

Modélisations quantochimiques des forces de  
dispersion de London par la méthode des  
phases aléatoires (RPA) : développements  
méthodologiques

Soutenance de thèse  
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► Dispersion

- polarisation dynamique mutuelle de nuages d'électrons
- fluctuations de la densité électronique qui se corrélient
- besoin d'une bonne description des corrélations longue portée

► RPA

- traitement de la longue-portée
- échange dans la fonction de réponse
- polarisabilité qui conduit au bon comportement à longue-portée ( $C_6/R^6$ )

► Performances

- moyennes, dû à la qualité de la corrélation courte-portée
- nettement meilleures dans un contexte de séparation de portée

formulation "matrice densité de corrélation"

formulation "matrice diélectrique"

formulation "de plasmon"

énergie de  
corrélation  $E_c$

formulation "de Riccati"

problème  
à  $N$ -corps

connexion  
adiabatique

HF    DFT

RSH

théorème de  
fluctuation-  
dissipation

Gradients

Lagrangien

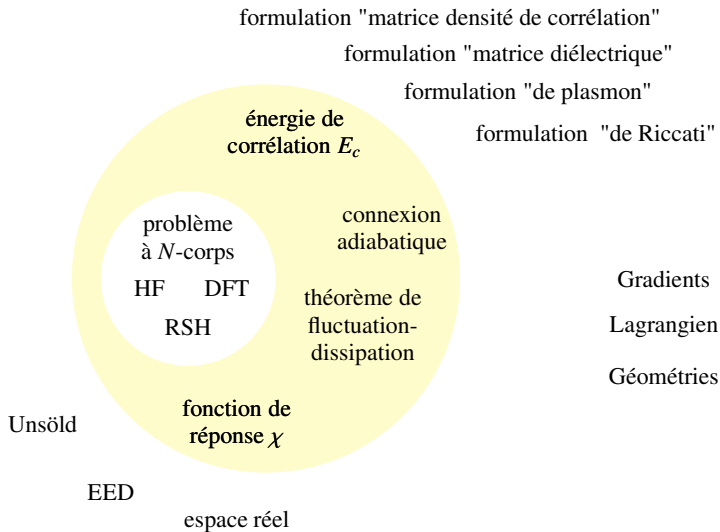
Géométries

Unsöld

fonction de  
réponse  $\chi$

EED

espace réel



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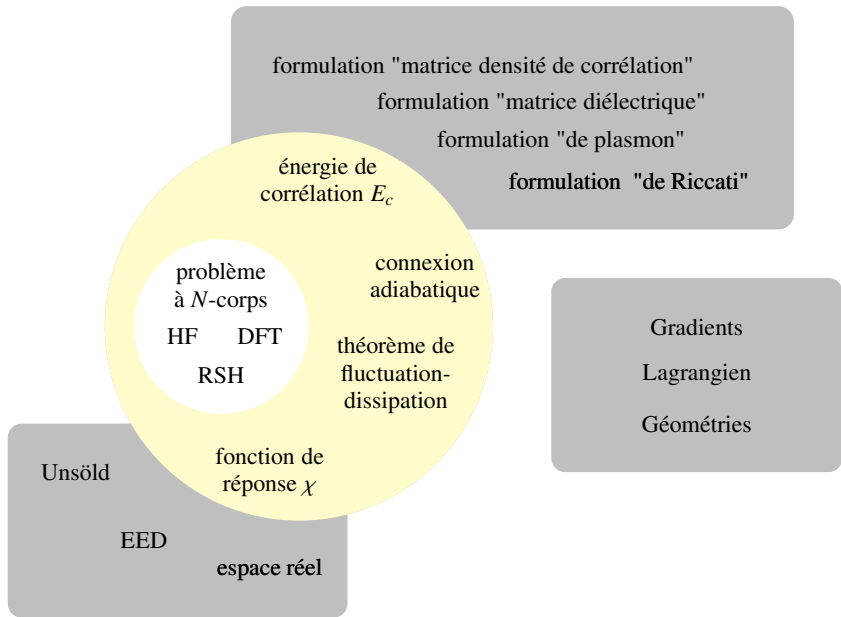
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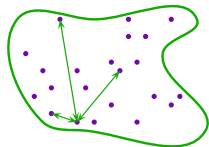


Contexte théorique

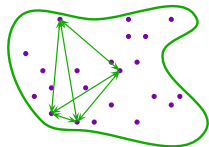
- ▶ montrer le contexte dans lequel sont dérivées les équations qui nous intéressent
- ▶ comprendre ce qu'est l'*énergie de corrélation*  $E_c$
- ▶ comment sont liées les notions telles que :
  - les fluctuations
  - la réponse d'un système (dissipation)

$$\left( \underbrace{\hat{T}_e + \hat{V}_{ne}}_{\sum_i \hat{h}_i} + \underbrace{\hat{V}_{ee}}_{\sum_{ij} \hat{g}_{ij}} \right) |\Psi\rangle = E |\Psi\rangle$$

$$\left( \underbrace{\hat{T}_e + \hat{V}_{ne}}_{\sum_i \hat{h}_i} + \underbrace{\hat{V}_{ee}}_{\sum_{ij} \hat{g}_{ij}} \right) |\Psi\rangle = E |\Psi\rangle$$



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## Approche "fonction d'onde"

$$\left( \underbrace{\hat{T}_e + \hat{V}_{ne}}_{\sum_i \hat{h}_i} + \underbrace{\hat{V}_{ee}}_{\sum_{ij} \hat{g}_{ij}} \right) |\Psi\rangle = E |\Psi\rangle$$

$$E = \min_{\Psi} \left\{ \hat{T} + \hat{V}_{ne} + \hat{V}_{ee} \right\}$$

$$E_{\text{HF}} = \langle \Phi_{\text{HF}} | \sum_i \hat{h}_i + \sum_{ij} \hat{g}_{ij} | \Phi_{\text{HF}} \rangle = \sum \langle \phi_i | \hat{h}_i + \hat{V}_{\text{Hx},i} | \phi_i \rangle \quad \frac{1}{r}$$

$$E = E_{\text{HF}} + E_c$$

## Approche "fonction d'onde"

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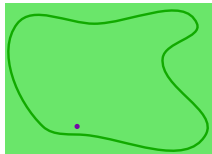
$$E_{\text{HF}} = \langle \Phi_{\text{HF}} | \sum_i \hat{h}_i + \sum_{ij} \hat{g}_{ij} | \Phi_{\text{HF}} \rangle = \sum \langle \phi_i | \hat{h}_i + \hat{V}_{\text{Hx},i} | \phi_i \rangle \quad \frac{1}{r}$$

$$E = E_{\text{HF}} + E_c$$

## Approche "fonctionnelle de la densité"

$$E = \min_n \left\{ T_e[n] + V_{ee}[n] + \int n(\mathbf{r}) v_{ne}(\mathbf{r}) \right\}$$

$$E = \min_n \left\{ T_s[n] + E_H[n] + E_{xc}[n] + \int n(\mathbf{r}) v_{ne}(\mathbf{r}) \right\}$$



## Approche "fonction d'onde"

$$\left( \underbrace{\hat{T}_e + \hat{V}_{ne}}_{\sum_i \hat{h}_i} + \underbrace{\hat{V}_{ee}}_{\sum_{ij} \hat{g}_{ij}} \right) |\Psi\rangle = E |\Psi\rangle$$

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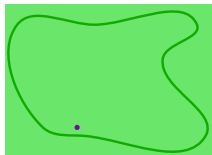
$$E_{\text{HF}} = \langle \Phi_{\text{HF}} | \sum_i \hat{h}_i + \sum_{ij} \hat{g}_{ij} | \Phi_{\text{HF}} \rangle = \sum \langle \phi_i | \hat{h}_i + \hat{V}_{\text{Hx},i} | \phi_i \rangle \quad \frac{1}{r}$$

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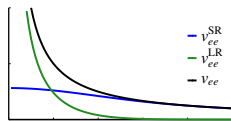
## Séparation de portée

$$E = \min_{\Psi} \left\{ \langle \Psi | \hat{T} + \hat{V}_{ee}^{\text{lr}} | \Psi \rangle + E_{\text{Hxc}}^{\text{sr}}[n_{\Psi}] + \int n_{\Psi}(\mathbf{r}) v_{ne}(\mathbf{r}) \right\}$$

$$E_{\text{RSH}} = \min_{\Phi} \left\{ \underbrace{\langle \Phi | \hat{T} + \hat{V}_{ne} + \hat{V}_{ee}^{\text{lr}} | \Phi \rangle}_{\text{blue}} + \underbrace{E_{\text{Hxc}}^{\text{sr}}[n_{\Phi}]}_{\text{green}} \right\}$$

énergie totale

$$E = E_{\text{RSH}} + E_c^{\text{lr}}$$



$$\frac{1}{r} = v_{ee}^{\text{lr}}(\mathbf{r}) + v_{ee}^{\text{sr}}(\mathbf{r})$$



# Connexion adiabatique

$$E_{\text{RSH}} = \min \left\{ \langle \Phi | \hat{T} + \hat{V}_{ne} + \hat{V}_{ee}^{\text{lr}} | \Phi \rangle + E_{\text{Hxc}}^{\text{sr}}[n_{\Phi}] \right\}$$

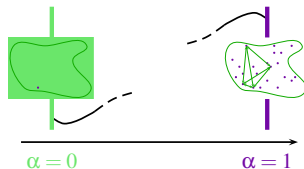
# Connexion adiabatique

$$E_{\text{RSH}} = \min \left\{ \langle \Phi | \hat{T} + \hat{V}_{ne} + \hat{V}_{ee}^{\text{lr}} | \Phi \rangle + E_{\text{Hxc}}^{\text{sr}}[n_{\Phi}] \right\}$$

$$E_{\alpha} = \min_{\Psi} \left\{ \langle \Psi | \hat{T} + \hat{V}_{ne} + (1 - \alpha) \hat{V}_{\text{Hx,HF}}^{\text{lr}} + \alpha \hat{V}_{ee}^{\text{lr}} | \Psi \rangle + E_{\text{Hxc}}^{\text{sr}}[n_{\Psi}] \right\}$$

$$\hat{H}_{\alpha} = \hat{T} + \hat{V}_{ne} + (1 - \alpha) \hat{V}_{\text{Hx,HF}}^{\text{lr}} + \alpha \hat{V}_{ee}^{\text{lr}} + \hat{V}_{\text{Hxc}}^{\text{sr}}[\Psi_{\alpha}]$$

- ▶ systèmes soumis à un potentiel intermédiaire
- ▶ connexion entre système non-interagissant et le système réel



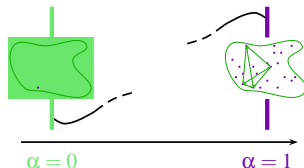
# Connexion adiabatique

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## Énergie de corrélation

$$\frac{dE_{\alpha}}{d\alpha} = \langle \Psi_{\alpha} | \hat{W}_{ee}^{\text{lr}} | \Psi_{\alpha} \rangle \quad ; \quad \hat{W}_{ee}^{\text{lr}} = \frac{d\hat{H}_{\alpha}}{d\alpha} = \hat{V}_{ee}^{\text{lr}} - \hat{V}_{\text{Hx,HF}}^{\text{lr}}$$

