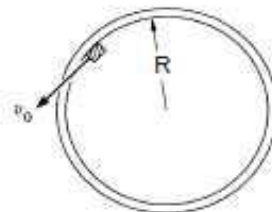
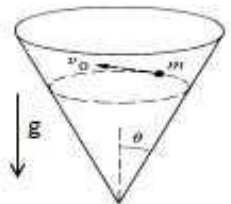
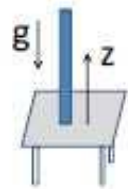
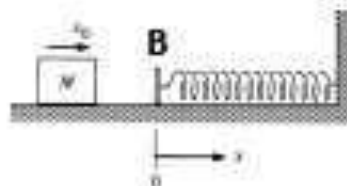

Marks will be deducted for illegible presentation.

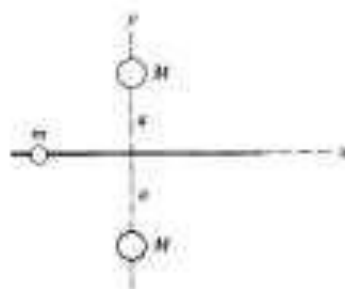
1. (a) A rod of uniform cross-sectional area is maintained vertically on a table. The linear density of the rod varies as $\rho(z) = \rho_0 z/l$, where l is the length of the rod and ρ_0 is a constant. Gravity acts along negative z -direction. Find the potential energy of the rod with respect to the table, in terms of ρ_0 , l , g . [3]
- (b) A particle of mass m slides without friction on the inside of a cone (see figure). The axis of the cone is vertical, and gravity acts downward. The apex half-angle of the cone is θ . The path of the particle happens to be a circle in a horizontal plane. The speed of the particle is v_0 . Find the radius of the circular path in terms of v_0 , g and θ . [3]
2. A block of mass m slides on a friction less table. It is constrained to move along the inner surface of a ring of radius R , which is fixed on the table. At $t = 0$ the block has a velocity v_0 as shown in the figure. The coefficient of friction between the block and the ring is μ .
- i. Find the angular frequency, $\omega(t)$, of the block as a function of time. [4]
- ii. Find the time taken by the block to make one full revolution of the ring (starting from $t = 0$). [2]



3. A block of mass M slides along a horizontal table. From $x = 0$ the mass starts experiencing forces due to a spring and the frictional force of the table (see figure). The spring has a spring-constant k . The coefficient of friction in this case is a function of x and is given by $\mu = bx$ (for $x > 0$), where b is a constant. The spring and board B are massless. At $x = 0$ the block has a speed v_0 . Find the loss of mechanical energy of the block during its motion from $x = 0$ to the point where it comes momentarily to rest, for the first time. [5]



4. (a) A block moves along the x -axis on a horizontal frictionless table. The table is covered with dust of linear density ρ . As the block moves it gathers all the dust in the path and gets heavier. At $t = 0$, the block's velocity is v_0 and its mass is M_0 .
- Find the velocity of the block as a function of distance, x . [3]
 - What horizontal force, $F(x)$, must be applied on the block, if the block has to move with constant velocity v_0 . [2]
- (b) A rubber ball of mass m dropped from a height H under gravity hits the ground and rebounds to a height h . Find the impulse it received from the ground. [2]
5. A bead of mass m slides without friction on a smooth rod along the X -axis. Two particles each of mass M are located at $(0, +a)$ and $(0, -a)$ as shown in the figure. The bead is under the gravitational attraction of the masses, M .
- Starting from the expression of the total potential energy obtain the force on the bead along the x -direction. [2]
 - Find the frequency of small oscillation of the bead about the equilibrium. [4]



a) Given $\rho(z) = \frac{\rho_0 z}{l}$ (linear density)



mass of the small element

$$dm = \rho(z) dz = \frac{\rho_0}{l} z dz$$

The Potential Energy of the element,

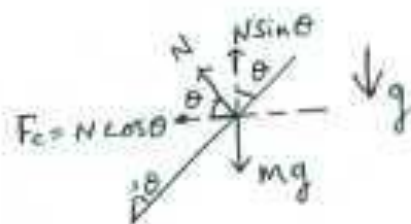
$$dU = dm g z = \frac{\rho_0}{l} g z^2 dz$$

$$\text{Total P.E.} = \int_{z=0}^{z=l} dU = \frac{\rho_0 g}{l} \int_0^l z^2 dz$$

$$\text{i.e. } U = \frac{\rho_0 g l^2}{3}$$

b)

From the force diagram with components resolved along horizontal & vertical directions,



Force Diagram
(Not required)

