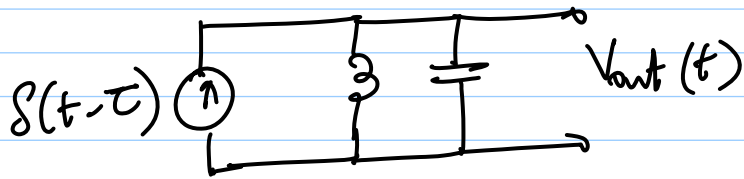
$$\Rightarrow \frac{\overline{v_n^2}}{\Delta f} = \frac{\overline{i_n^2}}{\Delta f} \cdot |Z|^2$$

$$= \frac{4kT}{R_p} \cdot \frac{(\omega_0 R_p)^2}{(2Q \Delta \omega)^2}$$

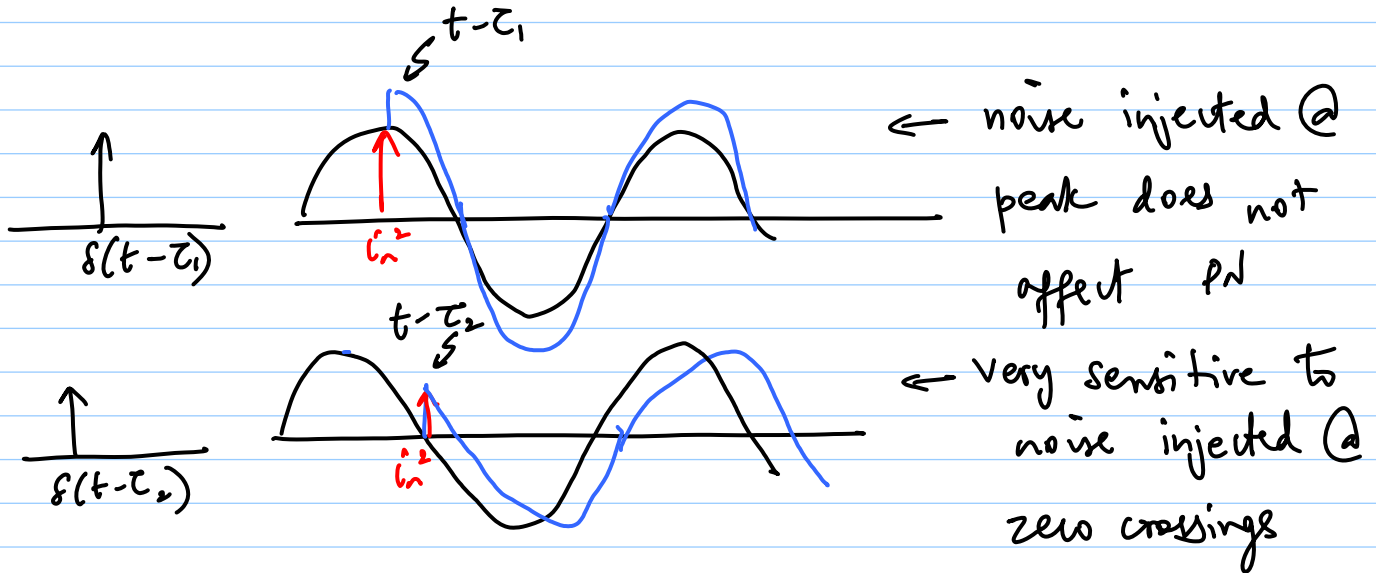
$$= 4kT R_p \cdot \left( \frac{\omega_0}{2Q \Delta \omega} \right)^2$$

rest of the analysis proceeds as before.

## Hajimiri - Lee PN model



$t$  = observation time  
 $\tau$  = excitation time



Impulse response (linearity still holds) for phase:

$$h_q(t, \tau) = \frac{\Gamma(\omega_0 \tau)}{q_{\max.}} u(t - \tau)$$

$q_{\max.}$  = max charge displacement across cap.  
 (to make  $\Gamma(\omega_0 \tau)$  ampl. independent)

$$\phi(t) = \frac{1}{q_{\max.}} \int_{-\infty}^t \Gamma(\omega_0 \tau) \cdot i(\tau) d\tau$$

$i(\tau)$  = noise current

$$\Gamma(\omega_0 \tau) = \text{ISF (normalised)}$$

$\phi(t)$  ← total phase @  $t$  = sum of all phase disturbances due to  $i_n(t)$  from  $-\infty$  to  $t$

