

Lecture 30 : Phase Noise

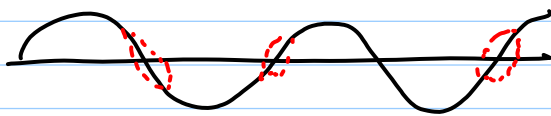
- * VCO is an autonomous clkt (i.e. its phase is not tied to any reference)
- * ideal VCO: $V(t) = V_0 \cos(\omega_0 t)$
- * Real VCO: random amplitude/phase changes due to noise

$$V(t) = V_0(t) \cos(\omega_0 t + \phi(t))$$

Amplitude
noise

Phase noise

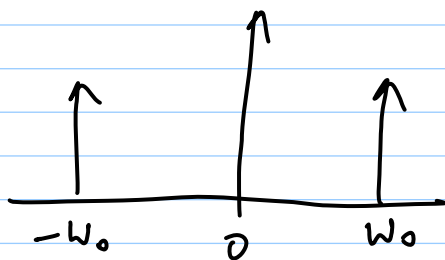
\equiv Jitter in time domain



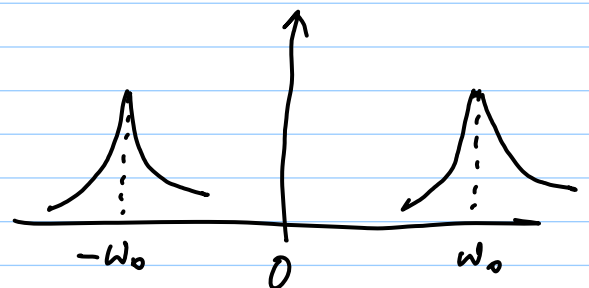
← zero crossings are modulated by noise

Instantaneous freq.

$$\omega = \omega_0 + \frac{d\phi(t)}{dt}$$

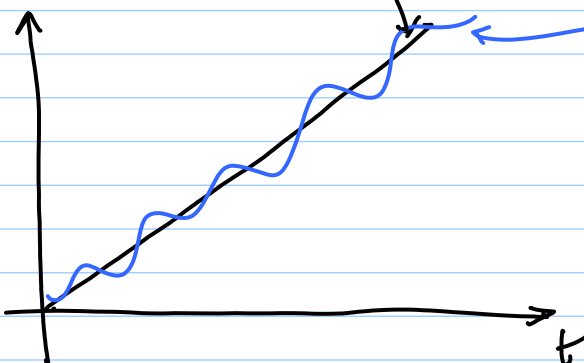


Ideal Osc.



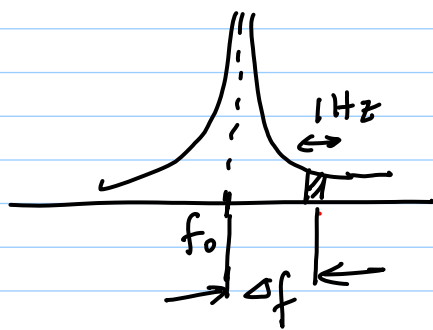
Real osc.

$\phi(t)$
(total phase)



$\phi(t)$ slope = ω_0

SSB phase noise



$$L(\Delta f) = 10 \cdot \log_{10} \left[\frac{P_{1\text{Hz}}(f_0 + \Delta f)}{P_s} \right]$$

Units : dBc/Hz

* $P_{1\text{Hz}}(f_0 + \Delta f) \equiv$ power in a 1 Hz BW at offset Δf from f_0

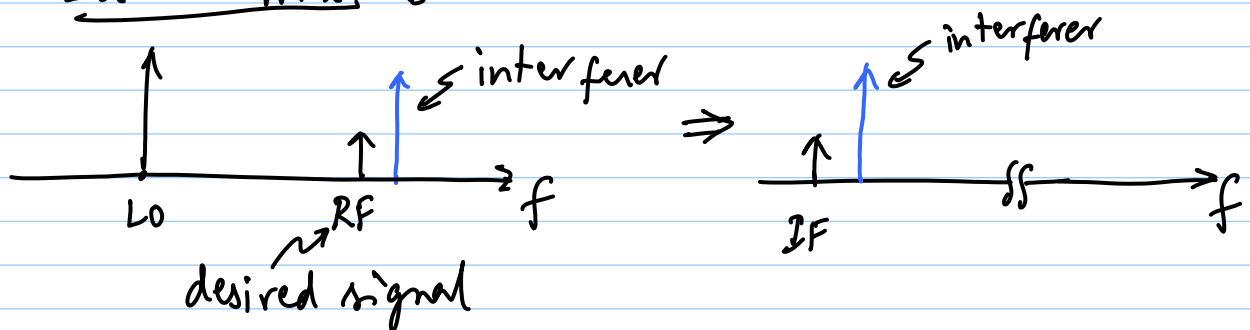
* $P_s =$ mean-square carrier power

* $L(\Delta f)$ is expressed in units of dB's below carrier per Hz {at offset Δf }