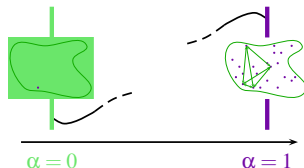
# Connexion adiabatique

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- ▶ systèmes soumis à un potentiel intermédiaire
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## Énergie de corrélation

$$\int_0^1 d\alpha \frac{dE_{\alpha}}{d\alpha} = \int_0^1 d\alpha \langle \Psi_{\alpha} | \hat{W}_{ee}^{\text{lr}} | \Psi_{\alpha} \rangle \quad ; \quad \hat{W}_{ee}^{\text{lr}} = \frac{d\hat{H}_{\alpha}}{d\alpha} = \hat{V}_{ee}^{\text{lr}} - \hat{V}_{\text{Hx,HF}}^{\text{lr}}$$

$$E_c^{\text{AC}} = \int_0^1 d\alpha \frac{1}{2} \int_{1,2} w(1,2) P_{c,\alpha}(1,2)$$

# Théorème de Fluctuation-Dissipation

$$E_c^{\text{AC}} = \frac{1}{2} \int_0^1 d\alpha \int_{1,2} w(1,2) P_{c,\alpha}(1,2)$$

$$\begin{aligned} P_{c,\alpha}(1,2) &= n_{2,\alpha}(1,2) - n_{2,0}(1,2) \\ &= \langle \Psi_\alpha | \delta \hat{n}_1(2) \delta \hat{n}_1(1) | \Psi_\alpha \rangle \\ &\quad - \langle \Psi_0 | \delta \hat{n}_1(2) \delta \hat{n}_1(1) | \Psi_0 \rangle + \Delta n_\alpha \end{aligned}$$

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## Fluctuations

$$E_c^{\text{AC}} = \frac{1}{2} \int_0^1 d\alpha \int_{1,2} w(1,2) \left[ \langle \Psi_\alpha | \delta \hat{n}(1) \delta \hat{n}(2) | \Psi_\alpha \rangle - \langle \Phi_0 | \delta \hat{n}(1) \delta \hat{n}(2) | \Phi_0 \rangle + \Delta n_\alpha \right]$$

# Théorème de Fluctuation-Dissipation

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## Réponse du système

$$\langle \Psi_\alpha | \delta \hat{n}(1) \delta \hat{n}(2) | \Psi_\alpha \rangle = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \chi_\alpha(1,2;\omega)$$

- système répond de la même manière quand il est mis hors équilibre par une force extérieure (dissipation) ou par des fluctuations quantiques

$$E_c^{AC-FDT} = \frac{1}{2} \int_0^1 d\alpha \int_{1,2} \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} w(1,2) \left[ \chi_\alpha(1,2;\omega) - \chi_0(1,2;\omega) + \Delta n_\alpha \right]$$

- expression exacte de l'énergie de corrélation

Approximation de la phase aléatoire



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$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{1,2;1',2'} w(1,2;1',2') \left[ \chi_\alpha(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_\alpha \right]$$

## “Flavors” de RPA dosage de l’inclusion de l’échange

- ▶ dans fonction de réponse  $\chi_\alpha$  : dRPA/RPA<sub>x</sub>
- ▶ dans l’interaction  $w$  : intégrales non-antisymétrisées (I) ou antisymétrisées (II)

## Intégrations analytiques/numériques

- ▶ intégrale analytique sur la fréquence  $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$   
formulation “matrice densité de corrélation”  $\mathbf{P}_{c,\alpha}$
- ▶ deux intégrales analytiques  $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$  et  $\left( \int_0^1 d\alpha \right)$   
formulation “de plasmon” et “de Riccati” ((d)rCCD)
- ▶ intégrale analytique sur la constante de couplage  $\left( \int_0^1 d\alpha \right)$   
formulation “matrice diélectrique”  $\varepsilon = \mathbf{1} - \mathbf{\Pi}_0 \mathbf{K}$

## Formulation “matrice densité de corrélation” $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1')\psi_a(1)\psi_a^*(2')\psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2')\psi_a(2)\psi_a^*(1')\psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - f_{\alpha}(1,2;1',2';\omega)$$

$$f_{\alpha}(1,2;1',2';\omega) = \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega)$$

## Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

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## Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

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$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

$$\mathbb{P}_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \mathbb{\Pi}_{\alpha}^{\text{RPA}} - \mathbb{\Pi}_0 \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1')\psi_a(1)\psi_a^*(2')\psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2')\psi_a(2)\psi_a^*(1')\psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

$$(\mathbb{\Pi}_0)^{-1} = \omega \mathbb{\Delta} - \mathbb{\Lambda}_0$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - \alpha f^{\text{RPA}}(1,2;1',2')$$

$$(\mathbb{\Pi}_{\alpha}^{\text{RPA}})^{-1} = (\mathbb{\Pi}_0)^{-1} - \alpha \mathbb{F}^{\text{RPA}}$$

$$\begin{aligned} f_{\alpha}(1,2;1',2';\omega) &= \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega) \\ &= \alpha f^{\text{RPA}}(1,2;1',2') = \alpha \left( \frac{1}{r_1 - r_2} + f_x(1,2;1',2') \right) \end{aligned}$$

# Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

$$\mathbb{P}_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \mathbb{P}_{\alpha}^{\text{RPA}} - \mathbb{P}_0 \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1')\psi_a(1)\psi_a^*(2')\psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2')\psi_a(2)\psi_a^*(1')\psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

$$(\mathbb{P}_0)^{-1} = \omega \Delta - \mathbb{A}_0$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - \alpha f^{\text{RPA}}(1,2;1',2')$$

$$(\mathbb{P}_{\alpha}^{\text{RPA}})^{-1} = (\mathbb{P}_0)^{-1} - \alpha \mathbb{F}^{\text{RPA}}$$

$$\begin{aligned} f_{\alpha}(1,2;1',2';\omega) &= \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega) \\ &= \alpha f^{\text{RPA}}(1,2;1',2') = \alpha \left( \frac{1}{r_1 - r_2} + f_x(1,2;1',2') \right) \end{aligned}$$

$$f^{\text{dRPA}} = w(1,2) [\delta(1,1') \delta(2,2')] + \dots$$

$$f^{\text{RPA}_x} = w(1,2) [\delta(1,1') \delta(2,2') - \delta(1,2') \delta(2,1')] + \dots$$

# Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

$$\mathbb{P}_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \mathbb{P}_{\alpha}^{\text{RPA}} - \mathbb{P}_0 \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1') \psi_a(1) \psi_a^*(2') \psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2') \psi_a(2) \psi_a^*(1') \psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

$$(\mathbb{P}_0)^{-1} = \omega \Delta - \mathbb{A}_0$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - \alpha f^{\text{RPA}}(1,2;1',2')$$

$$(\mathbb{P}_{\alpha}^{\text{RPA}})^{-1} = (\mathbb{P}_0)^{-1} - \alpha \mathbb{F}^{\text{RPA}}$$

$$f_{\alpha}(1,2;1',2';\omega) = \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega) \\ = \alpha f^{\text{RPA}}(1,2;1',2') = \alpha \left( \frac{1}{r_1 - r_2} + f_x(1,2;1',2') \right)$$

$$f^{\text{dRPA}} = w(1,2) [\delta(1,1') \delta(2,2')]$$

$$\mathbb{F}^{\text{dRPA}} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}$$

$$f^{\text{RPAx}} = w(1,2) [\delta(1,1') \delta(2,2') - \delta(1,2') \delta(2,1')]$$

$$\mathbb{F}^{\text{RPAx}} = \begin{pmatrix} \mathbf{A}' & \mathbf{B}' \\ \mathbf{B} & \mathbf{A}' \end{pmatrix}$$



# Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

$$\mathbb{P}_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \mathbb{\Pi}_{\alpha}^{\text{RPA}} - \mathbb{\Pi}_0 \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1') \psi_a(1) \psi_a^*(2') \psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2') \psi_a(2) \psi_a^*(1') \psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

$$(\mathbb{\Pi}_0)^{-1} = \omega \Delta - \mathbb{A}_0$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - \alpha f^{\text{RPA}}(1,2;1',2')$$

$$(\mathbb{\Pi}_{\alpha}^{\text{RPA}})^{-1} = (\mathbb{\Pi}_0)^{-1} - \alpha \mathbb{F}^{\text{RPA}}$$

$$f_{\alpha}(1,2;1',2';\omega) = \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega) \\ = \alpha f^{\text{RPA}}(1,2;1',2') = \alpha \left( \frac{1}{r_1 - r_2} + f_x(1,2;1',2') \right)$$

$$f^{\text{dRPA}} = w(1,2) [\delta(1,1') \delta(2,2')]$$

$$\mathbb{F}^{\text{dRPA}} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}$$

$$f^{\text{RPAx}} = w(1,2) [\delta(1,1') \delta(2,2') - \delta(1,2') \delta(2,1')]$$

$$\mathbb{F}^{\text{RPAx}} = \begin{pmatrix} \mathbf{A}' & \mathbf{B}' \\ \mathbf{B} & \mathbf{A}' \end{pmatrix}$$

## Énergie de corrélation

$$E_c^{\text{AC}} = \int_0^1 d\alpha \frac{1}{2} \int_{1,2} w(1,2) P_{c,\alpha}(1,2)$$

$$E_c^{\text{RPAx-I}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{I}} \cdot \mathbb{P}_{c,\alpha}^{\text{RPAx}})$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{I}} \cdot \mathbb{P}_{c,\alpha}^{\text{dRPA}})$$

$$E_c^{\text{RPAx-II}} = \frac{1}{4} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{II}} \cdot \mathbb{P}_{c,\alpha}^{\text{RPAx}})$$

$$E_c^{\text{dRPA-II}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{II}} \cdot \mathbb{P}_{c,\alpha}^{\text{dRPA}})$$

# Formulation "matrice densité de corrélation" $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$P_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega) - \chi_0(1,2;1',2';\omega) + \Delta n_{\alpha} \right]$$

$$\mathbb{P}_{c,\alpha}^{\text{RPA}} = \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \left[ \mathbb{P}_{\alpha}^{\text{RPA}} - \mathbb{P}_0 \right]$$

