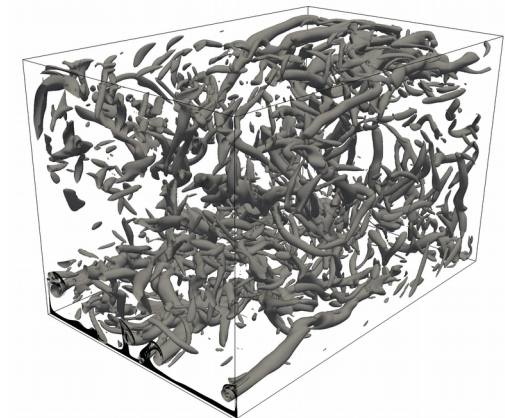
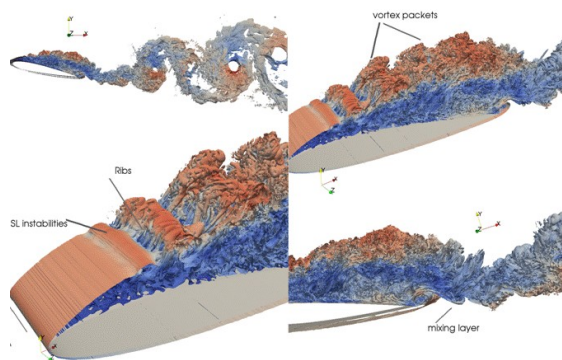


# COMPUTATIONAL ENGINEERING

## TACKLING TURBULENCE WITH (SUPER)COMPUTERS



# BURGULENCE

A GOOD TOY MODEL OF TURBULENCE

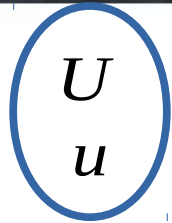
*On his death bed, Werner Heisenberg is reported to have said, «When I meet God, I am going to ask him two questions: Why relativity? And why turbulence? I really believe he will have an answer for the first.»*



# REMAINDER 1: TURBULENCE IS EXPENSIVE

$l$ : biggest eddies ( driving scale )

$\eta$ : smallest eddies ( Kolmogorov length scale )



(Computational) **Complexity of NS** grows very fast!!!

In space:  $dim=3 \Rightarrow (l/\eta)^3 \sim Re_l^{9/4}$

In time:  $t_{eddy}/t_\eta \sim Re_l^{1/2}$

In space  $\times$  time:  $(l/\eta)^3 t_{eddy}/t_\eta \sim Re_l^{11/4}$



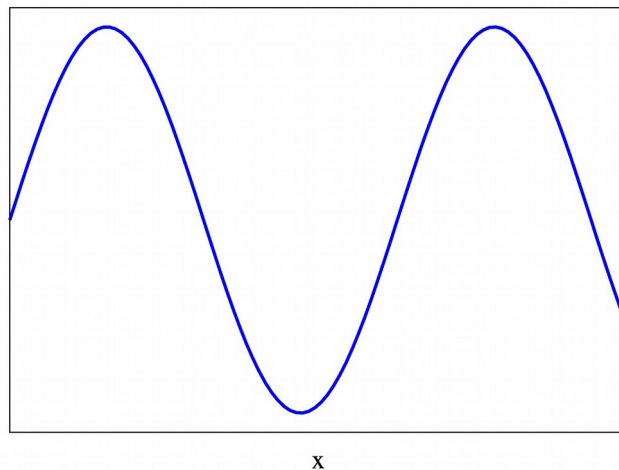
$$\begin{aligned}\eta/l &\sim Re_l^{-3/4} \\ u/U &\sim Re_l^{-1/4} \\ t_\eta/t_{eddy} &\sim Re_l^{-1/2}\end{aligned}$$

$$\begin{aligned}\eta/l &\sim (\nu^3/\varepsilon)^{1/4} \\ u &\sim (\nu\varepsilon)^{1/4} \\ t_\eta &\sim (\nu/\varepsilon)^{1/2}\end{aligned}$$

