

# DISCRETE PHYSICAL REALM

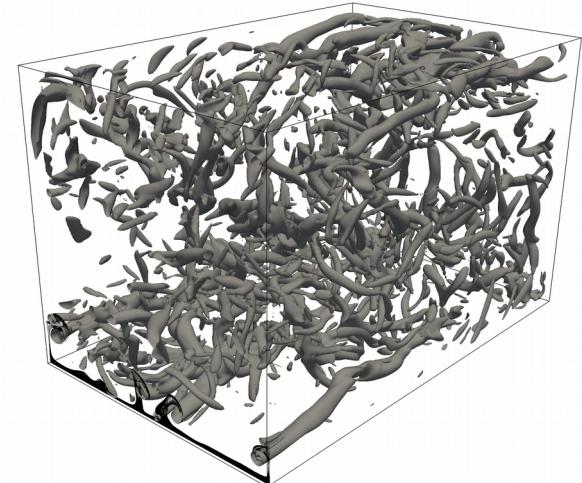
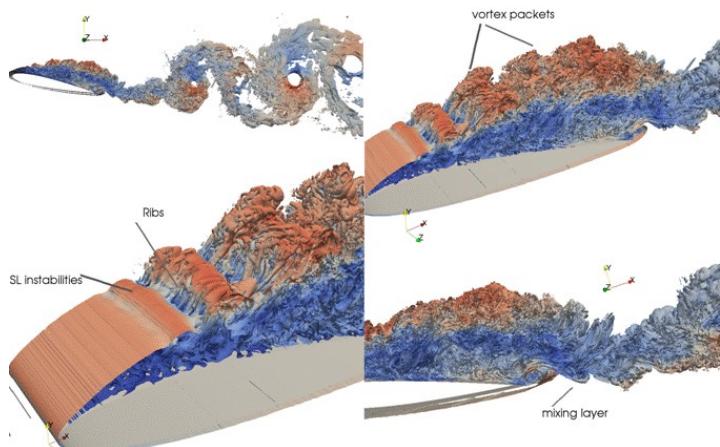
## ALGEBRAIC FORMULATION OF PHYSICS

A. BÁEZ, F.X. TRIAS AND N. VALLE



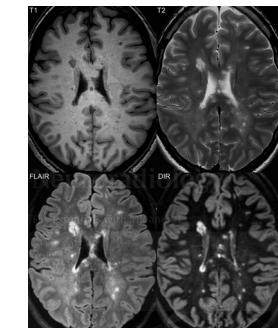
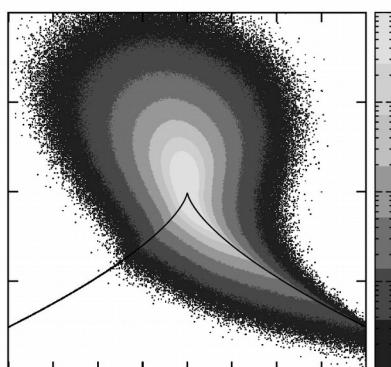
Centre Tecnològic de Transferència de Calor  
UNIVERSITAT POLITÈCNICA DE CATALUNYA





# INVARIANTS

## FLOW TOPOLOGY, VISUALIZATION AND TURBULENCE MODELING



*“First, solve the problem. Then, write the code.” John Johnson*



# WHAT IS AN INVARIANT?

**Def:** an invariant is a property, held by a class of mathematical objects, which **remains unchanged** when transformations of a certain type are applied to the objects. Namely,

Given an isomorphism  $T: U \rightarrow U$

$I(x)$  is an invariant respect to  $T \Leftrightarrow I(T(x)) = I(x), \forall x \in U$

$U$	$I(x)$	$T: U \rightarrow U$
$\mathbb{R}^n$	Euclidean distance between 2 points	{Space translations; Rotations; Reflexions;...}
$\mathbb{R}^{n \times n}$	{Determinant; Trace; Eigenvalues; Eigenvectors; conditioning number;...}	Changes of basis
$\mathbb{C}$	$\mathbb{R}(x)$	$\mathbb{R}(\bar{x}) = \mathbb{R}(x)$

Let us consider a  $3 \times 3$  tensor:  $\mathbf{A} \in \mathbb{R}^{3 \times 3}$

Three eigenvalues of  $\mathbf{A}$ :  $\lambda_1 \geq \lambda_2 \geq \lambda_3$

They are solutions of the characteristic equation:

$$\det(\lambda \mathbf{I} - \mathbf{A}) = \lambda^3 - P\lambda^2 + Q\lambda - R = 0$$

$$P = \lambda_1 + \lambda_2 + \lambda_3 \quad \text{← First invariant}$$

$$Q = \lambda_1 \lambda_2 + \lambda_1 \lambda_3 + \lambda_2 \lambda_3 \quad \text{← Second invariant}$$

$$R = \lambda_1 \lambda_2 \lambda_3 \quad \text{← Third invariant}$$

Caley-Hamilton theorem:

$$\lambda^3 - P\lambda^2 + Q\lambda - R = 0 \Rightarrow A^3 - PA^2 + QA - RI = 0$$