ISF  $\Gamma \rightarrow$  dimensionless,  
 $\rightarrow$  freq. & amplitude independent  
 $\rightarrow$  periodic w/ period  $2\pi$

$$\Gamma(\omega_0 t) = \frac{C_0}{2} + \sum_{n=1}^{\infty} C_n \cos(n\omega_0 t + \theta_n)$$

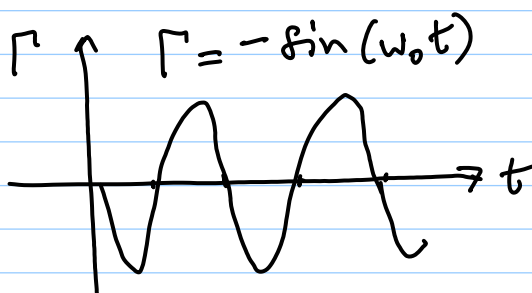
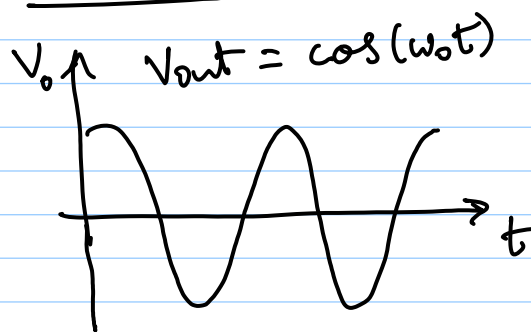
$C_n$  = Fourier Series coeffs.

$\theta_n$  = phase of  $n^{\text{th}}$  harmonic of ISF

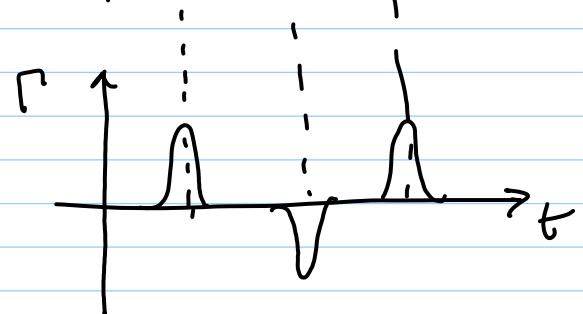
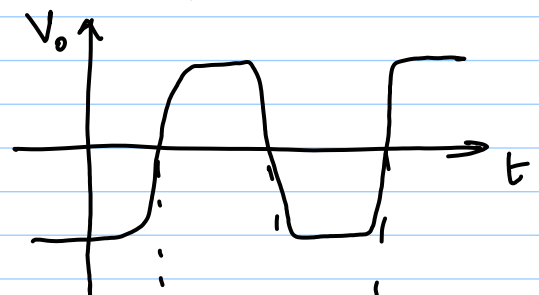
Note that from Parseval's Theorem,

$$\sum_{n=0}^{\infty} C_n^2 = \frac{1}{\pi} \int_0^{2\pi} |\Gamma(\varphi)|^2 d\varphi = 2 \Gamma_{\text{rms}}^2$$

LC osc.



Ring Osc.



\* In general, determine ISF accurately from simulation

$$\phi(t) = \frac{1}{q_{max.}} \left[ \frac{C_m}{2} \int_{-\infty}^t i(\tau) d\tau \right.$$

$$\left. + \sum_{n=1}^{\infty} C_n \int_{-\infty}^t i(\tau) \cos(n\omega_0 \tau) d\tau \right]$$