Représentation de Lehmann

$$\chi_0(1,2;1',2';\omega) = \sum_{ia} \frac{\psi_i^*(1') \psi_a(1) \psi_a^*(2') \psi_i(2)}{\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+} + \frac{\psi_i^*(2') \psi_a(2) \psi_a^*(1') \psi_i(1)}{-\omega - (\varepsilon_a - \varepsilon_i) + i\eta^+}$$

$$(\mathbb{P}_0)^{-1} = \omega \Delta - \mathbb{A}_0$$

Équation de Bethe-Salpeter

$$\chi_{\alpha}^{\text{RPA}}(1,2;1',2';\omega)^{-1} = \chi_0(1,2;1',2';\omega)^{-1} - \alpha f^{\text{RPA}}(1,2;1',2')$$

$$(\mathbb{P}_{\alpha}^{\text{RPA}})^{-1} = (\mathbb{P}_0)^{-1} - \alpha \mathbb{F}^{\text{RPA}}$$

$$f_{\alpha}(1,2;1',2';\omega) = \frac{\alpha}{r_1 - r_2} + f_{x,\alpha}(1,2;1',2';\omega) + f_{c,\alpha}(1,2;1',2';\omega) \\ = \alpha f^{\text{RPA}}(1,2;1',2') = \alpha \left( \frac{1}{r_1 - r_2} + f_x(1,2;1',2') \right)$$

$$f^{\text{dRPA}} = w(1,2) [\delta(1,1') \delta(2,2')]$$

$$\mathbb{F}^{\text{dRPA}} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}$$

$$f^{\text{RPAx}} = w(1,2) [\delta(1,1') \delta(2,2') - \delta(1,2') \delta(2,1')]$$

$$\mathbb{F}^{\text{RPAx}} = \begin{pmatrix} \mathbf{A}' & \mathbf{B}' \\ \mathbf{B} & \mathbf{A}' \end{pmatrix}$$

## Énergie de corrélation

$$E_c^{\text{AC}} = \int_0^1 d\alpha \frac{1}{2} \int_{1,2} w(1,2) P_{c,\alpha}(1,2)$$

$$E_c^{\text{RPAx-I}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{I}} \cdot \mathbb{P}_{c,\alpha}^{\text{RPAx}})$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{I}} \cdot \mathbb{P}_{c,\alpha}^{\text{dRPA}})$$

$$\mathbb{W}^{\text{I}} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}$$

$$E_c^{\text{RPAx-II}} = \frac{1}{4} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{II}} \cdot \mathbb{P}_{c,\alpha}^{\text{RPAx}})$$

$$E_c^{\text{dRPA-II}} = \frac{1}{2} \int_0^1 d\alpha \text{Tr}(\mathbb{W}^{\text{II}} \cdot \mathbb{P}_{c,\alpha}^{\text{dRPA}})$$

$$\mathbb{W}^{\text{II}} = \begin{pmatrix} \mathbf{A}' & \mathbf{B}' \\ \mathbf{B} & \mathbf{A}' \end{pmatrix}$$

Formulation “de plasmon”  $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$  et  $\left(\int_0^1 d\alpha\right)$

$$\Pi_{\alpha}^{-1}=\omega\Delta-\Lambda_0-\alpha\mathbb{F}$$

$$\mathbb{W}^{1/2}=\mathbb{F}^{d/x}$$

$$E_c=\int_0^1 d\alpha\,\mathrm{Tr}\left(\mathbb{W}\mathbb{P}_{c,\alpha}\right)=\int_0^1 d\alpha\,\sum\frac{d\omega_{\alpha,n}}{d\alpha}-\frac{d\omega_{\alpha,n}}{d\alpha}\Big|_{\alpha=0}=\sum\omega_{1,n}^{\mathrm{RPA}}-\omega_{1,n}^{\mathrm{TDA}}$$

Formulation “de plasmon”  $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$  et  $\left(\int_0^1 d\alpha\right)$

$$\Pi_{\alpha}^{-1} = \omega \Delta - \Lambda_0 - \alpha F$$

$$W^{I/II} = F^{d/x}$$

$$E_c = \int_0^1 d\alpha \operatorname{Tr} \left( \underline{W P}_{c,\alpha} \right) = \int_0^1 d\alpha \sum \frac{d\omega_{\alpha,n}}{d\alpha} - \frac{d\omega_{\alpha,n}}{d\alpha} \Big|_{\alpha=0} = \sum \omega_{1,n}^{\text{RPA}} - \omega_{1,n}^{\text{TDA}}$$

Formulation “de Riccati”  $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$  et  $\left(\int_0^1 d\alpha\right)$

$$\begin{pmatrix} \varepsilon + \mathbf{A}' & \mathbf{B} \\ \mathbf{B} & \varepsilon + \mathbf{A}' \end{pmatrix} \begin{pmatrix} \mathbf{X} \\ \mathbf{Y} \end{pmatrix} = \omega \begin{pmatrix} \mathbf{X} \\ -\mathbf{Y} \end{pmatrix} \quad \mathbf{T} = \mathbf{YX}^{-1}$$

$$\mathbf{R}[\mathbf{T}] = \mathbf{B} + [\mathbf{A}', \mathbf{T}]_+ + \mathbf{TBT} + [\varepsilon, \mathbf{T}]_+ = \mathbf{0}$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \operatorname{tr} \left\{ \mathbf{B}^{\text{dRPA}} \mathbf{T}^{\text{dRPA}} \right\}$$

$$E_c^{\text{RPAx-II}} = \frac{1}{4} \operatorname{tr} \left\{ \mathbf{B}^{\text{RPAx}} \mathbf{T}^{\text{RPAx}} \right\}$$

équivalent à (d)rCCD

## Variantes (formulation “matrice densité de corrélation”) $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$

$$E_c^{d/x-I} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{K} \right\}$$

$$E_c^{d/x-II} \rightsquigarrow E_c^{d/x-IIa} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{B} \right\} \quad (\mathbf{1} + \mathbf{P}_\alpha)^{-1} \approx (\mathbf{1} - \mathbf{P}_\alpha)$$

## Variantes (formulation “matrice densité de corrélation”) $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$E_c^{d/x-I} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{K} \right\}$$

$$E_c^{d/x-II} \rightsquigarrow E_c^{d/x-IIa} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{B} \right\} \quad (\mathbf{1} + \mathbf{P}_\alpha)^{-1} \approx (\mathbf{1} - \mathbf{P}_\alpha)$$

## Variantes (formulation “de Riccati”) $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$ et $\left( \int_0^1 d\alpha \right)$

$$\mathbf{0} = \mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+$$

$$E_c^{dRPA-I} = \operatorname{tr} \{ \mathbf{K} \mathbf{T} \}$$

$$E_c^{dRPA-I-SOSEX} = \operatorname{tr} \{ \mathbf{B} \mathbf{T} \}$$

$$\int_0^1 d\alpha \mathbf{P}_\alpha \text{ et } \mathbf{T} ?$$

Jansen, Liu, Ángyán; *J. Chem. Phys.* (2010)

## Variantes (formulation “matrice densité de corrélation”) $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$

$$E_c^{d/x-I} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{K} \right\}$$

$$E_c^{d/x-II} \rightsquigarrow E_c^{d/x-IIa} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{B} \right\} \quad (1 + \mathbf{P}_\alpha)^{-1} \approx (1 - \mathbf{P}_\alpha)$$

## Variantes (formulation “de Riccati”) $\left( \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \right)$ et $\left( \int_0^1 d\alpha \right)$

$$\mathbf{0} = \mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+$$

$$E_c^{\text{dRPA-I}} = \operatorname{tr} \{ \mathbf{K} \mathbf{T} \}$$

$$E_c^{\text{dRPA-I-SOSEX}} = \operatorname{tr} \{ \mathbf{B} \mathbf{T} \}$$

$$\int_0^1 d\alpha \mathbf{P}_\alpha \text{ et } \mathbf{T} ?$$

Hesselmann; *Phys. Rev. A* (2012)

Jansen, Liu, Ángyán; *J. Chem. Phys.* (2010)

$$\mathbf{0} = \mathbf{B} + [\mathbf{B}, \mathbf{T}]_+ + \mathbf{T} \mathbf{B} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+$$

$$E_c^{\text{RPAX2}} = \operatorname{tr} \{ \mathbf{K} \mathbf{T} \}$$

## Variantes (formulation “matrice densité de corrélation”) $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$

$$E_c^{d/x-I} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{K} \right\}$$

$$E_c^{d/x-II} \rightsquigarrow E_c^{d/x-IIa} = \int_0^1 d\alpha \operatorname{tr} \left\{ \mathbf{P}_\alpha^{d/x} \mathbf{B} \right\} \quad (1 + \mathbf{P}_\alpha)^{-1} \approx (1 - \mathbf{P}_\alpha)$$

## Variantes (formulation “de Riccati”) $\left(\int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i}\right)$ et $\left(\int_0^1 d\alpha\right)$

$$\mathbf{0} = \mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+$$

$$E_c^{dRPA-I} = \operatorname{tr} \{ \mathbf{K} \mathbf{T} \}$$

$$E_c^{dRPA-I-SOSEX} = \operatorname{tr} \{ \mathbf{B} \mathbf{T} \}$$

$$\int_0^1 d\alpha \mathbf{P}_\alpha \text{ et } \mathbf{T} ?$$

Hesselmann; *Phys. Rev. A* (2012)

Jansen, Liu, Ángyán; *J. Chem. Phys.* (2010)

$$\mathbf{0} = \mathbf{B} + [\mathbf{B}, \mathbf{T}]_+ + \mathbf{T} \mathbf{B} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+$$

$$E_c^{RPAX2} = \operatorname{tr} \{ \mathbf{K} \mathbf{T} \}$$

Comprendre et Unifier



## Formulation “matrice diélectrique” $\left(\int_0^1 d\alpha\right)$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \text{Tr}\{\mathbb{P}_\alpha^{d/x} \mathbb{W}^{l/ll} - \mathbb{P}_0 \mathbb{W}^{l/ll}\}$$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \sum_{n=2} \alpha^{n-1} \text{Tr}\{(\mathbb{P}_0 \mathbb{W}^{d/x})^{n-1} \mathbb{P}_0 \mathbb{W}^{l/ll}\}$$

$$\mathbb{P}_\alpha^{d/x} = (1 - \alpha \mathbb{P}_0 \mathbb{W}^{d/x})^{-1} \mathbb{P}_0$$

$$(1 - \alpha x)^{-1} = 1 + \sum_{n=2} \alpha^{n-1} x^{n-1}$$

$$-\sum_{n=2} \frac{x^n}{n} = \text{Log}(1-x) + x$$

## Formulation “matrice diélectrique” $\left(\int_0^1 d\alpha\right)$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \text{Tr}\{\Pi_\alpha^{d/x} \mathbb{W}^{l/II} - \Pi_0 \mathbb{W}^{l/II}\}$$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \sum_{n=2}^{\infty} \alpha^{n-1} \text{Tr}\{(\Pi_0 \mathbb{W}^{d/x})^{n-1} \Pi_0 \mathbb{W}^{l/II}\}$$

$$\begin{aligned} \Pi_\alpha^{d/x} &= (1 - \alpha \Pi_0 \mathbb{W}^{d/x})^{-1} \Pi_0 \\ (1 - \alpha x)^{-1} &= 1 + \sum_{n=2}^{\infty} \alpha^{n-1} x^{n-1} \\ - \sum_{n=2}^{\infty} \frac{x^n}{n} &= \text{Log}(1-x) + x \end{aligned}$$