Trace:  $tr(A^3) - P tr(A^2) + Q tr(A) - 3R = 0$

$$P = tr(A)$$

← First invariant

$$Q = \frac{1}{2} (tr^2(A) - tr(A^2)) \quad ← Second invariant$$

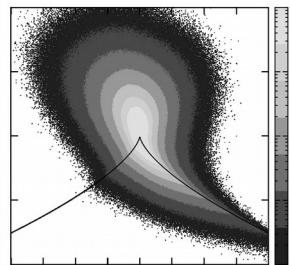
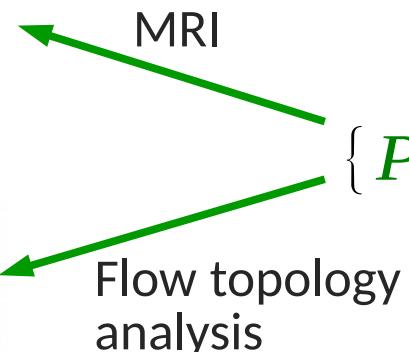
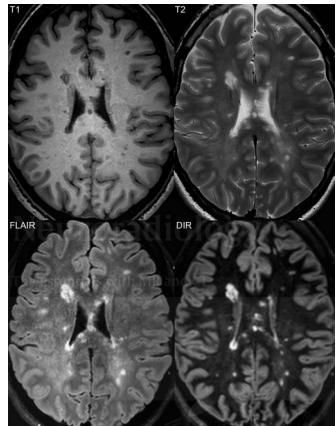
$$R = \det(A) = \frac{1}{6} (tr^3(A) - 3tr(A)tr(A^2) + 2tr(A^3))$$

← Third invariant

# TENSOR INVARIANTS: THEORY

$$\{\lambda_1, \lambda_2, \lambda_3\} \Leftrightarrow \{P, Q, R\} \Leftrightarrow \{tr(A), tr(A^2), tr(A^3)\}$$

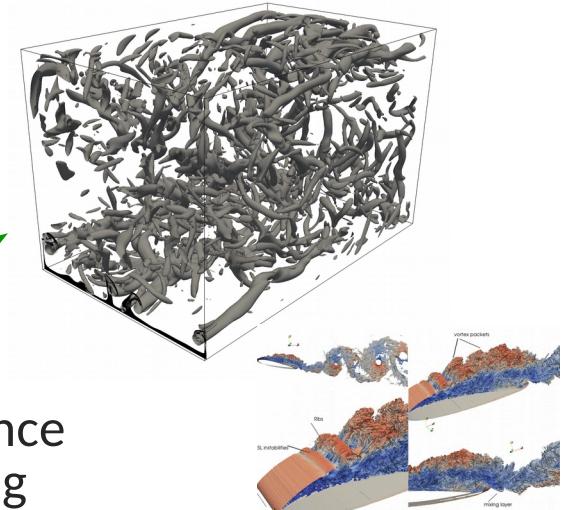
**P**, **Q** and **R** have a clear physical meaning and they are easier to compute



Flow visualization

Turbulence modeling

$$v_e^{S_{mag}} = (C_s \Delta)^2 |S(\bar{u})| = 2(C_s \Delta)^2 \sqrt{-Q_S}$$



$$a_{ij} = [\mathbf{A}]_{ij} \quad b_{ij} = [\mathbf{B}]_{ij} \quad c_{ij} = [\mathbf{C}]_{ij}$$

To know this  
is a MUST!



Products:

$$[\mathbf{AB}]_{ik} = a_{ij} b_{jk} \quad [\mathbf{AB}^T]_{ik} = a_{ij} b_{kj} \quad [\mathbf{ABC}]_{ik} = a_{ij} b_{jk} c_{km}$$

Tensor contraction (trace):

$$\text{tr}(\mathbf{A}) = a_{ii} \quad \text{tr}(\mathbf{AB}) = a_{ij} b_{ji} \quad \text{tr}(\mathbf{ABC}) = a_{ij} b_{jk} c_{ki}$$

So easy to show that  $\text{tr}(\mathbf{ABC}) = \text{tr}(\mathbf{BCA}) = \text{tr}(\mathbf{CAB})$

$$a_{ij} b_{jk} c_{ki} = b_{ij} c_{jk} a_{ki} = c_{ij} a_{jk} b_{ki}$$

Kronecker delta:  $\delta_{ij} = \begin{cases} 0 & i \neq j \\ 1 & i = j \end{cases}$

Levi-Civita tensor:  $\epsilon_{ijk} = \begin{cases} 1 & \text{3D cycle } (i,j,k) \text{ is even} \\ -1 & \text{3D cycle } (i,j,k) \text{ is odd} \\ 0 & \text{otherwise} \end{cases}$

$$\mathbf{e}_i \mathbf{e}_j = \delta_{ij} \quad [\nabla \times \mathbf{a}]_k = \epsilon_{ijk} \partial_i a_j \quad [\mathbf{a} \times \mathbf{b}]_k = \epsilon_{ijk} a_i b_j$$

$$\nabla \cdot \mathbf{a} = \partial_i a_i \quad \nabla \mathbf{a} = \partial_i a_j \quad \nabla^2 \mathbf{a} = \partial_{ii} a_j$$

To know this  
is a MUST!