consider sinusoidal current at  $(m\omega_0 + \Delta\omega)$

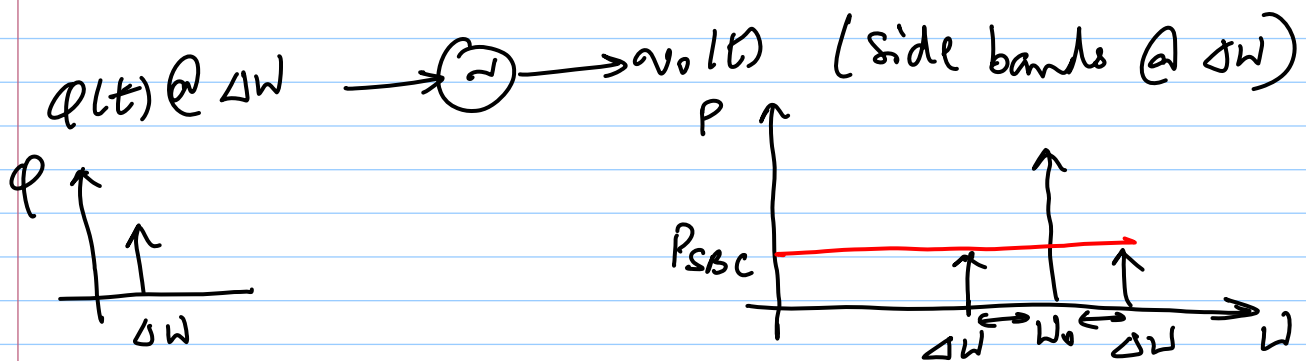
$$i(t) = I_m \cos[(m\omega_0 + \Delta\omega)t]$$

$$\Rightarrow \phi(t) \approx \frac{I_m C_m \sin(\Delta\omega t)}{2q_{max.} \Delta\omega} \quad \left\{ \begin{array}{l} \text{Even though noise} \\ \text{is @ } m\omega_0 + \Delta\omega \end{array} \right\}$$

$$v_o(t) = \cos(\omega_0 t + \phi(t)) \Rightarrow \text{phase to voltage converter}$$

(square  $V_o^2(t)$  to get power)

→ fundamentally nonlinear (PM)



$$P_{SBC}(\Delta\omega) \approx 10 \log \left[ \frac{I_m^2 C_m^2}{4q_{max.}^2 \Delta\omega^2} \right] \quad \begin{array}{l} \text{narrowband} \\ \text{FM approx.} \end{array}$$

\* see Razavi 3.2.2 for narrowband FM approx.

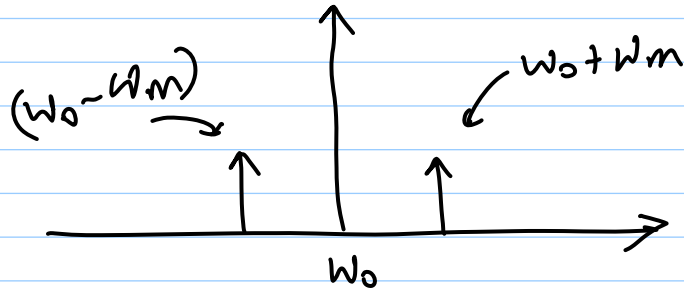
$V_c + V_m \cos \omega_m t$

$$v_o(t) = A_c \cos \left( \omega_0 t + \int_{-\infty}^t v_c(t) dt \right)$$

$\omega_m \ll \omega_0$

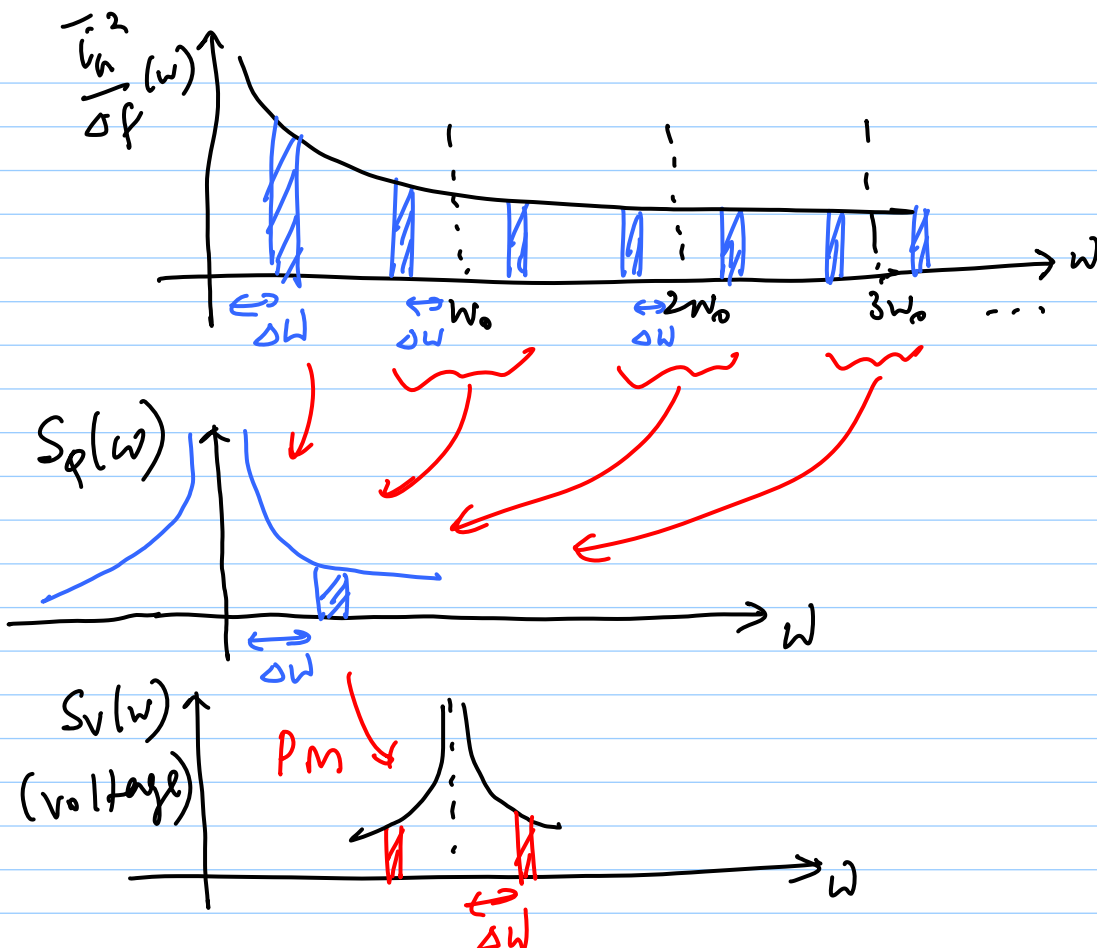
Use Bessel function approx.