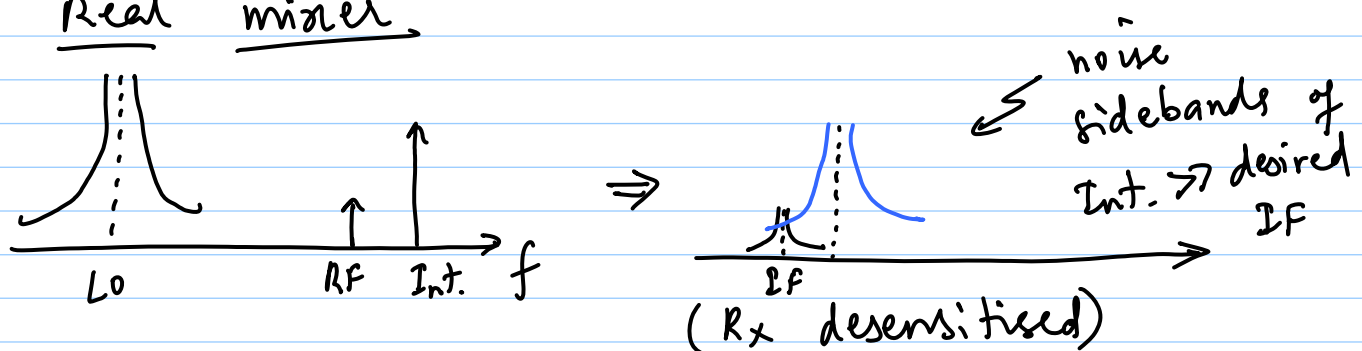
Effects of PN

1) Reciprocal Mixing

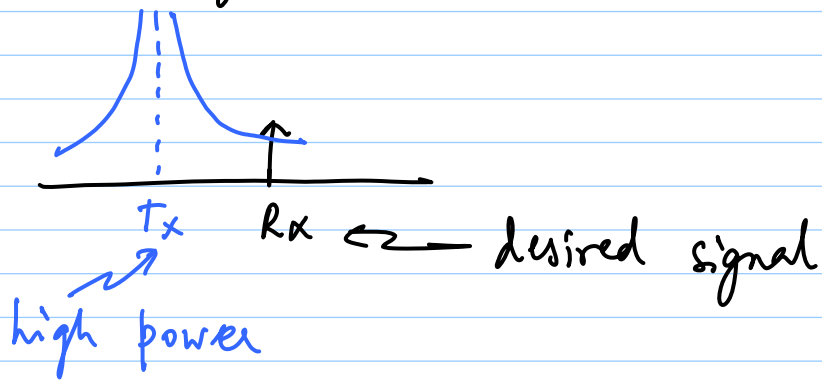
Ideal mixer :