It is easier to prove vector calculus identities, e.g.

$$\nabla \cdot (\nabla \mathbf{a})^T = \nabla \nabla \cdot \mathbf{a} \rightarrow \partial_i \partial_j a_i = \partial_j \partial_i a_i$$

$$\nabla \times \nabla \varphi = 0 \rightarrow \epsilon_{ijk} \partial_i \partial_j \varphi = \epsilon_{jik} \partial_i \partial_j \varphi = 0$$

$$\nabla \cdot \nabla \times \mathbf{a} = 0 \rightarrow \partial_k \epsilon_{ijk} \partial_i a_j = \epsilon_{ijk} \partial_k \partial_i a_j = \epsilon_{jik} \partial_k \partial_i a_j = 0$$

$$\nabla \times \nabla^2 \mathbf{a} = \nabla^2 \nabla \times \mathbf{a} \rightarrow \epsilon_{ijk} \partial_i \partial_{nn} a_j = \partial_{nn} \epsilon_{ijk} \partial_i a_j$$

# VELOCITY GRADIENT TENSOR

$$\mathbf{G} \equiv \nabla \vec{u}$$

Rate-of-strain tensor

$$\mathbf{S} = \frac{1}{2} (\mathbf{G} + \mathbf{G}^T)$$

Rate-of-rotation tensor

$$\boldsymbol{\Omega} = \frac{1}{2} (\mathbf{G} - \mathbf{G}^T)$$

$$P_G = \text{tr}(\mathbf{G}) = \nabla \cdot \vec{u} = 0$$

$$P_S = \text{tr}(\mathbf{S}) = 0$$

$$P_\Omega = \text{tr}(\boldsymbol{\Omega}) = 0$$

$$Q_G = -\frac{1}{2} \text{tr}(\mathbf{G}^2)$$

$$Q_S = -\frac{1}{2} \text{tr}(\mathbf{S}^2)$$

$$Q_\Omega = -\frac{1}{2} \text{tr}(\boldsymbol{\Omega}^2)$$

$$R_G = \frac{1}{3} \text{tr}(\mathbf{G}^3)$$

$$R_S = \frac{1}{3} \text{tr}(\mathbf{S}^3)$$

$$R_\Omega = \frac{1}{3} \text{tr}(\boldsymbol{\Omega}^3)$$

$$\text{tr}(\mathbf{S}^2) = S_{ij} S_{ji} = S_{ij} S_{ij} \rightarrow Q_S = -1/2 (\mathbf{S} : \mathbf{S}) \leq 0 \leftrightarrow \varepsilon = 2 \sqrt{\mathbf{S} : \mathbf{S}}$$

$$\text{tr}(\boldsymbol{\Omega}^2) = \Omega_{ij} \Omega_{ji} = -\Omega_{ij} \Omega_{ij} \rightarrow Q_\Omega = 1/2 (\boldsymbol{\Omega} : \boldsymbol{\Omega}) = 1/2 |\vec{\omega}|^2 \geq 0$$

$$\text{tr}(\mathbf{S} \boldsymbol{\Omega}) = S_{ij} \Omega_{ji} = -S_{ji} \Omega_{ij} = 0 \rightarrow$$

$$Q_G = Q_S + Q_\Omega$$

$$\mathbf{G} \equiv \nabla \vec{u}$$

Rate-of-strain tensor

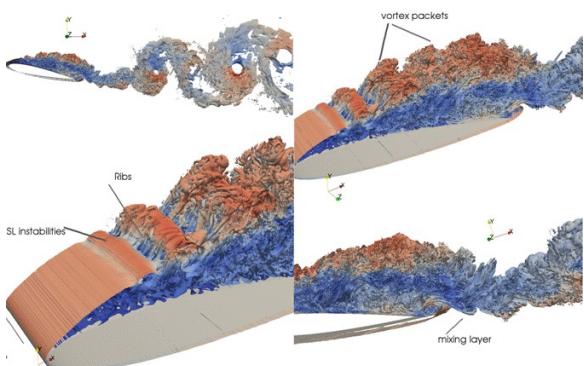
$$\mathbf{S} = \frac{1}{2} (\mathbf{G} + \mathbf{G}^T)$$

Rate-of-rotation tensor

$$\boldsymbol{\Omega} = \frac{1}{2} (\mathbf{G} - \mathbf{G}^T)$$

$$Q_S = -1/2 (\mathbf{S} : \mathbf{S}) \leq 0 \longleftrightarrow \varepsilon = 2 \nu \mathbf{S} : \mathbf{S}$$