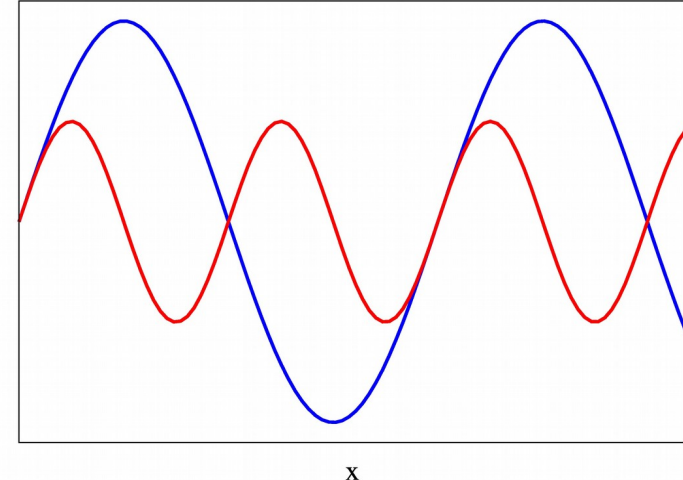
$$\frac{\partial \vec{u}}{\partial t} + \underbrace{(\vec{u} \cdot \nabla) \vec{u}}_{\text{nonlinear guy in NS}} = \nu \nabla^2 \vec{u} - \nabla p; \quad \nabla \cdot \vec{u} = 0$$

Let us consider a one single 1D wave,  $u(x) = \sin(kx)$

$$u \frac{\partial u}{\partial x} = k \sin(kx) \cos(kx) = \frac{k}{2} \sin(2kx)$$



$$u \frac{\partial u}{\partial x}$$



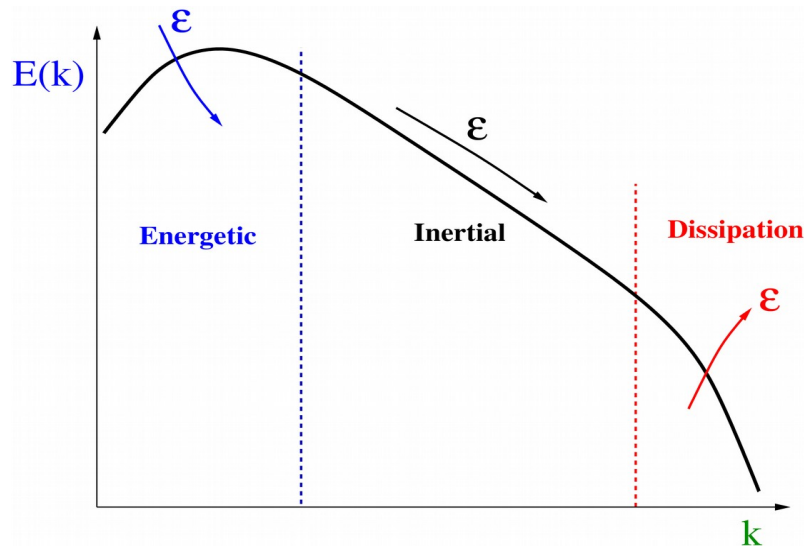
REMAINDER!!!

Double-angle formula:  $\sin(2x) = 2 \sin(x) \cos(x)$



# REMAINDER 3: ENERGY CASCADE

$$\begin{aligned}[k] &= [L^{-1}] & [\varepsilon] &= [L^2 T^{-3}] \\ [E_T] &= [L^2 T^{-2}] & [E(k)] &= [L^3 T^{-2}]\end{aligned}$$



Hypothesis (K41):  $E(k) \propto \varepsilon^a k^b$

$$\begin{aligned}2a - b &= 3 \\ -3a &= -2\end{aligned} \Rightarrow \begin{cases} a = 2/3 \\ b = -5/3 \end{cases}$$

Kolmogorov energy spectrum:

$$E(k) = C_K \varepsilon^{2/3} k^{-5/3}$$

**Formal proof** of the existence of the inertial subrange for the (incompressible) Navier-Stokes equations:  
C. Foias, O.P. Manley, R. Rosa and R. Temam: *Estimates for the energy cascade in three-dimensional turbulent flows*. C. R. Acad. Sci. Paris, Série I Math. 333, 499–504, (2001).

