

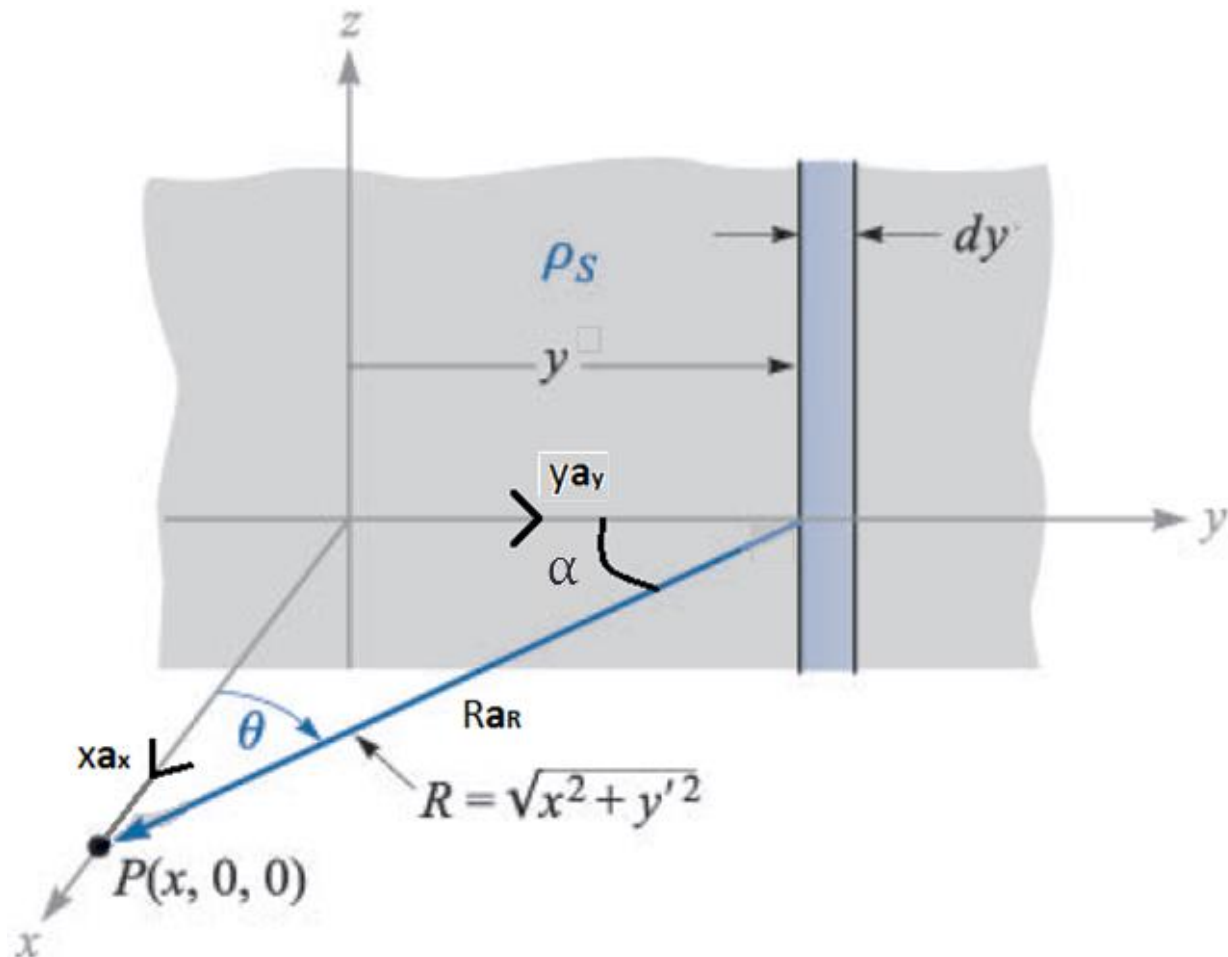
ELECTRIC FIELDS DUE TO CONTINUOUS CHARGE --- DISTRIBUTIONS - SURFACE CHARGE

Introduction

- Another basic charge configuration is the infinite sheet of charge having a uniform density of $\rho_s \text{ C/m}^2$
- Such a charge distribution may often be used to approximate the charge found on the plates of a parallel-plate capacitor
- For ease of derivation, the sheet of charge is considered to be **infinite**
- This is a good approximation since the distances involved in the measurement of fields are generally small compared to the dimensions of the sheet of charge