dRPA-I

$$\mathbb{W}^d = \mathbb{W}^l = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix} \quad ; \quad \text{Tr}\{(\Pi_0 \mathbb{W}^d)^n\} = \text{tr}\{(\Pi_0 \mathbf{K})^n\}$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{K}) + \Pi_0 \mathbf{K}\}$$

$$\epsilon = \mathbf{1} - \Pi_0 \mathbf{K}$$

# Formulation "matrice diélectrique" $\left(\int_0^1 d\alpha\right)$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \text{Tr}\{\Pi_\alpha^{d/x} \mathbb{W}^{I/II} - \Pi_0 \mathbb{W}^{I/II}\}$$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \sum_{n=2}^{\infty} \alpha^{n-1} \text{Tr}\{(\Pi_0 \mathbb{W}^{d/x})^{n-1} \Pi_0 \mathbb{W}^{I/II}\}$$

$$\begin{aligned} \Pi_\alpha^{d/x} &= (1 - \alpha \Pi_0 \mathbb{W}^{d/x})^{-1} \Pi_0 \\ (1 - \alpha x)^{-1} &= 1 + \sum_{n=2}^{\infty} \alpha^{n-1} x^{n-1} \\ - \sum_{n=2}^{\infty} \frac{x^n}{n} &= \text{Log}(1-x) + x \end{aligned}$$

dRPA-I

$$\mathbb{W}^d = \mathbb{W}^I = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix} ; \quad \text{Tr}\{(\Pi_0 \mathbb{W}^d)^n\} = \text{tr}\{(\Pi_0 \mathbf{K})^n\}$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{K}) + \Pi_0 \mathbf{K}\}$$

$$\epsilon = \mathbf{1} - \Pi_0 \mathbf{K}$$

dRPA-II

$$\mathbb{W}^d = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix} ; \quad \mathbb{W}^{II} = \begin{pmatrix} \mathbf{A}' & \mathbf{B} \\ \mathbf{B} & \mathbf{A}' \end{pmatrix} = \begin{pmatrix} \mathbf{B} & \mathbf{B} \\ \mathbf{B} & \mathbf{B} \end{pmatrix} + \begin{pmatrix} \mathbf{A}' - \mathbf{B} & \mathbf{0} \\ \mathbf{0} & \mathbf{A}' - \mathbf{B} \end{pmatrix}$$

$$E_c^{\text{dRPA-IIa}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{K}) \mathbf{K}^{-1} \mathbf{B} + \Pi_0 \mathbf{B}\}$$

# Formulation "matrice diélectrique" $\left(\int_0^1 d\alpha\right)$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \text{Tr}\{\Pi_\alpha^{d/x} \mathbb{W}^{I/II} - \Pi_0 \mathbb{W}^{I/II}\}$$

$$E_c^{\text{AC-FDT}} = \frac{1}{2} \int_0^1 d\alpha \int_{-\infty}^{\infty} \frac{-d\omega}{2\pi i} \sum_{n=2}^{\infty} \alpha^{n-1} \text{Tr}\{(\Pi_0 \mathbb{W}^{d/x})^{n-1} \Pi_0 \mathbb{W}^{I/II}\}$$

$$\begin{aligned} \Pi_\alpha^{d/x} &= (1 - \alpha \Pi_0 \mathbb{W}^{d/x})^{-1} \Pi_0 \\ (1 - \alpha x)^{-1} &= 1 + \sum_{n=2}^{\infty} \alpha^{n-1} x^{n-1} \\ -\sum_{n=2}^{\infty} \frac{x^n}{n} &= \text{Log}(1-x) + x \end{aligned}$$

dRPA-I

$$\mathbb{W}^d = \mathbb{W}^I = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix} ; \quad \text{Tr}\{(\Pi_0 \mathbb{W}^d)^n\} = \text{tr}\{(\Pi_0 \mathbf{K})^n\}$$

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{K}) + \Pi_0 \mathbf{K}\}$$

$$\epsilon = \mathbf{1} - \Pi_0 \mathbf{K}$$

dRPA-II

$$\mathbb{W}^d = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix} ; \quad \mathbb{W}^{II} = \begin{pmatrix} \mathbf{A}' & \mathbf{B} \\ \mathbf{B} & \mathbf{A}' \end{pmatrix} = \begin{pmatrix} \mathbf{B} & \mathbf{B} \\ \mathbf{B} & \mathbf{B} \end{pmatrix} + \begin{pmatrix} \mathbf{A}' - \mathbf{B} & \mathbf{0} \\ \mathbf{0} & \mathbf{A}' - \mathbf{B} \end{pmatrix}$$

$$E_c^{\text{dRPA-IIa}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{K}) \mathbf{K}^{-1} \mathbf{B} + \Pi_0 \mathbf{B}\}$$

RPAx-I

$$\mathbb{W}^x = \begin{pmatrix} \mathbf{A}' & \mathbf{B} \\ \mathbf{B} & \mathbf{A}' \end{pmatrix} = \begin{pmatrix} \mathbf{B} & \mathbf{B} \\ \mathbf{B} & \mathbf{B} \end{pmatrix} + \begin{pmatrix} \mathbf{A}' - \mathbf{B} & \mathbf{0} \\ \mathbf{0} & \mathbf{A}' - \mathbf{B} \end{pmatrix} ; \quad \mathbb{W}^{II} = \begin{pmatrix} \mathbf{K} & \mathbf{K} \\ \mathbf{K} & \mathbf{K} \end{pmatrix}$$

$$E_c^{\text{RPAx-Ia}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr}\{\text{Log}(\mathbf{1} - \Pi_0 \mathbf{B}) \mathbf{B}^{-1} \mathbf{K} + \Pi_0 \mathbf{K}\}$$

# Conclusion

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- ▶ différentes formulations pour écrire les équations RPA
  - formulation “matrice densité de corrélation”
  - formulation “matrice diélectrique”
  - formulation “de Riccati”
- ▶ des variantes émergent dans chacune de ces formulations
- ▶ comprendre les liens entre ces variantes
- ▶ systématiser les explorations (notamment “matrice diélectrique”)

formulation "matrice densité de corrélation"

formulation "matrice diélectrique"

formulation "de plasmon"

formulation "de Riccati"

énergie de  
corrélation  $E_c$

problème  
à  $N$ -corps

connexion  
adiabatique

HF DFT

RSH

théorème de  
fluctuation-  
dissipation

Unsöld

fonction de  
réponse  $\chi$

EED

espace réel

Gradients

Lagrangien

Géométries

Adaptation de l'approximation du  
dénominateur effectif à l'espace réel

$$E_c^{\text{dRPA-I}} = \frac{1}{2} \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \text{tr} \{ \text{Log}(\mathbf{1} - \mathbf{\Pi}_0 \mathbf{K}) + \mathbf{\Pi}_0 \mathbf{K} \}$$

- ▶ éviter la double sommation occ/vir ( $\chi_0 = \sum_{ia} \dots$ )
- ▶ éviter les états excités
- ▶ applications au calcul de :
  - polarisabilité dynamique
  - coefficients  $C_6$
  - énergie RPA



$$\chi(\mathbf{r}_1, \mathbf{r}_2; i\omega) = 2\text{Re}\left(\sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \Omega_\alpha}\right) \doteq 2\text{Re}\left(\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)\right)$$

## Approximation de Unsöld

$$\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \approx \frac{\sum n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \overline{\Omega}}$$

- résolution de l'identité
- $\overline{\Omega}_{\mathbf{r}_1, \mathbf{r}_2, \omega}$  dépend de  $(\mathbf{r}_1, \mathbf{r}_2, \omega)$

$\alpha$  états excités

$$\sum_{\alpha} \langle 0 | \hat{n}(\mathbf{r}_1) | \alpha \rangle \langle \alpha | \hat{n}(\mathbf{r}_2) | 0 \rangle$$

$$\sum_{\alpha} |\alpha\rangle \langle \alpha| + |0\rangle \langle 0| = 1$$

## Approximation de l'EED

- généralisation où  $\overline{\Omega}$  est une fonction
- moyennes sur l'état fondamental
- nouvelle énergie effective à approximer :  $\Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; \omega)$
- hiérarchie d'équations

$$\chi(\mathbf{r}_1, \mathbf{r}_2; i\omega) = 2\text{Re}\left(\sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \Omega_\alpha}\right) \doteq 2\text{Re}\left(\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)\right)$$

$$\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \doteq \frac{\sum n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)}$$

$$\Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) = \sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)\Omega_\alpha}{i\omega - \Omega_\alpha} \doteq \chi^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega)$$

$$\chi(\mathbf{r}_1, \mathbf{r}_2; i\omega) = 2\text{Re}\left(\sum_{\alpha \neq 0} \frac{n_{\alpha}(\mathbf{r}_1)n_{\alpha}(\mathbf{r}_2)}{i\omega - \Omega_{\alpha}}\right) \doteq 2\text{Re}\left(\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)\right)$$

$$\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \doteq \frac{\sum n_{\alpha}(\mathbf{r}_1)n_{\alpha}(\mathbf{r}_2)}{i\omega - \Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)}$$

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$$\chi^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \doteq \frac{\sum n_{\alpha}(\mathbf{r}_1)n_{\alpha}(\mathbf{r}_2)\Omega_{\alpha}}{i\omega - \Omega^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega)}$$

$$\Omega^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \chi^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega) = \sum_{\alpha \neq 0} \frac{n_{\alpha}(\mathbf{r}_1)n_{\alpha}(\mathbf{r}_2)\Omega_{\alpha}\Omega_{\alpha}}{i\omega - \Omega_{\alpha}} \doteq \chi^{jj}(\mathbf{r}_1, \mathbf{r}_2; i\omega)$$

$$\chi(\mathbf{r}_1, \mathbf{r}_2; i\omega) = 2\text{Re}\left(\sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \Omega_\alpha}\right) \doteq 2\text{Re}\left(\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)\right)$$

$$\chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \doteq \frac{\sum n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)}{i\omega - \Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega)}$$

$$\Omega^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \chi^{nn}(\mathbf{r}_1, \mathbf{r}_2; i\omega) = \sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)\Omega_\alpha}{i\omega - \Omega_\alpha} \doteq \chi^{nj}(\mathbf{r}_1, \mathbf{r}_2; i\omega)$$

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$$\Omega^{jj}(\mathbf{r}_1, \mathbf{r}_2; i\omega) \chi^{jj}(\mathbf{r}_1, \mathbf{r}_2; i\omega) = \sum_{\alpha \neq 0} \frac{n_\alpha(\mathbf{r}_1)n_\alpha(\mathbf{r}_2)\Omega_\alpha\Omega_\alpha\Omega_\alpha}{i\omega - \Omega_\alpha} \doteq \dots$$

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
- ▶ on cherche une expression explicite de  $\Omega^{nn}$

$$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$$

$$\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$$

$$\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$$

$$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$$

$$\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$$

$$\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$$

$$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$$

$$\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$$

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
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$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$ $\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$ $\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$	$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$ $\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$ $\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$	$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$ $\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$
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$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
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$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$ $\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$ $\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$	$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$ $\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$ $\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$	$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$ $\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$
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$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}}{i\omega - \frac{S^{jj}}{S^{nj}} \frac{i\omega - \Omega^{nj}}{i\omega - \Omega^{jj}}}$$

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
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$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$ $\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$ $\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$	$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$ $\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$ $\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$	$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$ $\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$
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$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}}{i\omega - \frac{S^{jj}}{S^{nj}} \frac{i\omega - \Omega^{nj}}{i\omega - \Omega^{jj}}}$$

...



# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
- ▶ on cherche une expression explicite de  $\Omega^{nn}$

$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$ $\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$ $\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$	$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$ $\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$ $\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$	$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$ $\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$
--	--	--

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

$$\Omega^{nn(1)} = \frac{S^{nj}}{S^{nn}}$$

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}}{i\omega - \frac{S^{jj}}{S^{nj}} \frac{i\omega - \Omega^{nj}}{i\omega - \Omega^{jj}}}$$

...

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
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--	--	--

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

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$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}}{i\omega - \frac{S^{jj}}{S^{nj}} \frac{i\omega - \Omega^{nj}}{i\omega - \Omega^{jj}}}$$

$$\Omega^{nn(2)} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}}}{i\omega - \frac{S^{jj}}{S^{nj}}}$$

...

# Approximations de $\Omega^{nn}$

- ▶ les numérateurs sont des règles de sommes
- ▶ on cherche une expression explicite de  $\Omega^{nn}$

$S^{nn} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2)$ $\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$ $\Omega^{nn} = \frac{\chi^{nj}}{\chi^{nn}}$	$S^{nj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha}$ $\chi^{nj} = \frac{S^{nj}}{i\omega - \Omega^{nj}}$ $\Omega^{nj} = \frac{\chi^{jj}}{\chi^{nj}}$	$S^{jj} = \sum n_{\alpha}(\mathbf{r}_1) n_{\alpha}(\mathbf{r}_2) \Omega_{\alpha} \Omega_{\alpha}$ $\chi^{jj} = \frac{S^{jj}}{i\omega - \Omega^{jj}}$
--	--	--

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}$$

$$\Omega^{nn(1)} = \frac{S^{nj}}{S^{nn}}$$

$$\Omega^{nn} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \Omega^{nn}}{i\omega - \Omega^{nj}}}{i\omega - \frac{S^{jj}}{S^{nj}} \frac{i\omega - \Omega^{nj}}{i\omega - \Omega^{jj}}}$$

$$\Omega^{nn(2)} = \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}}}{i\omega - \frac{S^{jj}}{S^{nj}}}$$

...

...

- ▶ dépendance explicite à la fréquence  $\omega$

# Approximations de $\chi^{nn}$

- **hierarchie d'expression** pour  $\chi^{nn}$
- uniquement en fonction de moyennes sur l'état fondamental
- dépendance à la fréquence

$$\chi^{nn} = \frac{S^{nn}}{i\omega - \Omega^{nn}}$$

$$\chi^{nn(1)} = \frac{S^{nn}}{i\omega - \frac{S^{nj}}{S^{nn}}}$$

$$\chi^{nn(2)} = \frac{S^{nn}}{i\omega - \frac{S^{nj}}{S^{nn}} \frac{i\omega - \frac{S^{nj}}{S^{nn}}}{i\omega - \frac{S^{jj}}{S^{nj}}}}$$

## Polarisabilité

$$\alpha_{\alpha\beta}(i\omega) = 2\text{Re} \sum_{ia} \frac{\langle i|\hat{r}_\alpha|a\rangle \langle a|\hat{r}_\beta|i\rangle}{i\omega - \Omega_{ia}} \quad \alpha/\beta = x, y, z$$

$$\alpha_{\alpha\beta}(i\omega)^{(1)} = 2\text{Re} \sum_i \frac{S_{\alpha\beta,i}^{nn}}{i\omega - \frac{S_{\alpha\beta,i}^{nj}}{S_{\alpha\beta,i}^{nn}}}$$

$$S_{\alpha\beta,i}^{nn} = 4C_{i\mu}^T (\mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta)_{\mu\nu} C_{\nu i}$$

$$S_{\alpha\beta,i}^{nj} = 4C_{i\mu}^T \left( \mathbf{r}_\alpha \mathbf{S}^{-1} \mathbf{F} \mathbf{Q} \mathbf{r}_\beta \right)_{\mu\nu} C_{\nu i} \\ - 4C_{i\mu}^T \left( \mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta \mathbf{S}^{-1} \mathbf{F} \right)_{\mu\nu} C_{\nu i}$$

## Polarisabilité

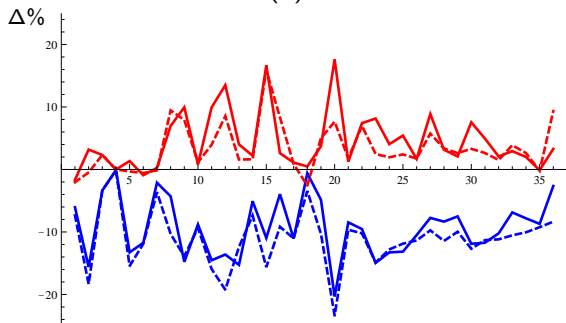
$$\alpha_{\alpha\beta}(i\omega) = 2\text{Re} \sum_{ia} \frac{\langle i|\hat{r}_\alpha|a\rangle \langle a|\hat{r}_\beta|i\rangle}{i\omega - \Omega_{ia}} \quad \alpha/\beta = x, y, z$$

$$\alpha_{\alpha\beta}(i\omega)^{(1)} = 2\text{Re} \sum_i \frac{S_{\alpha\beta,i}^{nn}}{i\omega - \frac{S_{\alpha\beta,i}^{nj}}{S_{\alpha\beta,i}^{nn}}}$$

$$S_{\alpha\beta,i}^{nn} = 4C_{i\mu}^\top (\mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta)_{\mu\nu} C_{\nu i}$$

$$S_{\alpha\beta,i}^{nj} = 4C_{i\mu}^\top \left( \mathbf{r}_\alpha \mathbf{S}^{-1} \mathbf{F} \mathbf{Q} \mathbf{r}_\beta \right)_{\mu\nu} C_{\nu i} \\ - 4C_{i\mu}^\top \left( \mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta \mathbf{S}^{-1} \mathbf{F} \right)_{\mu\nu} C_{\nu i}$$

► calcul de  $\alpha(0)$  sur 36 molécules



<span style="color: blue;">—</span>	$\alpha_{\text{moy}}^{\text{cano}(1)}$	9.3 %
<span style="color: blue;">- - -</span>	$\alpha_{\text{moy}}^{\text{loc}(1)}$	11.0 %
<span style="color: red;">—</span>	$\alpha_{\text{moy}}^{\text{cano}(2)}$	4.6 %
<span style="color: red;">- - -</span>	$\alpha_{\text{moy}}^{\text{loc}(2)}$	3.7 %

# Polarisabilité

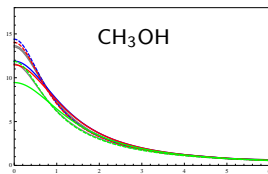
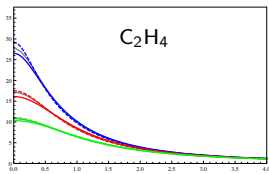
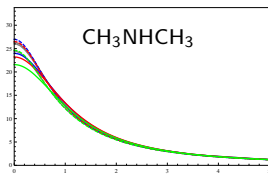
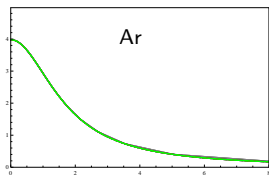
$$\alpha_{\alpha\beta}(i\omega) = 2\text{Re} \sum_{ia} \frac{\langle i|\hat{r}_\alpha|a\rangle \langle a|\hat{r}_\beta|i\rangle}{i\omega - \Omega_{ia}}$$

$\alpha/\beta = x, y, z$

$$\alpha_{\alpha\beta}(i\omega)^{(1)} = 2\text{Re} \sum_i \frac{S_{\alpha\beta,i}^{nn}}{i\omega - \frac{S_{\alpha\beta,i}^{nj}}{S_{\alpha\beta,i}^{nn}}}$$

$$S_{\alpha\beta,i}^{nn} = 4C_{i\mu}^\top (\mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta)_{\mu\nu} C_{\nu i}$$

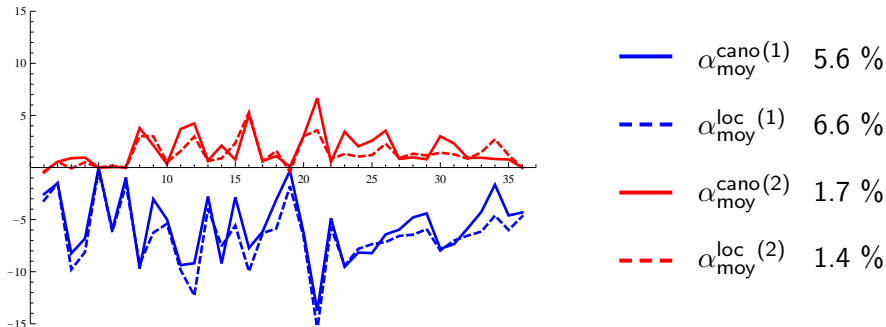
$$S_{\alpha\beta,i}^{nj} = 4C_{i\mu}^\top \left( \mathbf{r}_\alpha \mathbf{S}^{-1} \mathbf{F} \mathbf{Q} \mathbf{r}_\beta \right)_{\mu\nu} C_{\nu i} - 4C_{i\mu}^\top \left( \mathbf{r}_\alpha \mathbf{Q} \mathbf{r}_\beta \mathbf{S}^{-1} \mathbf{F} \right)_{\mu\nu} C_{\nu i}$$



—  $\alpha_{xx}^{(1)}$   
 - -  $\alpha_{xx}^{(2)}$   
 —  $\alpha_{yy}^{(1)}$   
 - -  $\alpha_{yy}^{(2)}$   
 —  $\alpha_{zz}^{(1)}$   
 - -  $\alpha_{zz}^{(2)}$

## Coefficients $C_6$

$$C_6^{ij} = \frac{3}{\pi} \int_0^\infty d\omega \bar{\alpha}^i(i\omega) \bar{\alpha}^j(i\omega)$$



- ▶ comme précédemment, ordre 1(2) sous-estime(sur-estime) systématiquement les  $C_6$
- ▶ comportement orbitales canoniques/localisées pas systématique
- ▶ nette amélioration dans la description des coefficients à l'ordre 2



- ▶ généralisation de l'approximation de Unsöld
- ▶ hierarchie d'expressions pour approximer  $\chi^{nn}$
- ▶ application au calcul de polarisabilité dynamique et de coefficients  $C_6$
- ▶ perspective du calcul de l'énergie de corrélation

formulation "matrice densité de corrélation"

formulation "matrice diélectrique"

formulation "de plasmon"

formulation "de Riccati"