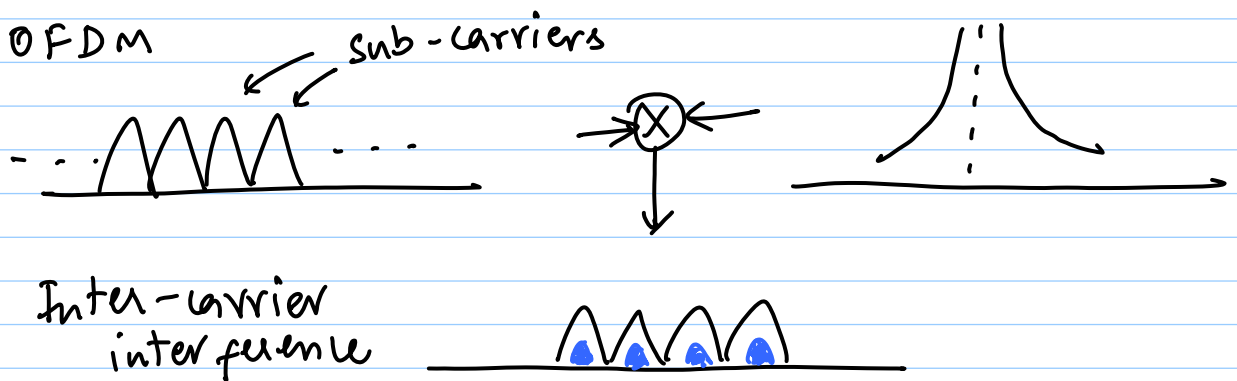
Real mixer



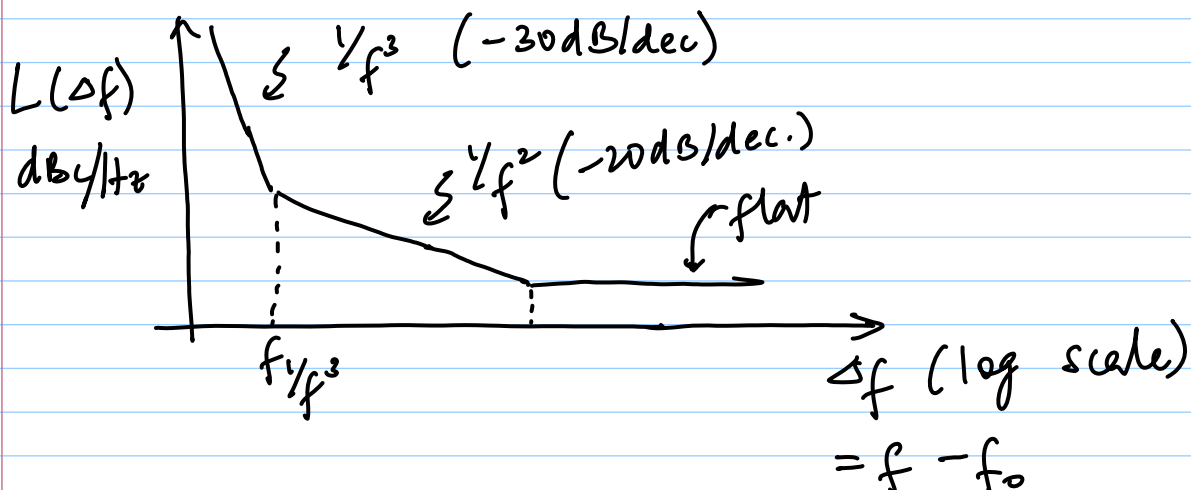
- 2) Powerful Tx with noisy LO can desensitize nearby Rx with weak signal



- 3) OFDM



Phase noise (zoomed in)



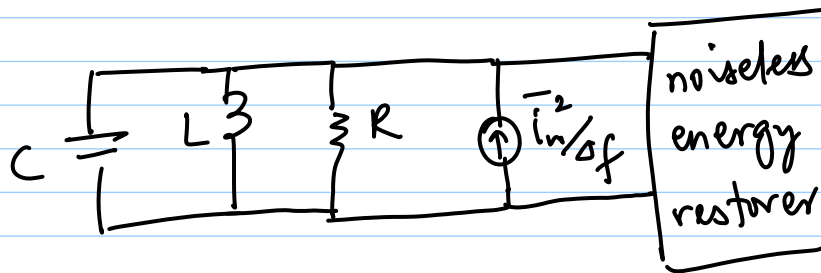
Leeson's phase-noise model

Ref: D. B. Leeson, "A simple model of feedback oscillator noise spectrum", Proceedings of the IEEE, pp. 320-330, Feb. 1966.

* LTI analysis of LC oscillator in steady-state using feedback theory

* Assumes noise from resistor only

→ excess noise compensated using empirical factor 'F'



idealised
osc. model

$$\overline{i_n^2} = 4kTG \quad A^2/Hz$$

$$\text{Energy } E = \frac{1}{2} CV_c^2 \Rightarrow V_c^2 = \frac{E}{C}$$

$$Q = \omega \cdot \frac{E}{P_D}$$

$\overline{v_n^2}$ = total mean squared noise over entire BW

$$= 4kTG \int_0^\infty |Z(f)|^2 df$$

$$\text{when } Z(s) = \frac{sL}{s^2 LC + \frac{sL}{R} + 1}$$

$$= \frac{kT}{C}$$

$$N/S = \frac{\overline{v_n^2}}{V_c^2} = \frac{kT}{E} = \frac{WkT}{Q \cdot P_D}$$