$$\frac{\partial \vec{u}}{\partial t} + \underbrace{(\vec{u} \cdot \nabla) \vec{u}}_{\text{nonlinear guy in NS}} = \nu \nabla^2 \vec{u} - \nabla p; \quad \nabla \cdot \vec{u} = 0$$

1D  No continuity eq.

$$\frac{\partial u}{\partial t} + \underbrace{u \frac{\partial u}{\partial x}}_{\text{nonlinear guy in Burgers'}} = \nu \frac{\partial^2 u}{\partial x^2} + f$$

This is simplified model that shares many of the features of the NS equations.

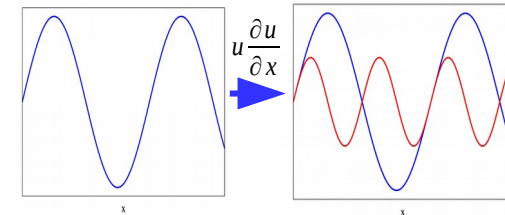
$$\frac{\partial \vec{u}}{\partial t} + \underbrace{(\vec{u} \cdot \nabla) \vec{u}}_{\text{nonlinear guy in NS}} = \nu \nabla^2 \vec{u} - \nabla p; \quad \nabla \cdot \vec{u} = 0$$

1D  $\downarrow$  No continuity eq.

$$\frac{\partial u}{\partial t} + \underbrace{u \frac{\partial u}{\partial x}}_{\text{nonlinear guy in Burgers'}}$$

$\downarrow$  Fourier Transform  $u(x) = \sum_{k=-N}^{k=+N} \hat{u}_k e^{ikx}$

$$\frac{\partial \hat{u}_k}{\partial t} + \underbrace{\sum_{p+q=k} \hat{u}_p i q \hat{u}_q}_{\text{nonlinear guy in Burgers'}}$$



We consider equation (6) on an interval  $\Omega$  with periodic boundary conditions. In Fourier space, the Burgers equation reads<sup>1</sup>

$$\partial_t \hat{u}_k + \sum_{k=p+q} \hat{u}_p i q \hat{u}_q = -\frac{k^2}{Re} \hat{u}_k + F_k \quad k = 0, \dots, N \quad (7)$$

where the forcing term is given by  $F_k = 0$  for  $k > 1$  and  $F_1$  such that  $\partial_t \hat{u}_1 = 0$  for  $t > 0$ . Here  $\hat{u}_k(t) \in \mathbb{C}$  denotes the  $k$ -th Fourier coefficient of  $u(x, t)$

$$u(x) \equiv \sum_{k=-N}^{k=+N} \hat{u}_k e^{ikx} \quad (8)$$

where  $N$  is the total number of Fourier modes. Note that since  $u(x, t) \in \mathbb{R}$  the following condition must be accomplished

$$\hat{u}_k = \overline{\hat{u}_{-k}} \quad (9)$$

where  $\overline{(\cdot)}$  denotes the complex conjugate. For details about the computation of the non-linear convective term the reader is referred to Appendix B.



$$\partial_t u + u \partial_x u = \frac{1}{Re} \partial_{xx} u + f \quad (\text{A.1})$$

each of the terms can be converted into Fourier space. Transient, diffusive and forcing terms are straightforward

$$\partial_t u = \sum_{k=-N}^{k=+N} (\partial_t \hat{u}_k) e^{ikx} \quad (\text{A.2})$$

$$\partial_{xx} u = \sum_{k=-N}^{k=+N} \hat{u}_k \partial_{xx} e^{ikx} = \sum_{k=-N}^{K=+N} (-k^2 \hat{u}_k) e^{ikx} \quad (\text{A.3})$$

$$f = \sum_{k=-N}^{k=+N} F_k e^{ikx} \quad (\text{A.4})$$

and non-linear convective term is given by

$$\partial_x u = \sum_{k=-N}^{k=+N} \hat{u}_k \partial_x e^{ikx} = \sum_{k=-N}^{k=+N} ik \hat{u}_k e^{ikx} \quad (\text{A.5})$$

