

# Quantum Computing Meets Condensed Matter Physics

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Alvarez Fellow Seminar  
Lawrence Berkeley National Laboratory  
04/05/2022



DMR-1752713  
PHY-1818914

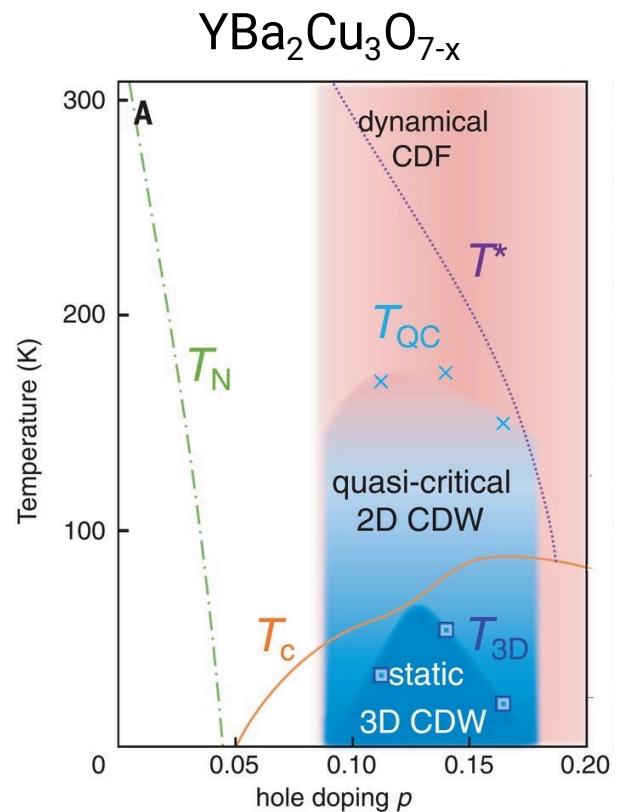


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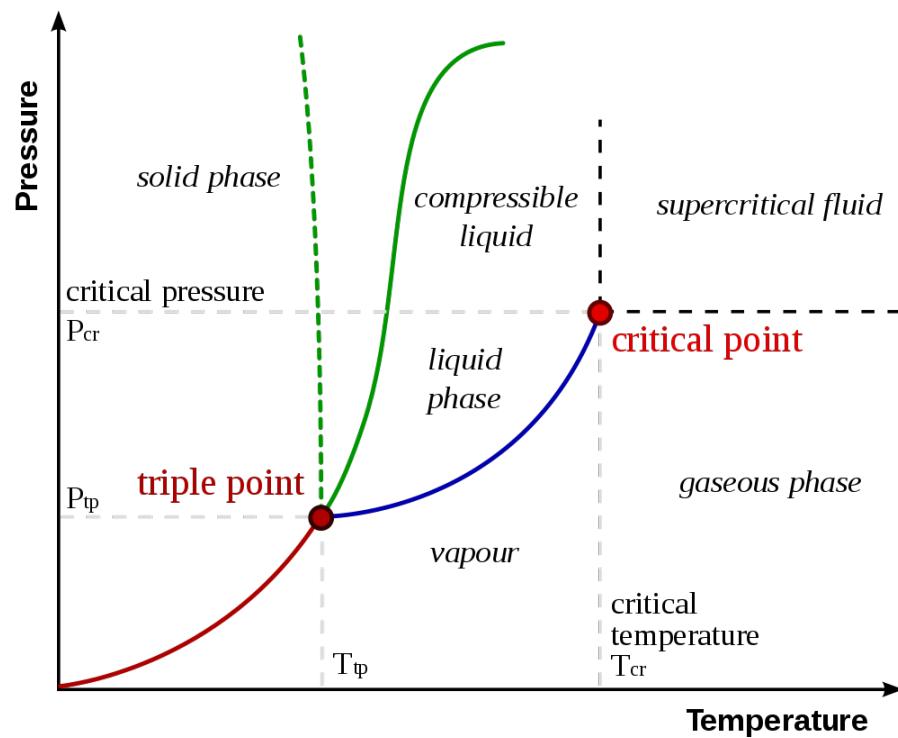
## Phase diagrams and Phase transitions

*“It’s not a cuprate talk unless it has a phase diagram...”*

-- Overheard at APS March Meeting



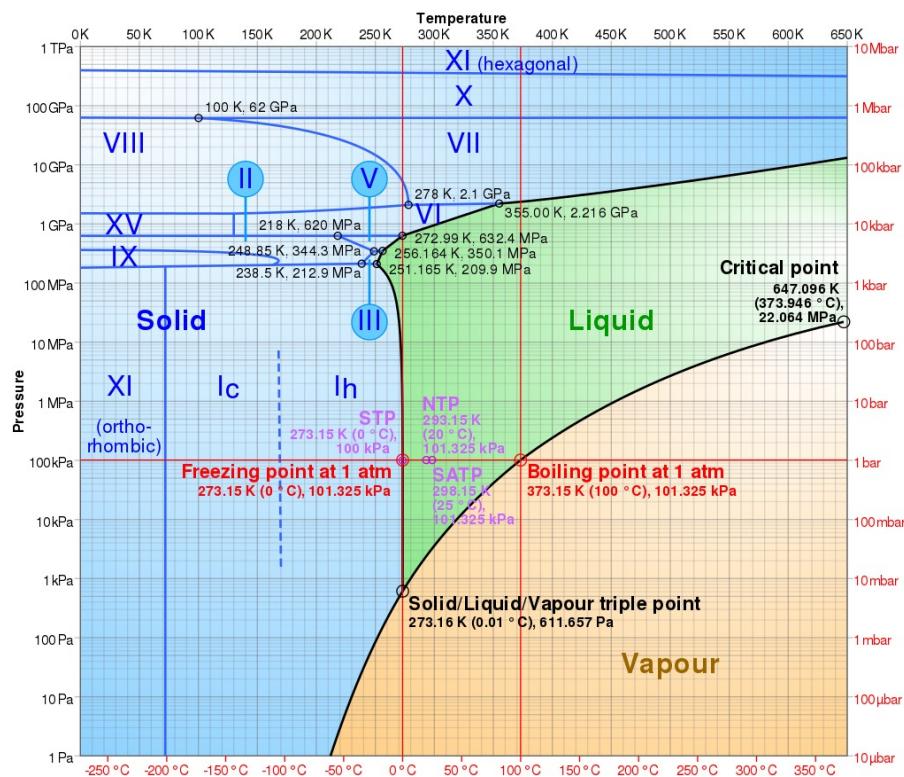
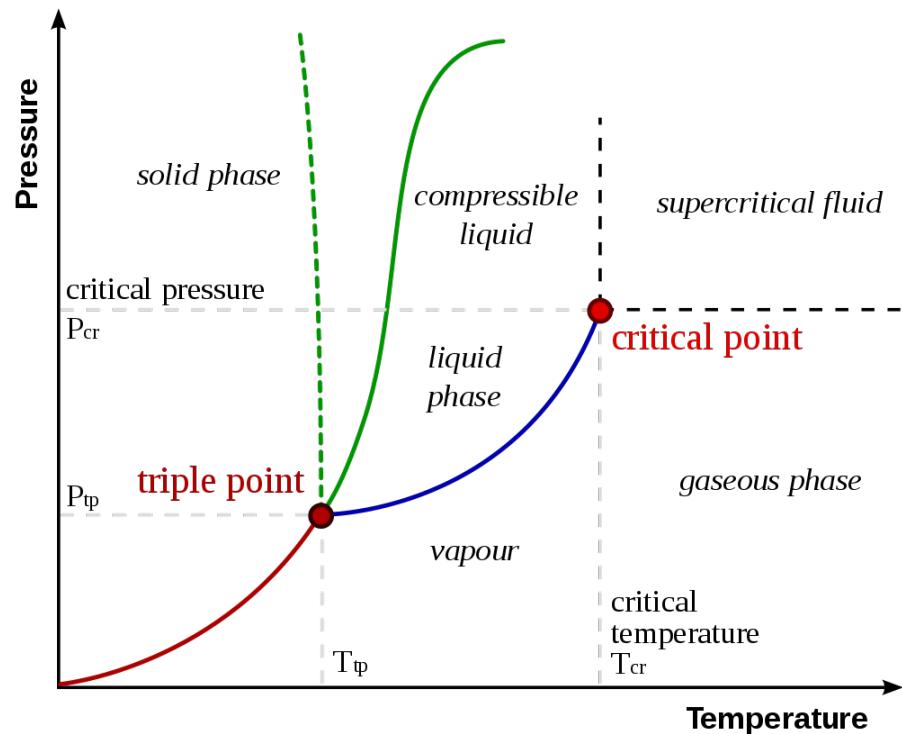
## Phase diagrams and Phase transitions



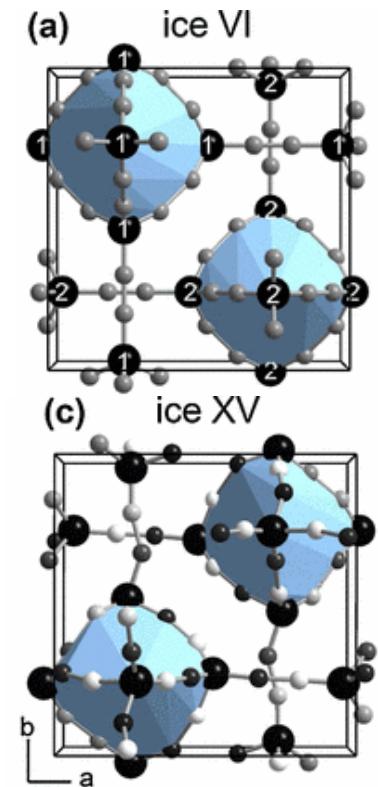
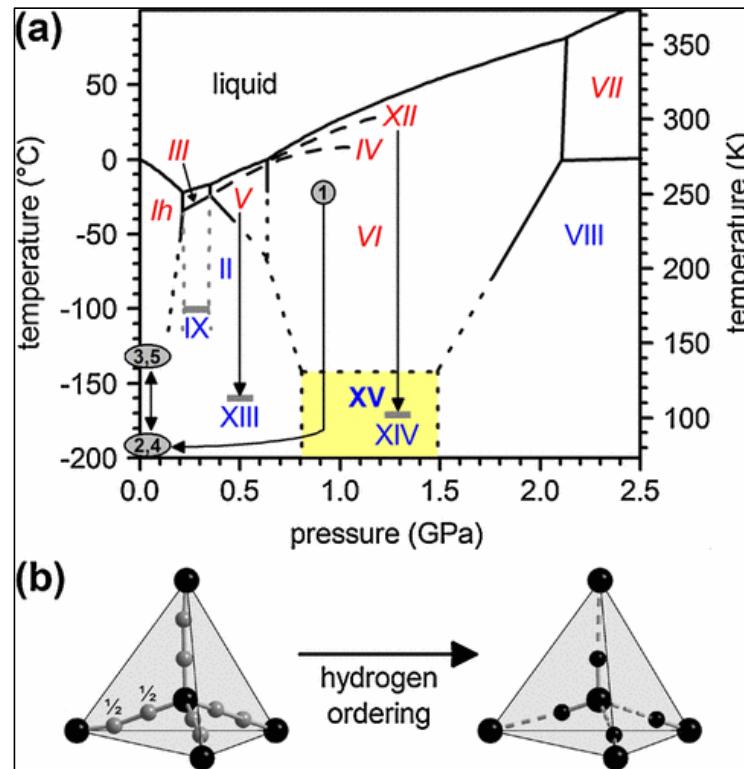
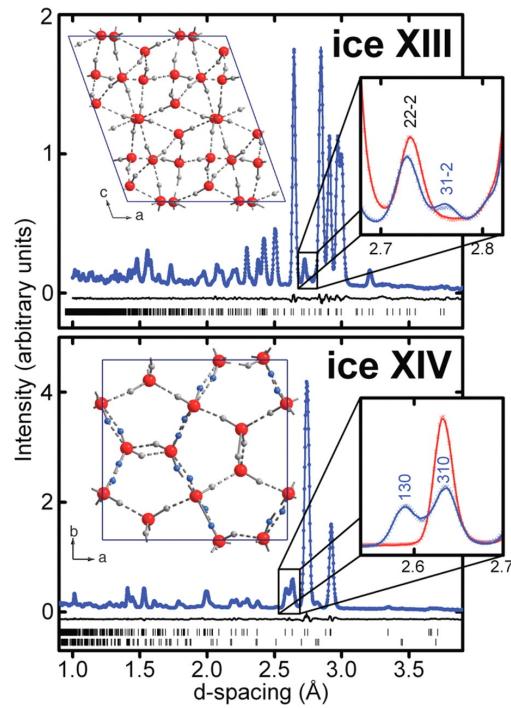
How many kinds (phases) of ice do we know about?



# Phase diagrams and Phase transitions



## Phase diagrams and Phase transitions



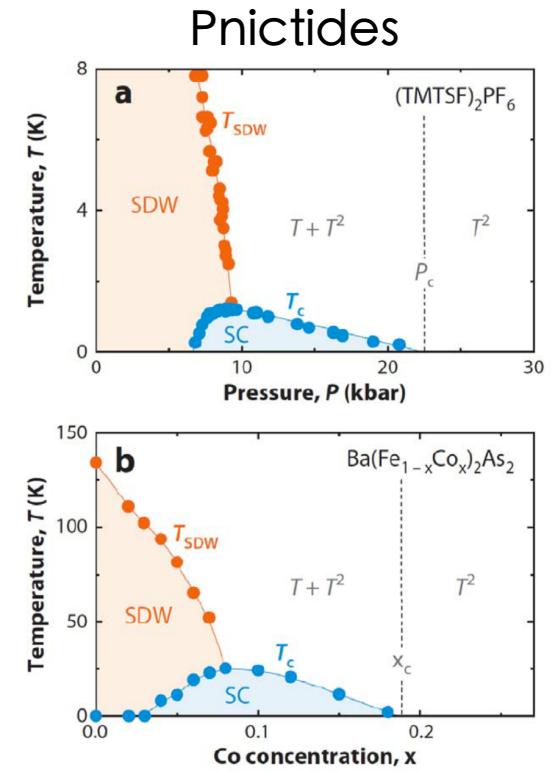
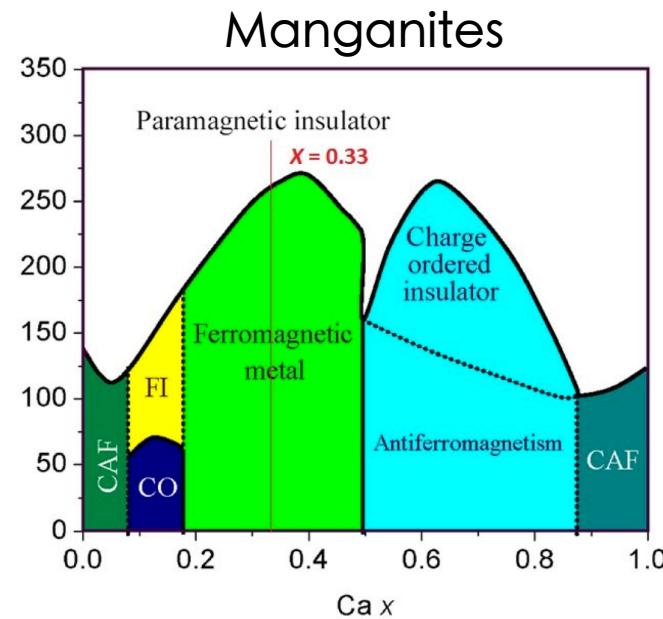
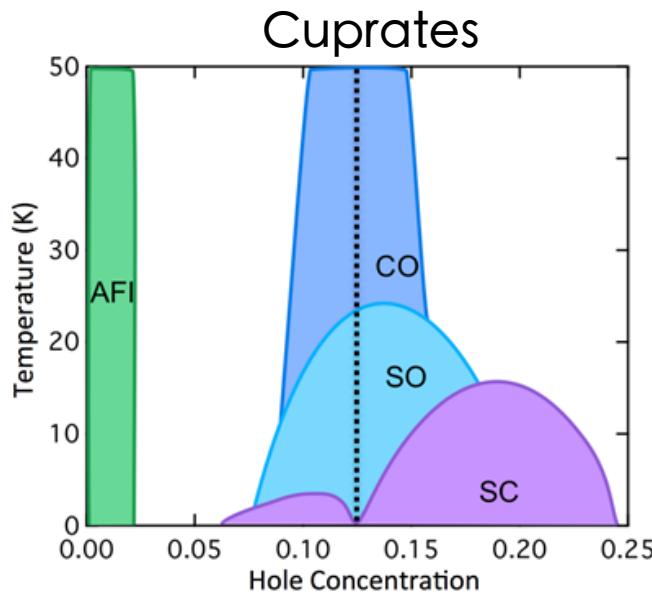
10.1126/science.1123896

10.1103/PhysRevLett.103.105701

## Phase diagrams and Phase transitions

	Group →	1	2	3	4	5	6	7	8	9	10	11	12	13	14	15	16	17	18	
↓ Period		1 H																2 He		
2		3 Li	4 Be															10 Ne		
3		11 Na	12 Mg															18 Ar		
4		19 K	20 Ca	21 Sc	22 Ti	23 V	24 Cr	25 Mn	26 Fe	27 Co	28 Ni	29 Cu	30 Zn	31 Ga	32 Ge	33 As	34 Se	36 Kr		
5		37 Rb	38 Sr	39 Y	40 Zr	41 Nb	42 Mo	43 Tc	44 Ru	45 Rh	46 Pd	47 Ag	48 Cd	49 In	50 Sn	51 Sb	52 Te	54 Xe		
6		55 Cs	56 Ba	57 La	*	72 Hf	73 Ta	74 W	75 Re	76 Os	77 Ir	78 Pt	79 Au	80 Hg	81 Tl	82 Pb	83 Bi	84 Po	86 Rn	
7		87 Fr	88 Ra	89 Ac	*	104 Rf	105 Db	106 Sg	107 Bh	108 Hs	109 Mt	110 Ds	111 Rg	112 Cn	113 Nh	114 Fl	115 Mc	116 Lv	117 Ts	118 Og
	*	58 Ce	59 Pr	60 Nd	61 Pm	62 Sm	63 Eu	64 Gd	65 Tb	66 Dy	67 Ho	68 Er	69 Tm	70 Yb	71 Lu					
	*	90 Th	91 Pa	92 U	93 Np	94 Pu	95 Am	96 Cm	97 Bk	98 Cf	99 Es	100 Fm	101 Md	102 No	103 Lr					