$$V_o(t) \approx A_o \cos \omega_o t + \frac{A_o V_m K_{VQ}}{2\omega_m} \cdot \left[ \cos(\omega_o + \omega_m)t - \cos(\omega_o - \omega_m)t \right]$$



for a white noise source

$$P_{SBc}(\Delta\omega) \approx 10 \log \left[ \frac{(\bar{i}_n^2 / \Delta f) \cdot \sum_{m=0}^{\infty} C_m^2}{4q^2 \Delta\omega^2} \right]$$



Thankfully  $|C_m|$  reduces as  $m \uparrow$ , so only the first few terms are significant  
 Spectrums in  $1/f^2$  region:

$$L(\Delta\omega) = 10 \log \left[ \frac{(\overline{i_n^2}/\Delta f) \cdot \Gamma_{rms}^2}{2q_{max}^2 \Delta\omega^2} \right]$$

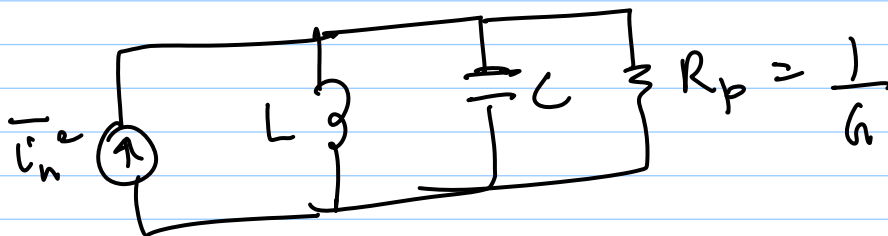
$\Rightarrow$  reduce  $\Gamma_{rms}$  to reduce  $PN$

$C = \text{tank cap.}$   
 $V_{pk} = \text{peak amplitude across tank}$  }  $q_{max} = C \cdot V_{pk}$

$$L(\Delta\omega) = 10 \log \left[ \frac{\overline{i_n^2} \Gamma_{rms}^2}{2(CV_{pk}^2) (\Delta\omega)^2} \right]$$

consider an LC oscillator

$$\Gamma(\omega_0 t) = -\sin(\omega_0 t) \Rightarrow \Gamma_{rms}^2 = \frac{1}{2}$$



$$L(\Delta\omega) = 10 \log \left[ \frac{4kT_G \cdot 1/2}{2(CV_{pk})^2 (\Delta\omega)^2} \right]$$

$$= 10 \log \left[ \frac{kT_G}{C^2 \cdot (2V_0^2)} \cdot \frac{1}{(\Delta\omega)^2} \right]$$

$$= 10 \log \left[ \frac{kT_G}{2\omega_0^2 C^2 V_0^2} \cdot \left( \frac{\omega_0}{\Delta\omega} \right)^2 \right]$$

$$= 10 \log \left[ \frac{kT \cdot 1/R}{2V_0^2 \cdot (Q^2/R^2)} \cdot \left( \frac{\omega_0}{\Delta\omega} \right)^2 \right]$$

$$= 10 \log \left[ \frac{2kTR}{V_0^2} \cdot \left( \frac{\omega_0}{2Q\Delta\omega} \right)^2 \right] \quad \leftarrow \text{Same as Leeson's linear result}$$