$$u \partial_x u = \sum_{p=-N}^{p=+N} \hat{u}_p e^{ipx} \sum_{q=-N}^{q=+N} iq \hat{u}_q e^{iqx} = \sum_{p=-N, q=-N}^{p=+N, q=+N} \hat{u}_p iq \hat{u}_q e^{i(p+q)x} \quad (\text{A.6})$$

and the condition that  $u(x, t) \in \mathbb{R}$  is automatically satisfied

$$\hat{u}_{-p} i(-q) \hat{u}_{-q} = -\overline{\hat{u}_p} iq \overline{\hat{u}_p} = \overline{\hat{u}_p iq \hat{u}_p} \quad (\text{A.7})$$

Therefore, the Burgers equation (A.1) in Fourier space results

$$\sum_{k=-N}^{k=+N} (\partial_t \hat{u}_k) e^{ikx} + \sum_{p=-N, q=-N}^{p=+N, q=+N} \hat{u}_p iq \hat{u}_q e^{i(p+q)x} = \frac{1}{Re} \sum_{k=-N}^{k=+N} (-k^2 \hat{u}_k) e^{ikx} + \sum_{k=-N}^{k=+N} F_k e^{ikx} \quad (\text{A.8})$$

The contribution of the convective term to the  $k^{th}$  mode is given by

$$\sum_{k=p+q} \hat{u}_p i q \hat{u}_q \in \mathbb{C}. \quad (\text{B.1})$$

$$\hat{u}_p = \overline{\hat{u}_{-p}}. \quad (\text{B.2})$$

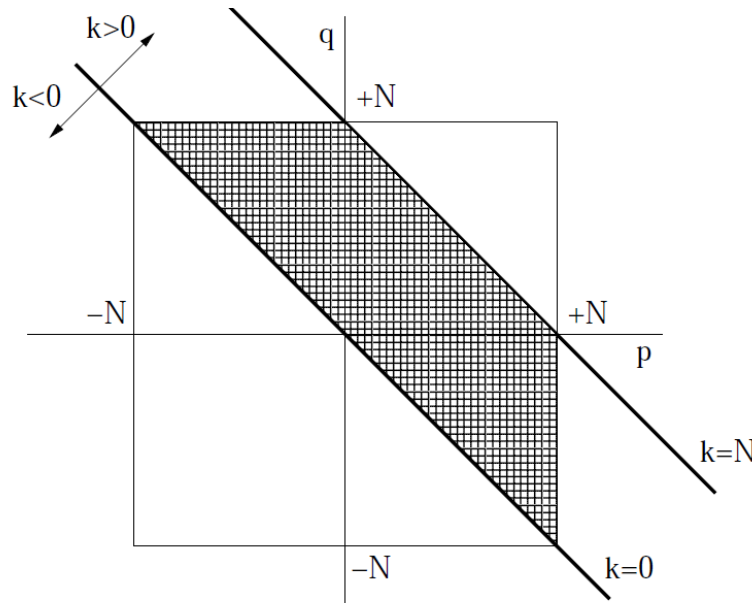


Figure 4: Representation of all possible triadic interactions between modes. Only the interactions between the straight lines  $k = 0$  and  $k = N$  need to be considered for the computation of the non-linear convective term.