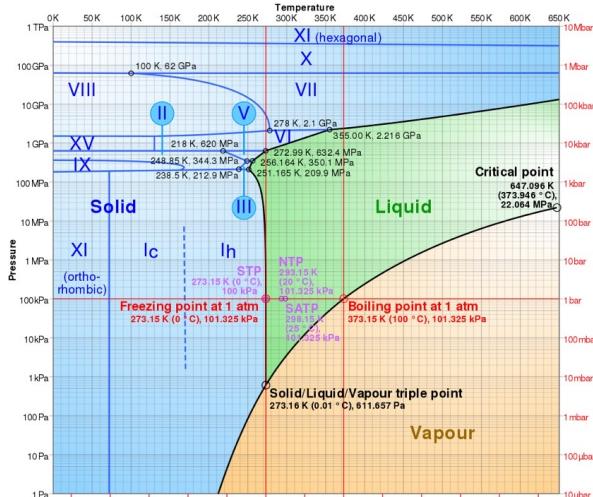
## Sample phase diagrams from “hard” condensed matter



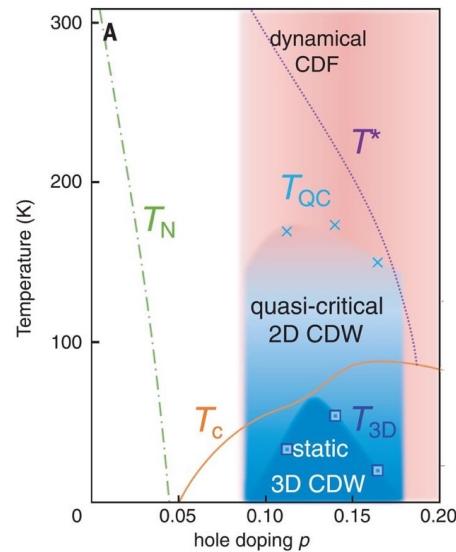
10.1146/annurev-conmatphys-070909-104117

# Why do we care about phase diagrams and transitions?

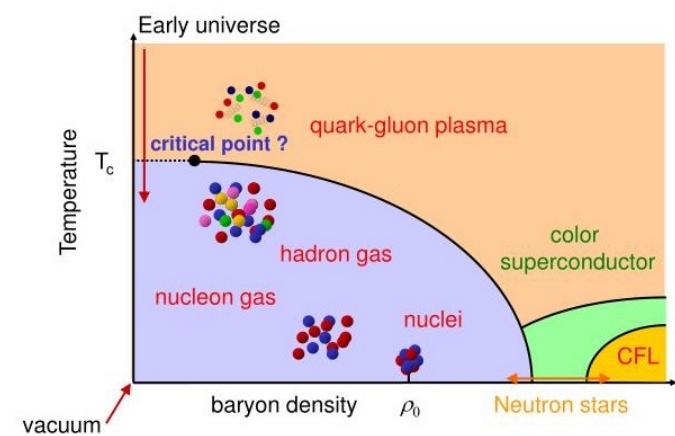
Fundamental science, of course!



Ice



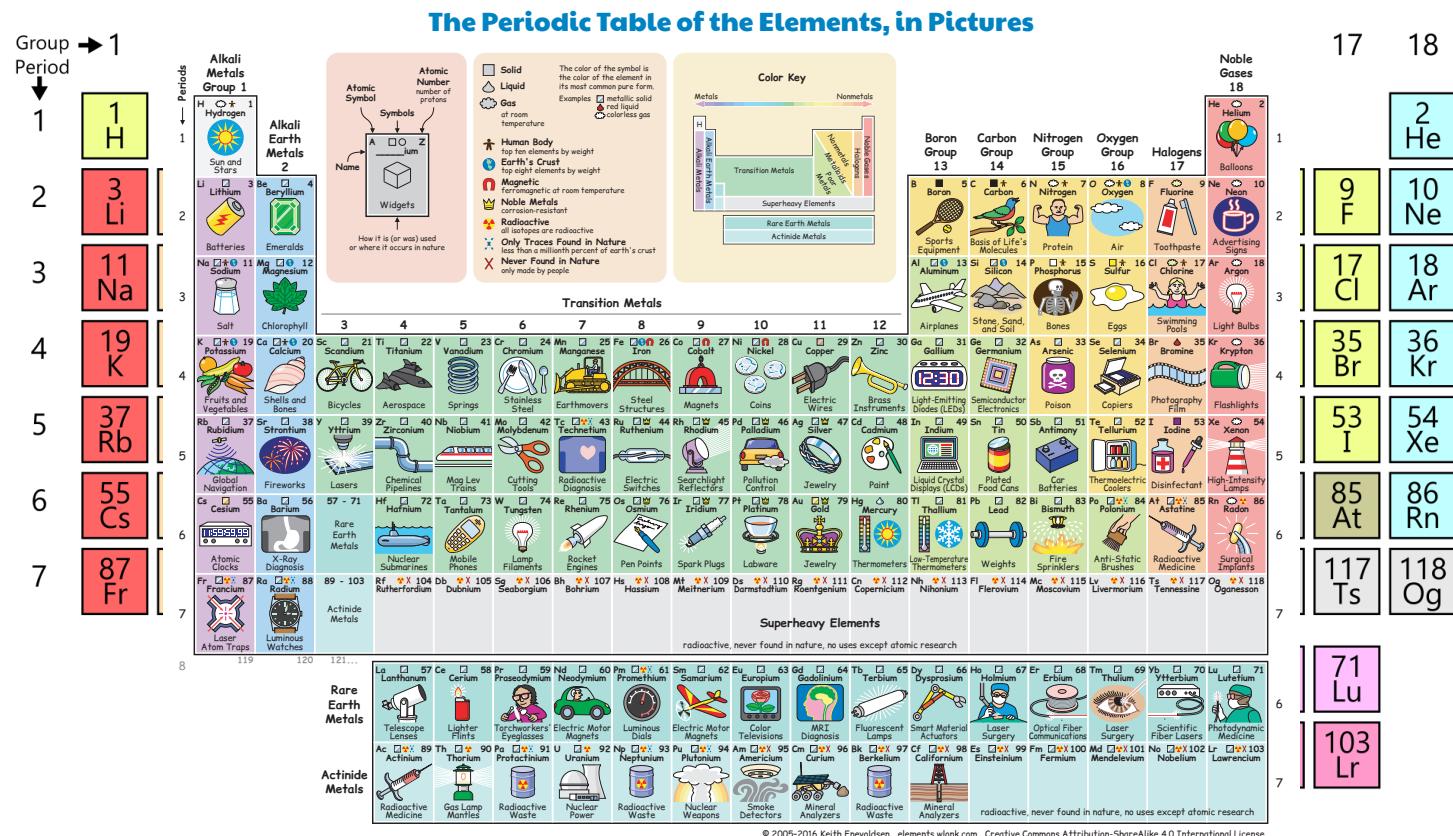
High- $T_c$  cuprate  
superconductor



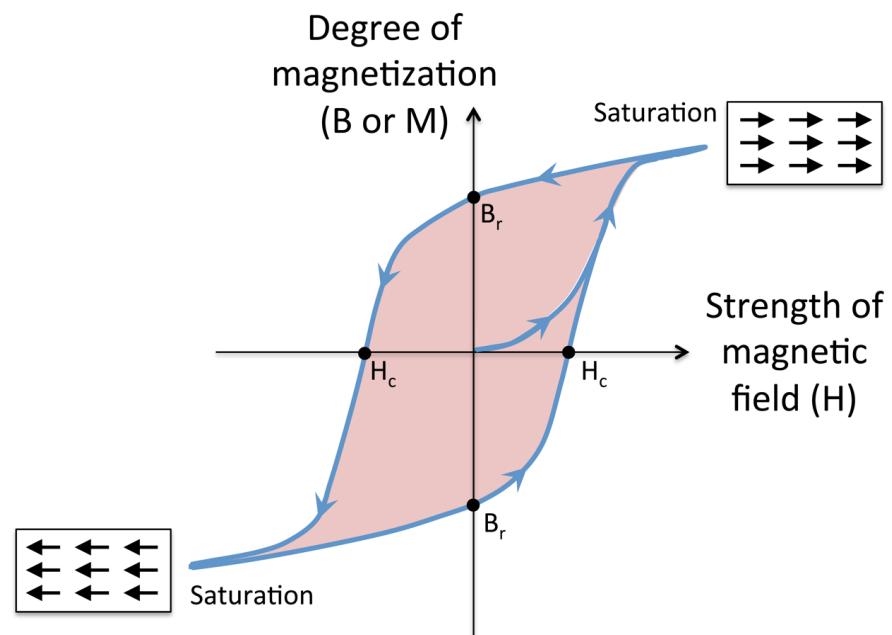
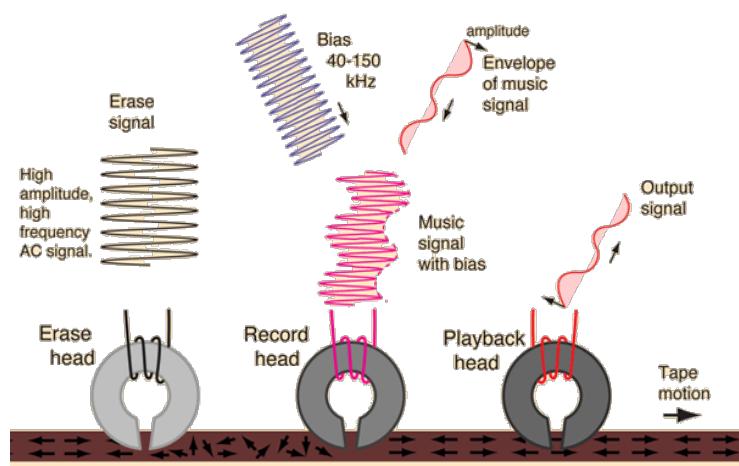
QCD

Why do we care about phase diagrams and transitions?

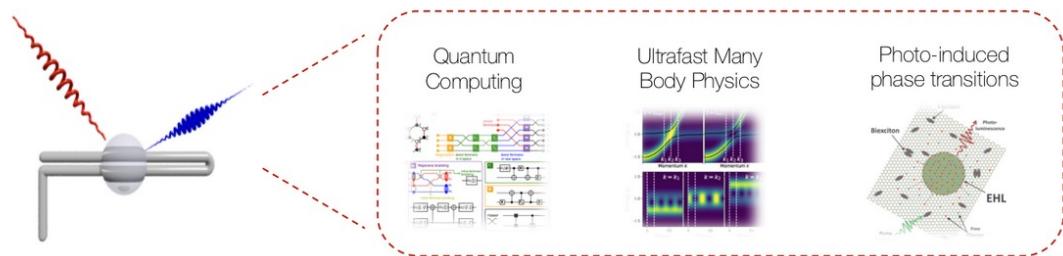
## But also...



## Why do we care about phase diagrams and transitions?



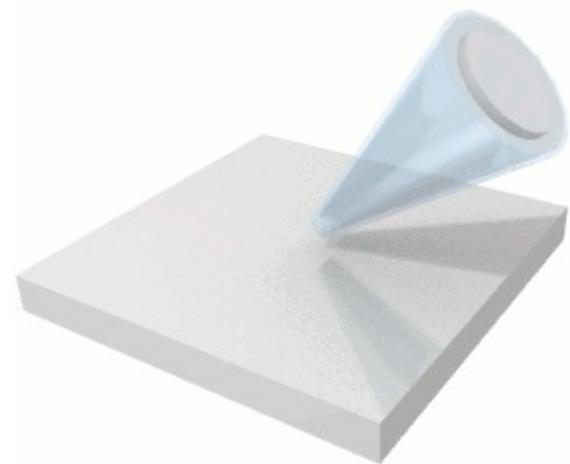
# Why quantum computing for condensed matter?



Kemper Lab

*Quantum materials in and out of equilibrium.*

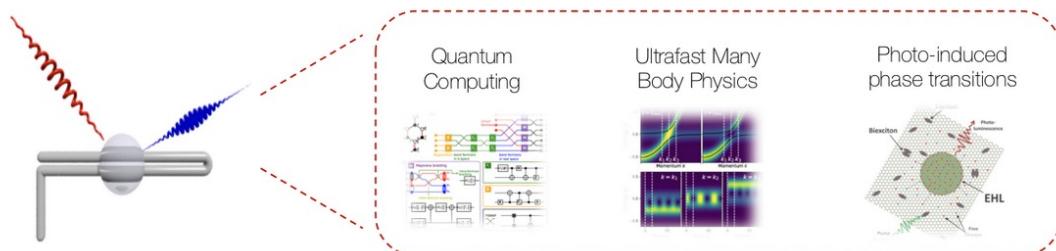
Time-resolved experiments



Shen group (Stanford)

11

# Why quantum computing for condensed matter?

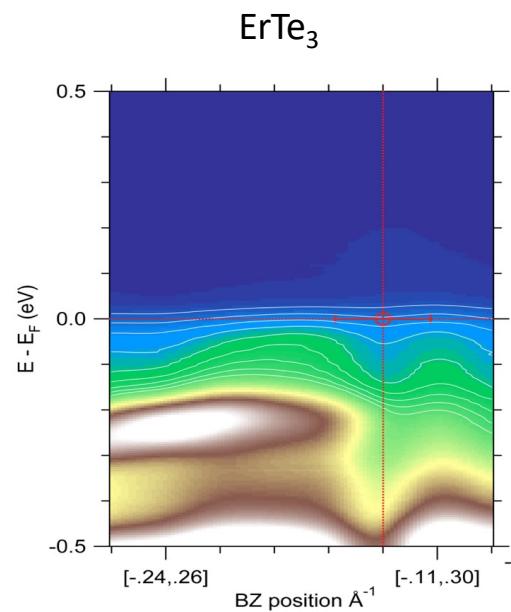


Kemper Lab

*Quantum materials in and out of equilibrium.*



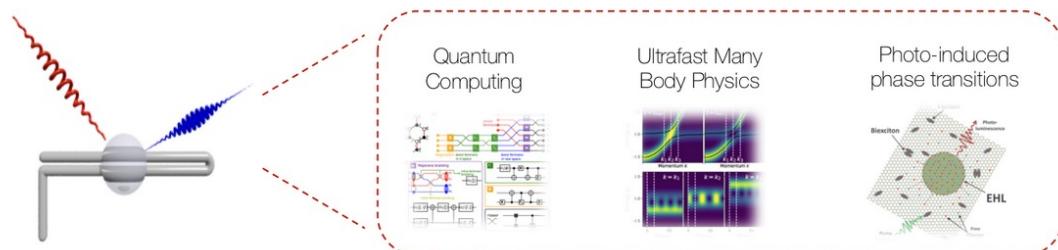
Time-resolved experiments



Shen group (Stanford)

12

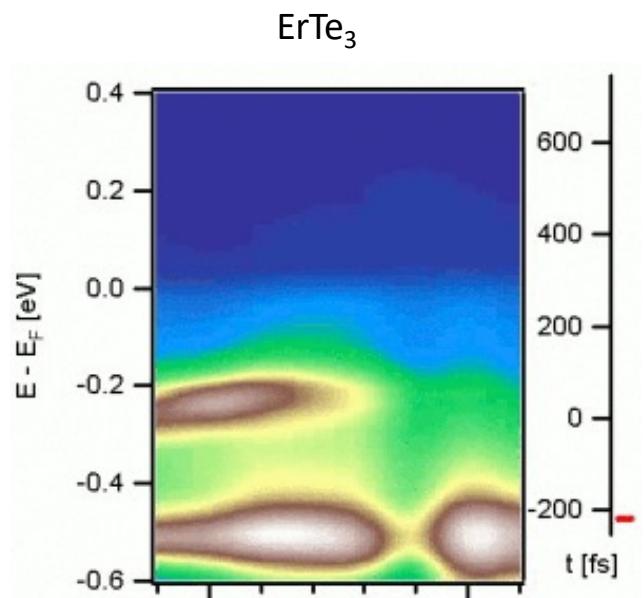
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Kemper Lab

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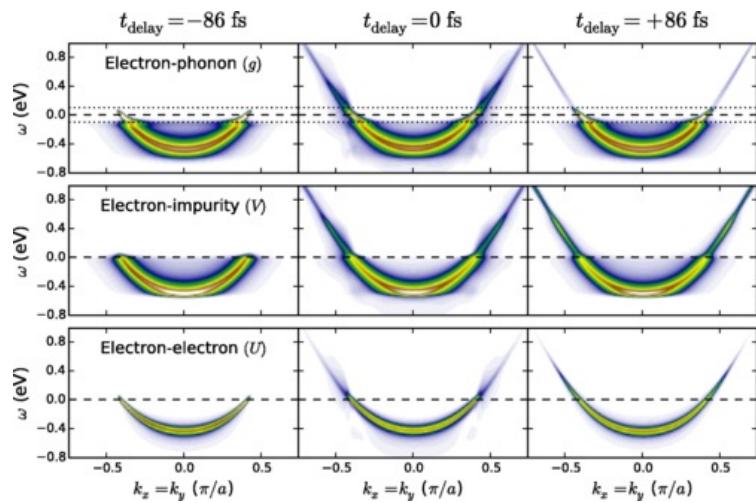
Time-resolved experiments



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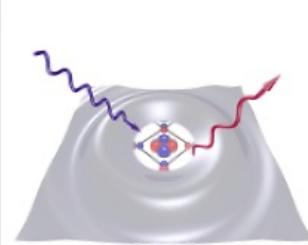
13

# Why quantum computing for condensed matter?



Non-Equilibrium Green's functions

*Phys. Rev. X 8, 041009 (2018)*



## Non-equilibrium physics of quantum materials

A. F. Kemper  
Alvarez Fellow  
Lawrence Berkeley National Laboratory

*NC State, January 14, 2015*



*This work is supported by Laboratory Directed Research and Development Program  
of Lawrence Berkeley National Laboratory under U.S. Department of Energy  
Contract No. DE-AC02-05CH11231.*