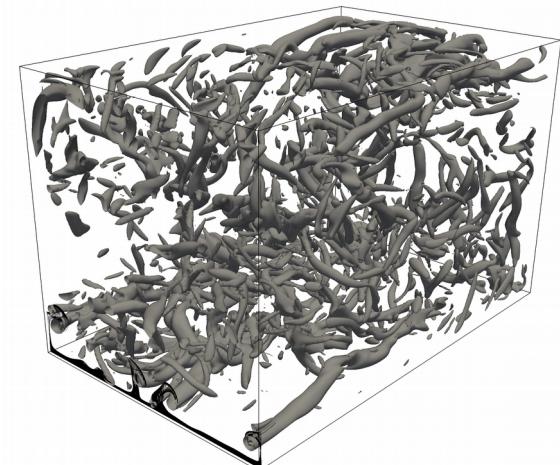
$$Q_\Omega = 1/2 (\boldsymbol{\Omega} : \boldsymbol{\Omega}) = 1/2 |\vec{\omega}|^2 \geq 0$$



$$Q_G = Q_S + Q_\Omega$$

$$Q_G > 0$$

Tube-like structures!



**Remark:** this is the well-known  $Q$ -criterion for flow visualization

$$\mathbf{G} \equiv \nabla \vec{u}$$

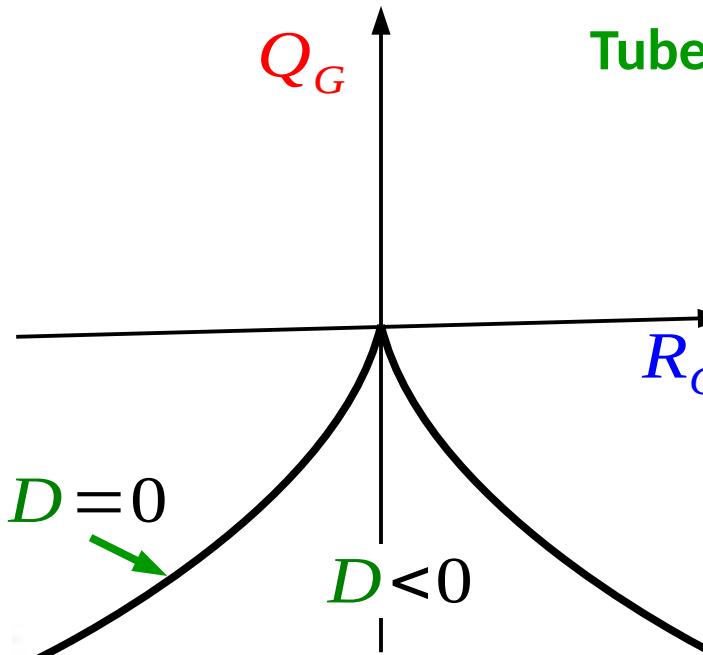
$$Q_G = Q_S + Q_\Omega$$

$$\lambda^3 + Q\lambda - R = 0$$

$$R = \lambda_1 \lambda_2 \lambda_3$$

$$D = (27/4) R^2 + Q^3$$

Discriminant



$Q_G > 0$   
Tube-like structures!

$$D < 0 \Rightarrow \lambda_1, \lambda_2, \lambda_3 \in \mathbb{R}$$

$$D > 0 \Rightarrow \lambda_1 \in \mathbb{R}, \quad \lambda_2 = \bar{\lambda}_3 \in \mathbb{C}$$

Vortex!