At a small freq offset $\Delta\omega$,

$$Y(\omega_0 + \Delta\omega) = G + j(\omega_0 + \Delta\omega) \cdot C + \frac{1}{j(\omega_0 + \Delta\omega)L}$$

$$= \frac{jG(\omega_0 + \Delta\omega) \cdot L - (\omega_0 + \Delta\omega)^2 LC + 1}{j(\omega_0 + \Delta\omega) \cdot L}$$

$$\approx \frac{jG(\omega_0 + \Delta\omega) - 2\omega_0 \Delta\omega C}{j(\omega_0 + \Delta\omega)}$$

$$= G + j \frac{2\omega_0 \Delta\omega}{(\omega_0 + \Delta\omega)} \cdot C$$

$$\begin{aligned} Z(\omega_0 + \Delta\omega) &= \frac{1}{G + j \frac{2\omega_0 \Delta\omega}{(\omega_0 + \Delta\omega)} \cdot C} \\ &= \frac{1}{G} \cdot \frac{1}{1 + 2j \frac{\omega_0 \cdot \Delta\omega}{(\omega_0 + \Delta\omega)} \cdot RC} \end{aligned}$$

Since $Q = \omega_0 RC$,

$$Z(\omega_0 + \Delta\omega) \approx R \cdot \frac{1}{1 + 2Qj \frac{\Delta\omega}{\omega_0}}$$

$$|Z(\omega_0 + \Delta\omega)| \approx \frac{R\omega_0}{2Q \Delta\omega}$$

$$\begin{aligned}
 \frac{\overline{v_n^2}}{\Delta f} &= \frac{\overline{i_n^2}}{\Delta f} \cdot |Z|^2 \\
 &= 4kT/R \cdot R^2 \cdot \frac{\omega_0^2}{(2Q\Delta\omega)^2} \\
 &= 4kTR \cdot \left(\frac{\omega_0}{2Q\Delta\omega}\right)^2
 \end{aligned}$$

* includes both amplitude & phase noise

→ Equipartition theorem of Thermodynamics states:

Noise $\begin{cases} \rightarrow \text{Ampl. Noise} \\ \rightarrow \text{Phase Noise} \end{cases} \begin{cases} \text{equally} \\ \text{split} \end{cases}$

→ assume amplitude limiter eliminates all amplitude noise

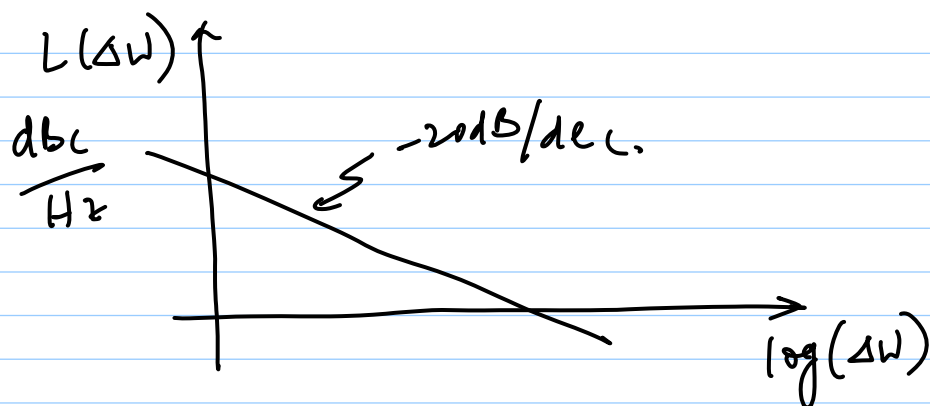
Phase component of Noise

$$= \frac{1}{2} \cdot \frac{\overline{v_n^2}}{\Delta f} = 2kTR \cdot \left(\frac{\omega_0}{2Q\Delta\omega}\right)^2$$

$$L(\Delta\omega) = 10 \log_{10} \left[\frac{\frac{1}{2} \cdot \overline{v_n^2} / \Delta f}{v_{sig}^2} \right]$$

$$= 10 \log_{10} \left[\frac{2kTR \left(\frac{\omega_0}{2Q\Delta\omega}\right)^2}{v_{sig}^2} \right]$$

$$= 10 \log_{10} \left[\frac{2kT}{P_{sig.}} \cdot \left(\frac{\omega_0}{2Q\Delta\omega}\right)^2 \right]$$



* we know actual VCO PN has different slopes not accounted for in this simplistic model

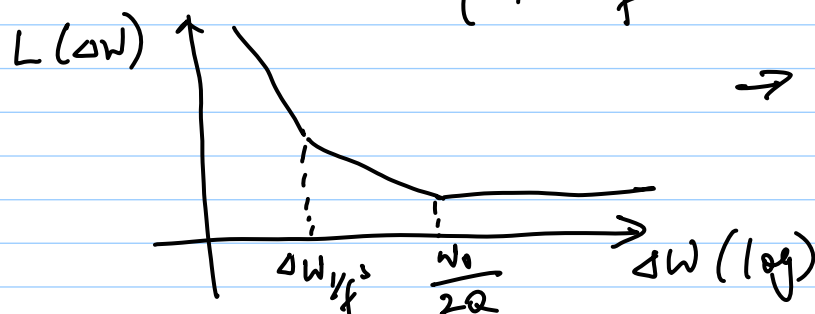
- $1/f^3$ region
- flat noise floor
- Device noise