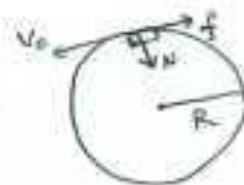
$$N \sin \theta = mg \quad \text{--- (1)}$$

$$N \cos \theta = \frac{m v_0^2}{r} \quad \text{--- (2)}$$

$$\frac{(1)}{(2)} \Rightarrow \tan \theta = \frac{mg}{m v_0^2 / r} = \frac{r g}{v_0^2}$$

$$\text{So } r = \frac{v_0^2 \tan \theta}{g}$$

(X.2)
i) Given at $t=0$; $v=v_0$
or $\omega = v_0/R$



Total force on the particle

$$\vec{F} = N(\hat{r}) + f(\hat{\theta}) \quad \begin{matrix} (f - \text{friction} \\ N - \text{normal} \\ \text{reaction}) \end{matrix}$$

The general form of acceleration

$$\vec{a} = (\ddot{r} - r\dot{\theta}^2)\hat{r} + (2\dot{r}\dot{\theta} + r\ddot{\theta})\hat{\theta}$$

Since $r=R$; $\dot{r} = \ddot{r} = 0$.

$$\therefore \vec{a} = -R\dot{\theta}^2\hat{r} + R\ddot{\theta}\hat{\theta} \quad \text{--- (1)}$$

Comparing forces and accelerations
along \hat{r} & $\hat{\theta}$ directions

$$\left. \begin{aligned} N &= mR\dot{\theta}^2 \quad (\text{along } \hat{r}) \quad \text{--- (2)} \\ f &= -mR\ddot{\theta} \quad (\text{along } \hat{\theta}) \quad \text{--- (3)} \end{aligned} \right\}$$

Since $f = \mu N$, $\mu mR\dot{\theta}^2 = -mR\ddot{\theta}$

$$\dot{\theta} = \omega, \quad \mu \omega^2 = -\frac{d\omega}{dt}$$

$$\text{or} \quad -\frac{d\omega}{\omega^2} = \mu dt$$

Integrating $\frac{1}{\omega} = \mu t + C$

Using Initial condition, $t=0$, $\omega = v_0/R$

$$\frac{1}{\omega} = \mu t + R/v_0$$

$$\text{Or } \omega(t) = \frac{v_0}{(R + \mu v_0 t)}$$

Q.2 ii)

$$\omega(t) = \frac{d\theta}{dt} = \frac{v_0}{(R + \mu v_0 t)}$$

$$\text{or } d\theta = \frac{v_0 dt}{(R + \mu v_0 t)}$$

Integrating between $t=0$ to $t=T$
where T is the time taken for the 1st full revolution.

$$\int_0^T d\theta = v_0 \int_0^T \frac{dt}{(R + \mu v_0 t)}$$

$$\text{i} \quad 2\pi = v_0 \frac{1}{\mu v_0} \ln(R + \mu v_0 t) \Big|_0^T$$

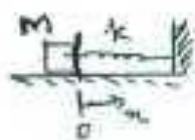
$$\text{or } 2\pi \mu = \ln\left(\frac{R + \mu v_0 T}{R}\right)$$

$$\text{or } 1 + \frac{\mu v_0 T}{R} = e^{2\pi \mu}$$

$$\text{i} \quad T = \frac{R(e^{2\pi \mu} - 1)}{\mu v_0}$$

Q.3. Using Work energy theorem

$$E_i - E_f = -W^{nc} \quad \text{--- (1)}$$



Where E_i - initial mechanical energy (total)

E_f - final mechanical energy (total)

W^{nc} - Work done by non-conservative forces (here friction, $f = \mu Mg$)

$$E_i = \frac{1}{2} M v_0^2 ; E_f = \frac{1}{2} k x_0^2 \quad \text{--- (2)}$$

x_0 - unknown, the distance at which the block comes to rest.

$$W^{nc} = - \int_0^{x_0} f dx = - \int_0^{x_0} \mu Mg dx$$

Since $\mu = b x$

$$W^{nc} = - Mgb \int_0^{x_0} x dx = - \frac{Mgb x_0^2}{2} \quad \text{--- (3)}$$

Using (1) & (2), $Mgb \frac{x_0^2}{2} = \frac{1}{2} M v_0^2 - \frac{1}{2} k x_0^2$

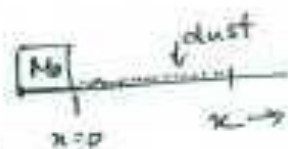
$$\therefore x_0^2 = \frac{M v_0^2}{(k + Mgb)}$$

Mechanical energy lost is $-W^{nc}$, say ΔE .

