# Excitations, Emergent Facilitation and Glassy Dynamics in Supercooled Liquids

 $\mathsf{Muhammad}\ \mathsf{R}.\ \mathsf{Hasyim}^1,\ \mathsf{Kranthi}\ \mathsf{K}.\ \mathsf{Mandadapu}^2$ 

APS March Meeting 2024, March 4, 2024

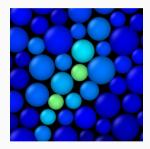
 $<sup>^{1}</sup>$ Department of Chemical and Biomolecular Engineering, University of California, Berkeley, CA 94720

<sup>&</sup>lt;sup>2</sup>Chemical Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

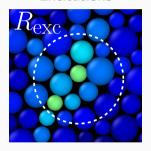
#### **Excitations**



Below the onset temperature  $T_{\rm o}$ , localized excitations drive glassy dynamics! (Keys, et al. *Phys. Rev. X* 2011)

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

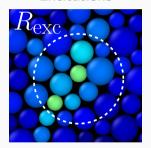
#### **Excitations**



Below the onset temperature  $T_{\rm o}$ , localized excitations drive glassy dynamics! (Keys, et al. *Phys. Rev. X* 2011)

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

#### **Excitations**

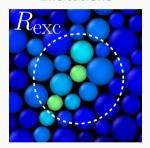


Energy Cost  $J_{\sigma}$ 

Below the onset temperature  $T_{\rm o}$ , localized excitations drive glassy dynamics! (Keys, et al. *Phys. Rev. X* 2011)

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

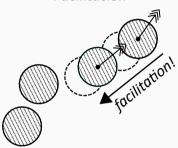
#### **Excitations**



Energy Cost  $J_{\sigma}$ 

Below the onset temperature  $T_{\rm o}$ , localized excitations drive glassy dynamics! (Keys, et al. *Phys. Rev. X* 2011)

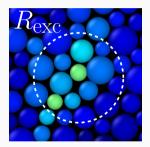
#### **Facilitation**



Excitations facilitate the creation and relaxation of another excitation close by.

DF theory (Chandler and Garrahan, *Annu. Rev. Phys. Chem*, 2010) can explain glassy dynamics without knowledge of local structure with two key ideas:

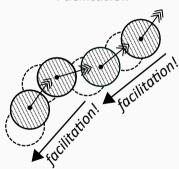
#### **Excitations**



Energy Cost  $J_{\sigma}$ 

Below the onset temperature  $T_{\rm o}$ , localized excitations drive glassy dynamics! (Keys, et al. *Phys. Rev. X* 2011)

### **Facilitation**



#### A Theory of Onset Temperature in 2D

#### Inherent-State Melting and the Onset of Glassy Dynamics in Two-Dimensional Supercooled Liquids

Dimitrios Fraggedakis, <sup>1</sup>, <sup>1</sup> Muhammad R. Hasyim, <sup>1</sup>, <sup>1</sup> and Kranthi K. Mandadapu<sup>1</sup>, <sup>2</sup>, <sup>†</sup>

<sup>1</sup>Department of Chemical & Biomolecular Engineering, University of California, Berkeley, CA 94720

<sup>2</sup>Chemical Sciences Division, Lawrence Berkeley National Laboratory, Berkeley, CA 94720

(Datet: April 18, 2022)

Below the onset temperature T<sub>o</sub>, the equilibrium relaxation time of most glass-forming liquids exhibits glassy dynamics characterized by super-Arrhenius temperature dependence. In this supercooled regime, the relaxation dynamics also proceeds through localized elastic excitations corresponding to hopping events between inherent states, i.e., potential-energy minimizing configurations of the liquid. Despite its importance in distinguishing the supercooled regime from the hightemperature regime, the microscopic origin of T<sub>c</sub> is not yet known. Here, we construct a theory for the onset temperature in two dimensions and find that inherent-state melting transition, described by the binding-unbinding transition of dipolar elastic excitations, delineates the supercooled regime from the high-temperature regime. The corresponding melting transition temperature is in good agreement with the onset temperature found in various two-dimensional atomistic models of glass formers. We discuss the predictions of our thory on the displacement and density correlations of two-dimensional appercooled liquids, which are consistent with observations of the Mermin-Wagner

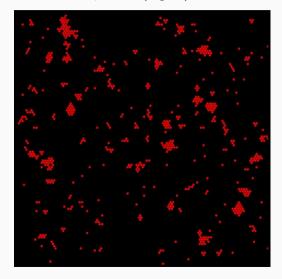
Keywords: Two-dimensional glassy dynamics, Kosterlitz-Thouless transition, excitations, geometric charges

D. Fraggedakis, M.R. Hasyim, K.K. Mandadapu, PNAS 120 (14) e2209144120 (2023)

$$\beta J_{\sigma} = 1$$
 (High-T)

Origin of Glassy Behavior
A combination of emergent constraints
and revertibility of dynamics.

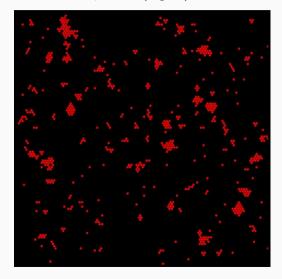
$$\beta J_{\sigma} = 1$$
 (High-T)



# Origin of Glassy Behavior A combination of emergent constraints and revertibility of dynamics.

 Emergent Constraints: Dynamics only persist in a few spots of the liquid.

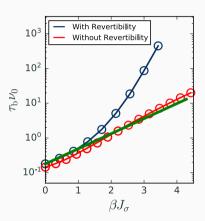
$$\beta J_{\sigma} = 1$$
 (High-T)



# Origin of Glassy Behavior A combination of emergent constraints and revertibility of dynamics.

- Emergent Constraints: Dynamics only persist in a few spots of the liquid.
- Revertibility: Within an active region, excitations can go back and forth, wasting time in the process.

$$\beta J_{\sigma} = 1$$
 (High-T)

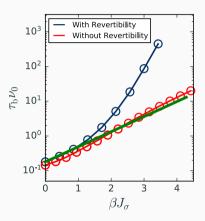


Relaxation time data when revertibility is turned off (red markers) along with an Arrhenius timescale  $au_{\rm arr}\sim e^{\beta J_\sigma}$ .

# Origin of Glassy Behavior A combination of emergent constraints and revertibility of dynamics.

- Emergent Constraints: Dynamics only persist in a few spots of the liquid.
- Revertibility: Within an active region, excitations can go back and forth, wasting time in the process.

$$\beta J_{\sigma} = 1$$
 (High-T)



Relaxation time data when revertibility is turned off (red markers) along with an Arrhenius timescale  $au_{\rm arr}\sim e^{\beta J_\sigma}$ .

# Origin of Glassy Behavior A combination of emergent constraints and revertibility of dynamics.

- Emergent Constraints: Dynamics only persist in a few spots of the liquid.
- Revertibility: Within an active region, excitations can go back and forth, wasting time in the process.

No revertibility, no super-Arrhenius behavior.