### Flicker noise

$$\overline{i_n^2}_{1/f} = \overline{i_n^2} \cdot \frac{\omega_{1/f}}{\Delta\omega} \quad \text{in the } 1/f \text{ noise region}$$

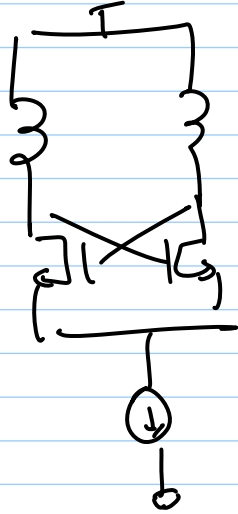
$$\Rightarrow L(\Delta\omega) = 10 \log \left[ \frac{(\overline{i_n^2}/\omega_f) \cdot C^2}{8q_{\text{min}}^2 \cdot \Delta\omega^2} \cdot \frac{\omega_{1/f}}{\Delta\omega} \right] \quad \leftarrow \text{flicker noise only around DC!}$$

$$\Rightarrow \Delta\omega_{1/f^3} = \omega_{1/f} \cdot \frac{C^2}{4\overline{i_{\text{rms}}}^2} = \omega_{1/f} \cdot \left( \frac{\Gamma_{\text{dc}}}{\Gamma_{\text{rms}}} \right)^2$$

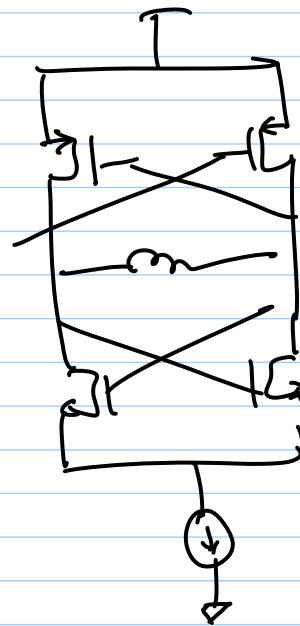
$\Rightarrow 1/f^3$  noise corner can be reduced by decreasing  $\Gamma_{\text{dc}}$  (also  $\neq \omega_{1/f}$ )



$\Gamma_{ac} = 0 \Rightarrow$  Symmetrical rise & fall times!

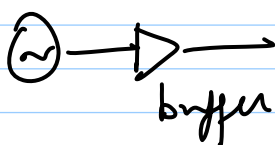
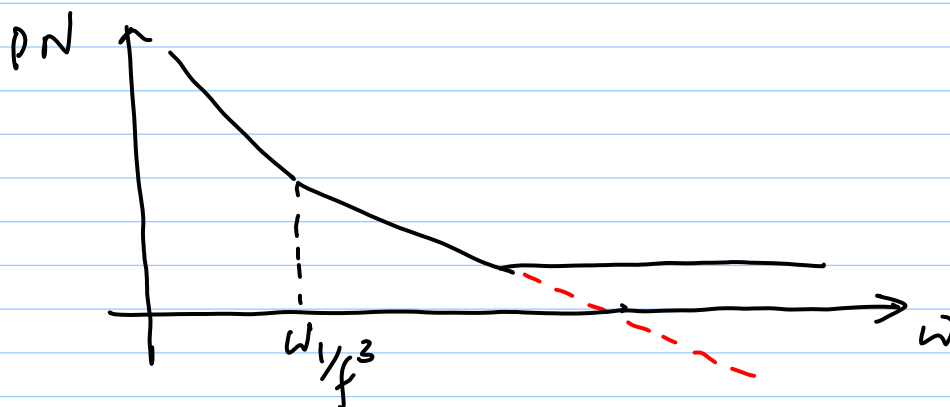


$t_r \neq t_f$



$t_r = t_f$   
is possible  
by careful  
Design!

\* One last point: why do you have a flat noise region?



\* noise of buffer

\* noise floor of meas. equipment  
(if it is high)