

PHYSICAL FORMULATION

MATHEMATICS AS THE LANGUAGE OF PHYSICS

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EQUATIONS

CONSTRAINING SOLUTIONS

- Concept of operator
- Link between algebra and calculus
- Operator properties
- Function spaces as vector spaces

1)Definition

2)Physical space

3)Analytical solutions

Equations impose an equality condition to the solution.

They can be expressed with one or more operators.

$$Ax = b$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (u \rho) = 0$$

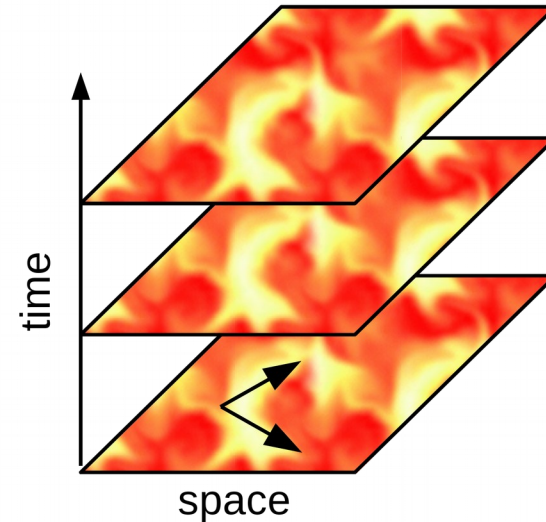
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When approaching physical problems, our solution will depend, at least, on space and time.

$$\frac{\partial \theta}{\partial t} - \nabla^2 \theta = 0$$

Solution is the combination of spatial and temporal evolution.

$$S = T \otimes X$$



- 1) Relativistic effects are neglected. Measures are independent!
- 2) Operators are naturally split. Temporal and spatial operators.

Let's take a diffusion equation as an example:

$$\frac{\partial \theta}{\partial t} - \nabla^2 \theta = 0$$

$$\theta|_{x=\partial\Omega} = 0 \quad \forall t \geq 0$$

$$\theta|_{t=0} = u_0(x) \quad \forall x \in \Omega$$

BC: Boundary Conditions

IC: Initial conditions

Important:

- BC: Defined for all time at space boundaries only.
- IC: Defined for all space at time boundaries (i.e., initial time) only.
- More details in

[Introduction to PDE | MIT](#)

Separation of variables:

$$\theta = T(t) X(x)$$

$$X \frac{\partial T}{\partial t} - T \nabla^2 X = 0$$

$$X \partial_t T = T \nabla^2 X$$

$$\frac{\partial_t T}{T} = \frac{\nabla^2 X}{X} = \lambda$$

Important:

- λ must be a scalar.
- We end up with 2 separate eigenvalue problems:

$$\partial_t T = \lambda T$$

$$\nabla^2 X = \lambda X$$

- So we are diagonalizing the operators!

Eigenvalue problem:

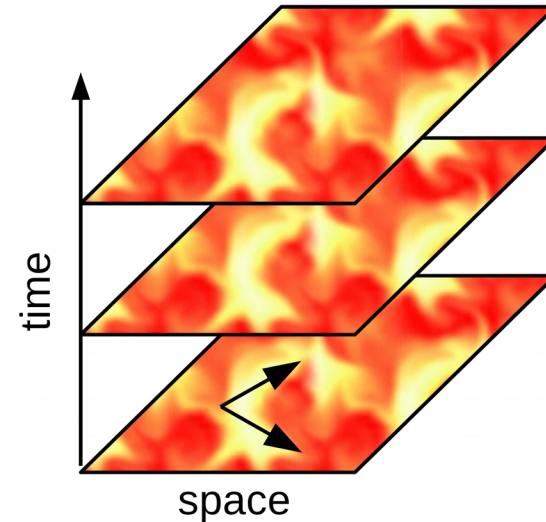
$$\partial_t T = \lambda T$$

$$\nabla^2 X = \lambda X$$

As with algebraic operators, we try to express it in a diagonal form.

$$Ax = P \Lambda P^{-1} x$$

Subject to a base P.



Every operator has an optimal way to look at it. It is the diagonal form.

Eigenvalue problem:

$$\partial_t T = \lambda T$$
$$\nabla^2 X = \lambda X$$

As with algebraic operators, we try to express it in a diagonal form.

$$Ax = P \Lambda P^{-1} x$$

Subject to a base P.