$$D > 0, R > 0 \Rightarrow \lambda_1 > 0$$

vortex stretching

$$D > 0, R < 0 \Rightarrow \lambda_1 < 0$$

vortex compression

In these figures



$$\lambda^3 + Q\lambda + R = 0$$

$$R = -\lambda_1 \lambda_2 \lambda_3$$

...instead of

$$\lambda^3 + Q\lambda - R = 0$$

$$R = \lambda_1 \lambda_2 \lambda_3$$

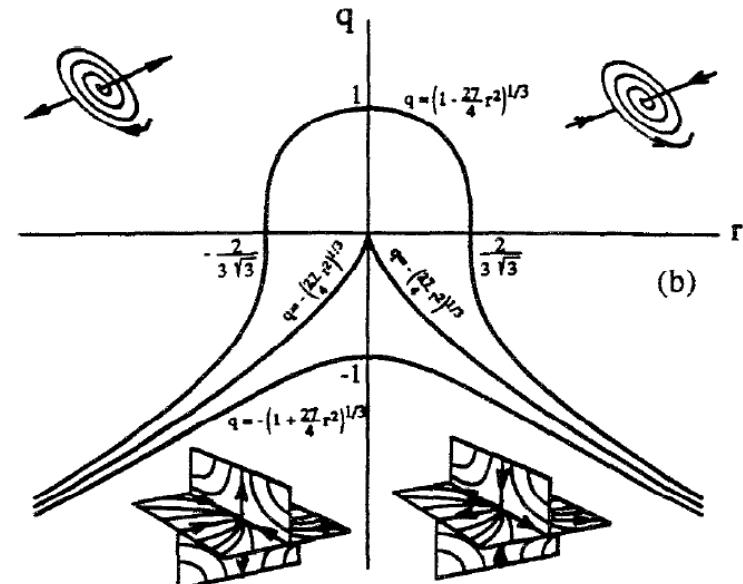
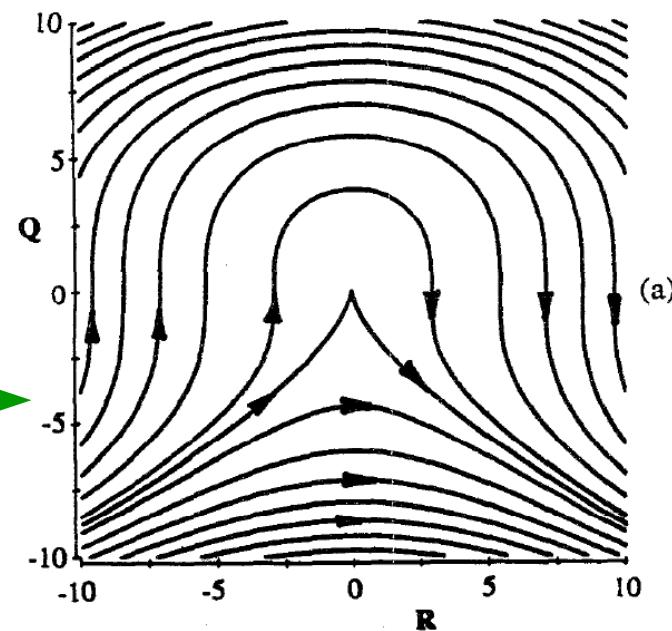


FIG. 1. Trajectories of the invariants of the velocity gradient tensor: (a) dimensioned  $Q$  vs  $R$ , (b) dimensionless  $q$  vs  $r$  with schematic diagrams of local streamlines.

Figure extracted from: B. J. Cantwell, "Exact solution of the restricted Euler equation for the velocity gradient tensor," Phys. Fluids A4, 782 (1992).

In these figures



$$\lambda^3 + Q\lambda + R = 0$$

$$R = -\lambda_1 \lambda_2 \lambda_3$$

...instead of

$$\lambda^3 + Q\lambda - R = 0$$

$$R = \lambda_1 \lambda_2 \lambda_3$$

Flow dynamics for NS equations

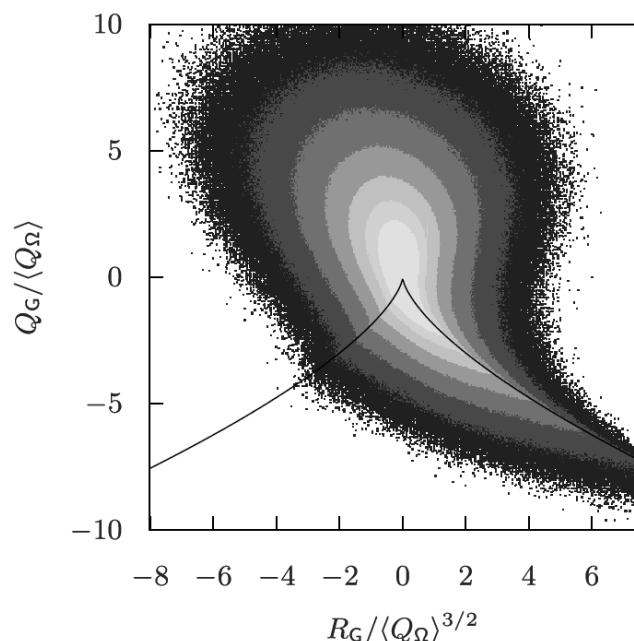
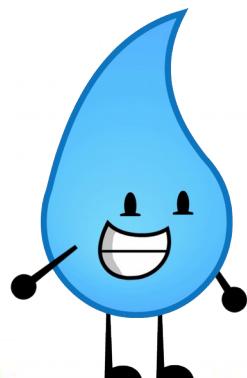


Figure extracted from: F. Dabbagh, F.X. Trias, A. Gorobets, and A. Oliva. "On the evolution of flow topology in turbulent Rayleigh-Bénard convection", Physics of Fluids, 28:115105, 2016.