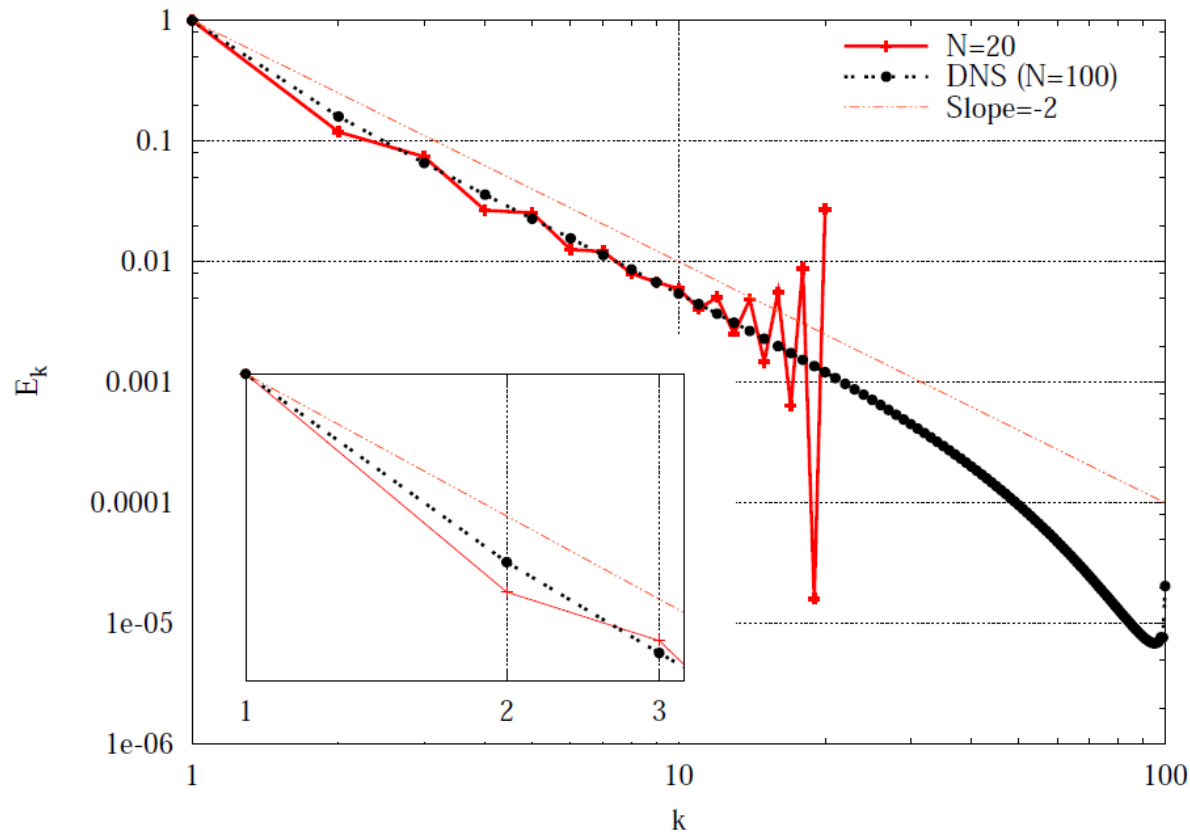