 BERKELEY LAB  
LAWRENCE BERKELEY NATIONAL LABORATORY

## Theory background 2: Equations of motion

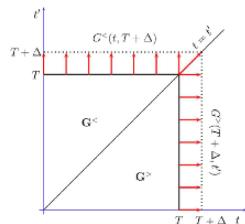
$$(i\partial_t - \varepsilon_k) G_k(t, t') = \delta(t - t') + \int d\bar{t} \Sigma(t, \bar{t}) G_k(\bar{t}, t')$$

Green's function  
for quantum state  $k$

Interaction kernel,  
depends on  $G_k$

- Non-linear initial value problem
- Results depend on the history

Time stepping method:  
[A.Stan, N.E.Dahlen, R. van Leeuwen, J.Chem.Phys.130, 224101 \(2009\)](#)

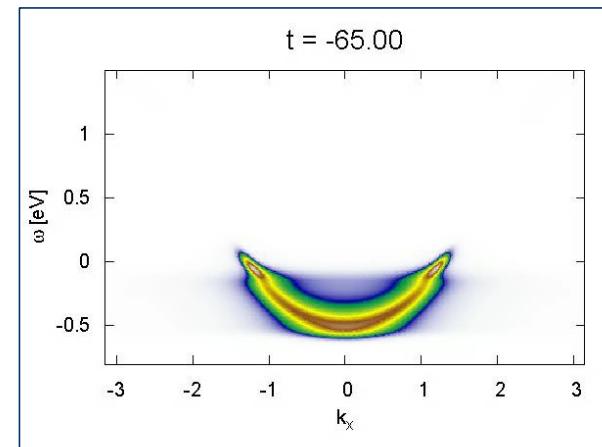
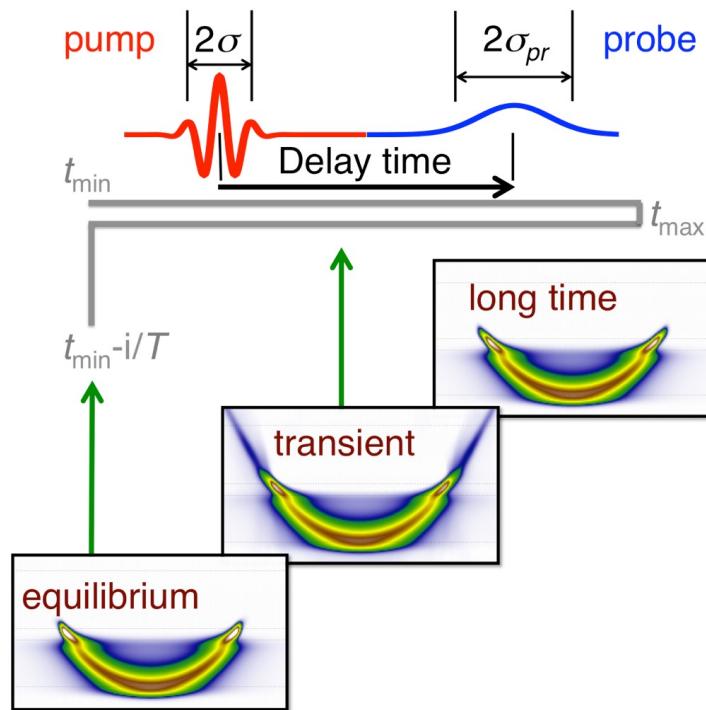


## (Some) computational details

$$(i\partial_t - \varepsilon_k) G_k(t, t') = \delta(t - t') + \int d\bar{t} \Sigma(t, \bar{t}) G_k(\bar{t}, t')$$

- 6,400 quantum  $k$  states are needed (somewhat reducible by symmetry)
- For each  $k$ , the  $G$  matrix must be stored and one copy of  $\Sigma$ 
  - Roughly 20 million unique entries
  - Entries are complex double 1x1 or 2x2 matrices (0.5 Gb or 2 Gb)
- Hybrid parallelism
  - MPI used to distribute set of  $G_k$  among MPI tasks
  - OpenMP used to parallelize integrations and other time loops

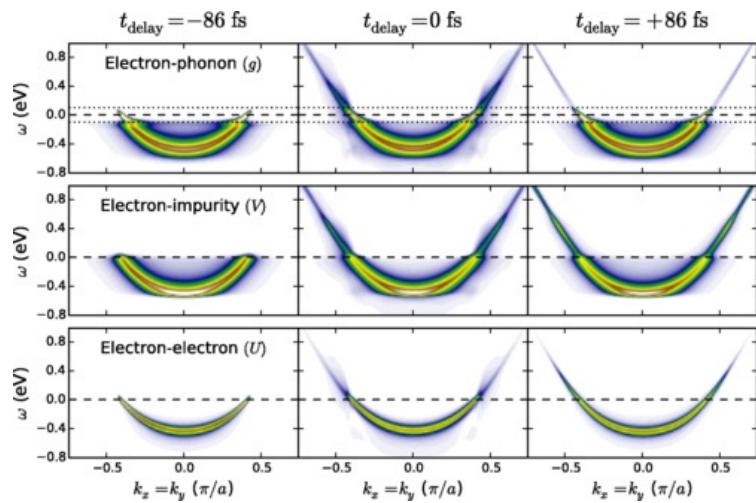
Pump-probe photoemission simulations



M. Sentef et al, PRX (2013)

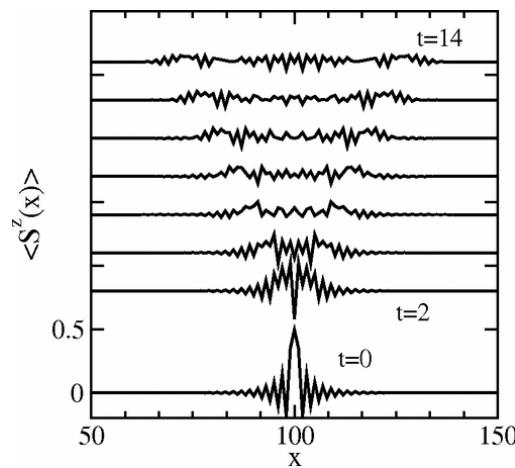
$$A_{\mathbf{k}}(\omega, t_0) = \text{Im} \frac{1}{2\pi\sigma^2} \int dt dt' G_{\mathbf{k}}^{<}(t, t') e^{-(t-t_0)^2/2\sigma^2} e^{-(t'-t_0)^2/2\sigma^2} e^{i\omega(t-t')}$$

# Why quantum computing for condensed matter?



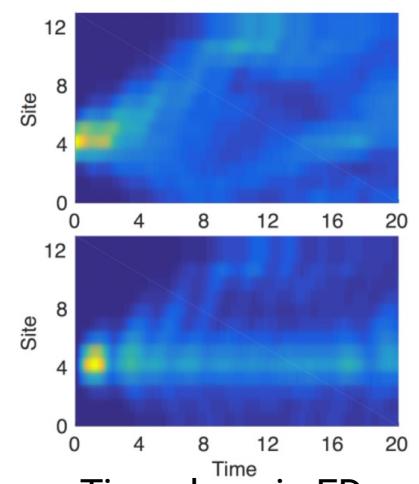
Non-Equilibrium Green's functions

*Phys. Rev. X 8, 041009 (2018)*



Time domain DMRG

*Phys. Rev. Lett. 93, 076401 (2004)*



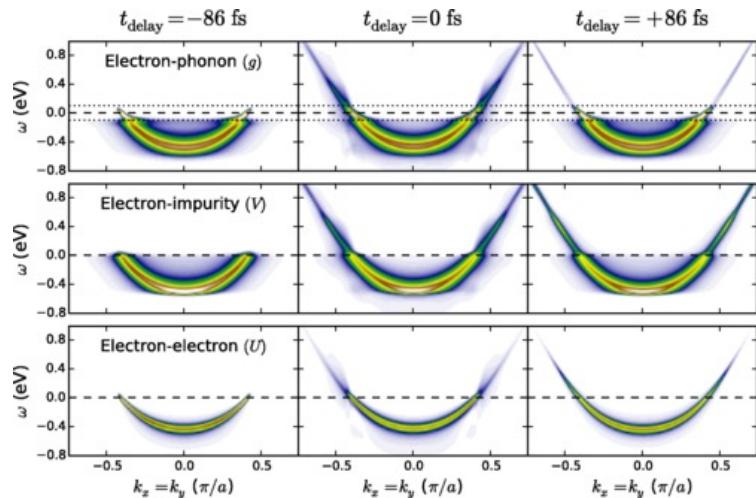
Time domain ED

Johnston & Kemper, unpublished

# Why quantum computing for condensed matter?

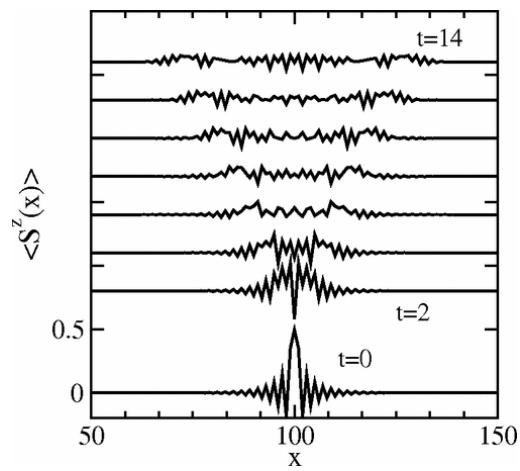


All these techniques eventually reach a barrier.



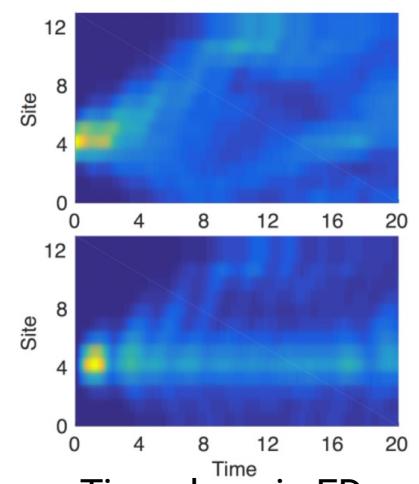
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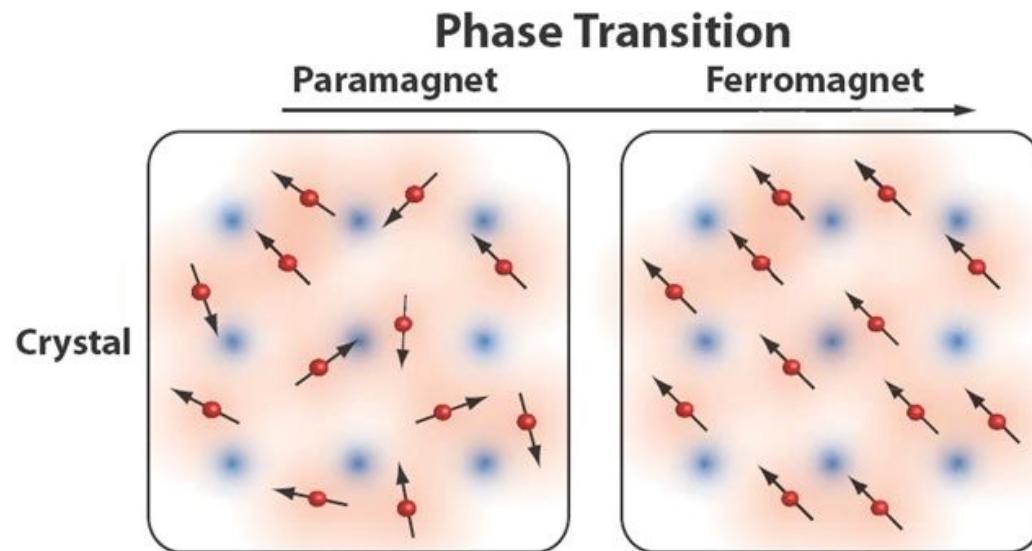
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# Why quantum computing for condensed matter?



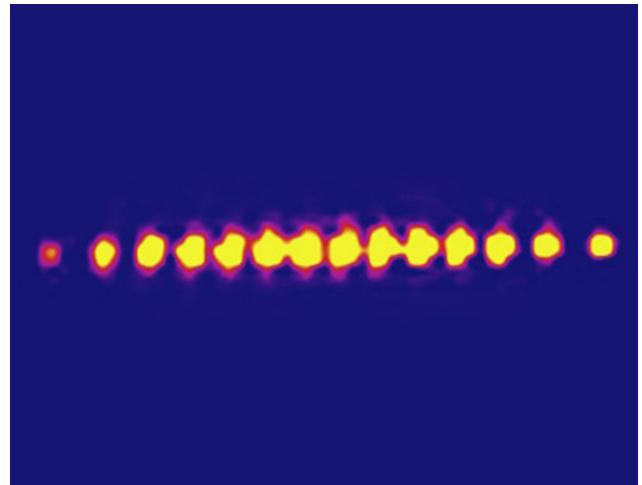
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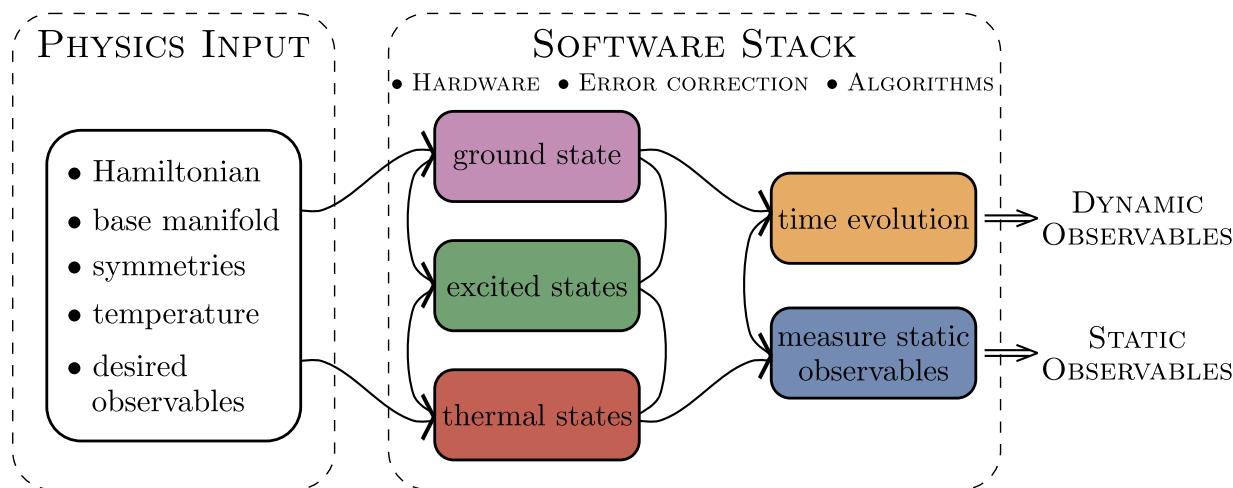
50 spins = 1,125,899,906,842,624 states

# Quantum Computers

- Superconducting (IBM, Google)
- Ion (Ionq)
- Topological? (Microsoft)
- Photonic (Xanadu)



## Quantum Matter meets Quantum Computing

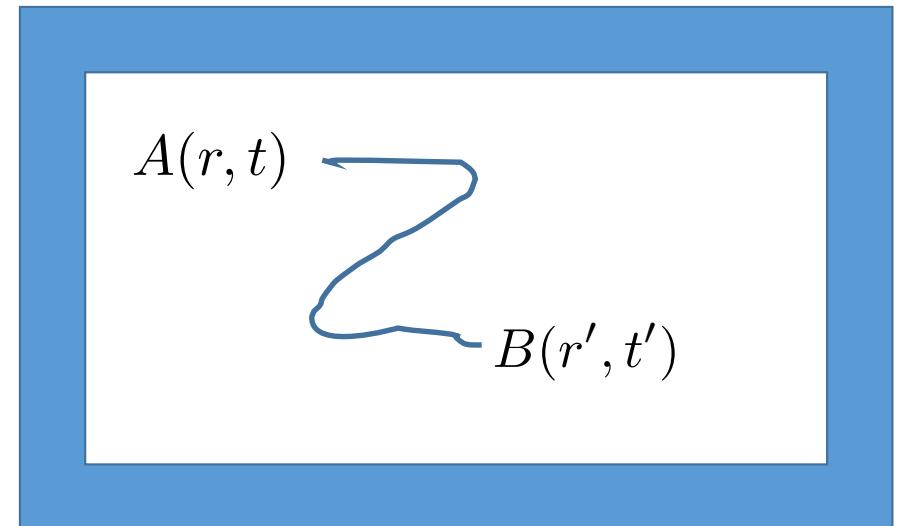


- Measuring correlation functions
- Preparing/measuring topological states
- Modeling driven/dissipative systems
- Time evolution via Lie algebraic decomposition/compression
- Thermodynamics

## Low-energy excitations: correlation functions

$$\langle A(r, t)B(r', t') \rangle$$

*Given some (observable) operator B at (r',t'), what is the likelihood of some (observable) operator A at (r,t)?*