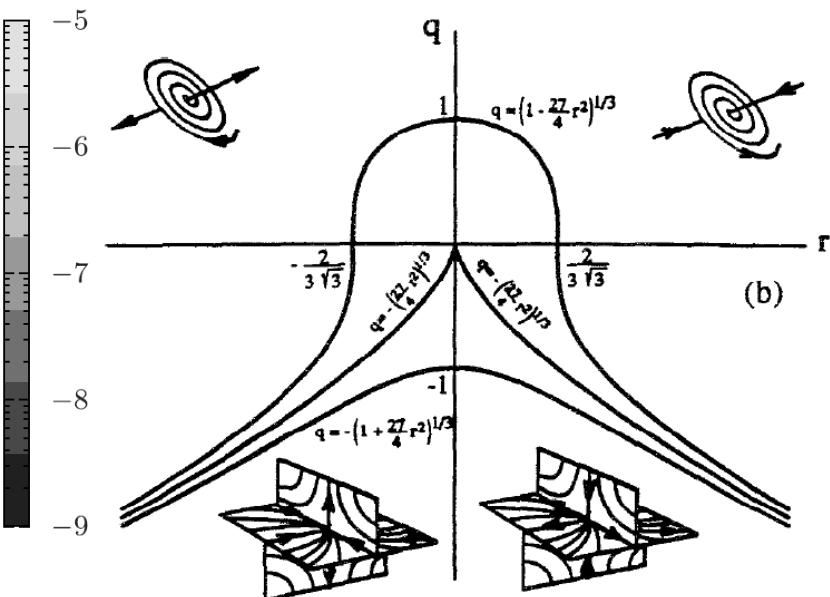
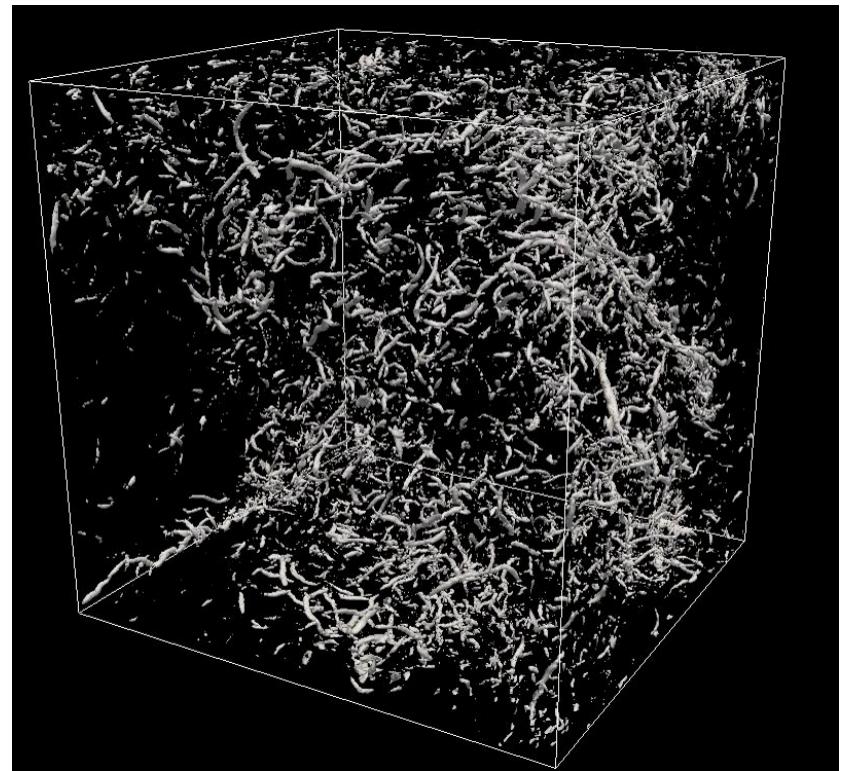
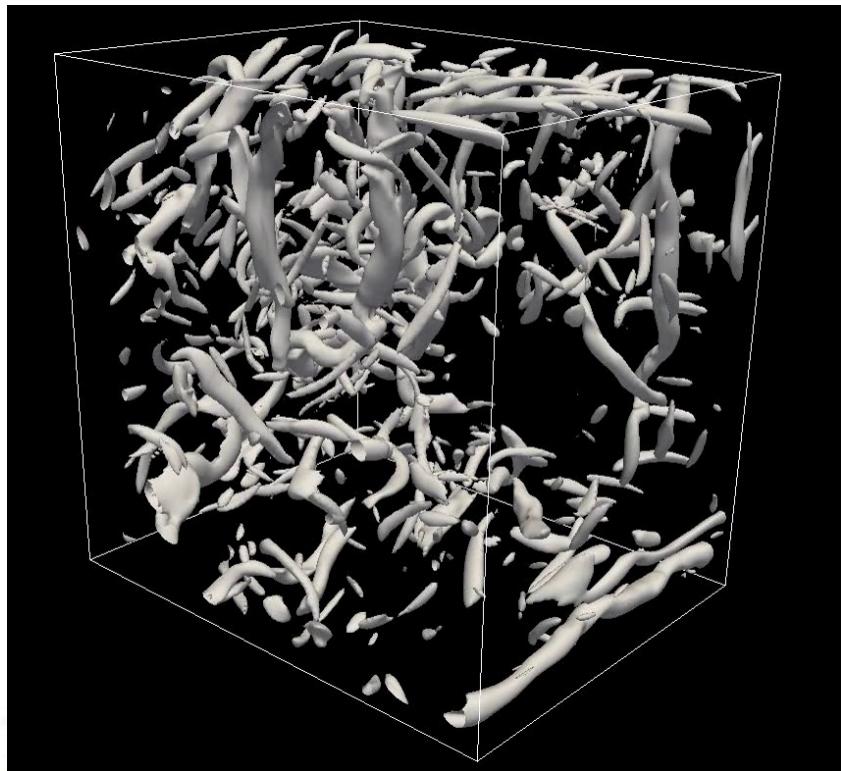


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$$Q_G = Q_S + Q_\Omega$$

$Q_G > 0$  Tube-like structures!