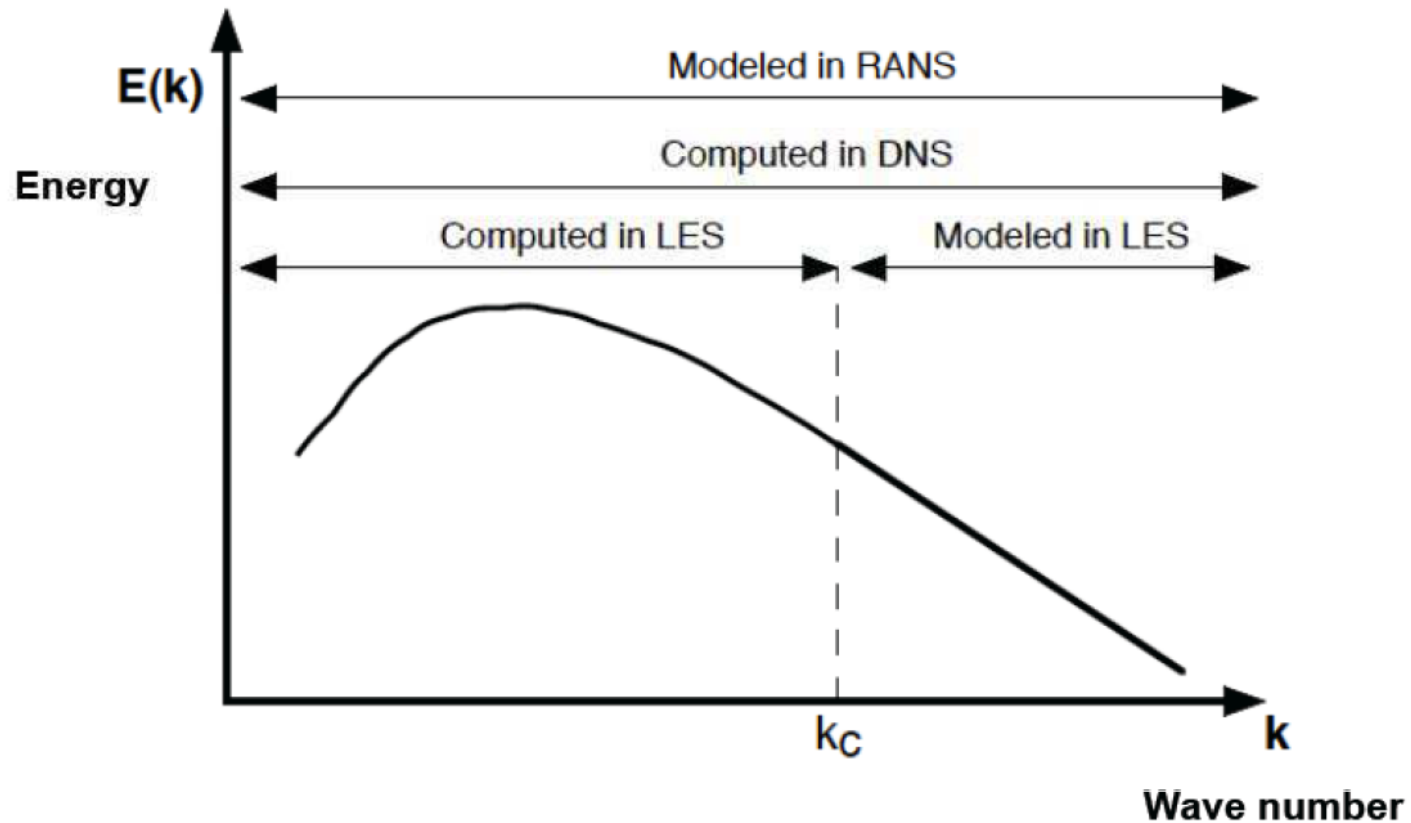


