# ELECTRIC POTENTIAL AND POTENTIAL ENERGY

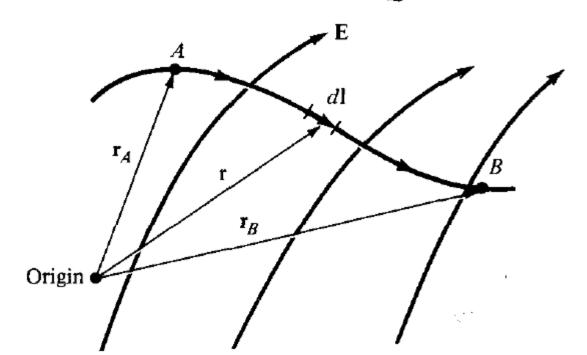
#### Introduction

- The electric field intensity **E** due to a charge distribution can be obtained from Coulomb's law in general or from Gauss's law when the charge distribution is symmetric
- It would be desirable if we could find some as yet undefined scalar function with a single integration and then determine the electric field from this scalar by some simple straightforward procedure, such as differentiation
- This scalar function does exist and is known as the *potential* or scalar potential field
- ➤The method of obtaining E from the electric scalar potential V is discussed in this lecture

#### **Electric Potential**

- ➤ Suppose we wish to move a point charge Q from point A to point B in an electric field E
- From Coulomb's law, the force on Q is F = QE so that the work done in displacing the charge by dl is:

$$dW = -\mathbf{F} \cdot d\mathbf{l} = -Q\mathbf{E} \cdot d\mathbf{l}$$



#### **Electric Potential**

- The negative sign indicates that the work is being done by an external agent
- Thus, the total work done, or the potential energy required, in moving Q from A to B is:

$$W = -Q \int_A^B \mathbf{E} \cdot d\mathbf{l}$$

Polividing W by Q in the above equation gives the potential energy per unit charge:

$$V_{AB} = \frac{W}{Q} = -\int_{A}^{B} \mathbf{E} \cdot d\mathbf{l}$$

This quantity, denoted by  $V_{AB}$ , is known as the potential difference between points A and B

#### Electric Potential - Important Points

- 1. In determining  $V_{AB}$ , A is the initial point while B is the final point
- 2. If  $V_{AB}$  is negative, there is a loss in potential energy in moving Q from A to B; this implies that the work is being done by the field
- 3. If  $V_{AB}$  is positive, there is a gain in potential energy in the movement; an external agent performs the work
- 4.  $V_{AB}$  is independent of the path taken (to be proved later)
- 5.  $V_{AB}$  is measured in joules per coulomb, commonly referred to as volts (V)

## Electric Potential - Point Charge

As an example, if the **E** field in the figure is due to a point charge Q located at the origin, then:

$$\mathbf{E} = \frac{Q}{4\pi\varepsilon_0 r^2} \mathbf{a}_r$$

Then the equation for potential difference between point A and B becomes:

$$V_{AB} = -\int_{r_A}^{r_B} \frac{Q}{4\pi\varepsilon_0 r^2} \mathbf{a}_r \cdot dr \, \mathbf{a}_r$$
$$= \frac{Q}{4\pi\varepsilon_0} \left[ \frac{1}{r_B} - \frac{1}{r_A} \right]$$

≻Or:

$$V_{AB} = V_B - V_A$$

 $\triangleright$  Here  $V_A$  and  $V_B$  are potentials or absolute potentials at A and B

## Electric Potential - Point Charge

- The potential difference  $V_{AB}$  may be regarded as the potential at B with reference to A
- In problems involving point charges, it is customary to choose infinity as reference where the potential is zero
- Thus, if  $V_A = 0$  as  $r_A \to \infty$ , the potential at any point  $(r_B \to r)$  due to a point charge Q located at the origin is:

$$V = \frac{Q}{4\pi\varepsilon_0 r}$$

If the point charge Q is not located at the origin but at a point whose position vector is  $\mathbf{r}'$ , the potential V(x, y, z) or simply V(r) at r becomes:

$$V(\mathbf{r}) = \frac{Q}{4\pi\varepsilon_{\rm o}|\mathbf{r} - \mathbf{r}'|}$$

## Electric Potential - Point Charge

Note that because **E** points in the radial direction, any contribution from a displacement in the  $\theta$  or  $\emptyset$  direction is wiped out by the dot product, i.e.

$$\mathbf{E}.d\mathbf{l} = E \cos\theta \ dl = E dr$$

- $\succ$  Hence the potential difference  $V_{AB}$  is independent of the path as mentioned earlier (work done depends upon displacement)
- The potential at any point is the potential difference between that point and a chosen point at which the potential is zero
- ➤ By assuming zero potential at infinity, the potential at a distance r from the point charge is the work done per unit charge by an external agent in transferring a test charge from infinity to that point