# TURBULENCE MODELING

$$V^2 = |S \vec{\omega}|^2$$

$$\{ Q_G, R_G, Q_S, R_S, V^2 \}$$

← Unified framework!

**Smagorinsky (1963)**

J. Smagorinsky, Mon. Weather Rev. 91, 99–164 (1963)

$$v_e^{Smag} = (C_S \Delta)^2 |S(\bar{u})| = 2(C_S \Delta)^2 (-Q_S)^{1/2}$$

**WALE (1999)**

F. Nicoud and F. Ducros, Flow, Turbul. Combust. 62(3), 183–200 (1999)

$$v_e^W = (C_W \Delta)^2 \frac{(V^2 + 2Q_G^2/3)^{3/2}}{(-2Q_S)^{5/2} + (V^2/2 + 2Q_G^2/3)^{5/4}}$$

**Vreman's (2004)**

A.W. Vreman, Phys. Fluids 16(10), 3670–3681 (2004)

$$v_e^{Vr} = (C_{Vr} \Delta)^2 \left( \frac{V^2 + Q_G^2}{2(Q_G - 2Q_S)} \right)^{1/2}$$

**QR-model (2011)**

R. Verstappen, J. Sci. Comput. 49(1), 94–110 (2011)

$$v_e^{Ve} = (C_{Ve} \Delta)^2 \frac{|R_S|}{-Q_S}$$

**Sigma (2011)**

F. Nicoud et al., " Phys. Fluids 23(8), 085106 (2011)

$$v_e^\sigma = (C_\sigma \Delta)^2 \frac{\sigma_3(\sigma_1 - \sigma_2)(\sigma_2 - \sigma_3)}{\sigma_1^2}; \quad \sigma_i = \sqrt{\lambda_i^{GG}}$$

**S3QR (2015)**

F.X.Trias et al. Phys. Fluids, 27: 065103, (2015)

$$v_e^{S3QR} = (C_{S3QR} \Delta)^2 \frac{R_G^{5/3}}{Q_G - 2Q_S}$$

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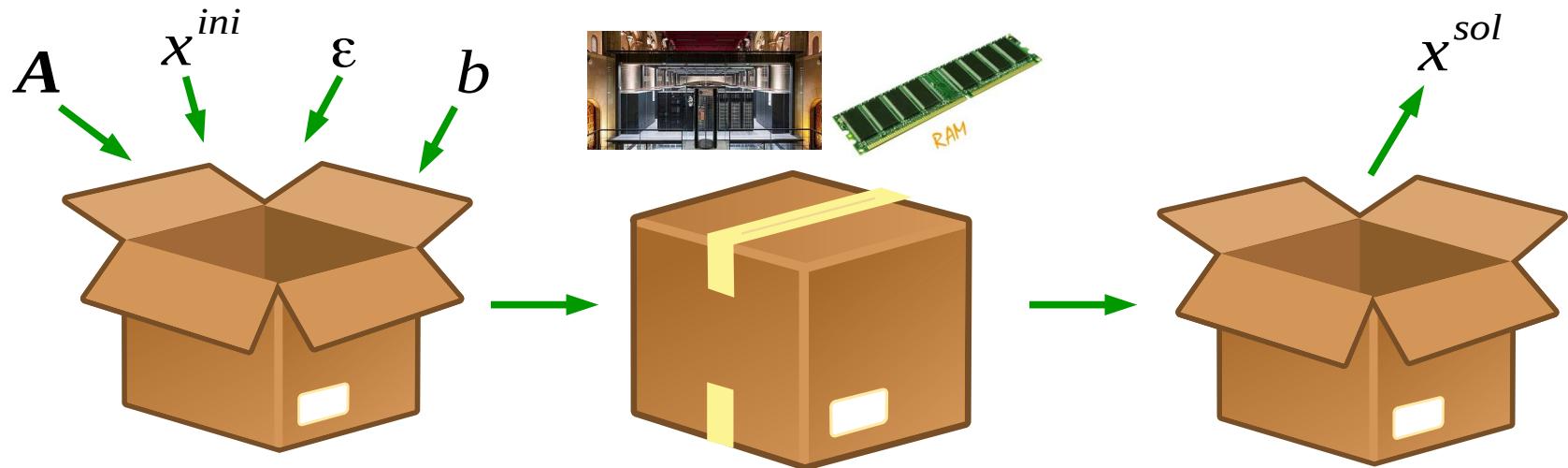
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**BONUS**

# BONUS: LINEAR SOLVERS

$$\boxed{A x = b} \rightarrow x^{sol} \approx A^{-1} b \quad \text{where } \|A x^{sol} - b\| \leq \varepsilon$$