Figure 2: Energy spectrum of the steady-state solution of the Burgers equation for  $N = 20$  and  $N = 100$  (DNS). The steady state is reached at  $t = 3$  approximately.



## DNS vs LES vs RANS



# TURBULENCE MODELING: A PREVIEW

DNS

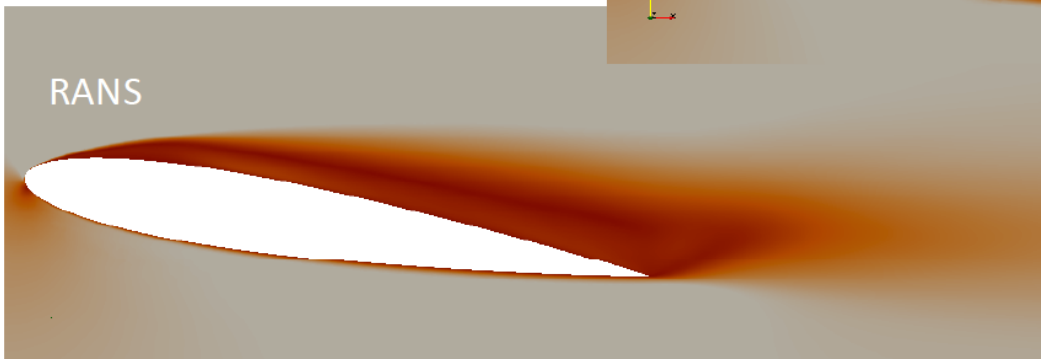


Flow past a NACA0012  
airfoil at  $\text{AoA}=9^\circ$   $\text{Re}=5\text{e}4$

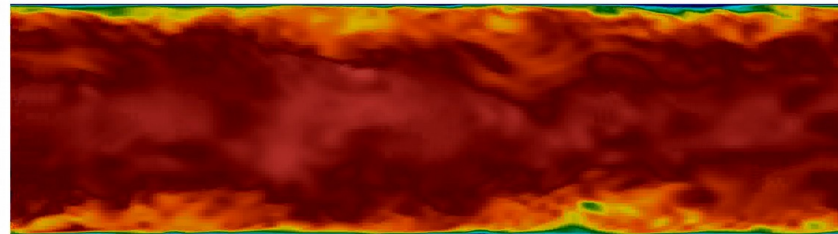
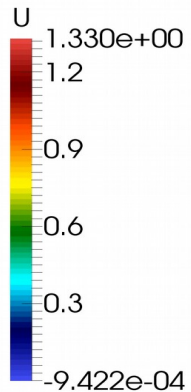
LES



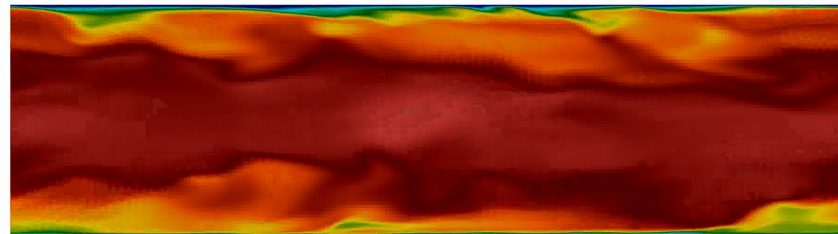
RANS



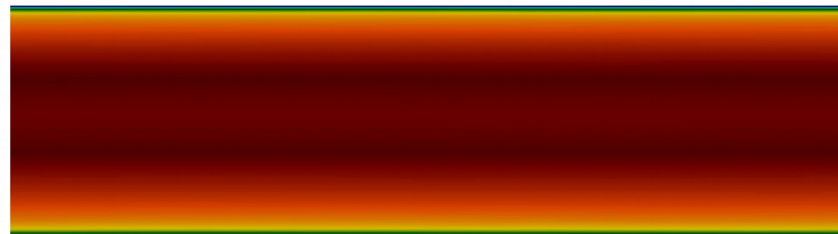
# TURBULENCE MODELING: A PREVIEW



DNS



LES



RANS

Turbulent channel flow at  $Re_\tau = 395$  ( $Re_{bulk} = 13760$ )

$$\nu_t(k/k_N) = \nu_t^{+\infty} \left( \frac{E_{k_N}}{k_N} \right)^{1/2} \nu_t^* \left( \frac{k}{k_N} \right) \quad (19)$$

with

$$\nu_t^{+\infty} = 0.31 \frac{5-m}{m+1} \sqrt{3-m} C_K^{-3/2} \quad (20)$$

where  $m$  is the slope of the energy spectrum, that is  $k^{-m}$ ,  $E_{k_N}$  is the energy at the cutoff frequency,  $k_N$ , and  $C_K$  is the Kolmogorov constant.  $\nu_t^*$  is a non-dimensional eddy-viscosity equal to 1 for small values of  $k/k_N$  and with a strong increase for higher  $k$  up to  $k/k_N = 1$ ; it reads

$$\nu_t^* \left( \frac{k}{k_N} \right) = 1 + 34.5 e^{-3.03(k_N/k)} \quad (21)$$

Note that the classical  $\nu_t(k/k_N)$  is recovered for  $m = 5/3$ . In our case, the energy spectrum is approximately  $m = 2$  (see Figure 2) and the Kolmogorov constant (for 1D Burgers equation) is  $C_K \approx 0.4523$ . In our case,  $\nu = Re^{-1}$ , and therefore, in practice the only modification in the code implies to modify  $\nu$  by  $\nu_{eff}(k) = \nu + \nu_t(k)$ . Try to reproduce the results reported in Figure 3.



