# Quantum Circuits and Quantum Machine Learning

Mar 6th, Spring 2023

Bozhang Chen, Linzhe Teng, Yun Wang, Shuzhen Zhang Ivor Yidong Chen (Team Leader) Marius Junge (Faculty mentor)

University of Illinois at Urbana-Champaign





## Linear Algebra

 Hilbert Space is a (finite-dimensional) complex vector space H equipped with an inner product:

$$\langle v|(|w\rangle) \equiv \langle v|w\rangle \equiv (|v\rangle, |w\rangle)$$

- A matrix U is **Unitary** if  $UU^{\dagger} = U^{\dagger}U = I$
- A matrix X is **Self-adjoint** if  $X^{\dagger} = X$
- A matrix Y is symmetric positive semi-definite if it is Self-adjoint and its eigenvalues are positive.
- Trace of an n × n matrix A:  $Tr(A) = \sum_{i=1}^{n} a_{ii}$

#### **Tensor Products**

- Suppose V and W are Hilbert spaces of dimension m and n respectively. Then V  $\otimes$  W has dimension m  $\times$  n.
- Calculation rule of tensor product:

$$A \otimes B \equiv \left[ \begin{array}{cccc} A_{11}B & A_{12}B & \dots & A_{1n}B \\ A_{21}B & A_{22}B & \dots & A_{2n}B \\ \vdots & \vdots & \vdots & \vdots \\ A_{m1}B & A_{m2}B & \dots & A_{mn}B \end{array} \right] \right\} mp \,.$$

#### Basic Quantum mechanics I

- To describe the state space(Hilbert Space) in the quantum physic system, we use the unit vector in the state space.
- The unit vector in the simplest state space  $\mathbb{C}^2$  is called a **qubit**:

$$|\psi\rangle = a|0\rangle + b|1\rangle, \ \langle\psi|\psi\rangle = 1$$

 $|0\rangle$  and  $|1\rangle$  are the orthonormal basis for  $\mathbb{C}^2,~a,b$  are complex numbers. Hence, we could represent every state by this equation.

#### Basic Quantum mechanics II

The evolution of a closed quantum system is described by the Schrödinger Equation:

$$i\hbar \frac{\mathsf{d}|\psi\rangle}{\mathsf{d}t} = \mathsf{H}|\psi\rangle$$

 $\hbar$  is the *Planck's constant*, and *H* is a Hermitian (self-adjoint) operator.

### Quantum Density& Measurement

• The density operator for the system is defined as:

$$\rho = \sum_{i} p_{i} |\psi_{i}\rangle\langle\psi_{i}|$$

- $Tr(\rho) = 1$  and  $\rho$  is a positive operator.
- Generalized measurement(**POVM**): **POVM** is a set of positive semi-definite matrices  $\{T_i\}$ , i is the outcome that may occur in the experiment, and it satisfies  $\sum_{i=1}^{n} T_i = \mathbb{I}$ .
- The probability of the measurement result i, P(i) is:

$$P(i) = \langle \psi | T_i | \psi \rangle$$

#### Quantum Gates and Quantum Circuit I

- Quantum Logic Gates could operate on qubit.
- Quantum Gates are unitary matrices and reversible.
- A gate which acts on n qubits:  $|\psi_1\rangle \otimes |\psi_2\rangle \otimes \cdots \otimes |\psi_n\rangle$  is a  $2^n \times 2^n$  matrix.
- Pauli Gates:

$$old X = egin{bmatrix} 0 & 1 \\ 1 & 0 \end{bmatrix} \quad old Y = egin{bmatrix} 0 & -i \\ i & 0 \end{bmatrix} \quad old Z = egin{bmatrix} 1 & 0 \\ 0 & -1 \end{bmatrix}$$

•

Control X : 
$$\begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \quad \text{SWAP} : \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \end{bmatrix}$$

#### Quantum Gates and Quantum Circuit II

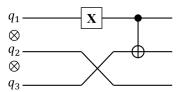
 The rotation operators are generated by the exponentiation of Pauli matrices:

$$\exp(iP\theta) = \cos(\theta)I + i\sin(\theta)P$$

• A rotation gate along z-axis by  $\theta$  radians is:

$$\begin{bmatrix} 1 & 0 \\ 0 & e^{i\theta} \end{bmatrix}$$

• An example quantum circuit composed by basic gates:



## Representation Theory

- Definition and Tensor Product Representation
- Decomposition Theory
- Character
- $H \otimes H \otimes H$  and  $S_3$

### Definition and Tensor Product Representation

 Definition: Let V be a finite-dimensional complex vector space and G be a group. We say V is a representation of G, if ρ : G → GL(V) is homomorphism

- Tensor Product Representation:
- Let V, W be a representation of G
- V ⊗ W is a new representation of G

# **Decomposition Theory**

- Let V be a representation of G
- $V = U_1 \oplus U_2 \oplus ... \oplus U_n$
- **Note:** Every *U* is a irreducible representation
- A representation is called irreducible if it has no non-trivial sub-representation
- Abelian group only has one dimensional representation

#### Character

- **Definition:** The character of a representation  $\rho$  is a function  $\chi_{\rho}$  on G defined by  $\chi_{\rho}(g) = Tr(\rho(g))$ .
- Another property: The character is the same for elements related by conjugation.  $\chi_{\rho}(g) = \chi_{\rho}(h^{-1}gh)$

	e	$(1\ 2)$	$(1\ 2\ 3)$
$\chi_{\rm triv}$	1	1	1
$\chi_{ m sign}$	1	-1	1
$\chi V$	2	0	-1

Table 1. Character table of  $S_3$ .

# $H \otimes H \otimes H$ and $S_3$

- Flip:  $f: V \otimes W \rightarrow W \otimes V$
- $S_3 = \{e, (12), (13), (23), (123), (132)\}$
- H ⊗ H ⊗ H is the irreducible representation of S<sub>3</sub>
- suppose  $h_1 \otimes h_2 \otimes h_3 \in H \otimes H \otimes H$ , Then

$$\forall \sigma \in S_3 : \sigma(h_1 \otimes h_2 \otimes h_3) = h_{\sigma(1)} \otimes h_{\sigma(2)} \otimes h_{\sigma(3)}$$

• **EX**:  $\sigma_{(132)} \in S_3$ 

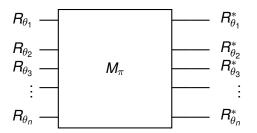
$$\sigma_{(132)}(h_1\otimes h_2\otimes h_3)=h_2\otimes h_3\otimes h_1$$

• It just a flip of  $h_1 \otimes (h_2 \otimes h_3) \rightarrow (h_2 \otimes h_3) \otimes h_1$ 



### Project goal - Permutation Matrix $M_{\pi}$

• Our goal is to find the permutation matrix  $M_{\pi}$  with  $\pi \in S_n$  by a black-box quantum gate.



- Known inputs:  $(R_{\theta_1}, R_{\theta_2}, ..., R_{\theta_n})$ , each  $R_{\theta_i}$  represents a rotation.
- Known outputs after measurement:  $(R_{\theta_1}^*, R_{\theta_2}^*, ..., R_{\theta_n}^*)$