Conductivity

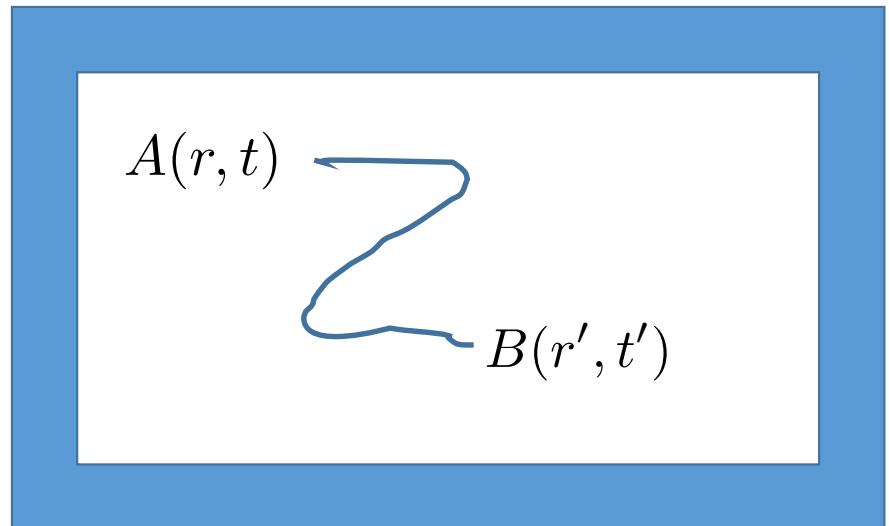
$$\langle j(r, t)j(r', t') \rangle$$

Single-particle spectra (ARPES)

$$\langle c(r, t)c^\dagger(r', t') \rangle$$

Spin-resolved neutron scattering

$$\sigma_{\alpha\beta}^{x,y,z} \langle S_\alpha(r, t)S_\beta(r', t') \rangle$$

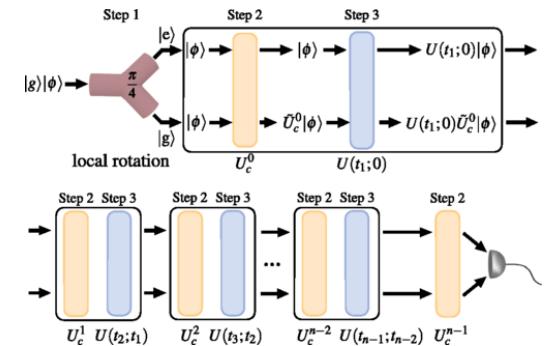
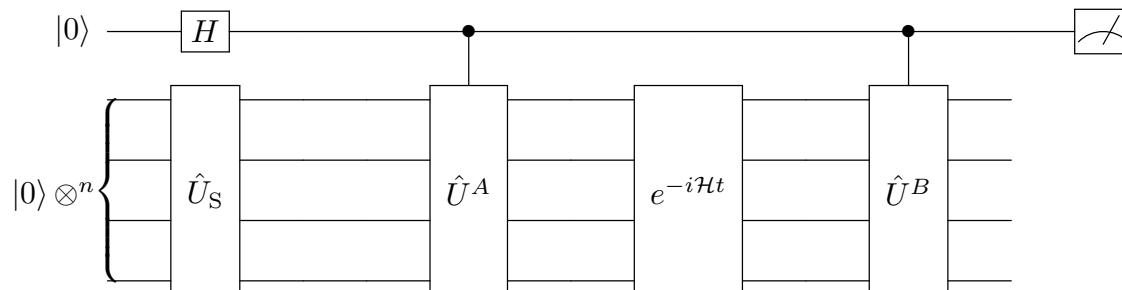


# Low-energy excitations: correlation functions

Express the correlation function through the Lehmann representation:

$$C(t) = \langle \Phi | \hat{U}^B(t) \hat{U}^A(0) | \Phi \rangle = \sum_m e^{-i(E_m - E_0)t} \langle \phi_0 | U^B | m \rangle \langle m | U^A | \phi_0 \rangle.$$

Quantum circuit:



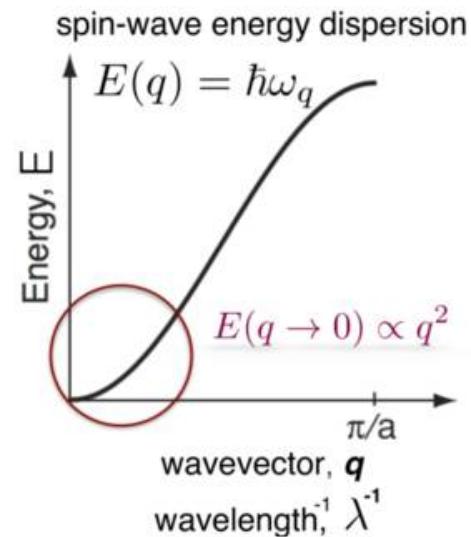
# Low-energy excitations: 2-site magnons

Spin-spin correlation function for periodic Heisenberg model: Magnons!

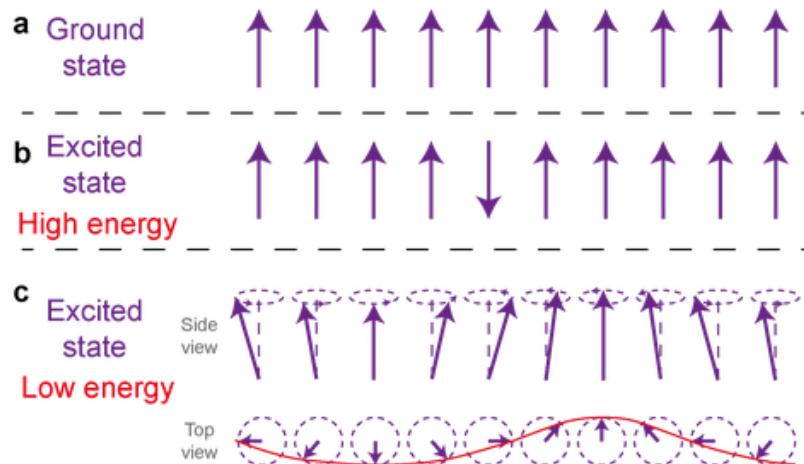
a Ground state



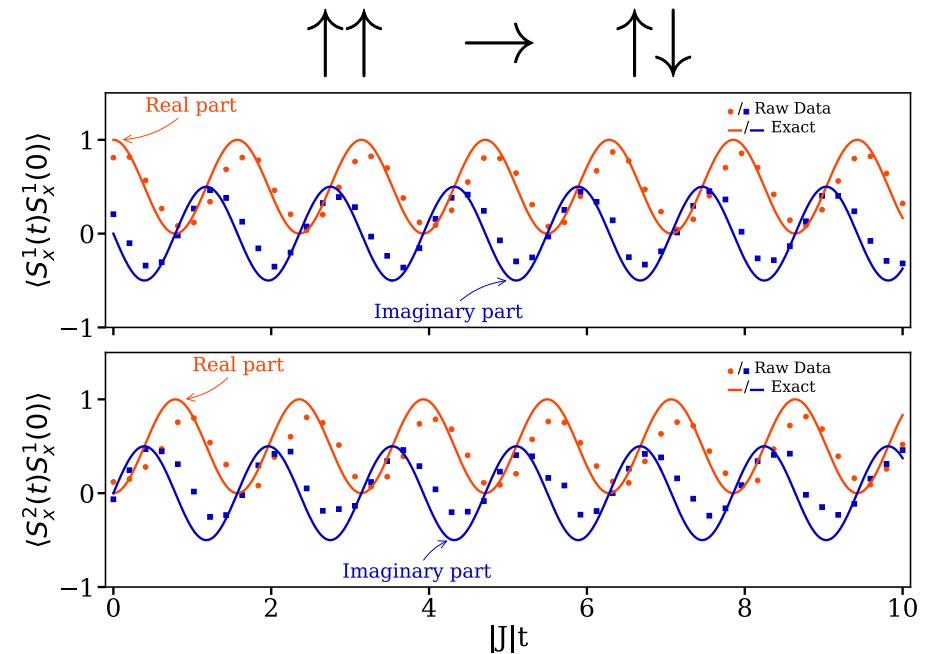
$$\mathcal{H} = -J \sum_i \vec{S}_i \cdot \vec{S}_{i+1}$$



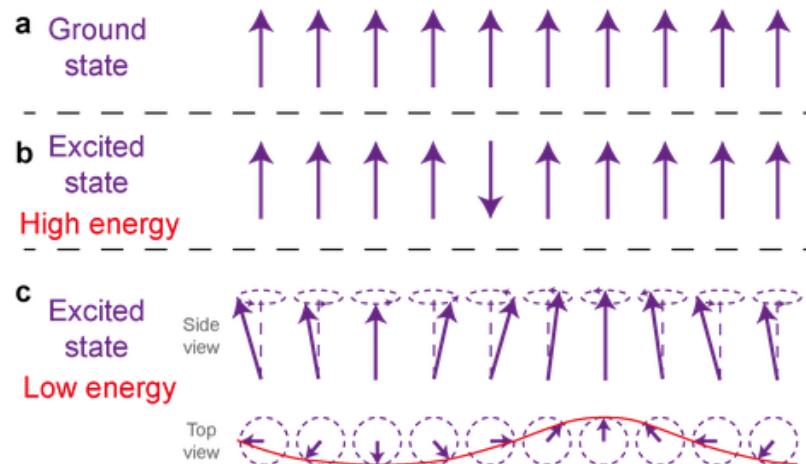
# Low-energy excitations: 2-site magnons



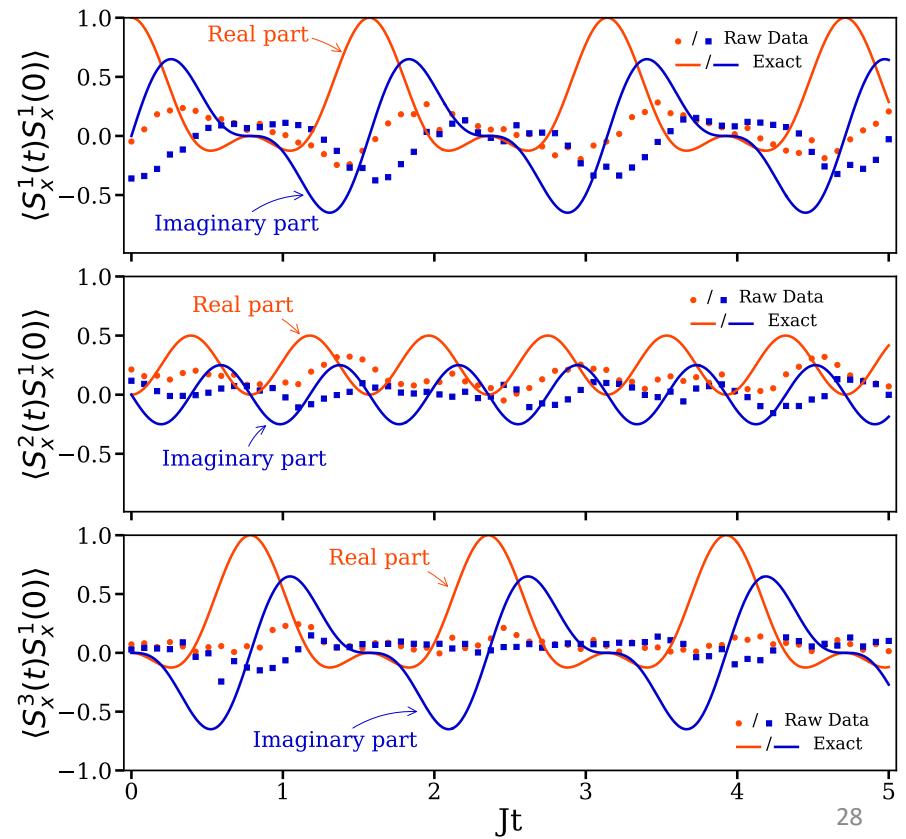
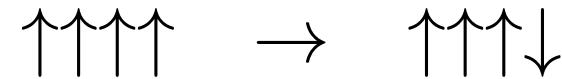
$$\hat{H} = 2SJ \sum_k (1 - \cos(k)) \hat{c}_k^\dagger \hat{c}_k$$



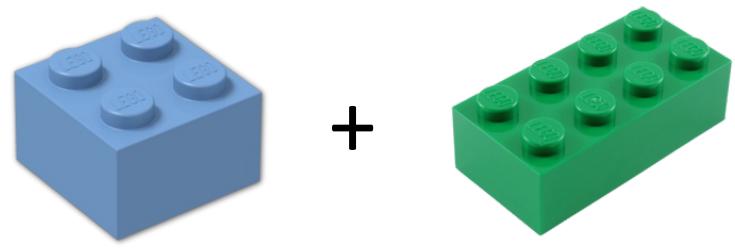
# Low-energy excitations: 4-site magnons



$$\hat{H} = 2SJ \sum_k (1 - \cos(k)) \hat{c}_k^\dagger \hat{c}_k$$



## Low-energy excitations: 4-site magnons

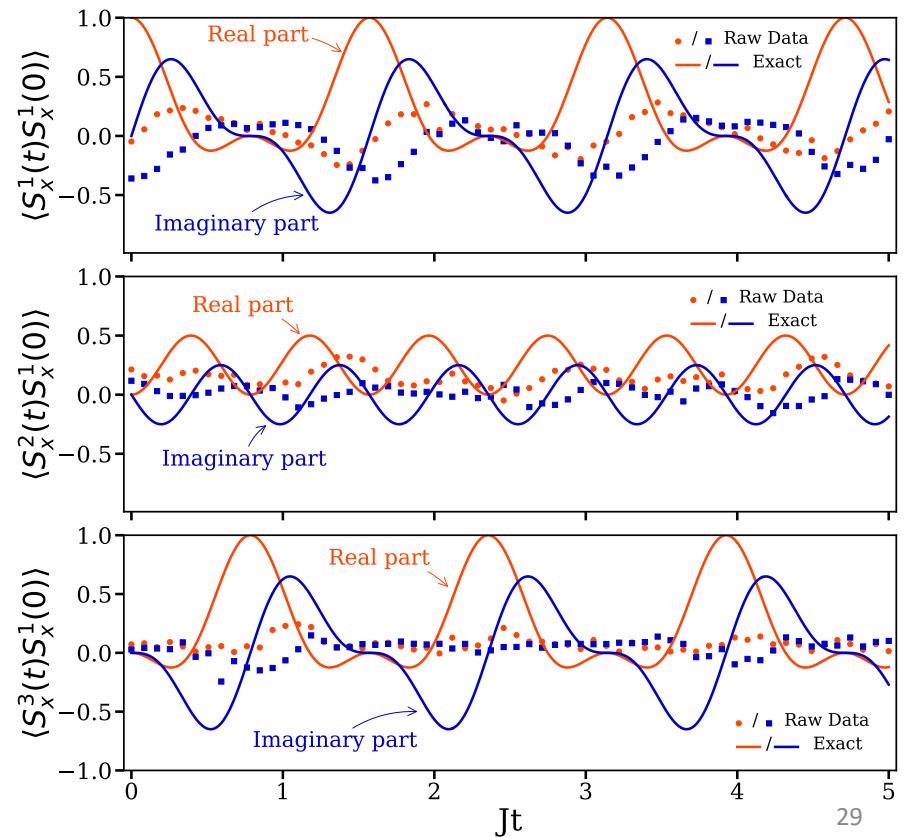
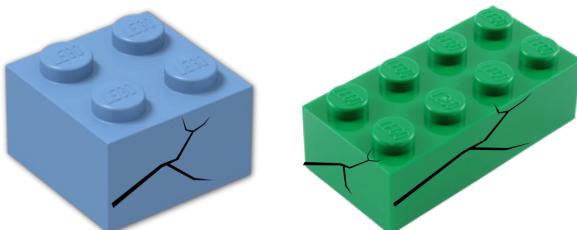


Single Qubit Gates

Two Qubit Gate

= Anything!\*

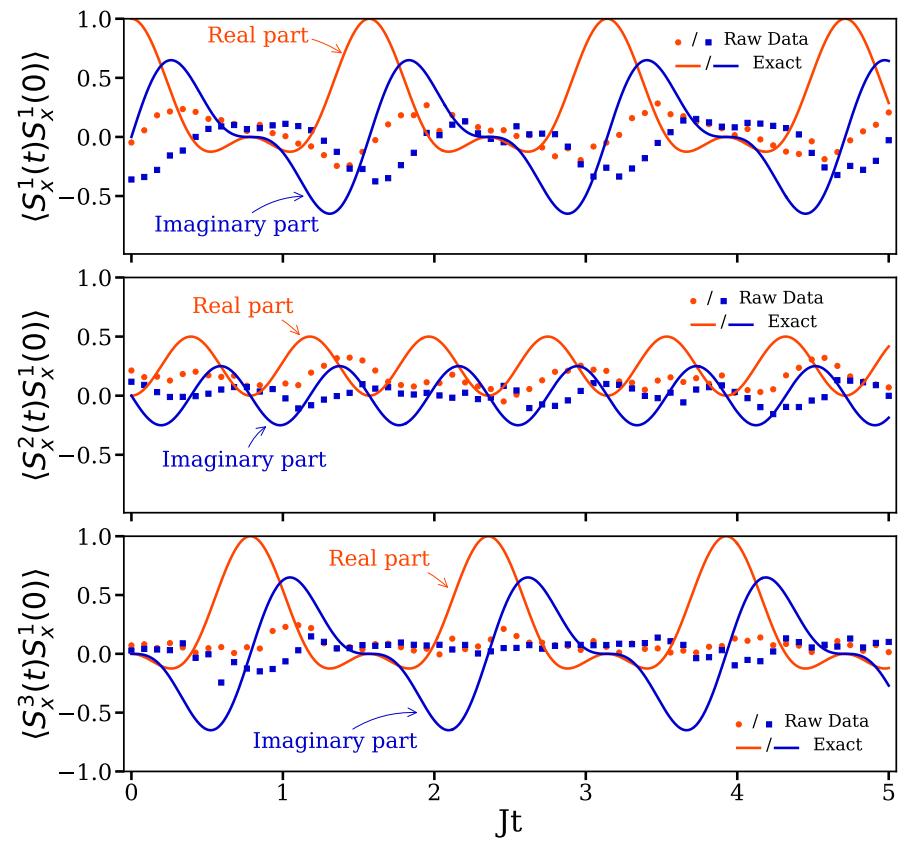
\*Actual Gates



# Low-energy excitations: 4-site magnons

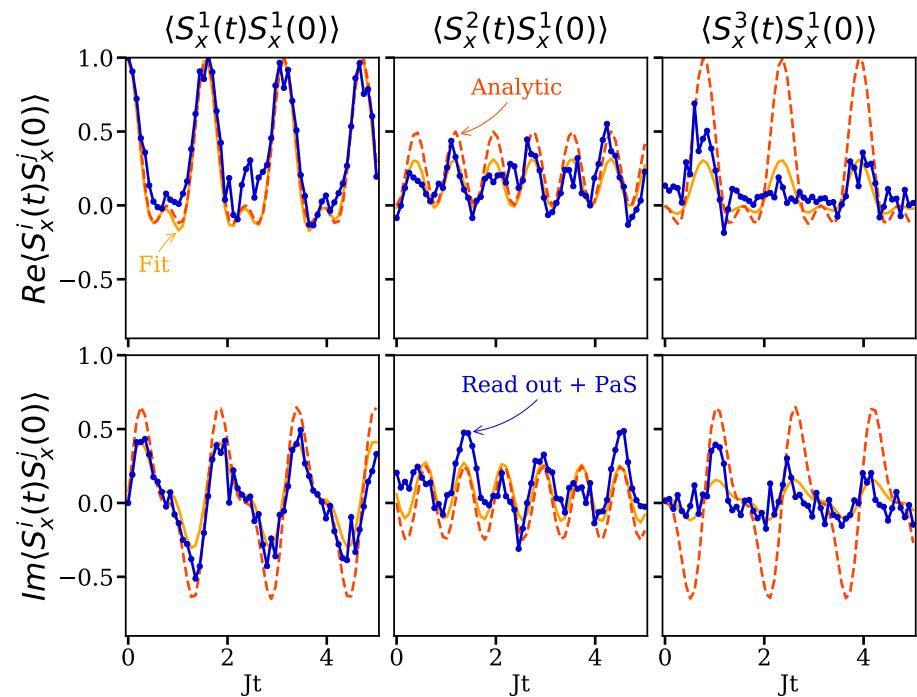
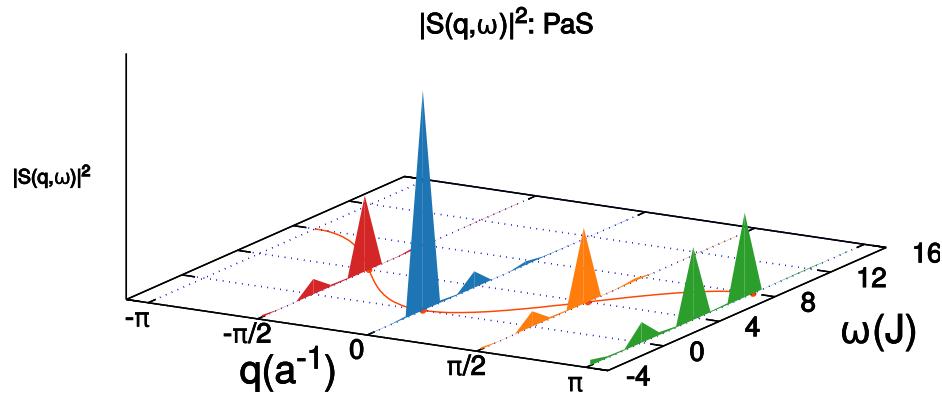
Error mitigation strategies:

1. Read out error correction
2. Phase-and-scale
  - Multiply by complex number
3. Fourier fitting
  - Assume only 4 frequencies



# Low-energy excitations: 4-site magnons

Spin-spin correlation function for periodic Heisenberg model: Magnons!

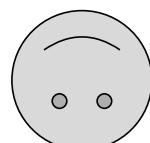


[10.1103/PhysRevB.101.014411](https://arxiv.org/abs/101.014411)

# Quantum computing for condensed matter?



Quantum computers are promising tools for solving problems in CMT...



... but care needs to be taken to ameliorate their noisiness in the NISQ era.

Noisy  
Intermediate  
Scale  
Quantum

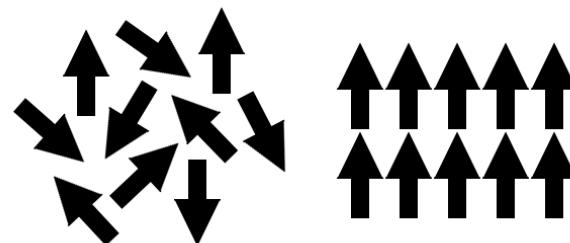


## Quantum Phase Transition

Phase transitions:



Quantum phase transitions:  $T=0$ , realized by quantum fluctuation controlled by physical quantities e.g. magnetic fields:

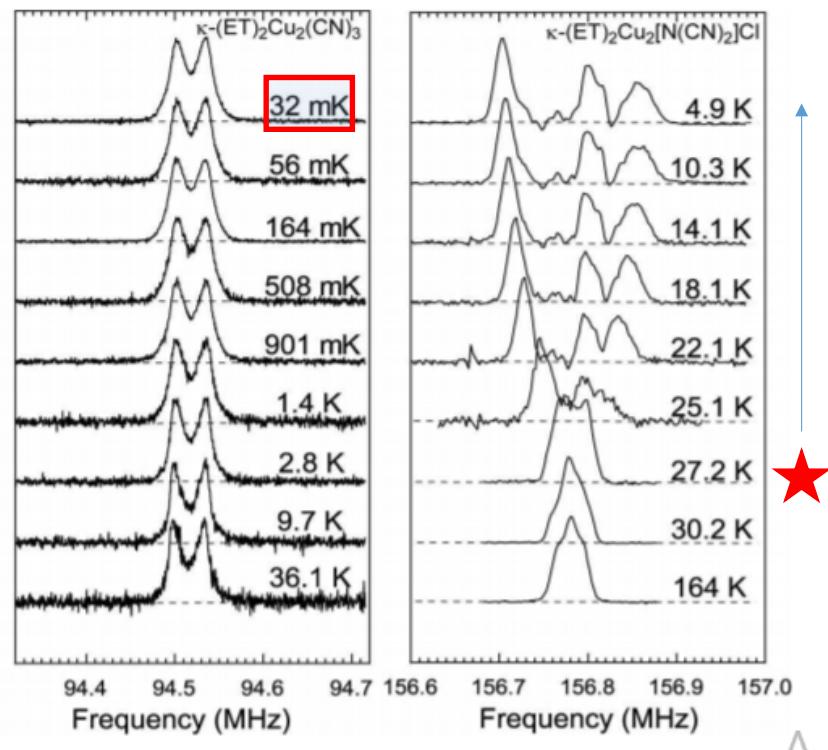


Applied Magnetic  
Field Absent

Applied Magnetic  
Field Present

# Topological Quantum Phase Transitions

The important feature of topological quantum phase transition is **featureless** in the sense that there is **no local order parameter** and no symmetry breaking.



Spin Liquid State in an Organic Mott Insulator with a Triangular Lattice

Y. Shimizu, K. Miyagawa, K. Kanoda, M. Maesato, and G. Saito  
Phys. Rev. Lett. **91**, 107001 – Published 4 September 2003



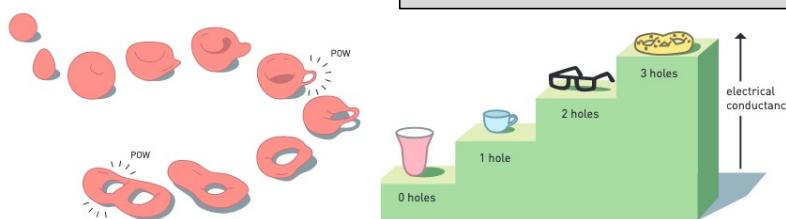
## Emergence of topological matters (after 1980)

In 1980, the first ordered phase beyond symmetry breaking was discovered.

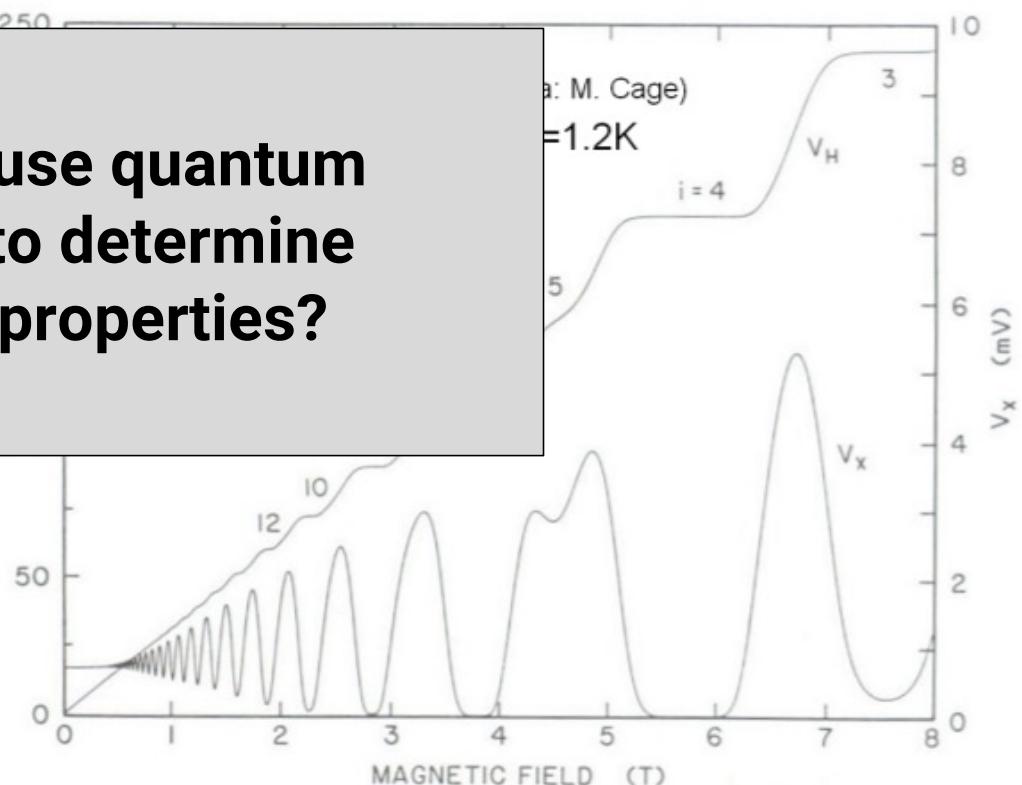
*Electrons confined to magnetic field show, temperature, plateau conductance" appears*

$$\sigma_{xy} = n$$

**How do we use quantum computers to determine topological properties?**



**Fig 3. Topology.** This branch of mathematics is interested in properties that change step-wise, like the number of holes in the above objects. Topology was the key to the Nobel Laureates' discoveries, and it explains why electrical conductivity inside thin layers changes in integer steps.



Nobel Prize in Physics 2016

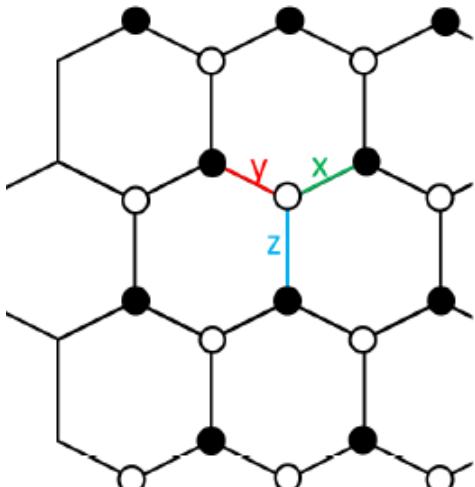
# The Kitaev model

PHYSICAL REVIEW B **95**, 045112 (2017)

## Kitaev honeycomb tensor networks: Exact unitary circuits and applications

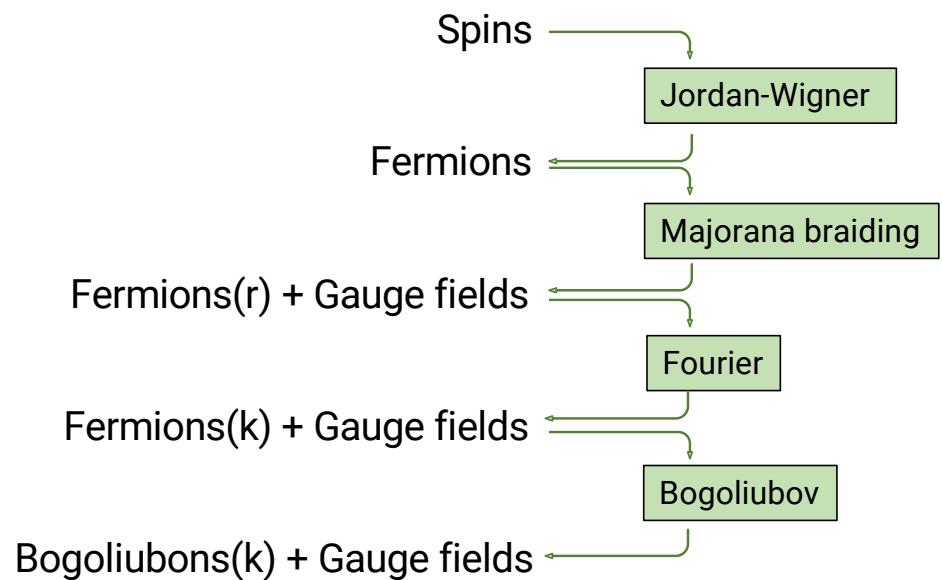
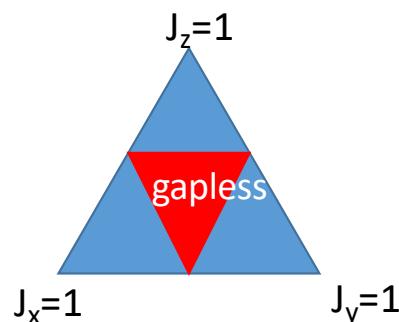
Philipp Schmoll and Román Orús

*Institute of Physics, Johannes Gutenberg University, 55099 Mainz, Germany*

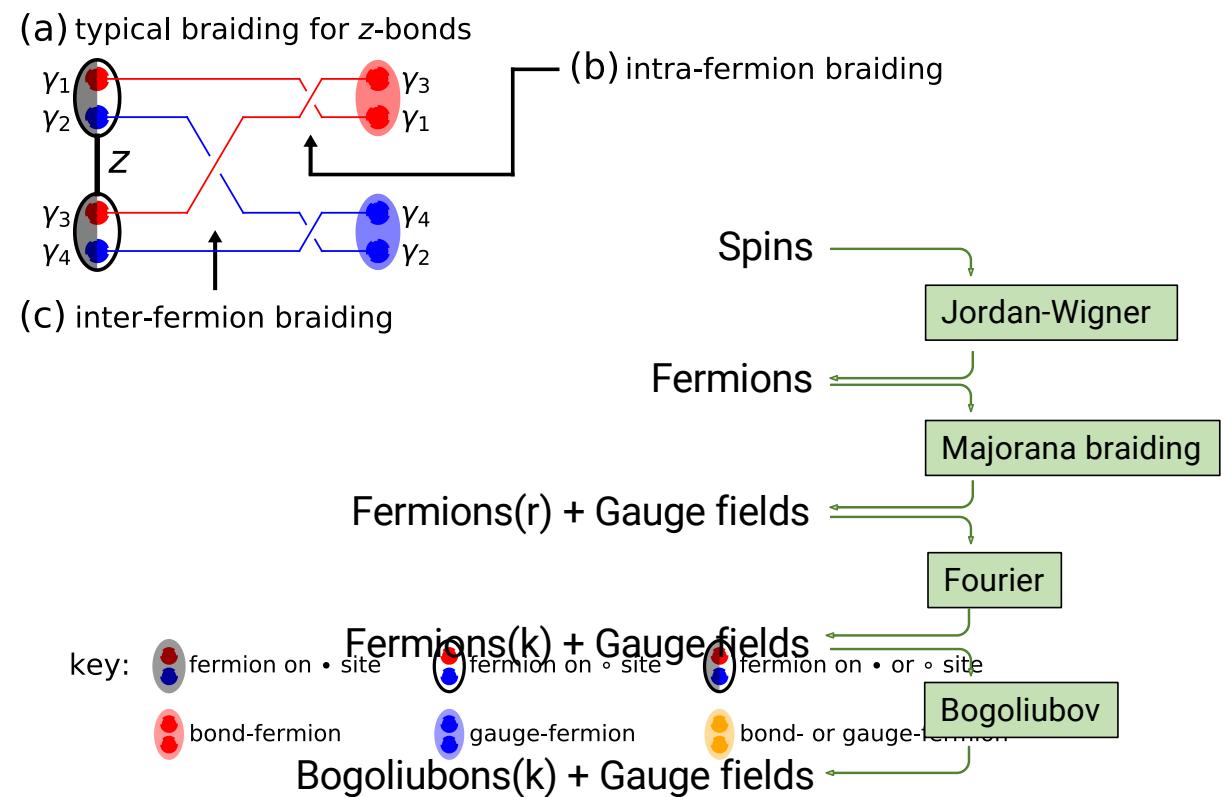
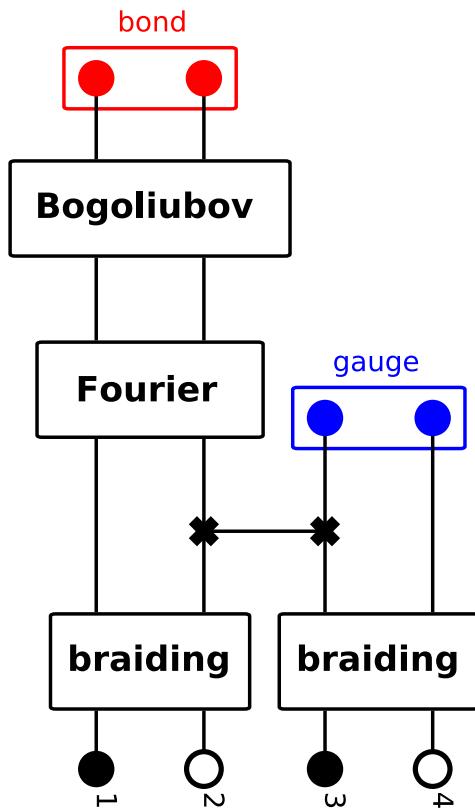


No local order  
parameter!