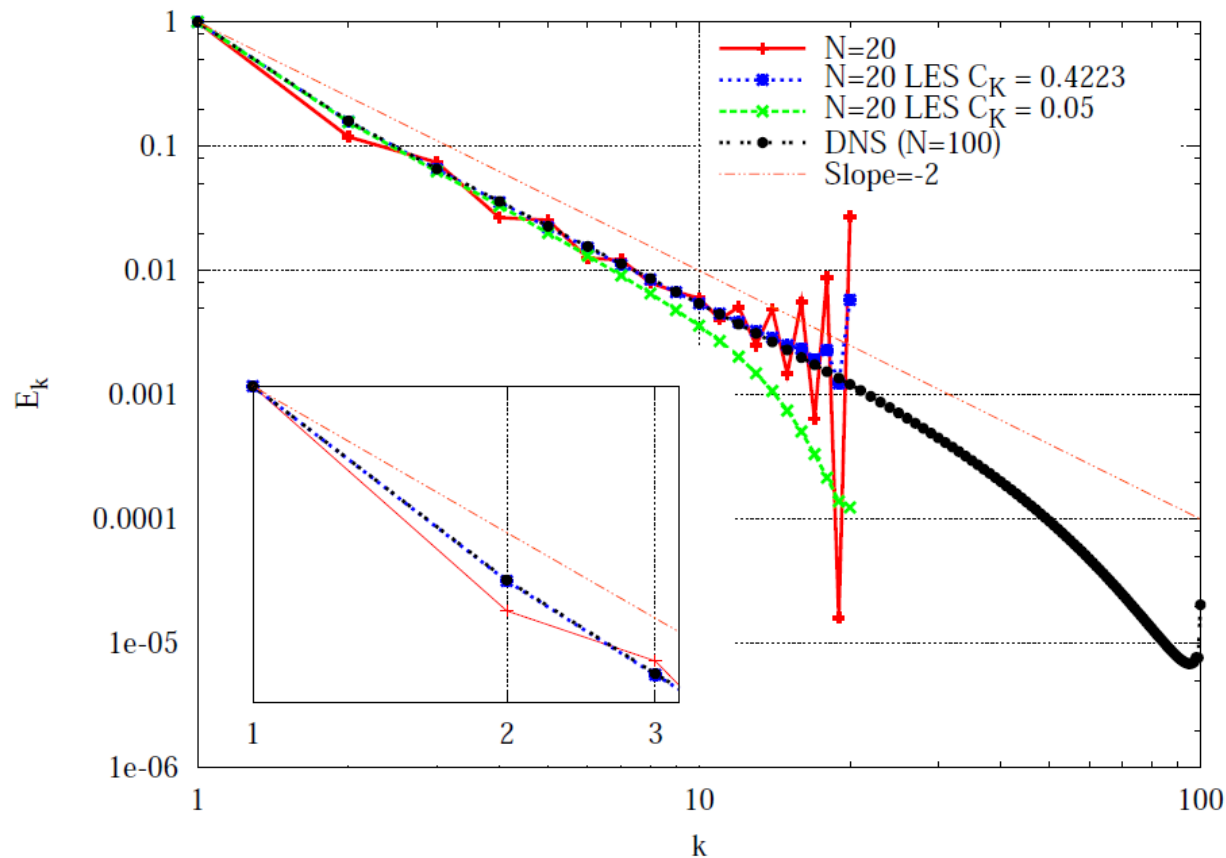


Figure 3: Energy spectrum of the steady-state solution of the Burgers equation for  $N = 20$  (with spectral eddy-viscosity model taking  $m = 2$  and  $C_K = 0.4523$ ) and  $N = 100$  (DNS).

- Note that although  $E_k \in \mathbb{R}$ ,  $\hat{u}_k \in \mathbb{C}$ . Therefore, your code must be able to perform basic operations with complex numbers.
- You can use a fully explicit time-integration scheme. Be careful with  $\Delta t$ !! A CFL-like condition must be imposed<sup>4</sup>

$$\Delta t < C_1 \frac{Re}{N^2} \quad (22)$$

- All the results reported here correspond to the steady state solution. You can also visualize the temporal evolution, play with other  $Re$ -numbers, different initial conditions, etc... It would help you to understand better the role of each term.