## Electric Potential - Multiple Charges

- The superposition principle, which we applied to electric fields, applies to potentials
- For *n* point charges  $Q_1, Q_2,...,Q_n$  located at points with position vectors  $\mathbf{r}_1, \mathbf{r}_2,..., \mathbf{r}_n$ , the potential at  $\mathbf{r}$  is:

$$V(\mathbf{r}) = \frac{Q_1}{4\pi\varepsilon_0|\mathbf{r}-\mathbf{r}_1|} + \frac{Q_2}{4\pi\varepsilon_0|\mathbf{r}-\mathbf{r}_2|} + \cdots + \frac{Q_n}{4\pi\varepsilon_0|\mathbf{r}-\mathbf{r}_n|}$$

>Or:

$$V(\mathbf{r}) = \frac{1}{4\pi\varepsilon_0} \sum_{k=1}^{n} \frac{Q_k}{|\mathbf{r} - \mathbf{r}_k|}$$
 (point charges)

#### Electric Potential - Charge Distributions

For continuous charge distributions, we replace  $Q_k$  with charge element  $\rho_L dl$ ,  $\rho_s dS$ , or  $\rho_v dv$  and the summation becomes an integration, so the potential at r becomes:

$$V(\mathbf{r}) = \frac{1}{4\pi\varepsilon_{o}} \int_{L} \frac{\rho_{L}(\mathbf{r}')dl'}{|\mathbf{r} - \mathbf{r}'|} \qquad \text{(line charge)}$$

$$V(\mathbf{r}) = \frac{1}{4\pi\varepsilon_{o}} \int_{S} \frac{\rho_{S}(\mathbf{r}')dS'}{|\mathbf{r} - \mathbf{r}'|} \qquad \text{(surface charge)}$$

$$V(\mathbf{r}) = \frac{1}{4\pi\varepsilon_{o}} \int_{S} \frac{\rho_{\nu}(\mathbf{r}')d\nu'}{|\mathbf{r} - \mathbf{r}'|} \qquad \text{(volume charge)}$$

➤ Where the primed coordinates are used to denote source point location and the unprimed coordinates refer to field point

As discussed, the potential difference between points A and B is independent of the path taken, hence:

$$V_{BA} = -V_{AB}$$

>Therefore:

$$V_{BA} + V_{AB} = \oint \mathbf{E} \cdot d\mathbf{l} = 0$$

≻Or:

$$\oint \mathbf{E} \cdot d\mathbf{l} = 0$$

- This shows that the line integral of **E** along a closed path must be zero
- ➤ Physically, this implies that no net work is done in moving a charge along a closed path in an electrostatic field

>Applying Stokes's theorem, we get:

$$\oint \mathbf{E} \cdot d\mathbf{l} = \int (\mathbf{\nabla} \times \mathbf{E}) \cdot d\mathbf{S} = 0$$

 $\triangleright$  Or:  $\nabla \times \mathbf{E} = 0$ 

- Any vector field that satisfies the above equations is said to be conservative
- >Thus, an electrostatic field is a conservative field
- The above equation is referred to as Maxwell's second equation for static electric fields

From the way we defined potential,  $V = -\int \mathbf{E} \cdot d\mathbf{l}$ , it follows that:

$$dV = -\mathbf{E} \cdot d\mathbf{l} = -E_x dx - E_y dy - E_z dz$$

The differential of a scalar quantity may be written as:

$$dV = \frac{\partial V}{\partial x} dx + \frac{\partial V}{\partial y} dy + \frac{\partial V}{\partial z} dz$$

➤ By comparing we get:

$$E_x = -\frac{\partial V}{\partial x}, \qquad E_y = -\frac{\partial V}{\partial y}, \qquad E_z = -\frac{\partial V}{\partial z}$$

>As:

$$\mathbf{E} = E_{x} \mathbf{a}_{x} + E_{y} \mathbf{a}_{y} + E_{z} \mathbf{a}_{z}$$

➤So we get:

$$\mathbf{E} = -\nabla V$$

- $\triangleright$ Therefore, the electric field intensity is the gradient of V
- The negative sign shows that the direction of **E** is opposite to the direction in which *V* increases or **E** is directed from higher to lower levels of *V*
- Since the curl of the gradient of a scalar function is always zero  $(\nabla \times \nabla V = 0)$ , this implies that **E** must be a gradient of some scalar function
- $\triangleright$ So by using the potential difference V, a vector problem is reduced to a scalar problem

#### Problem-1

- $\triangleright$  An infinite line charge having charge density  $\rho_L$  C/m<sup>2</sup> is placed along the z-axis. Find the work done in moving a point charge Q:
- a) From point b to a on the y-axis, where a is closer to the line charge compared to b.
- b) In a circle around the line charge.