

1. An electron is injected with a velocity  $\vec{u}_0 = \hat{y} u_0$  <sup>at origin (0,0,0)</sup> into a region where both an electric field  $\vec{E}$  and a magnetic field  $\vec{B}$  exist. Describe the motion of the electron if  $\vec{E} = \hat{z} E_0$ , and  $\vec{B} = \hat{x} B_0$ . Discuss the effect of the relative magnitudes of  $|\vec{E}|$ , and  $|\vec{B}|$ , on the electron paths.

Sol:  $q = -|e|$        $\left. \vec{V} \right|_{t=0} = (0, u_0, 0)$        $\vec{E} = (0, 0, E_0)$   
 $\vec{B} = (B_0, 0, 0)$

Lorentz force equation  $F = m \frac{d\vec{V}}{dt} = q (\vec{E} + \vec{V} \times \vec{B})$

$$m \frac{dv_z}{dt} = 0 \quad \Rightarrow \quad v_z(t) = 0$$

$$\frac{dv_y}{dt} = -\frac{|e|}{m} B_0 v_z \quad \Rightarrow \quad v_y(t) = (u_0 - E_0/B_0) \cos \omega_c t + E_0/B_0$$

$$\frac{dv_z}{dt} = -\frac{|e|}{m} (E_0 - B_0 v_y) \quad \Rightarrow \quad v_z(t) = \left( \frac{E_0}{B_0} - u_0 \right) \sin \omega_c t$$

Cyclotron frequency  $= \omega_c = \frac{|e|}{m} B_0$

The electron is injected at the origin ( $x(0)=0$   $y(0)=0$   $z(0)=0$ )

$$x(t) = 0 \quad \text{Since } v_x(t) = 0$$

$$y(t) = \frac{1}{\omega_c} (u_0 - E_0/B_0) \sin \omega_c t + \frac{E_0 t}{B_0}$$

$$z(t) = -\frac{1}{\omega_c} (u_0 - E_0/B_0) (1 - \cos \omega_c t)$$

Equation of motion

$$\left( y(t) - \frac{E_0 t}{B_0} \right)^2 + \left( z(t) + \frac{(u_0 - E_0/B_0)}{\omega_c} \right)^2 = \frac{(u_0 - E_0/B_0)^2}{\omega_c^2}$$

cycloid motion

if  $E_0/B_0 = u_0$

$$v_z = 0 \quad ; \quad v_y = u_0 \quad \& \quad v_x = 0$$

$$y(t) = u_0 t \quad \text{with } x(t) = 0 \quad \& \quad z(t) = 0$$

is line motion

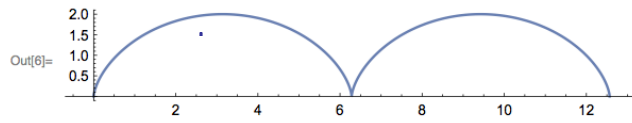
$$\Rightarrow a = \frac{E_0}{u_0 B_0}$$

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In[1]:= y = (1 - a) Sin[t] + t;  
z = - (1 - a) (1 - Cos[t]);
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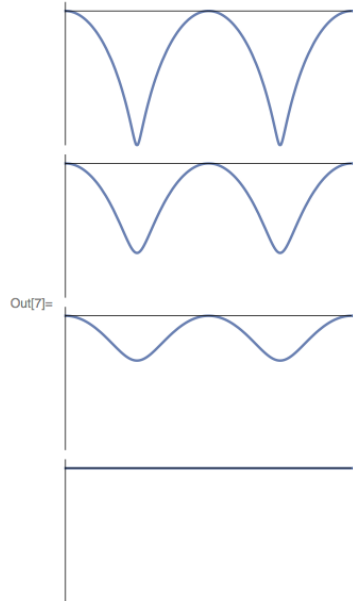
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In[3]:= a = 0.5;  
ParametricPlot[{y, z}, {t, 0, 10 π}]
```



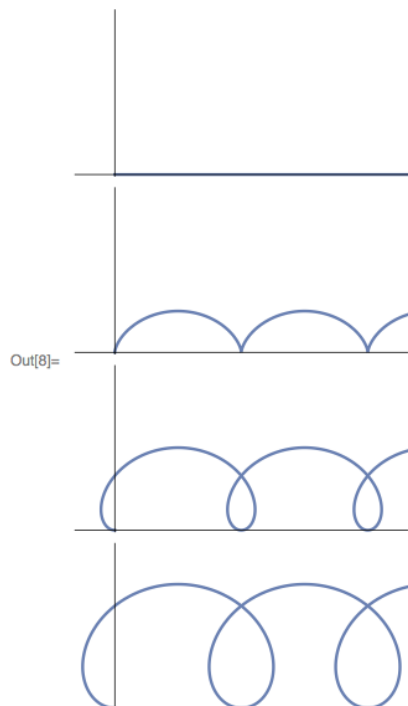
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In[5]:= a = 2;  
ParametricPlot[{y, z}, {t, 0, 4 π}]
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In[7]:= g1 = Grid[Table[  
  {ParametricPlot[{y, z}, {t, 0, 10 π}, Ticks -> None, AspectRatio -> 1 / 2,  
    PlotRange -> {{0, 4 π}, {-1.5, 0.1}}}], {a, 0.25, 1, 0.25}]]
```