Using (3), $\Delta E = \frac{Mgb}{2} \frac{M v_0^2}{(k + Mgb)}$

Q.4.

a) i) Given at $t=0$, $M = M_0$
and $V = V_0$



Since there is no external forces along the x -direction the total linear momentum is conserved.

That is, $M_0 V_0 = M(x) V(x)$

$$V(x) = \frac{M_0 V_0}{M(x)}$$

The mass of the block at $M(x) = M_0 + \rho \cdot x$
 $\rho \cdot x$ - being the mass accumulated over the distance x .

Thus $V(x) = \frac{M_0 V_0}{(M_0 + \rho \cdot x)}$

a) i) 2nd solution:

$$F_x = \frac{dP_x}{dt} = V \frac{dM}{dt} + M \frac{dV}{dt} = 0$$

or $\frac{dM}{M} + \frac{dV}{V} = 0$

Integrating $\ln M = -\ln V + \ln C$ some constant

or $M = C/V$

at $x=0$; $M = M_0$ and $V = V_0$

Hence $C = M_0 V_0$

that $M(x) = \frac{M_0 V_0}{V(x)}$

or $V(x) = \frac{M_0 V_0}{(M_0 + \rho \cdot x)}$

(24 a) ii)

$$F(x) = \frac{d(M(x)V(x))}{dt}$$

$$= M(x) \frac{dV(x)}{dt} + V(x) \frac{dM(x)}{dt}$$

but $V(x) = V_0$ In this case

$$\text{So } F(x) = V_0 \frac{dM(x)}{dt}$$

$$= V_0 \frac{dM(x)}{dt} \frac{dx}{dt} \Rightarrow V_0$$

$$F(x) = V_0^2 \rho$$

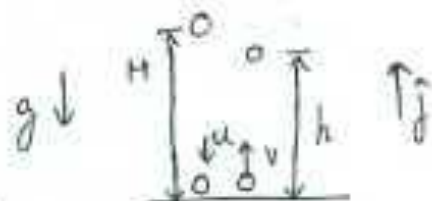
$$\text{ii } \underline{F(x) = \rho V_0^2 \text{ (a const.)}}$$

b)

The velocity with the ball hits the ground

$$u = \sqrt{2gH}$$

— (1)



The velocity with which it rebounds

$$v = \sqrt{2gh} \quad \text{— (2)}$$

Impulse it received from the ground

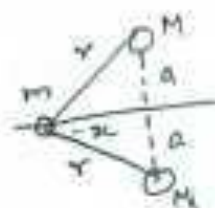
$$= \int_0^T F dt = p_f - p_i \quad \text{— (} p_i \text{ + } p_f \text{ are the initial + final momenta)}$$

$$= mv\hat{j} - m(-u\hat{j}) = m(v+u)\hat{j}$$

$$\underline{\underline{\text{Impulse} = m(\sqrt{2gH} + \sqrt{2gh})\hat{j}}}$$

Q.5

- i) The total potential energy of the mass "m" due to the two masses "M" is



$$U = +2 \frac{GMm}{r} \quad \text{where } r = \sqrt{a^2 + x^2}$$

$$F_x = -\frac{\partial U}{\partial x} = 2GMm \frac{d(1/r)}{dx}$$

$$= -\frac{2GMm}{r^2} \frac{dr}{dx}$$

$$\therefore F_x = -\frac{2GMm \cdot x}{r^3} \quad \text{--- (1)}$$

$$\begin{aligned} r^2 &= x^2 + a^2 \\ 2r \frac{dr}{dx} &= 2x \\ \frac{dr}{dx} &= x/r \end{aligned}$$

- ii) At equilibrium $F_x = 0$

From (1) equilibrium point is, $x_0 = 0$.
The effective force constant for small oscillation of "m" about $x = 0$ is given by

$$k = \left. \frac{d^2 U}{dx^2} \right|_{x=0}$$

$$\begin{aligned} \therefore k &= 2GMm \frac{d}{dx} \left(\frac{x}{r^3} \right) = 2GMm \left(\frac{1}{r^3} + x \cdot \frac{-3}{r^4} \cdot \frac{dr}{dx} \right) \bigg|_{x=0} \\ &= 2GMm \left(\frac{1}{r^3} - \frac{3x^2}{r^5} \right) \bigg|_{x=0} \end{aligned}$$

$$\therefore k = \frac{2GMm}{a^3} \quad \left(\text{Since at } x=0; r=a \right)$$

$$\text{Frequency } \omega = \sqrt{\frac{k}{m}} = \sqrt{\frac{2GM}{a^3}}$$