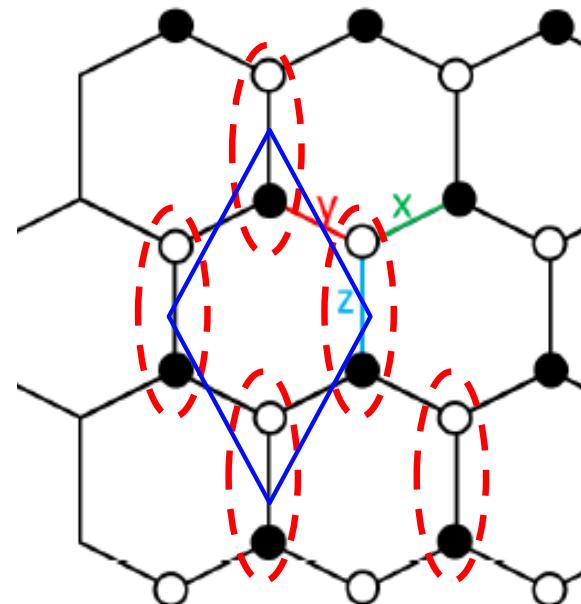
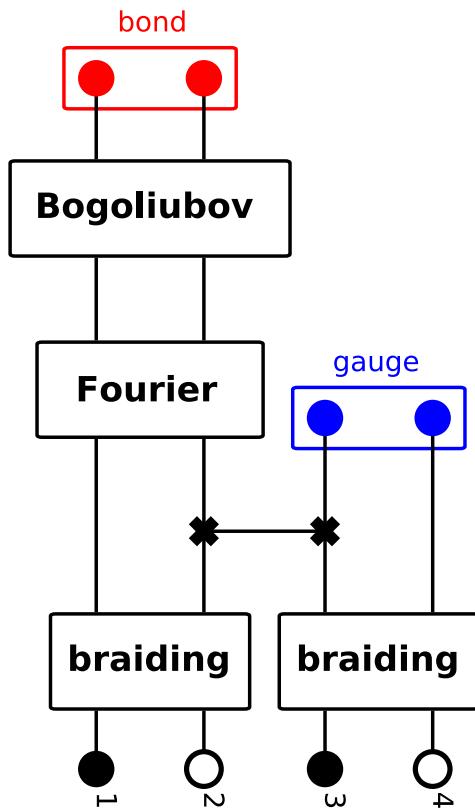
Topological QPT



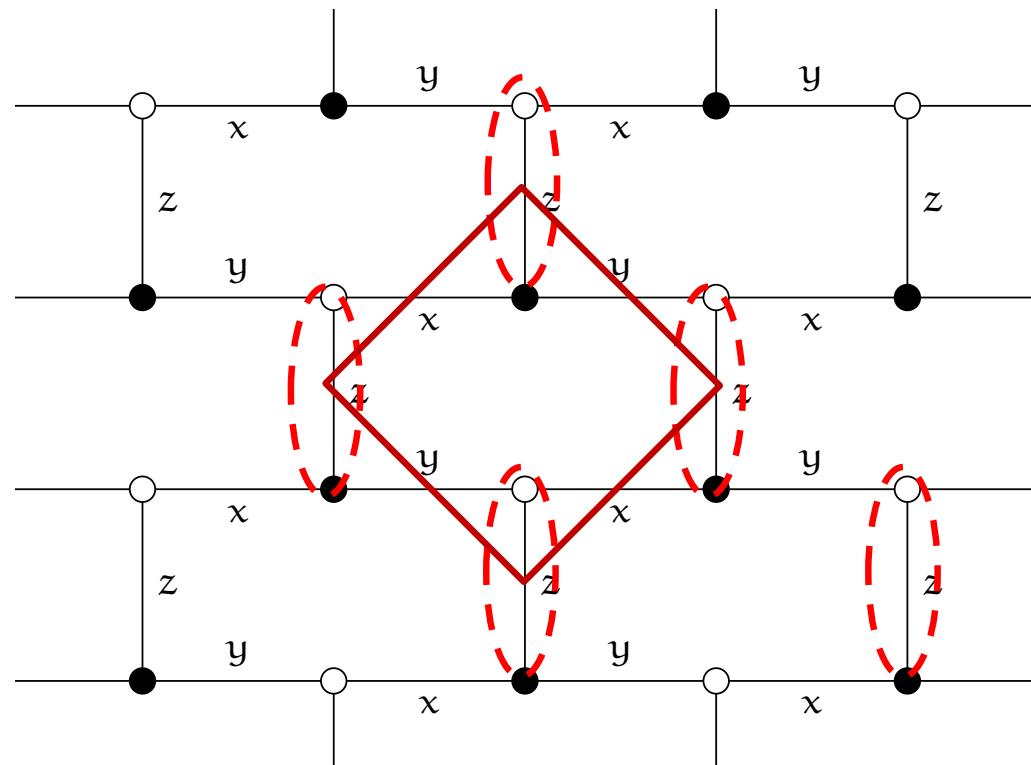
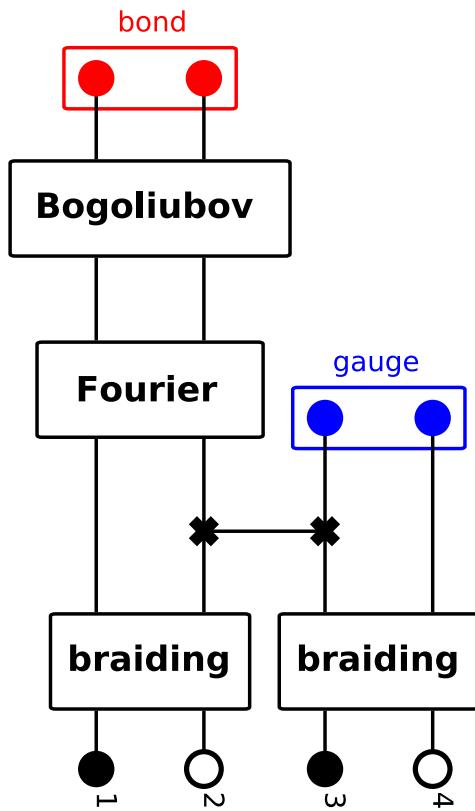
# Majorana Braiding



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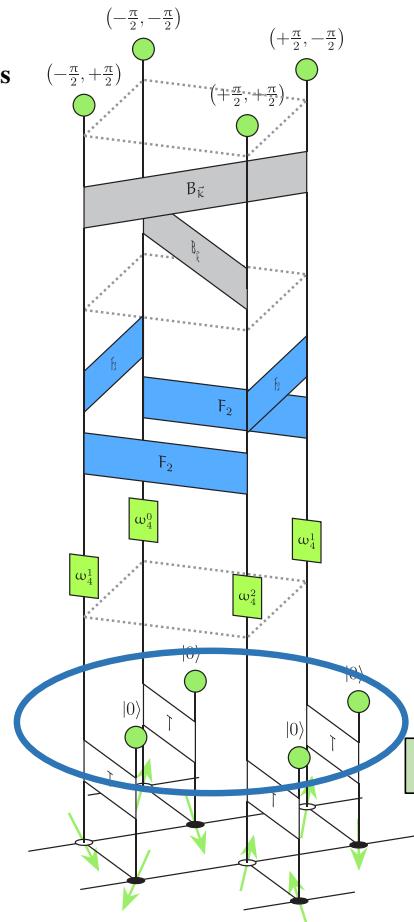
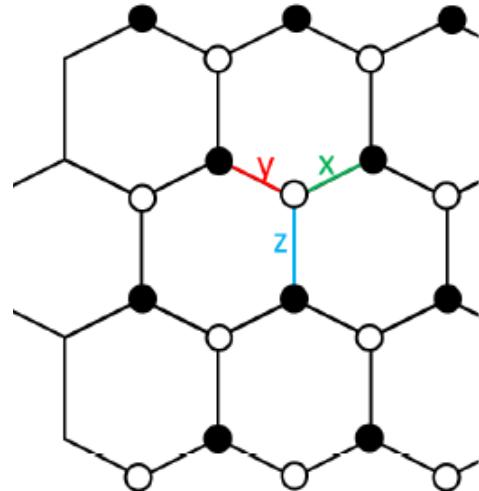
# The Kitaev model

PHYSICAL REVIEW B **95**, 045112 (2017)

## Kitaev honeycomb tensor networks: Exact unitary circuits and applications

Philipp Schmoll and Román Orús

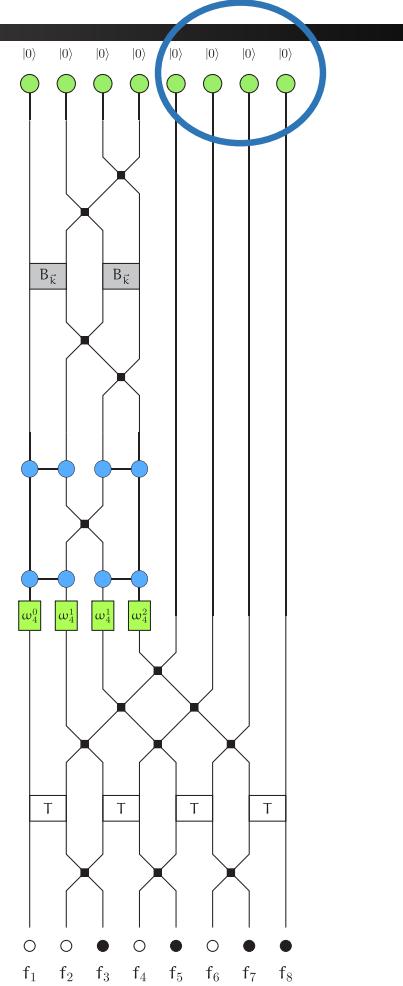
*Institute of Physics, Johannes Gutenberg University, 55099 Mainz, Germany*



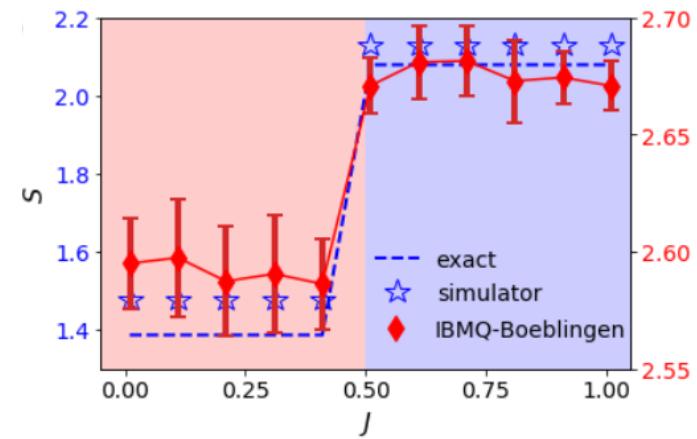
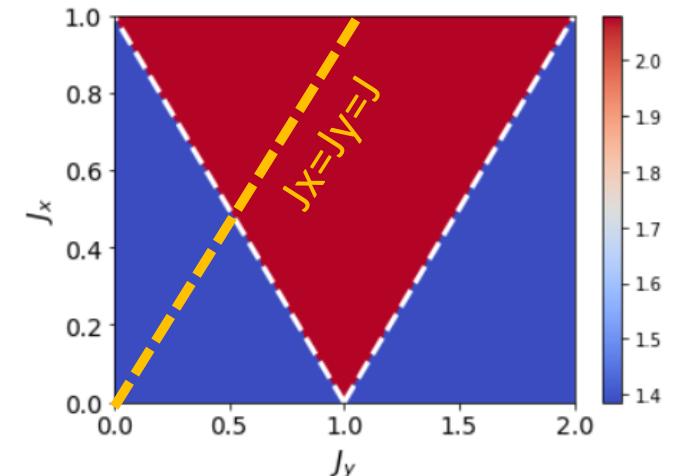
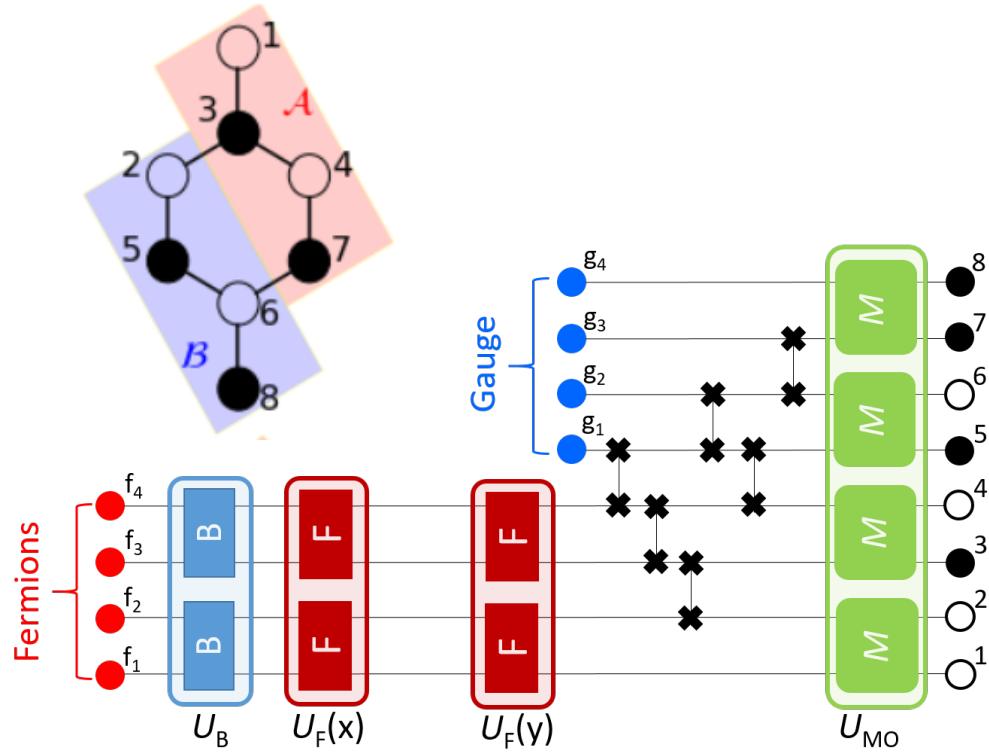
Bogoliubov

Fourier

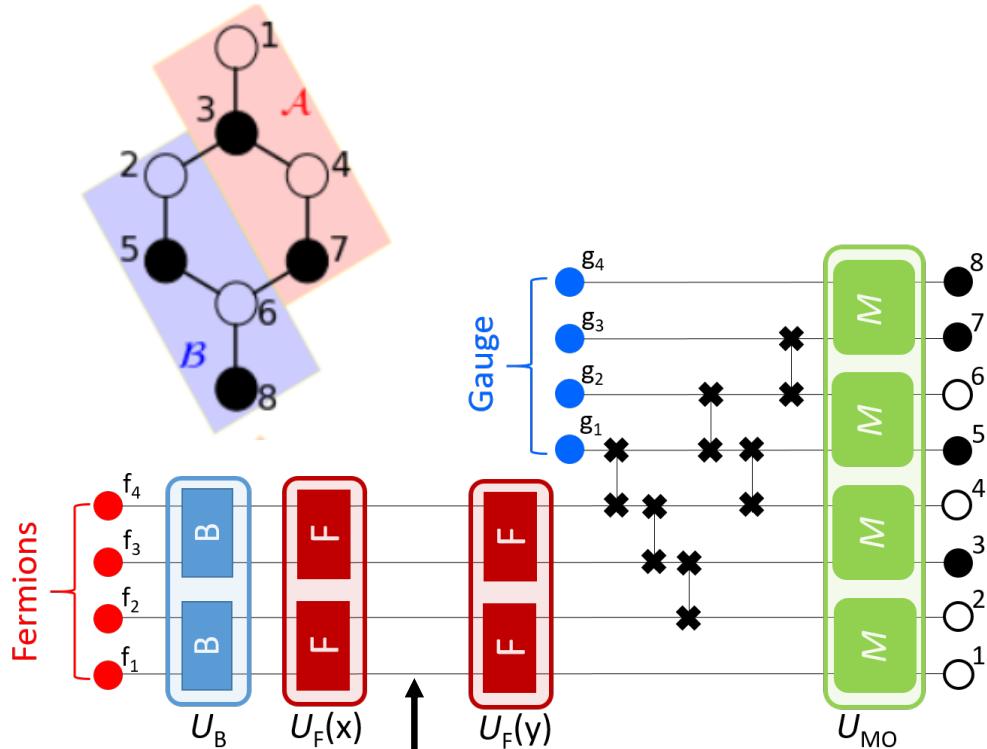
Majorana braiding



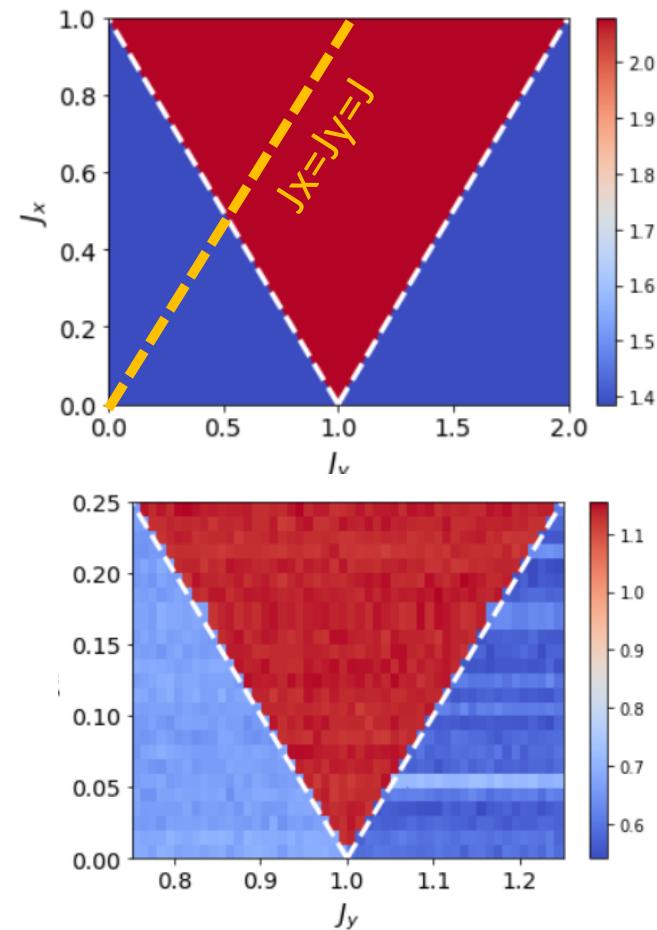
# Characterize QPT of Kitaev spin model by quantum computer