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In[8]:= g2 = Grid[Table[  
  {ParametricPlot[{y, z}, {t, 0, 10 π}, Ticks -> None,  
    AspectRatio -> 1 / 2, PlotRange -> {{-2, 4 π + 2}, {-0.1, 8}}}], {a, 1, 4, 1}]]
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2. The surface current density in the  $xy$  plane is given to be

$$\vec{K}(x, y) = K_0 \left[ \frac{x^2}{a^2} \hat{i} + \frac{y^2}{b^2} \hat{j} \right]$$

Evaluate  $\nabla \cdot \vec{K}$  and hence calculate the time rate of change of charge contained in the first quadrant of a circle of unit radius centred on the origin. Show explicitly that for this quadrant, the continuity equation is satisfied in its integral form.

Sol:  $\nabla \cdot \vec{K} = -\frac{\partial \rho}{\partial t} \Rightarrow \frac{\partial \rho}{\partial t} = -2K_0 \left( \frac{x}{a^2} + \frac{y}{b^2} \right)$

The charge  $q$  is contained in 1<sup>st</sup> quadrant of a circle of unit radius.

$$\frac{dq}{dt} = \int_{r=0}^1 \int_0^{\pi/2} \frac{\partial \rho}{\partial t} r dr d\theta$$

$$= -\frac{2K_0}{3} \left( \frac{1}{a^2} + \frac{1}{b^2} \right)$$

Let us evaluate the total charge  $q$  flowing out of the boundaries of the 1<sup>st</sup> quadrant of unit circle.

Along  $OP$ ,  $\vec{K}$  is  $\hat{i}$  direction and along  $OQ$ ,  $\vec{K}$  is along  $\hat{j}$  direction. No charge is flowing across  $OP$  and  $OQ$ .

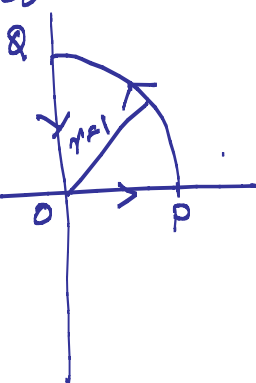
$PQ$  is the part where charge leaves from circle.

$$d\vec{l} = r \hat{\theta} d\theta \Big|_{r=1} = \hat{\theta} d\theta$$

$$\vec{K} = K_0 \left[ \frac{r^2 \cos^2 \theta}{a^2} (\hat{r} \cos \theta - \hat{\theta} \sin \theta) + \frac{r^2 \sin^2 \theta}{b^2} (\hat{r} \sin \theta + \hat{\theta} \cos \theta) \right]$$

$$\vec{K} \times d\vec{l} = K_0 \left[ \frac{r^2 \cos^2 \theta}{a^2} (\hat{r} \cos \theta - \hat{\theta} \sin \theta) + \frac{r^2 \sin^2 \theta}{b^2} (\hat{r} \sin \theta + \hat{\theta} \cos \theta) \right] \times (\hat{r} dr + r \hat{\theta} d\theta)$$

$$|\vec{K} \times d\vec{l}| = K_0 \left( \frac{\cos^3 \theta}{a^2} + \frac{\sin^3 \theta}{b^2} \right) d\theta \int_0^{\pi/2} |\vec{K} \times d\vec{l}| = \frac{2K_0}{3} \left( \frac{1}{a^2} + \frac{1}{b^2} \right)$$



3. A uniformly charged thin disk of charge density  $\sigma$  radius  $R$  and thickness  $t < R$  rotates with an angular velocity  $\omega$  about the  $z$  axis of symmetry as shown in Fig. Find out the magnetic field  $B$  at distance  $z$  in the vertical axis.

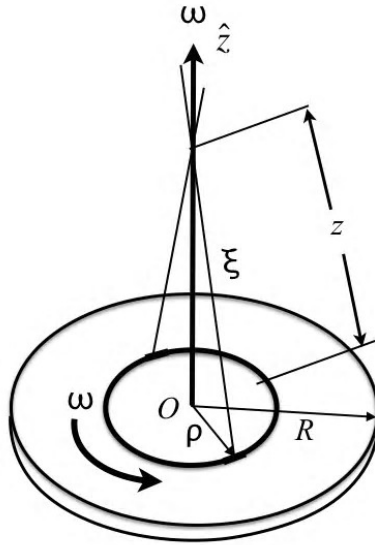


Figure 1: Problem 3

Sol: The current density in the disk is given by  $\rho \vec{v} = \sigma \omega \rho \hat{\phi}$  where  $\rho$  is the distance from the axis of the disk.