énergie de  
corrélation  $E_c$

problème  
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théorème de  
fluctuation-  
dissipation

Unsöld

fonction de  
réponse  $\chi$

EED

espace réel

Gradients

Lagrangien

Géométries

Gradients analytiques d'énergies de type hybride à séparation de portée électronique (RSH) mêlant théorie de la fonctionnelle de la densité (DFT) et approximation de la phase aléatoire (RPA)

## On dérive ici les gradients analytiques d'énergies RSH+RPA

- ▶ **nouveau:** gradients analytiques RPA

depuis : gradients analytiques HF+RPA

Rekkedal, Coriani, Iozzi, Teale, **Helgaker**, Pedersen; *J. Chem. Phys.* (2013)

et : gradients analytiques PBE+dRPA(DF)

Burow, Bates, **Furche**, Eshuis; *J.C.T.C.* (Just Accepted Manuscript)

- ▶ **nouveau:** gradients analytiques d'énergies sr+lr

en fait : gradients RSH+MP2

Chabbal, **Stoll**, Werner, **Leininger**; *Mol. Phys.* (2010)

ici : dérivation “tout-en-un”

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Budapest, Hongrie

# Motivation

- ▶ obtenir forces sur les noyaux  $\frac{\partial E}{\partial \kappa}$  (géométries, états de transition)
- ▶ obtenir des constantes de forces  $\frac{\partial^2 E}{\partial \kappa_1 \partial \kappa_2}$  (fréquence de vibration)
- ▶ toute propriété monoélectronique définie comme  $\frac{\partial E}{\partial x}$  (dipôles, ...)
- ▶ gradients numériques (inefficaces, imprécis)



# Gradients analytiques

## Paramètres

$$E \doteq E[\kappa, \mathbf{V}(\kappa), \mathbf{T}(\kappa)]$$



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## Gradients

$$\frac{\partial E}{\partial \kappa} = \frac{\partial E}{\partial \kappa} + \frac{\partial E}{\partial \mathbf{V}} \frac{\partial \mathbf{V}}{\partial \kappa} + \frac{\partial E}{\partial \mathbf{T}} \frac{\partial \mathbf{T}}{\partial \kappa}$$



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$$\frac{\partial E}{\partial \mathbf{h}} \mathbf{h}^{(\kappa)} + \frac{\partial E}{\partial (\mu\nu|\sigma\rho)} (\mu\nu|\sigma\rho)^{(\kappa)} + \frac{\partial E}{\partial \mathbf{S}} \mathbf{S}^{(\kappa)}$$

# Gradients analytiques

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$$E \doteq E[\kappa, \mathbf{V}(\kappa), \mathbf{T}(\kappa)]$$



## Gradients

$$\frac{\partial E}{\partial \kappa} = \frac{\partial E}{\partial \kappa} + \cancel{\frac{\partial E}{\partial \mathbf{V}} \frac{\partial \mathbf{V}}{\partial \kappa}} + \cancel{\frac{\partial E}{\partial \mathbf{T}} \frac{\partial \mathbf{T}}{\partial \kappa}}$$

$$-\frac{\partial E}{\partial \mathbf{h}} \mathbf{h}^{(\kappa)} + \frac{\partial E}{\partial (\mu\nu|\sigma\rho)} (\mu\nu|\sigma\rho)^{(\kappa)} + \frac{\partial E}{\partial \mathbf{S}} \mathbf{S}^{(\kappa)}$$

## Méthodes variationnelles

$$\frac{\partial E_{HF}}{\partial \kappa} = \delta \underline{h_{\alpha\beta}} P_{\alpha\beta} + \frac{1}{2} \delta (\underline{\mu\lambda|\nu\sigma}) (P_{\mu\lambda} P_{\nu\sigma} - P_{\mu\sigma} P_{\nu\lambda}) + \delta \underline{S_{\mu\nu}} S_{\nu\lambda}^{-1} F_{\lambda\sigma} P_{\sigma\mu}$$

# Gradients analytiques

## Paramètres

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## Méthodes non-variationnelles

# Gradients analytiques

## Paramètres

$$E \doteq E[\kappa, \mathbf{V}(\kappa), \mathbf{T}(\kappa)]$$



## Gradients

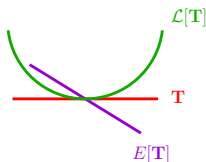
$$\frac{\partial E}{\partial \kappa} = \frac{\partial E}{\partial \kappa} + \cancel{\frac{\partial E}{\partial \mathbf{V}} \frac{\partial \mathbf{V}}{\partial \kappa}} + \frac{\partial E}{\partial \mathbf{T}} \frac{\partial \mathbf{T}}{\partial \kappa}$$

$$\frac{\partial E}{\partial \mathbf{h}} \mathbf{h}(\kappa) + \frac{\partial E}{\partial (\mu\nu|\sigma\rho)} (\mu\nu|\sigma\rho)(\kappa) + \frac{\partial E}{\partial \mathbf{S}} \mathbf{S}(\kappa)$$

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## Méthodes non-variationnelles



Travailler avec un **objet alternatif** qui **est** variationnel

# Technique du Lagrangien

**Rappel :**

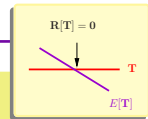
pour non-variationnelle

énergie

$E[\mathbf{V}, \mathbf{T}]$

règle pour  $\mathbf{T}$

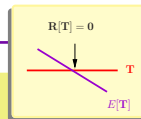
$\mathbf{R}[\mathbf{T}] = 0$



# Technique du Lagrangien

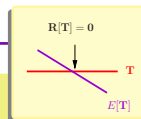
**Rappel :**  
pour non-variationnelle

énergie      règle pour  $\mathbf{T}$   
 $E[\mathbf{V}, \mathbf{T}]$        $\mathbf{R}[\mathbf{T}] = 0$



on introduit le Lagrangien

$$\mathcal{L}[\mathbf{V}, \mathbf{T}, \boldsymbol{\lambda}] = E[\mathbf{V}, \mathbf{T}] + \langle \boldsymbol{\lambda} \mathbf{R}[\mathbf{T}] \rangle$$



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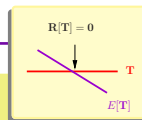
conditions stationnaires pour  $\mathcal{L}$

$$\frac{\partial \mathcal{L}}{\partial \mathbf{T}} = \frac{\partial E}{\partial \mathbf{T}} + \langle \boldsymbol{\lambda} \frac{\partial \mathbf{R}}{\partial \mathbf{T}} \rangle = 0$$

$$\frac{\partial \mathcal{L}}{\partial \boldsymbol{\lambda}} = \mathbf{R}[\mathbf{T}] = 0$$



# Technique du Lagrangien



**Rappel :**

pour non-variationnelle

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$$\mathcal{L}[\mathbf{V}, \mathbf{T}, \lambda] = E[\mathbf{V}, \mathbf{T}] + \langle \lambda \mathbf{R}[\mathbf{T}] \rangle$$

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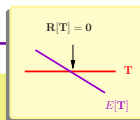
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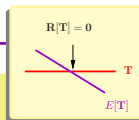
pour non-variationnelle

énergie

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$\mathbf{R}[\mathbf{T}] = 0$



on introduit le Lagrangien

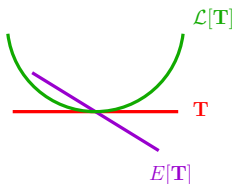
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# Énergie et Lagrangien RSH+RPA

**Rappel:**

$$E = E_{\text{RSH}} + E_c^{\text{lr}} = \langle \Phi | \hat{T} + \hat{V}_{ne} + \hat{V}_{ee}^{\text{lr}} | \Phi \rangle + E_{Hxc}^{\text{sr}}[n_\Phi] + E_c^{\text{lr}}$$

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## Notation avec fockiennes

$$\begin{aligned} E &= \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + E_c^{\text{lr}} \\ &= \langle \mathbf{d}^{(0)} \mathbf{f}^{\text{lr}} \rangle + \Delta_{\text{DC}}^{\text{lr}} + E_c^{\text{lr}} \\ &\quad + \langle \mathbf{d}^{(0)} \mathbf{g}^{\text{sr}} \rangle + \Delta_{\text{DC}}^{\text{sr}} \end{aligned}$$

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$$\mathbf{f}^{\text{lr}} = \mathbf{h} + \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}]$$

$$\Delta_{\text{DC}}^{\text{lr}} = -\frac{1}{2} \langle \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rangle$$

$$g^{\text{lr}}[\mathbf{d}^{(0)}]_{pq} = d_{rs}^{(0)} \left( (pq|rs)^{\text{lr}} - \frac{1}{2} (pr|qs)^{\text{lr}} \right)$$

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$$\mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}]_{pq} = d_{rs}^{(0)} ((pq|rs)^{\text{lr}} - \frac{1}{2} (pr|qs)^{\text{lr}})$$

$$E_{\text{Hxc}}^{\text{sr}}[n] = \int dr F[\xi]$$

$$\xi = \{\xi_A\} = \{n, n_\alpha, \nabla n_\alpha, \dots\}$$

$$\mathbf{g}_{ab}^{\text{sr}} = \int dr \sum_A \frac{\partial F}{\partial \xi_A} \frac{\partial \xi_A}{\partial d_{ab}^{(0)}}$$

$$\Delta_{\text{DC}}^{\text{sr}} = E_{\text{Hxc}}^{\text{sr}}[n] - \langle \mathbf{d}^{(0)} \mathbf{g}^{\text{sr}} \rangle$$

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## Lagrangien RSH+RPA (sr+lr)

- ▶ 2 paramètres non-varia. : amplitudes  $\mathbf{T}$ , coefficients orbitaux  $\mathbf{C}$
- ▶ 3 contraintes:  $\mathbf{R}[\mathbf{T}, \mathbf{C}] = 0$ ,  $(\mathbf{f})_{ai} = 0$ ,  $(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) = 0$



# Énergie et Lagrangien RSH+RPA

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$$\mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}]_{pq} = d_{rs}^{(0)} \left( (pq|rs)^{\text{lr}} - \frac{1}{2} (pr|qs)^{\text{lr}} \right)$$

$$\begin{aligned} E_{Hxc}^{\text{sr}}[n] &= \int dr F[\xi] \quad \xi = \{\xi_A\} = \{n, n_\alpha, \nabla n_\alpha, \dots\} \\ \mathbf{g}_{ab}^{\text{sr}} &= \int dr \sum_A \frac{\partial F}{\partial \xi_A} \frac{\partial \xi_A}{\partial d_{ab}^{(0)}} \\ \Delta_{\text{DC}}^{\text{sr}} &= E_{Hxc}^{\text{sr}}[n] - \langle \mathbf{d}^{(0)} \mathbf{g}^{\text{sr}} \rangle \end{aligned}$$

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$$\mathcal{L}[\mathbf{T}, \lambda, \mathbf{C}, \mathbf{z}, \mathbf{x}] = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + E_c^{\text{lr}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