Sheet of Charge

- Consider an infinite sheet of charge with uniform charge density ρ_s in the yz -plane
- We consider a small portion of the sheet having height dz and width dy
- Let the point of observation be at $P(x, 0, 0)$



Sheet of Charge

- The charge associated with an elemental area dS is:

$$dQ = \rho_s dS \quad \text{where} \quad dS = dydz$$

- The sheet of charge may be assumed to consist of **infinite line charges** that extend from $-\infty$ to $+\infty$ along the z-axis

- We have for an infinite line charge:

$$\mathbf{E} = \frac{\rho_L \mathbf{a}_R}{2\pi\epsilon_0 R}$$

- Therefore, the differential intensity due to the line charge is:

$$d\mathbf{E} = \frac{\rho_s dy \mathbf{a}_R}{2\pi\epsilon_0 R}; \quad \text{where } \rho_L = \rho_s dy$$

Sheet of Charge

➤ It may be observed from the figure that vector \mathbf{R} is perpendicular to the line charge

➤ Therefore: $y\mathbf{a}_y + R\mathbf{a}_R = x\mathbf{a}_x$ or $R\mathbf{a}_R = x\mathbf{a}_x - y\mathbf{a}_y$
and

$$\mathbf{a}_R = \frac{x\mathbf{a}_x - y\mathbf{a}_y}{\sqrt{x^2 + y^2}}$$

➤ Substituting values in the equation for $d\mathbf{E}$, we get:

$$d\mathbf{E} = \frac{\rho_S dy (x\mathbf{a}_x - y\mathbf{a}_y)}{2\pi\epsilon_o (x^2 + y^2)}$$

Sheet of Charge

- The intensity due to all the line charges will be their vector sum
- When we take the vector sum, the **y-component will cancel** due to symmetry
- For every line charge along the positive y-axis, there is a line charge at the same distance along the negative y-axis
- So in the summation, the y-axis is ignored and limits are taken along z-axis from $-\infty$ to $+\infty$

Sheet of Charge

➤ Therefore, we have the total electric field as:

$$\mathbf{E} = \int_{-\infty}^{\infty} \frac{\rho_S dy x \mathbf{a}_x}{2\pi\epsilon_o(x^2 + y^2)}$$

➤ We use change of variables to solve the above equation

$$\tan\alpha = \frac{x}{y}$$

$$\Rightarrow y = x \cot\alpha \quad \text{and} \quad dy = -x \operatorname{cosec}^2\alpha d\alpha$$

$$\text{when } y = -\infty \quad \Rightarrow \cot\alpha = -\infty \quad \text{so } \alpha = \pi$$

$$\text{when } y = \infty \quad \Rightarrow \cot\alpha = \infty, \quad \text{so } \alpha = 0$$

Sheet of Charge

➤ So, we get:

$$\mathbf{E} = - \int_{\pi}^0 \frac{\rho_S x^2 \operatorname{cosec}^2 \alpha d\alpha \mathbf{a}_x}{2\pi\epsilon_o(1 + \cot^2 \alpha)x^2}$$

Or

$$\mathbf{E} = \int_0^{\pi} \frac{\rho_S d\alpha \mathbf{a}_x}{2\pi\epsilon_o}$$

$$\mathbf{E} = \frac{\rho_S \mathbf{a}_x}{2\epsilon_o}$$

- We see that for an infinite sheet of charge, the intensity does not depend upon any coordinate
- So if we take P anywhere, the intensity will remain the same but the **direction** will change

Sheet of Charge

- If the point of observation is located at the back of the sheet, then we have:

$$\mathbf{E} = -\frac{\rho_S \mathbf{a}_x}{2\epsilon_0}$$

Or generally:

$$\mathbf{E} = \frac{\rho_S \mathbf{a}_n}{2\epsilon_0}$$

- Where \mathbf{a}_n is a vector normal to the sheet

Problem-1

- A circular disk of radius a is uniformly charged with ρ_s C/m². If the disk lies on the $z = 0$ plane with its axis along the z -axis,
- (a) Show that at point $(0, 0, h)$

$$\mathbf{E} = \frac{\rho_s}{2\epsilon_0} \left\{ 1 - \frac{h}{[h^2 + a^2]^{1/2}} \right\} \mathbf{a}_z$$

- (b) From this, derive the \mathbf{E} field due to an infinite sheet of charge on the $z = 0$ plane