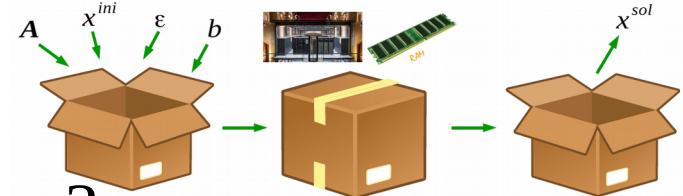
Easy to test!

$$x^{rand} \rightarrow b = A x^{rand} \xrightarrow[\varepsilon]{A} \boxed{\text{?}} \rightarrow x^{sol} \rightarrow x^{sol} \approx x^{rand}$$

## BONUS: LINEAR SOLVERS

$$\boxed{\mathbf{A}x = b} \rightarrow x^{sol} \approx \mathbf{A}^{-1}b \quad \text{where } \|\mathbf{A}x^{sol} - b\| \leq \varepsilon$$

How to choose a solver? 



- What type of matrix/problem do I have?  
dense or sparse?  $\mathbf{A} = \mathbf{A}^T$ ? diagonal dominant?...  
circulant? block diagonal? condition number?...  
 $\mathbf{A}x_i = b_i, \quad i=1\dots N$ ; Is  $N$  large or small?
- What is  $\varepsilon$ ?  
do I need machine accuracy or  $\varepsilon=0.001$  is enough?
- How are  $x^{ini}$  and  $b$ ?  
is  $x^{ini}$  a good guess of  $x^{sol}$ ? what is the spectrum of  $b$ ?
- How is my computer?  
parallel? memory shared or distributed?

# BONUS: LINEAR SOLVERS

Typical linear solvers in CFD:

To know this  
is a MUST!



Direct

Diagonalizable:

e.g. circulant matrix (easily solved with a FFT )

LU-based:  $Ax = LUx = b \Rightarrow y = L^{-1}b \Rightarrow x = U^{-1}y$

e.g. TDMA, band-LU, Cholesky (if  $A = A^T$ ),

Schur complement (for block/partitioned  $A$ )

Stationary:  $A = \tilde{A} + (A - \tilde{A}) \Rightarrow \tilde{A}x_{i+1} = b - (A - \tilde{A})x_i$

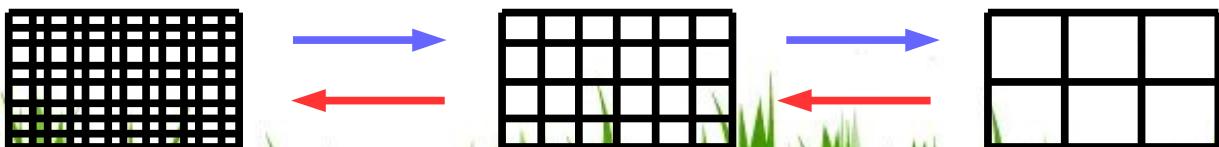
e.g. Jacobi ( $\tilde{A} = D$ ), Gauss-Seidel ( $\tilde{A} = D + L$ ),...

Non-stationary: e.g. Krylov subspace (CG if  $A = A^T$ ),  
Multigrid solvers:

**Smoothen** (e.g. Jacobi) easily solves high freqs

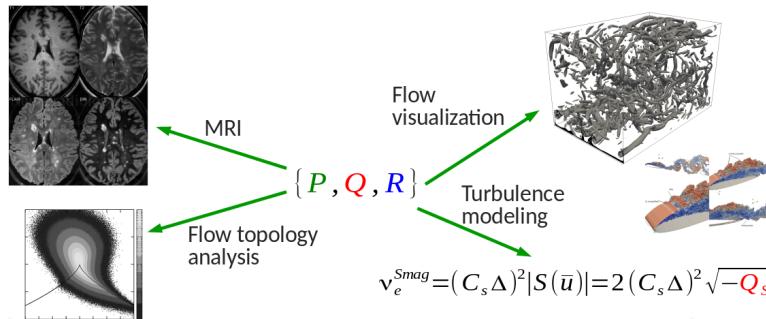


Idea: transform low freqs into high freqs ( **restriction** )  
solve them and come back ( **prolongation** )



# Take-away messages

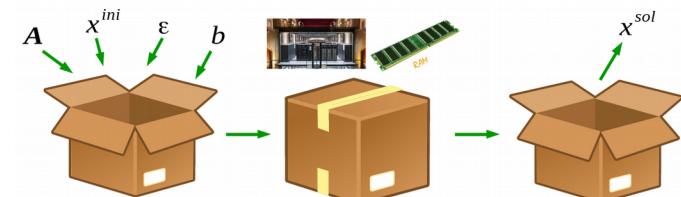
- **Invariants** play a very important role in CFD: e.g. flow visualization (**Q**-criterion), turbulence modeling, flow topology analysis...



- (Sparse) **linear algebra** forms the lowest level of a CFD code. Hence, **few highly optimized basic operators** can be the basis of your code.

$$y \leftarrow Ax; \quad \alpha \leftarrow x \cdot y; \quad y \leftarrow \alpha x + y$$

**SpMV**      **dot**      **axpy**



# Open questions