## Characterize QPT of Kitaev spin model by quantum computer

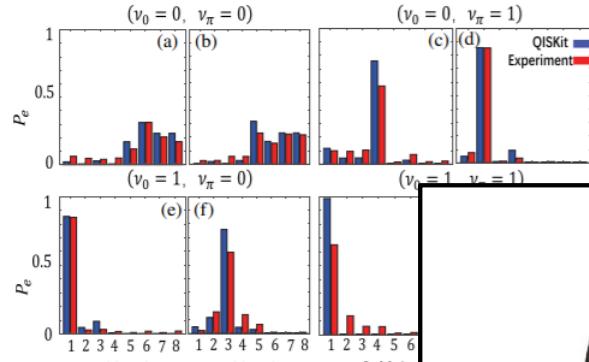


But we can do better! The gauge part has a trivial contribution, so we can instead compute a correlation function  $\langle \eta_x(-k_y) \eta_{x'}(k_y) \rangle$

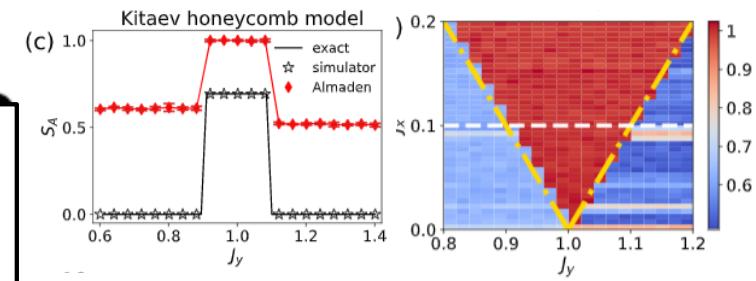


# Topological states are robust... or are they?

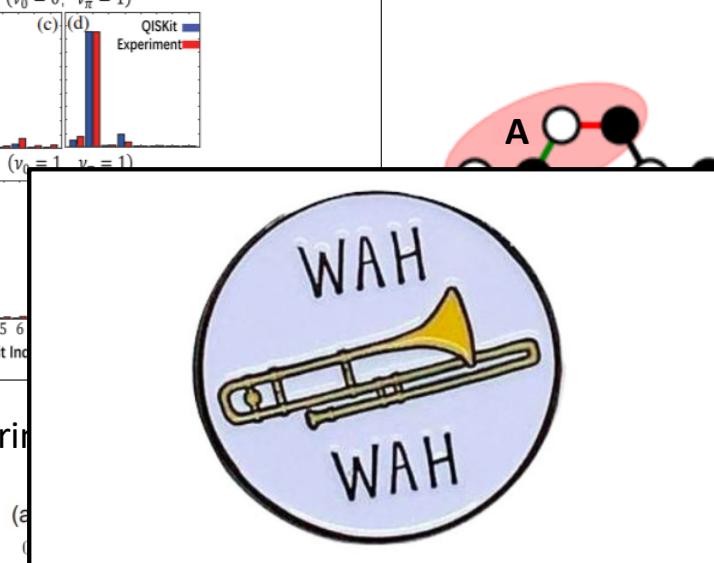
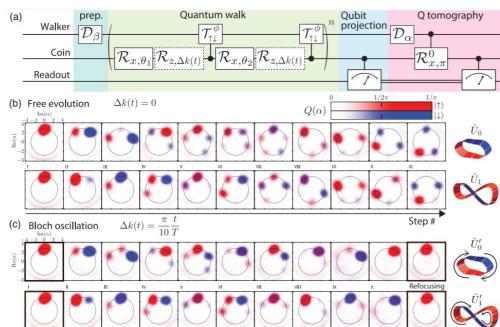
Floquet topological states, Mei et al (PRL 2020)



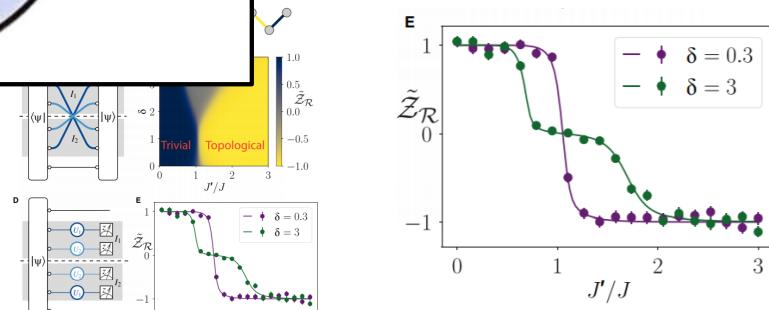
Kitaev Honeycomb Model, Xiao et al (arXiv:2006.05524)



1D topological quantum walk, Flurin et al (PRL 2019)



Protected Topological (SPT) state, Elben (Sci Adv 2020)

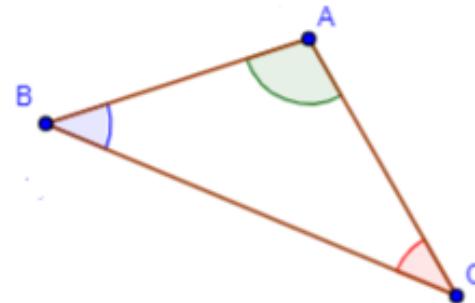


## A Brief Detour Into Geometry



The earth is round. [[citation needed](#)]

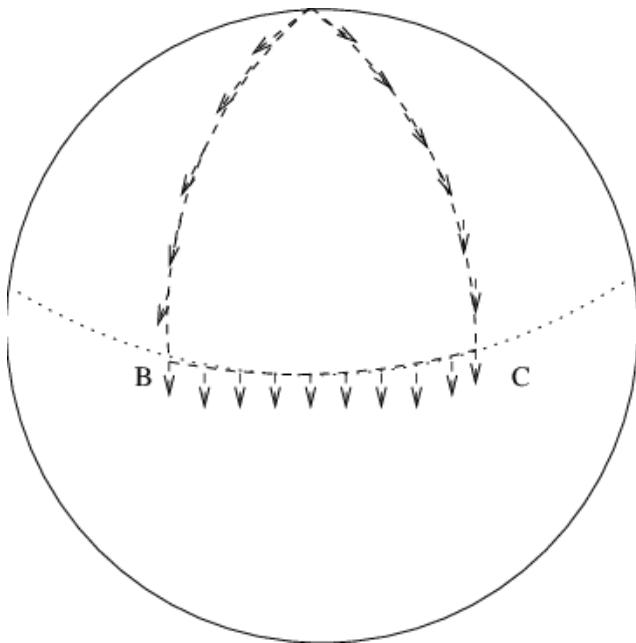
How can we tell?



$$\angle A + \angle B + \angle C = 180^\circ$$

unless you're on a curved space...

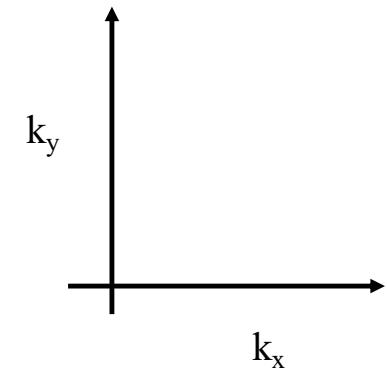
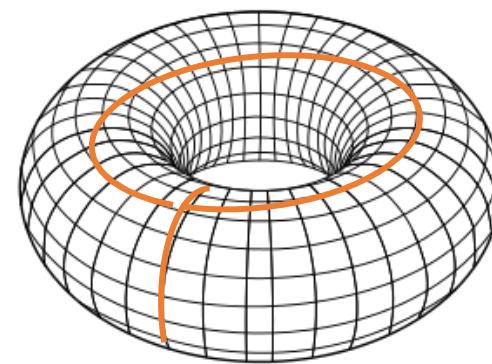
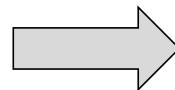
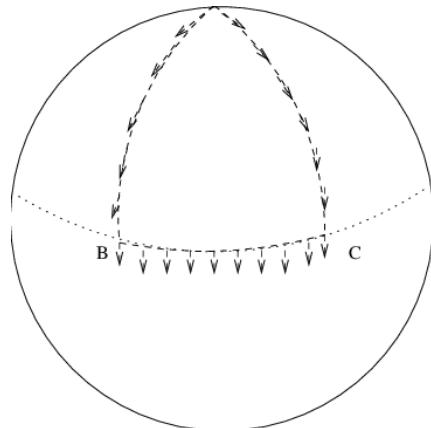
## A Brief Detour Into Geometry



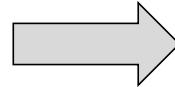
We could also use *parallel transport*.

$$\frac{\partial}{\partial \vec{s}} \vec{v}$$

If you are on a curved space, as you make a loop your vector may point a different way than you started.



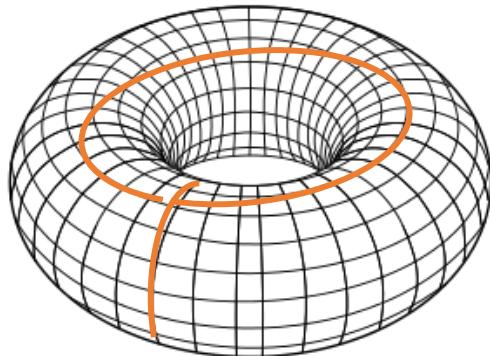
$$\frac{\partial}{\partial \vec{s}} \vec{v}$$



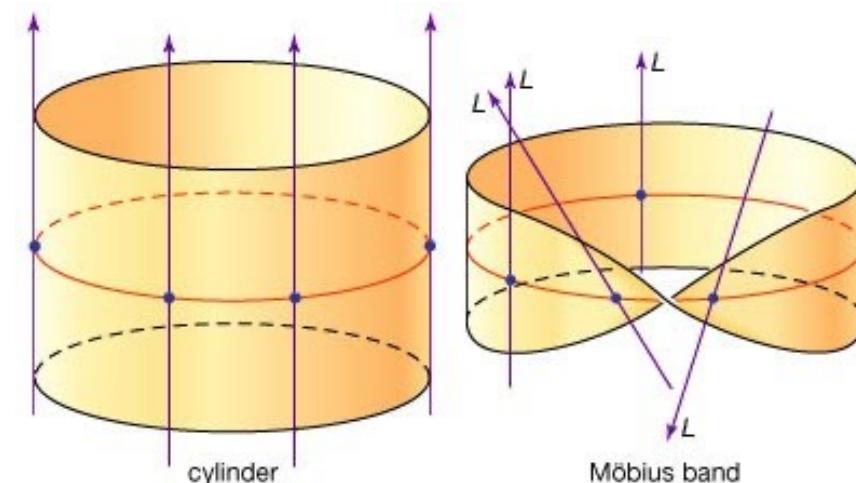
$$\langle \Psi_\Theta | \partial_\Theta | \Psi_\Theta \rangle$$

Q: If you walk a closed loop around the torus, what phase do you pick up?

A: Some **integer** multiple of  $2\pi$



$$\langle \Psi_\Theta | \partial_\Theta | \Psi_\Theta \rangle$$



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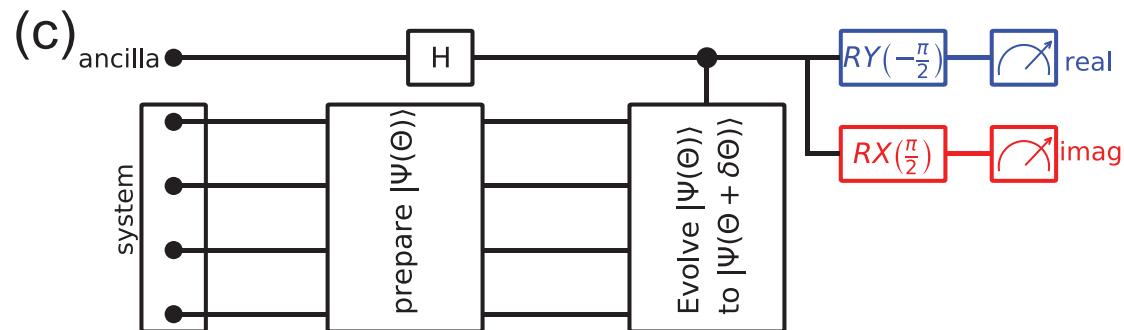
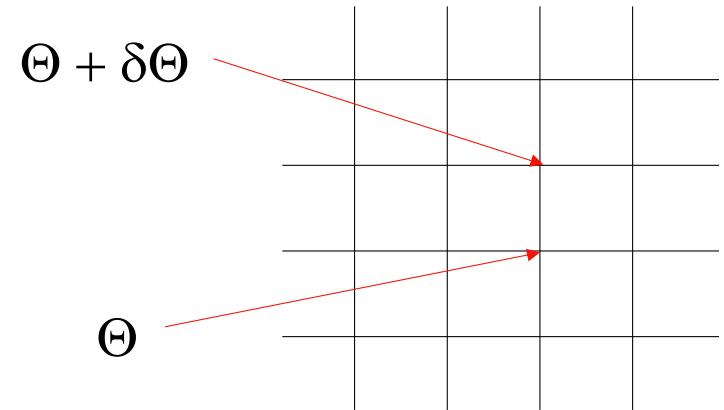
Q: If you walk a closed loop around the torus, what phase do you pick up?

A: Some **integer** multiple of  $2\pi$

How can we measure the topology of a wave function (bundle)?

$$\mathcal{C}_{Berry} = \langle \Psi_\Theta | \partial_\Theta | \Psi_\Theta \rangle$$

$$\delta\Theta \langle \Psi_\Theta | \partial_\Theta | \Psi_\Theta \rangle \approx \ln [1 + \delta\Theta \langle \Psi_\Theta | \partial_\Theta | \Psi_\Theta \rangle] \approx \ln \langle \Psi_\Theta | \Psi_{\Theta+\delta\Theta} \rangle$$

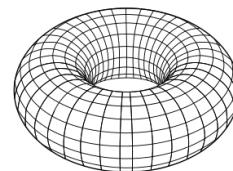


# Quantum Geometry

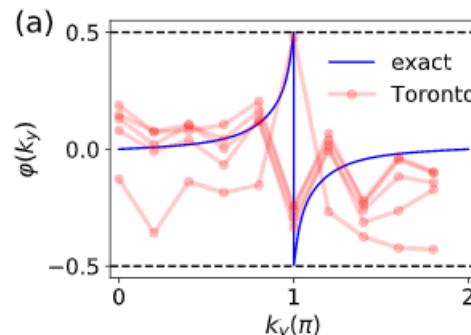
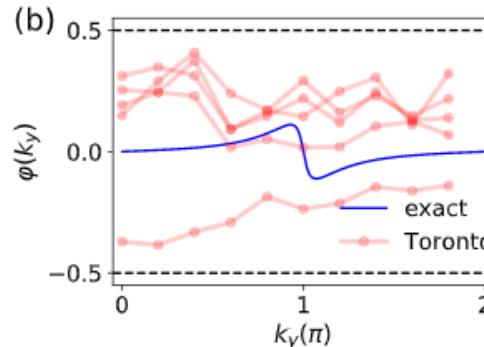
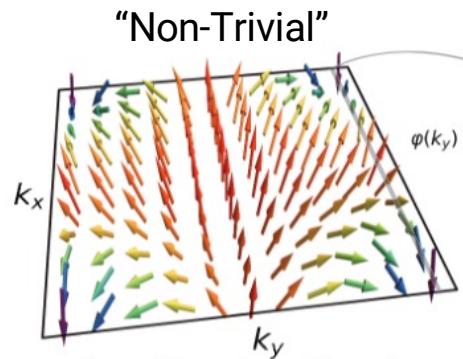
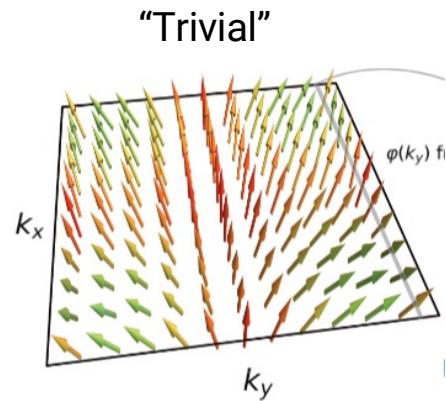
Chiral  $p$ -wave  
superconductor

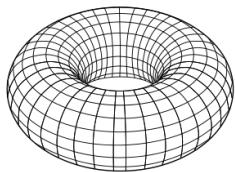
$$\mathcal{H}(\mathbf{k}) = \Delta(\sin k_y \sigma_x + \sin k_x \sigma_y) - \epsilon(\mathbf{k}) \sigma_z$$

$$\epsilon(\mathbf{k}) = t(\cos k_x + \cos k_y) + \mu$$

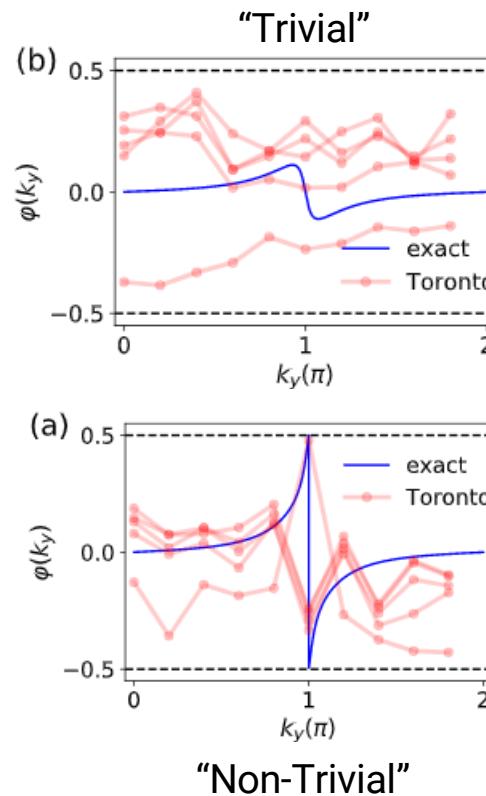
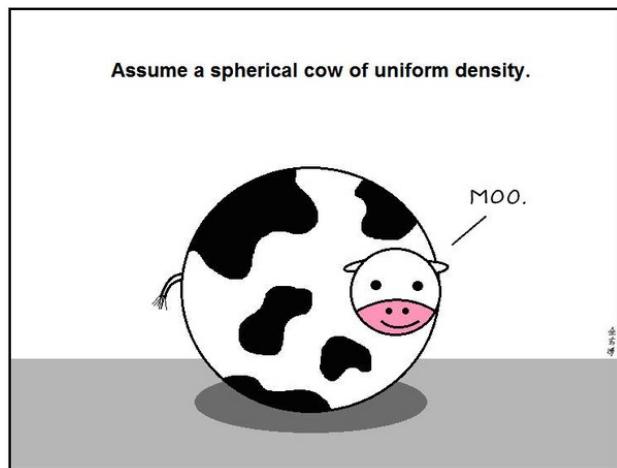