Leeson's model:

$$L(\Delta W) = 10 \cdot \log_{10} \left[\frac{2KTF}{P_{\text{sig.}}} \left\{ 1 + \left(\frac{W_0}{2Q\Delta W} \right)^2 \right\} \left\{ 1 + \frac{\Delta W_{1/f^2}}{|\Delta W|} \right\} \right]$$

→ F is an empirical factor {device excess noise number}

→ $\Delta W_{1/f^3} = 1/f^3$ noise corner
($\neq 1/f$ corner of device)



→ Q = Quality factor
with loading

Leeson's model

- accounts for trends in all regions of $L(\Delta\omega)$
- provides some insight
 - e.g. $\omega_0 \uparrow \Rightarrow \text{PN} \uparrow$ (6dB for every 2x)
 - $Q \uparrow \Rightarrow \text{PN} \downarrow$ (?? for every 2x)
- provides no info on F & $\Delta\omega_{1/f^3}$
- We know that oscillators
 - have non-linearity
 - are time-variant

Rewrite:

$$L(\Delta\omega) = 10 \log_{10} \left[\frac{2FKT}{P_{\text{sig.}}} \left(\frac{\omega_0}{2Q\Delta\omega} \right)^2 \right]$$

$$= 10 \log_{10} \left[\frac{2FKT}{\left(\frac{2I_T}{\pi} \cdot R_P \right)^2} \cdot \left(\frac{\omega_0}{2Q\Delta\omega} \right)^2 \right]$$

$$= 10 \log_{10} \left[\frac{(\pi^2/2)FKT}{I_T^2} \cdot \left(\frac{\omega_0}{QL} \right) \cdot \left(\frac{1}{2Q\Delta\omega} \right)^2 \right]$$

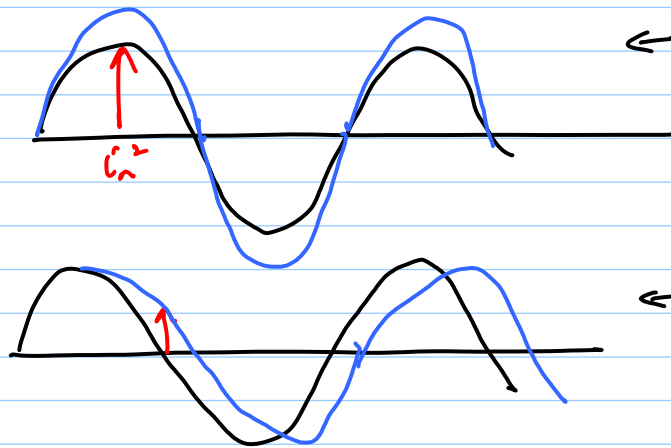
$$\text{i.e. } \text{PN} \propto 10 \log_{10} \left[\frac{1}{Q^3} \cdot (\dots) \right]$$

More sophisticated PN models

- 1) A. Hajimiri & T.H. Lee, "A General Theory of phase-noise in oscillators"

* Oscillators are time-variant (as far as noise is concerned)

e.g.



← noise injected @ peak does not affect PN

← very sensitive to noise injected @ zero crossings

* $\overline{v_n^2} \ll v_{sig}^2 \Rightarrow$ Linearity assumption is valid for noise

\Rightarrow LTV model

* "Impulse Sensitivity Function" or ISF can be determined for an oscillator
 \rightarrow PN is determined based on this

- 2) J. Rael & A. Abidi, "Physical processes of phase noise in differential LC oscillators"

IEEE Custom Integrated Circuits Conf., 2000

* Non-linear, time-invariant analysis

* Similar to mixer noise analysis

* Expression for F in Leeson's Model :

$$F = 2 + \frac{8\gamma R I_T}{\pi V_o} + \gamma \cdot \frac{8}{9} g_m R$$

where $V_o = \frac{2}{\pi} I_T R = \text{peak diff. ampl.}$

$$P_{sig.} = \frac{V_o^2}{R}$$

* If I_T is increased while keeping g_m constant $\Rightarrow F$ remains constant

$$\Rightarrow L(\omega) = 10 \log_{10} \left[\frac{4FkT}{V_o^2} (\dots) \right] \downarrow$$

P_N decreases due to increase in V_o

* If I_T is increased beyond supply rail limitation, V_o does not increase but $F \uparrow$ & $P_N \uparrow$