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par rapport à  $\mathbf{T}$

$$\frac{\partial}{\partial \mathbf{T}} \left( \langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\boldsymbol{\varepsilon}, \mathbf{T}]_+) \rangle \right) = 0$$

$$-\mathbf{P} = \mathbf{Q}[\mathbf{T}] \lambda + \lambda \mathbf{Q}[\mathbf{T}]^T$$

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par rapport à  $\mathbf{T}$

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par rapport à  $\mathbf{C}$

$$\langle \mathbf{K}(\mathbf{T} + \lambda + [\lambda, \mathbf{T}]_+ + \mathbf{T} \lambda \mathbf{T}) \rangle + \langle \epsilon[\lambda, \mathbf{T}]_+ \rangle \doteq \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{d}_\lambda^{(2)} \mathbf{f} \rangle$$

# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

par rapport à  $\mathbf{T}$   $\frac{\partial}{\partial \mathbf{T}} (\langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\epsilon, \mathbf{T}]_+) \rangle) = 0$   
 $-\mathbf{P} = \mathbf{Q}[\mathbf{T}] \lambda + \lambda \mathbf{Q}[\mathbf{T}]^T$

par rapport à  $\mathbf{C}$   $\langle \mathbf{K}(\mathbf{T} + \lambda + [\lambda, \mathbf{T}]_+ + \mathbf{T} \lambda \mathbf{T}) \rangle + \langle \epsilon[\lambda, \mathbf{T}]_+ \rangle \doteq \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{d}_\lambda^{(2)} \mathbf{f} \rangle$   
 $\mathcal{L} = \langle (\mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z}) \mathbf{f} \rangle + \Delta_{\text{DC}} + \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle$

# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

par rapport à  $\mathbf{T}$

$$\frac{\partial}{\partial \mathbf{T}} \left( \langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\epsilon, \mathbf{T}]_+) \rangle \right) = 0$$

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$$\langle \mathbf{K}(\mathbf{T} + \lambda + [\lambda, \mathbf{T}]_+ + \mathbf{T} \lambda \mathbf{T}) \rangle + \langle \epsilon[\lambda, \mathbf{T}]_+ \rangle \doteq \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{d}_\lambda^{(2)} \mathbf{f} \rangle$$

$$\mathcal{L} = \langle (\mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z}) \mathbf{f} \rangle + \Delta_{\text{DC}} + \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle$$

$$\sum_{kc,b} (ib|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,b} (ab|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (ij|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (aj|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

par rapport à  $\mathbf{T}$

$$\frac{\partial}{\partial \mathbf{T}} (\langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\epsilon, \mathbf{T}]_+) \rangle) = 0$$

$$-\mathbf{P} = \mathbf{Q}[\mathbf{T}] \lambda + \lambda \mathbf{Q}[\mathbf{T}]^T$$

par rapport à  $\mathbf{C}$

$$\mathcal{L} = \langle \mathbf{K}(\mathbf{T} + \lambda + [\lambda, \mathbf{T}]_+ + \mathbf{T} \lambda \mathbf{T}) \rangle + \langle \epsilon[\lambda, \mathbf{T}]_+ \rangle \doteq \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{d}_\lambda^{(2)} \mathbf{f} \rangle$$

$$\mathcal{L} = \langle (\mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z}) \mathbf{f} \rangle + \Delta_{\text{DC}} + \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle$$

$\mathbf{d} \mathbf{h} \rightarrow \mathbf{d} \mathbf{h}$

$$\mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow \mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}]$$

$$-\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}]$$

$\mathbf{d} \mathbf{f}^{\text{lr}} +$

$\mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(2)} + \mathbf{z}]$

$$\sum_{kc,b} (ib|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,b} (ab|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (ij|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (aj|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

par rapport à  $\mathbf{T}$

$$\frac{\partial}{\partial \mathbf{T}} (\langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\epsilon, \mathbf{T}]_+) \rangle) = 0$$

$$-\mathbf{P} = \mathbf{Q}[\mathbf{T}] \lambda + \lambda \mathbf{Q}[\mathbf{T}]^T$$

par rapport à  $\mathbf{C}$

$$\mathcal{L} = \langle \mathbf{K}(\mathbf{T} + \lambda + [\lambda, \mathbf{T}]_+ + \mathbf{T} \lambda \mathbf{T}) \rangle + \langle \epsilon[\lambda, \mathbf{T}]_+ \rangle \doteq \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{d}_\lambda^{(2)} \mathbf{f} \rangle$$

$$\mathcal{L} = \langle (\mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z}) \mathbf{f} \rangle + \Delta_{\text{DC}} + \langle \mathbf{K} \mathbf{M}_\lambda \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle$$

$\mathbf{d} \mathbf{h} \rightarrow \mathbf{d} \mathbf{h}$

$$\mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow \mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}]$$

$$-\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}]$$

$$\mathbf{d} \mathbf{g}^{\text{sr}} \rightarrow \mathbf{d} \mathbf{g}^{\text{sr}} + \mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}]$$

$$\Delta_{\text{DC}}^{\text{sr}} \rightarrow -\mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}^{(0)}]$$

$$\mathbf{d} \mathbf{f}^{\text{lr}} +$$

$$\mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(2)} + \mathbf{z}]$$

$$\mathbf{d} \mathbf{g}^{\text{sr}} +$$

$$\mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}^{(2)} + \mathbf{z}]$$

$$\sum_{kc,b} (ib|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,b} (ab|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (ij|kc) (\mathbf{M}_\lambda)_{kc,jb}$$

$$\sum_{kc,j} (aj|kc) (\mathbf{M}_\lambda)_{kc,jb}$$



# Conditions Stationnaires

$$\mathcal{L} = \langle \mathbf{d}^{(0)} \mathbf{f} \rangle + \Delta_{\text{DC}} + \underline{E_c^{\text{lr}}} + \langle \lambda \mathbf{R}[\mathbf{T}, \mathbf{C}] \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle + \langle \mathbf{z} \mathbf{f} \rangle$$

par rapport à  $\mathbf{T}$

$$\frac{\partial}{\partial \mathbf{T}} \left( \langle \mathbf{K} \mathbf{T} \rangle + \langle \lambda (\mathbf{K} + [\mathbf{K}, \mathbf{T}]_+ + \mathbf{T} \mathbf{K} \mathbf{T} + [\epsilon, \mathbf{T}]_+) \rangle \right) = 0$$

$$-\mathbf{P} = \mathbf{Q}[\mathbf{T}] \lambda + \lambda \mathbf{Q}[\mathbf{T}]^T$$

par rapport à  $\mathbf{C}$

$$\mathcal{L} = \langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f} \rangle + \Delta_{\text{DC}} + \langle \mathbf{K} \mathbf{M}_{\lambda} \rangle + \langle \mathbf{x}(\mathbf{C}^T \mathbf{S} \mathbf{C} - \mathbf{1}) \rangle$$

$\mathbf{d} \mathbf{h} \rightarrow \mathbf{d} \mathbf{h}$

$$\mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow \mathbf{d} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}]$$

$$-\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] \rightarrow -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(0)}] + -\frac{1}{2} \mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(2)} + \mathbf{z}]$$

$$\mathbf{d} \mathbf{g}^{\text{sr}} \rightarrow \mathbf{d} \mathbf{g}^{\text{sr}} + \mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}]$$

$$\Delta_{\text{DC}}^{\text{sr}} \rightarrow -\mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}^{(0)}]$$

$\mathbf{d} \mathbf{f}^{\text{lr}} +$

$\mathbf{d}^{(0)} \mathbf{g}^{\text{lr}}[\mathbf{d}^{(2)} + \mathbf{z}]$

$\mathbf{d} \mathbf{g}^{\text{sr}} +$

$\mathbf{d}^{(0)} \mathbf{W}^{\text{sr}}[\mathbf{d}^{(2)} + \mathbf{z}]$

$$\sum_{kc,b} (ib|kc) (\mathbf{M}_{\lambda})_{kc,jb}$$

$$\sum_{kc,b} (ab|kc) (\mathbf{M}_{\lambda})_{kc,jb}$$

$$\sum_{kc,j} (ij|kc) (\mathbf{M}_{\lambda})_{kc,jb}$$

$$\sum_{kc,j} (aj|kc) (\mathbf{M}_{\lambda})_{kc,jb}$$

$$\begin{cases} \left( \Theta - \Theta^T + \mathbf{f} \mathbf{z} - \mathbf{z} \mathbf{f} + 4 \mathbf{g}^{\text{lr}}(\mathbf{z}) + 4 \mathbf{W}^{\text{sr}}[\mathbf{z}] \right)_{ai} = 0 & (\text{CP-RPA}) \\ (1 + \tau_{pq}) \left( \Theta + \tilde{\Theta}(\mathbf{z}) \right)_{pq} = -4(\mathbf{x})_{pq} \end{cases}$$

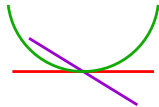
## Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr} + \text{lr}} \right\rangle + \Delta_{\text{DC}}^{\text{sr} + \text{lr}} + \langle \mathbf{K} \mathbf{M}_{\lambda} \rangle + \left\langle \mathbf{x} (\mathbf{C}^{\text{T}} \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

# Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr} + \text{lr}} \right\rangle + \Delta_{\text{DC}}^{\text{sr} + \text{lr}} + \langle \mathbf{K} \mathbf{M}_{\lambda} \rangle + \left\langle \mathbf{x} (\mathbf{C}^{\text{T}} \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

simplement la dérivée d'un objet variationnel

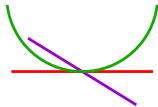


# Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr} + \text{lr}} \right\rangle + \Delta_{\text{DC}}^{\text{sr} + \text{lr}} + \langle \mathbf{K} \mathbf{M}_{\lambda} \rangle + \left\langle \mathbf{x} (\mathbf{C}^{\text{T}} \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

simplement la dérivée d'un objet variationnel

$$\mathcal{L}^{(\kappa)} =$$



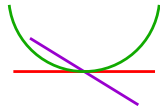
# Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr}} + \text{lr} \right\rangle + \Delta_{\text{DC}}^{\text{sr}} + \text{lr} + \langle \mathbf{K} \mathbf{M}_{\lambda} \rangle + \left\langle \mathbf{x} (\mathbf{C}^{\text{T}} \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

simplement la dérivée d'un objet variationnel

$$\mathcal{L}^{(\kappa)} = D_{\mu\nu}^1 H_{\mu\nu}^{(\kappa)} + D_{\mu\nu, \rho\sigma}^2 (\mu\nu | \rho\sigma)^{\text{lr}(\kappa)}$$

$$\begin{aligned} (\mathbf{D}^1)_{\mu\nu} &= C_{\mu\rho} \left( \mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z} \right)_{\rho q} C_{q\nu}^{\dagger} = \left( \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\nu} \\ (\mathbf{D}^2)_{\mu\nu, \sigma\rho} &= \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\nu} D_{\rho\sigma}^{(0)} - \frac{1}{2} \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\rho} D_{\nu\sigma}^{(0)} \end{aligned}$$



# Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr}} + \text{lr} \right\rangle + \Delta_{\text{DC}}^{\text{sr}} + \text{lr} + \langle \mathbf{KM}_{\lambda} \rangle + \left\langle \mathbf{x}(\mathbf{C}^{\text{T}} \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

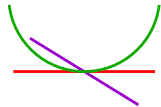
simplement la dérivée d'un objet variationnel

$$\mathcal{L}^{(\kappa)} = D_{\mu\nu}^1 H_{\mu\nu}^{(\kappa)} + D_{\mu\nu,\rho\sigma}^2 (\mu\nu|\rho\sigma)^{\text{lr}(\kappa)} + \Gamma_{\mu\nu,\rho\sigma}^2 (\mu\nu|\rho\sigma)^{\text{lr}(\kappa)}$$

$$(\mathbf{D}^1)_{\mu\nu} = C_{\mu\rho} \left( \mathbf{d}^{(0)} + \mathbf{d}_{\lambda}^{(2)} + \mathbf{z} \right)_{\rho q} C_{q\nu}^{\dagger} = \left( \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\nu}$$

$$(\mathbf{D}^2)_{\mu\nu,\sigma\rho} = \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\nu} D_{\rho\sigma}^{(0)} - \frac{1}{2} \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_{\lambda}^{(2)} + \mathbf{Z} \right)_{\mu\rho} D_{\nu\sigma}^{(0)}$$

$$(\Gamma^2)_{\mu\nu,\sigma\rho} = C_{\mu k} C_{\nu j} C_{c\rho}^{\dagger} C_{b\sigma}^{\dagger} (\mathbf{M}_{\lambda})_{ia,kc}$$



# Expression du gradient

$$\mathcal{L} = \left\langle (\mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z}) \mathbf{f}^{\text{sr} + \text{lr}} \right\rangle + \Delta_{\text{DC}}^{\text{sr} + \text{lr}} + \langle \mathbf{KM}_\lambda \rangle + \left\langle \mathbf{x}(\mathbf{C}^\top \mathbf{S} \mathbf{C} - \mathbf{1}) \right\rangle$$

simplement la dérivée d'un objet variationnel

$$\mathcal{L}^{(\kappa)} = D_{\mu\nu}^1 H_{\mu\nu}^{(\kappa)} + D_{\mu\nu,\rho\sigma}^2 (\mu\nu|\rho\sigma)^{\text{lr}(\kappa)} + \Gamma_{\mu\nu,\rho\sigma}^2 (\mu\nu|\rho\sigma)^{\text{lr}(\kappa)} + \text{SR}^{(\kappa)} + X_{\mu\nu} S_{\mu\nu}^{(\kappa)}$$

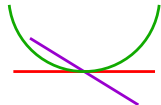
$$(\mathbf{D}^1)_{\mu\nu} = C_{\mu\rho} \left( \mathbf{d}^{(0)} + \mathbf{d}_\lambda^{(2)} + \mathbf{z} \right)_{\rho q} C_{q\nu}^\dagger = \left( \mathbf{D}^{(0)} + \mathbf{D}_\lambda^{(2)} + \mathbf{Z} \right)_{\mu\nu}$$

$$(\mathbf{D}^2)_{\mu\nu,\sigma\rho} = \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_\lambda^{(2)} + \mathbf{Z} \right)_{\mu\nu} D_{\rho\sigma}^{(0)} - \frac{1}{2} \left( \frac{1}{2} \mathbf{D}^{(0)} + \mathbf{D}_\lambda^{(2)} + \mathbf{Z} \right)_{\mu\rho} D_{\nu\sigma}^{(0)}$$

$$(\Gamma^2)_{\mu\nu,\sigma\rho} = C_{\mu k} C_{\nu j} C_{c\rho}^\dagger C_{b\sigma}^\dagger (\mathbf{M}_\lambda)_{ia,kc}$$

$$\text{SR}^{(\kappa)} = \omega_\lambda^{(\kappa)} \left( F(\xi_A) + \frac{\partial F}{\partial \xi_A} \left( \xi_A^{\mathbf{d}^{(2)}} + \xi_A^{\mathbf{z}} \right) \right)$$

$$+ \omega_\lambda \frac{\partial F}{\partial \xi_B} \left( \xi_B^{\mathbf{d}^{(0)}(x)} + \xi_B^{\mathbf{d}_\lambda^{(2)}(x)} + \xi_B^{\mathbf{z}(x)} \right) + \omega_\lambda \frac{\partial^2 F}{\partial \xi_B \partial \xi_A} \left( \xi_A^{\mathbf{d}^{(2)}} + \xi_A^{\mathbf{z}} \right) \xi_B^{(\kappa)}$$

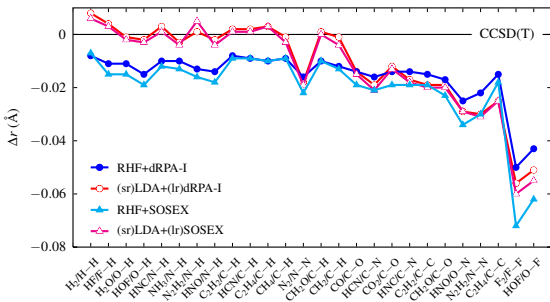


# Implémentation

- ▶ implémenté dans MOLPRO, dans le “coeur”
- ▶ utilise un parallèle avec les gradients RSH+MP2
- ▶ validation par la correspondance aux gradients numériques
- ▶ temps de calcul est double du calcul d'une énergie
- ▶ croissance en  $N^6$
- ▶ optimisation de géométrie
- ▶ densité corrélée
- ▶ dipôles
- ▶ de manière générale : meilleure convergence avec la base



# Longueurs de liaisons de simples molécules



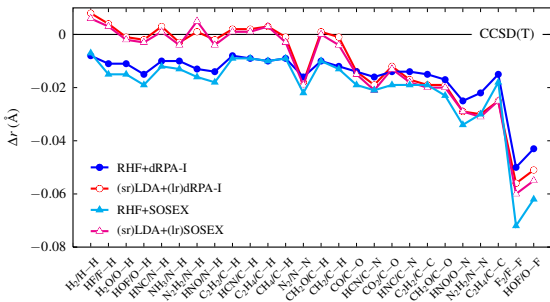
RHF+dRPA-I : 0.016

(sr)LDA+(lr)dRPA-I : 0.013

RHF+SOSEX : 0.021

(sr)LDA+(lr)SOSEX : 0.014

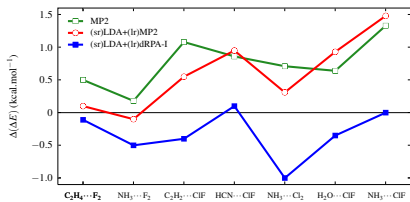
# Longueurs de liaisons de simples molécules



RHF+dRPA-I :	0.016
(sr)LDA+(lr)dRPA-I :	0.013
RHF+SOSEX :	0.021
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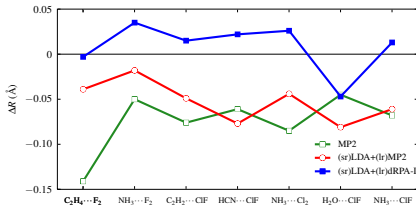
- ▶ à la limite  $\mu = 0$ , résultats sans séparation de portée sont bien reproduits
- ▶ convergence avec la base utilisée meilleure avec séparation de portée
- ▶ RSH atténue les différences de performances entre MP2, dRPA et SOSEX
- ▶ erreurs sont inférieures à 0.1 Å
- ▶ déviations moyennes sont meilleures avec la séparation de portée
- ▶ gain surtout sur les liaisons X—H
- ▶ F—X sont pathologiques

## Énergies d'interaction

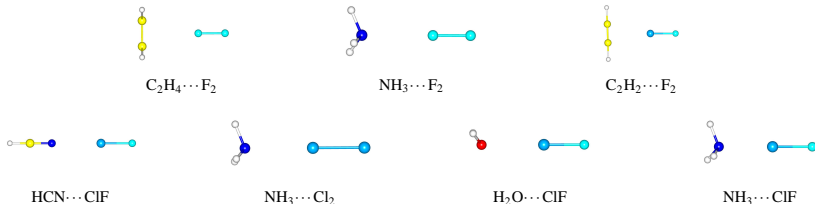


MP2 : 0.76 (sr)LDA+(lr)MP2 : 0.63  
 (sr)LDA+(lr)dRPA-I : 0.35

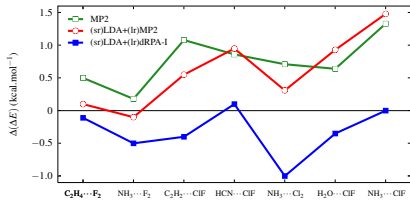
## Distances inter-monomères



MP2 : 0.075 (sr)LDA+(lr)MP2 : 0.053  
 (sr)LDA+(lr)dRPA-I : 0.023

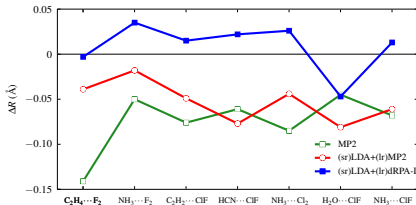


## Énergies d'interaction



MP2 : 0.76 (sr)LDA+(lr)MP2 : 0.63  
 (sr)LDA+(lr)dRPA-I : 0.35

## Distances inter-monomères



MP2 : 0.075 (sr)LDA+(lr)MP2 : 0.053  
 (sr)LDA+(lr)dRPA-I : 0.023

- ▶ après optimisation sans correction *counterpoise*
- ▶  $NH_3 \cdots F_2$  : faibles magnitudes des valeurs  
pas un cas problématique
- ▶  $H_2O \cdots ClF$  : bonne énergie d'interaction, longueur inter-monomère moyenne
- ▶  $NH_3 \cdots Cl_2$  : bonne longueur inter-monomère, énergie d'interaction moyenne  
liaisons plus déformées lors de la dimérisation (sr)LDA+(lr)dRPA-I

## Formulations RPA

- ▶ intégration numérique sur la fréquence (formulation “matrice diélectrique”)
- ▶ implémentation du *Density Fitting*

## EED

- ▶ clarifier question des orbitales
- ▶ calculs de l'énergie de corrélation

## Gradients analytiques

- ▶ implémentation additionnelle de variantes telles que SO1 et SO2 difficile pour des raisons techniques (contraction de  $\mathbf{JM}_\lambda$ )
- ▶ applications supplémentaires (densité corrélée, dipôle, ...)
- ▶ *Density Fitting*

# Remerciements

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- ▶ membres du jury
- ▶ toute l'équipe du CRM<sup>2</sup>
- ▶ le laboratoire de chimie théorique de Eötvös Loránd University
- ▶ Virginie Pichon, Faculté de Pharmacie
- ▶ Sébastien Lebègue
- ▶ János Ángyán