Phase:  $2\pi\mathcal{C}$





Phase:  $2\pi\mathcal{C}$



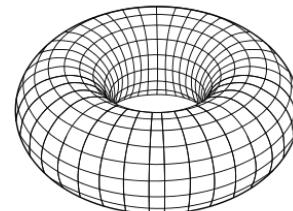
*Solution: Rather than taking lines across the manifold, consider the whole manifold.*

## Quantum Geometry

Chiral  $p$ -wave  
superconductor

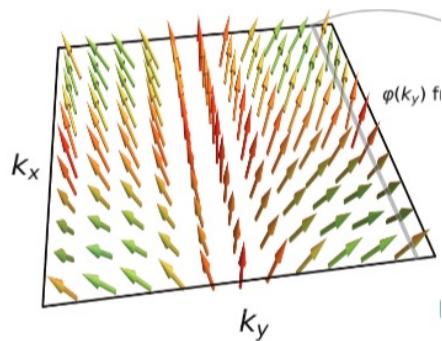
$$\mathcal{H}(\mathbf{k}) = \Delta(\sin k_y \sigma_x + \sin k_x \sigma_y) - \epsilon(\mathbf{k}) \sigma_z$$

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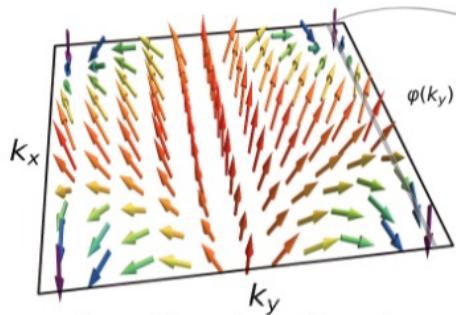


Phase:  $2\pi\mathcal{C}$

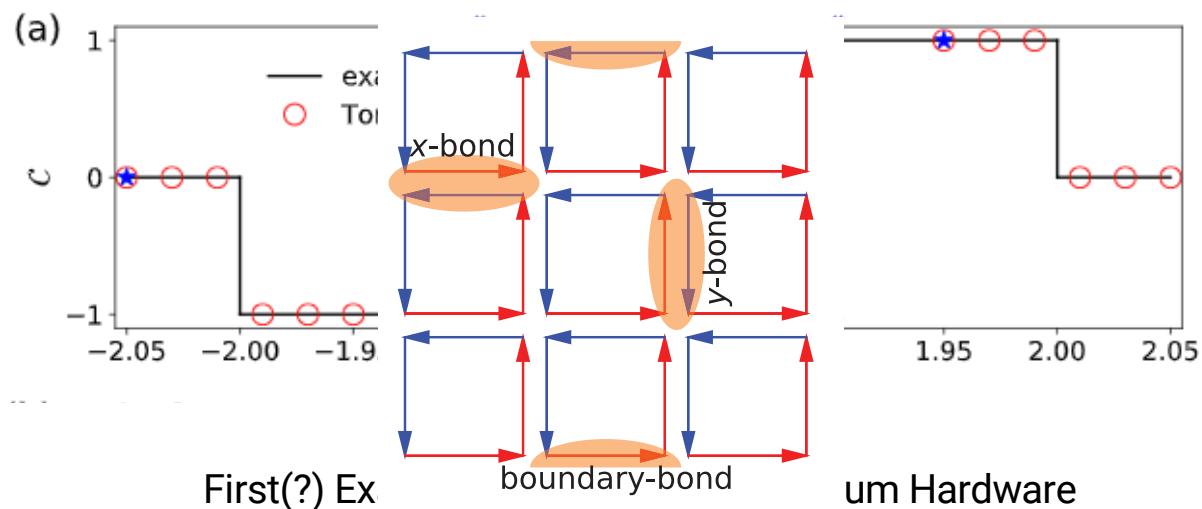
"Trivial"



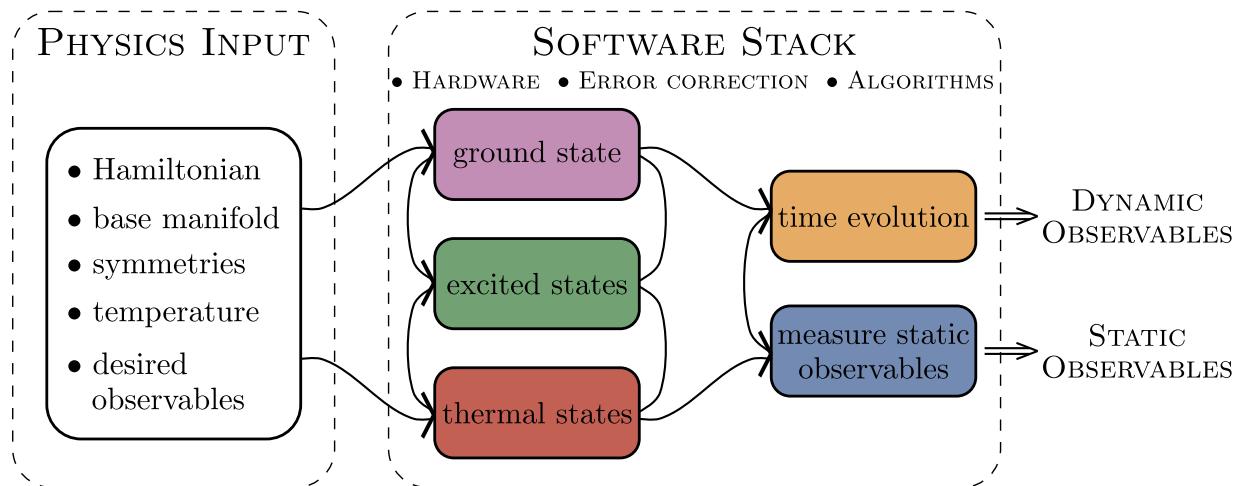
"Non-Trivial"



$$\mathcal{C} = \oint d\mathbf{k} \cdot \mathcal{A} = \int d^2k [\nabla \times \mathbf{A}(\mathbf{k})] \cdot \mathbf{n}_z = \int d^2k \mathcal{F}_{xy}(\mathbf{k}) \xrightarrow{\text{discretize}} \frac{1}{i2\pi} \sum_{\mathbf{k}} \mathcal{F}(\mathbf{k})$$

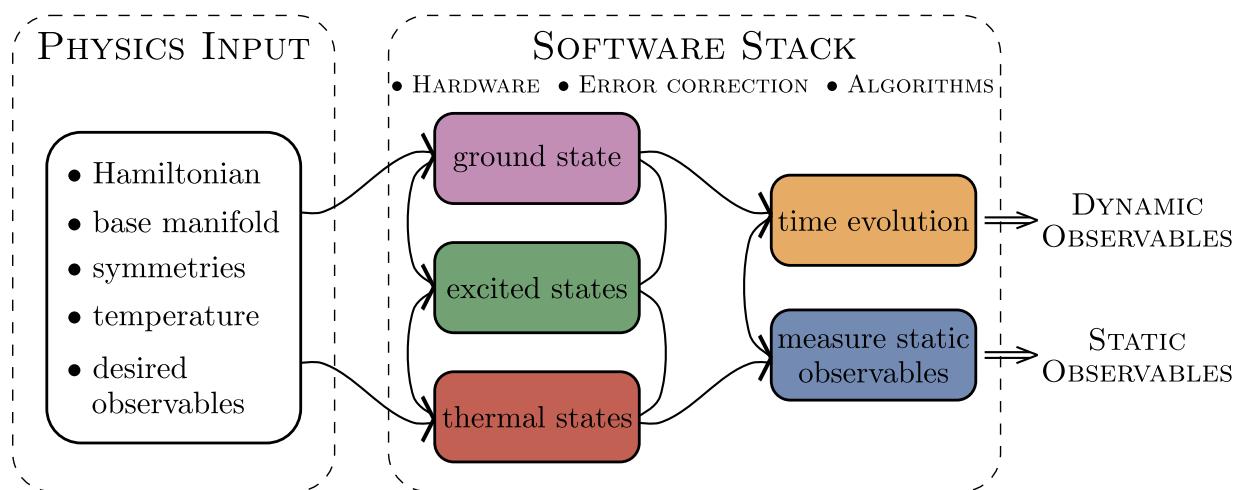


## Quantum Matter meets Quantum Computing

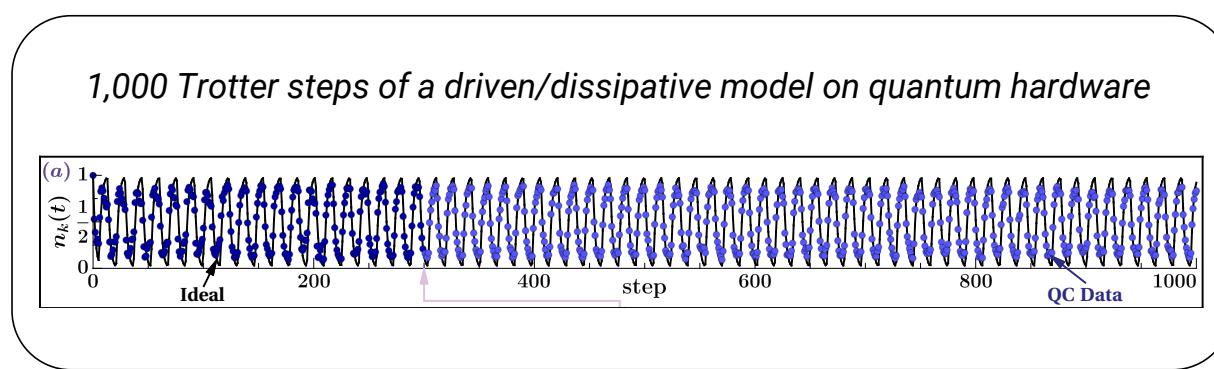


- Measuring correlation functions
- Preparing/measuring topological states
- Modeling driven/dissipative systems
- Time evolution via Lie algebraic decomposition/compression
- Thermodynamics

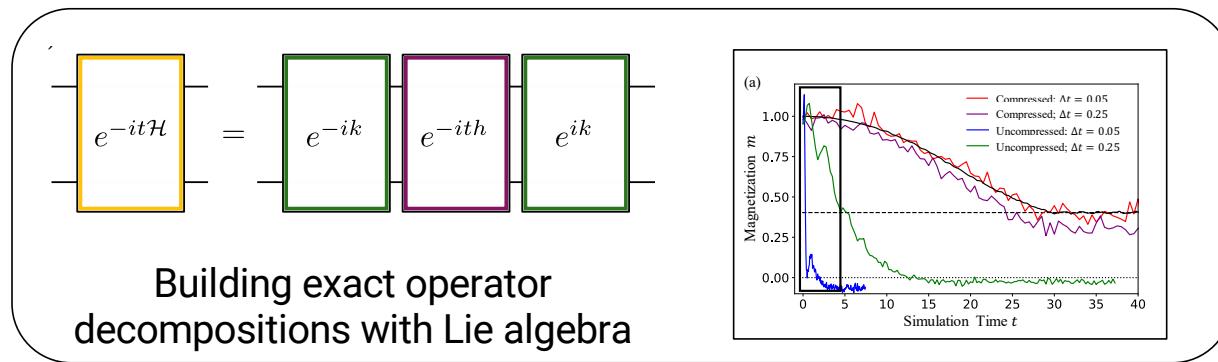
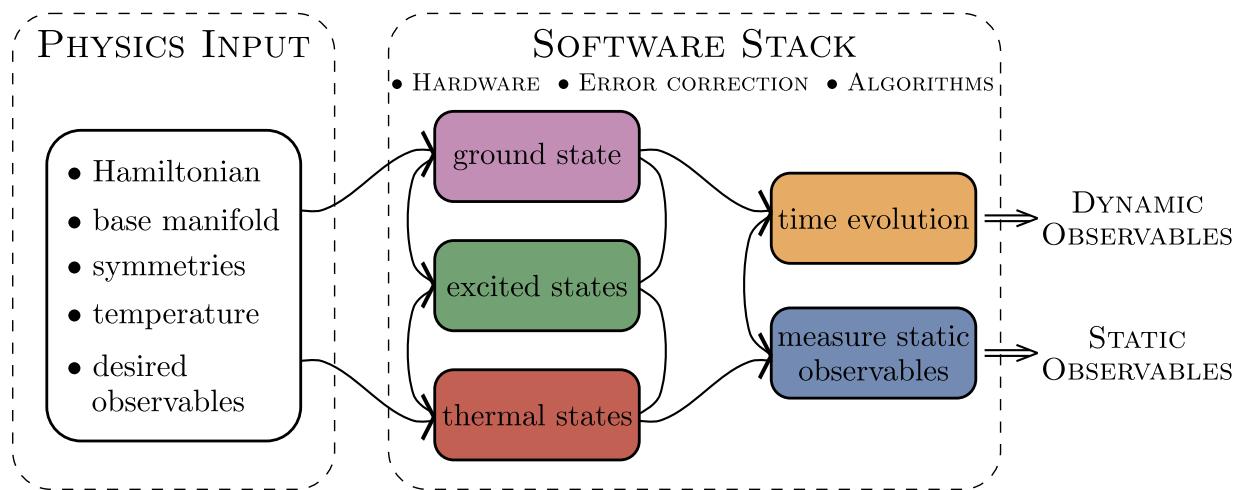
## Quantum Matter meets Quantum Computing



- Measuring correlation functions
- Preparing/measuring topological states
- **Modeling driven/dissipative systems**
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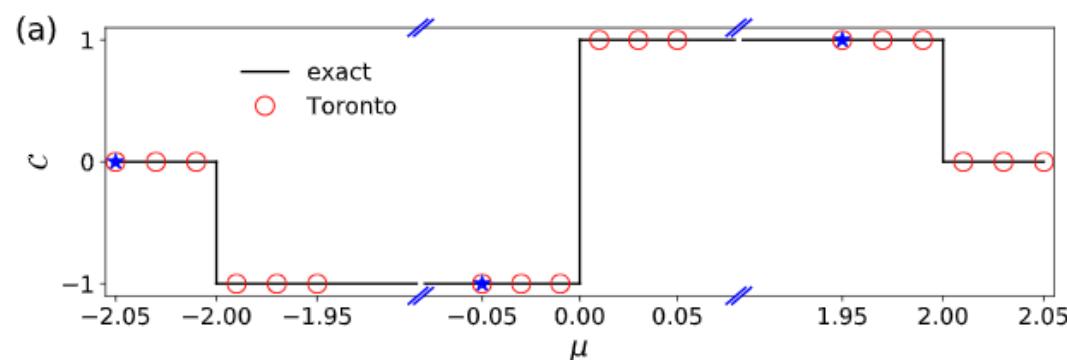
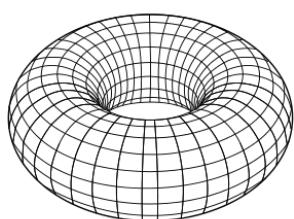
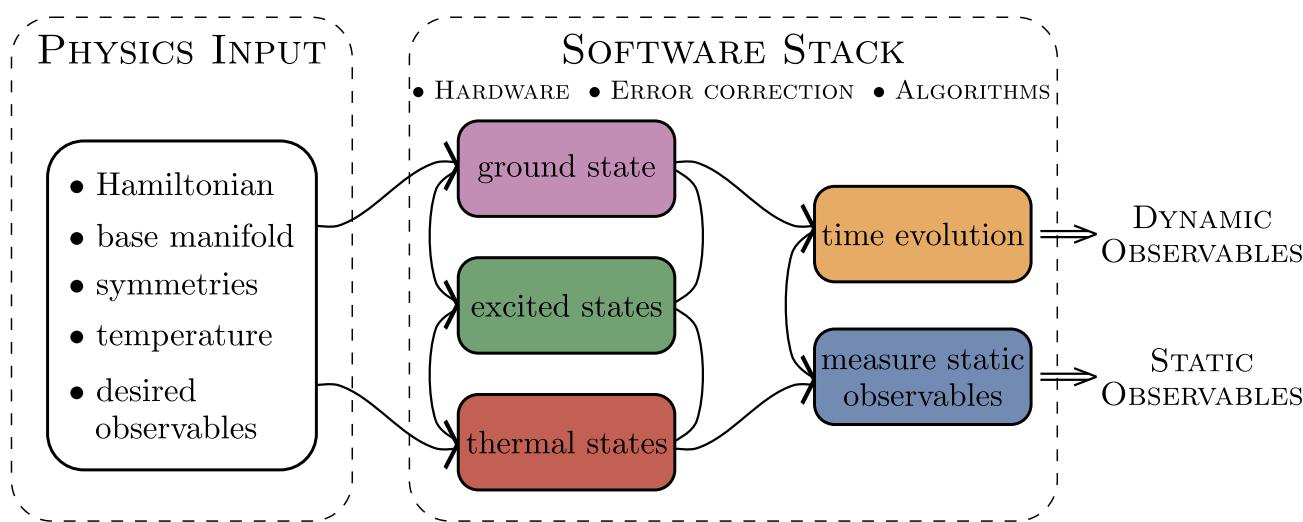


# Quantum Matter meets Quantum Computing



- Measuring correlation functions
- Preparing/measuring topological states
- Modeling driven/dissipative systems
- Time evolution via Lie algebraic decomposition/compression
- Thermodynamics

## Conclusion



Topological physics on QC  
Quantum 5, 553 (2021)  
arXiv:2101.07283

Thermodynamics on QC  
Science Advances Vol. 7, no.  
34, eabf2447  
arXiv:2112.05688

Exact Time Evolution  
arXiv:2104.00728  
Phys. Rev. A 105, 032420 (2022)  
arXiv:2108.03283

Computing Spectra  
10.1103/PhysRevB.101.014411

Quantum Matter on QC Review  
Quantum Sci. Technol. 6  
043002