Homework Assignment 5
Particles Want To Be Free

Practice Problems

These problems are graded on effort only.

Griffiths: P2.13, P2.18, P2.21
Hints on P2.21 are at the end of the problem set.

Additional Problems

Problem 1: Because it is so widely applicable, the Fourier series is written in many different ways, and this can be quite confusing. But you never need to be told more than the Fourier series, as you can deduce the rest for yourself. Suppose we write the Fourier series for a function $f(t)$ with period T in this way:

$$f(t) = \sum_m c_m \exp(i\omega_m t)$$

where:

$$\omega_m \equiv \frac{2\pi m}{T}$$

This form obscures the role of orthonormal basis functions in the series, but it is neater to write.

(A) Calculate:

$$\int_{-T/2}^{+T/2} \exp(-i\omega_m t) \, \exp(i\omega_m t) \, dt$$

(B) Without calculating integrals you already know are zero, write the general solution to this integral:

$$\int_{-T/2}^{+T/2} \exp(-i\omega_n t) \, \exp(i\omega_m t) \, dt$$

(C) Use Fourier's trick to determine the formulas for the coefficients c_m. That is, calculate:

$$\int_{-T/2}^{+T/2} \exp(-i\omega_n t) f(t) dt$$

and use the results to write a formula for calculating c_n.

Problem 2: At $t = 0$ a free particle has the initial wave function:

$$\Psi(x, 0) = \begin{cases} A & -a \leq x \leq a \\ 0 & \text{otherwise} \end{cases}$$

(A) Find the value of A that normalizes $\Psi(x, 0)$ appropriately.

(B) Find $\widetilde{\Psi}(k, 0)$ the Fourier transform of $\Psi(x, 0)$.

(C) Write down $\Psi(x, t)$ as an integral.

(D) Demonstrate explicitly that $\Psi(x, t)$ is a solution to the time dependent Schrödinger Equation.

Problem 3: For an operator $\hat{O}$ we defined its Hermetian adjoint $\hat{O}^\dagger$ by the defining equation:

$$\langle f | \hat{O}^\dagger g \rangle = \langle \hat{O} f | g \rangle$$

.

(A) Find $\hat{p}^\dagger$ by writing $\langle \hat{p} f | g \rangle$ as an integral, and doing whatever it takes to move the action over to g.

(B) Show that this definition is equivalent to the definition used by Griffiths:

$$\langle f | \hat{O} g \rangle = \langle \hat{O}^\dagger f | g \rangle$$

(C) Show that $\left(\hat{O}^\dagger \right)^\dagger = \hat{O}$

(D) Show that $\left(\hat{A}\hat{B} \right)^\dagger = \hat{B}^\dagger \hat{A}^\dagger$

(E) Is $(i\hat{x})$ a Hermetian operator?

(F) Is $(\hat{x}\hat{p})$ a Hermetian operator?

(G) Is $\left(\frac{\partial^2}{\partial x^2} \right)$ a Hermetian operator?

Problem 4: This problem is broken into a lot of steps but each step is very doable. Recall our beloved ladder operators for the simple harmonic oscillator:

$$\hat{a}_- \equiv \frac{1}{\sqrt{2}}\left(i\frac{\hat{p}}{p_0} + \frac{\hat{x}}{x_0}\right)$$

$$\hat{a}_+ \equiv \frac{1}{\sqrt{2}}\left(-i\frac{\hat{p}}{p_0} + \frac{\hat{x}}{x_0}\right)$$

where:

$$x_0 \equiv \sqrt{\frac{\hbar}{m\omega}}, \qquad \text{and} \qquad p_0 \equiv \frac{\hbar}{x_0} = \sqrt{\hbar m\omega}.$$

(A) Show that:

$$\hat{a}_+^\dagger = \hat{a}_-$$

and:

$$\hat{a}_-^\dagger = \hat{a}_+$$

(B) For a particle in state ψ_n show that the expectation value:

$$\langle \hat{a}_+ \hat{a}_- \rangle = n$$

and:

$$\langle \hat{a}_- \hat{a}_+ \rangle = n + 1$$

You may use the known results from lecture:

$$\hat{a}_+ \hat{a}_- \psi_n = n\,\psi_n$$
$$\hat{a}_- \hat{a}_+ \psi_n = (n+1)\,\psi_n$$

(C) We showed that these operators act as raising and lowering operators:

$$\hat{a}_+ \psi_n = c_n \psi_{n+1}$$
$$\hat{a}_- \psi_n = d_n \psi_{n-1}$$

with unknown constants c_n and d_n. To determine the constant d_n, we calculate $\langle \hat{a}_+ \hat{a}_- \rangle$ in a second way:

$$\begin{aligned}
\langle \hat{a}_+ \hat{a}_- \rangle &= \langle \psi_n | \hat{a}_+ \hat{a}_- \psi_n \rangle \\
&= \langle \hat{a}_+^\dagger \psi_n | \hat{a}_- \psi_n \rangle \\
&= \langle \hat{a}_- \psi_n | \hat{a}_- \psi_n \rangle \\
&= \langle d_n \psi_n | d_n \psi_n \rangle \\
&= d_n^* d_n \langle \psi_n | \psi_n \rangle \\
&= |d_n|^2
\end{aligned}$$

so we see that:

$$|d_n|^2 = n.$$

Calculate $\langle \hat{a}_- \hat{a}_+ \rangle$ in similar fashion to determine $|c_n|^2$.

(D) We conclude that:

$$|d_n|^2 = n$$

and

$$|c_n|^2 = n + 1$$

But there is (one) arbitrary phase factor per stationary state ψ_n. We can certainly pick:

$$d_n = \sqrt{n}$$

Show that this choice obligates us to choose:

$$c_n = \sqrt{n+1}$$

Hint: calculate $\langle \hat{a}_+ \hat{a}_- \rangle$ a third way.

(E) Suppose we picked instead:

$$d_n = i\sqrt{n}$$

what would c_n be? Would the spatial wave functions $\psi_n(x)$ be real functions?

(F) Determine the ground state wave function from the simple differential equation:

$$\hat{a}_- \psi_0 = 0$$

(G) Determine $\psi_1(x)$ and $\psi_2(x)$ from:

$$\psi_n(x) = \frac{1}{\sqrt{n!}} (\hat{a}_+)^n \psi_0$$

Additional Hints

For P2.21:
You may use:

$$\int_{-\infty}^{+\infty} e^{-x^2} dx = \sqrt{\pi}$$

Use change of variables to handle additional constants. (A) You should get:

$$A = \left(\frac{2a}{\pi}\right)^{\frac{1}{4}}$$

if you choose a positive real normalization factor.
(B) Calculate the Fourier transform of $\Psi(x,0)$ which I write as $\widetilde{\Psi}(k,0)$. You should find:

$$\widetilde{\Psi}(k,0) \propto \exp\left(-\frac{k^2}{4a}\right)$$

Check dimensionality of $\widetilde{\Psi}$ you should find:

$$[\widetilde{\Psi}] = \left[\frac{1}{\sqrt{k}}\right]$$

Then calculate:

$$\widetilde{\Psi}(k,t) = \widetilde{\Psi}(k,0)\exp(-i\omega t) = \widetilde{\Psi}(k,0)\exp\left(-i\frac{\hbar k^2 t}{2m}\right)$$

make sure to identify the γ term immediately... this will save you a lot of effort! You should find:

$$\widetilde{\Psi}(k,t) \propto \exp\left(-\frac{k^2}{4a}\gamma^2\right)$$

Find $\Psi(x,t)$ as the inverse Fourier transform of $\widetilde{\Psi}(k,t)$.