A ring-shaped portion of the disk of radial thickness  $dp$  thus constitutes a current ring  $dI = \sigma \omega \rho dp$

The  $\vec{B}$  field on the  $z$ -axis due to this current ring is

$$d\vec{B} = \frac{\mu_0}{2} \frac{dI \rho^2}{\xi^3} \hat{z}$$

Note that  $\vec{B}$  field on axis of circular loop is

$$\vec{B} = B_z \hat{z} = \frac{\mu_0 I R^2}{2 \xi^3} \hat{z} \quad R = \text{circular loop radius.}$$

$$\Rightarrow d\vec{B} = \frac{\mu_0}{2} \frac{\sigma \omega \rho^3 dp}{\xi^3} \hat{z} \quad \xi^2 = \rho^2 + z^2$$

$$\vec{B} = \int_0^R \left( \frac{\mu_0 \omega \sigma}{2} \right) \frac{\rho^3 dp}{(\rho^2 + z^2)^{3/2}} \hat{z} = \hat{z} \frac{\mu_0 \omega \sigma}{2} \left[ \frac{R^2 + 2z^2}{\sqrt{R^2 + z^2}} - 2|z| \right]$$

4. A thin conducting wire is bent into the shape of a regular polygon of  $N$  sides. A current  $I$  flows in the wire. Show that the magnetic flux density at the center is

$$\vec{B} = \frac{\mu_0 N I}{2\pi b} \tan \frac{\pi}{N} \hat{a}_n$$

where  $b$  is the radius of the circle circumscribing the polygon and  $\hat{a}_n$  is a unit vector normal to the plane of the polygon. Show also that as  $N$  becomes very large this result reduces to

$$\vec{B} = \frac{\mu_0 I b^2}{2(z^2 + b^2)^{3/2}} \hat{a}_n \quad \text{from Griffiths equation 5.35}$$

with  $z = 0$ .

Sol:

$$\tan \alpha = l/a$$



polygon with  $N$  side  $\alpha = \pi/N$

$$\begin{aligned} \vec{B} &= \frac{\mu_0 N I}{2\pi a} \sin \alpha \hat{e}_n \\ &= \frac{\mu_0 N I}{2\pi b} \tan \alpha \hat{e}_n \end{aligned}$$

$$\vec{B} = \lim_{N \rightarrow \infty} \frac{\mu_0 N I}{2\pi b} \tan(\pi/N) \hat{e}_n$$

$$= \frac{\mu_0 N I}{2\pi b} \frac{\pi}{N} \hat{e}_n = \frac{\mu_0 I}{2b} \hat{e}_n$$

$$\vec{B} = \left. \frac{\mu_0 I b^2}{2(b^2 + z^2)^{3/2}} \right|_{z=0} \hat{e}_n = \frac{\mu_0 I}{2b} \hat{e}_n$$

$\hat{e}_n$  is  $\perp$  to the plane of polygon.

5. A curved wire kept in  $yz$  plane, carrying a current  $I_1$  subtends an angle  $2\theta_0$  at the location of another long straight wire carrying a current  $I_2$  as shown in the figure. Find the force exerted by the straight wire on the curved wire.

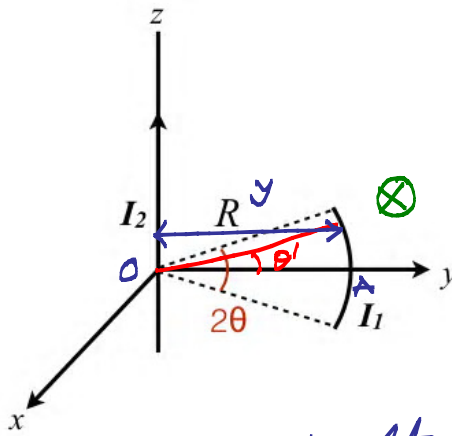


Figure 2:

Sol: The field due to the straight wire at a distance  $y$  is  $\vec{B} = -\frac{\mu_0 I_2}{2\pi y} \hat{i} = -\frac{\mu_0 I_2}{2\pi R \cos \theta'} \hat{i}$

Force on curved wire

$$\vec{F} = \int I_1 d\vec{l} \times \vec{B} = -\frac{\mu_0 I_1 I_2}{2\pi R} \int \frac{d\vec{l} \times \hat{i}}{\cos \theta'}$$

Due to symmetry about  $OA$ , only the  $z$ -component of  $d\vec{l}$  will contribute to the force.

$$d\vec{l} = R d\theta \hat{\theta} = R \cos \theta' d\theta' \hat{k}$$

$$\vec{F} = -\frac{\mu_0 I_1 I_2}{2\pi} \int_{-\theta}^{\theta} d\theta' \hat{j} = -\frac{\mu_0 I_1 I_2}{2\pi} 2\theta \hat{j} = -\frac{\mu_0 I_1 I_2 \theta}{\pi} \hat{j}$$

Note that  $y = R \cos \theta'$   $z = R \sin \theta'$

$$d\vec{l} = -R \sin \theta' d\theta' \hat{y} + R \cos \theta' d\theta' \hat{z} = R d\theta' (-\sin \theta' \hat{y} + \cos \theta' \hat{z})$$

$\hat{\theta} = -\sin \theta' \hat{y} + \cos \theta' \hat{z}$  is not  $\hat{\theta}$  in cylindrical co-ordinate system