Important:

- P is the base spanned by eigenvectors.
- Λ components are the eigenvalues.
- Differential operators always diagonalize with base e.

$$\frac{d^n}{dt^n} \exp(v t) = v^n \exp(v t)$$

X eigenvalue problem:

$$\nabla^2 X = \lambda X$$

Eigenfunctions look like:

$$X(x) = a_n e^{(\sqrt{\lambda_n} x)} + b_n e^{(-\sqrt{\lambda_n} x)}$$

T eigenvalue problem:

$$\partial_t T = \lambda T$$

Eigenfunctions looks like:

$$T(t) = d_n e^{(\lambda_n t)}$$

Finally, the **solution space** will look like:

$$\theta(x, t) = T(t) X(x) = \sum_{n=-\infty}^{+\infty} d_n e^{(\lambda_n t)} (a_n e^{(\sqrt{\lambda_n} x)} + b_n e^{(-\sqrt{\lambda_n} x)})$$

Important:

- We have composed time with space to obtain our solution.
- The solution is an *infinite-dimensional vector space!*
- Values of λ are, potentially, the full set of reals.
- We still have to determine the coefficients and the eigenvalues.

Eigenvalues

Boundary Conditions: $\theta|_{x=\partial\Omega}=0 \forall t$

In 1D: $\theta|_{x=0}=\theta|_{x=L}=0 \forall t$

By expanding solution: $X(0)=X(L)=0$

For convenience: $\pm\sqrt{\lambda_n}=\pm\mu_n$

We end up with:

$$\begin{pmatrix} X(0) \\ X(L) \end{pmatrix} = \begin{bmatrix} e^{(\mu_n 0)} & e^{(-\mu_n 0)} \\ e^{(\mu_n L)} & e^{(-\mu_n L)} \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$
$$\begin{bmatrix} B \end{bmatrix} \quad \begin{bmatrix} c_n \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix}$$

Eigenvalues

To find a compatible solution:

$$0 = \exp(-\mu_n L) - \exp(\mu_n L) \rightarrow$$

$$\mu_n = n \frac{\pi}{L} i$$

Which, substituting back in B gives:

$$\begin{bmatrix} 1 & 1 \\ e^{(n\pi i)} & e^{(-n\pi i)} \end{bmatrix} \begin{bmatrix} a_n \\ b_n \end{bmatrix} = \begin{bmatrix} 0 \\ 0 \end{bmatrix} \rightarrow a_n = -b_n$$

And so:

$$X(x) = a_n (e^{\mu_n x} - e^{-\mu_n x}) = 2 a_n i \sin(\mu_n x)$$

Real case: $\mu_n \in \mathbb{R}$

$$\begin{pmatrix} 1 & 1 \\ e^{\mu_n L} & e^{-\mu_n L} \end{pmatrix} \begin{pmatrix} a_n \\ b_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$a_n = b_n = 0$$

Trivial solution!

We end up with:

$$X(x) = 0$$

Imaginary case: $i\mu_n \in \mathfrak{I}$

$$\begin{pmatrix} \cos(0) & \sin(0) \\ \cos(\mu_n L) & \sin(\mu_n L) \end{pmatrix} \begin{pmatrix} a_n \\ b_n \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

$$a_n = 0 \quad b_n \neq 0$$

$$\mu_n = n \frac{\pi}{L} \quad \forall n \in \mathbb{N} > 0$$

We end up with:

$$X(x) = \sum_{n=1}^{+\infty} b_n \sin\left(n\pi \frac{x}{L}\right)$$

At this point the eigenvectors look like:

$$\theta(x, t) = \sum_{n=1}^{+\infty} d_n e^{-\left(\frac{n\pi}{L}\right)^2 t} \sin\left(n\pi \frac{x}{L}\right) \quad \forall n \in \mathbb{N} > 0$$

We can obtain qualitative information at this point:

- Solution is going to smooth down with time.
- Higher spatial frequencies will dissipate faster.
- Characteristic time is $\tau = \left(\frac{L}{n\pi}\right)^2$
- There is no interaction between frequencies

Coefficients

Initial conditions: $\theta_0(x) = \theta(x, 0) = \sum_{n=1}^{+\infty} d_n \sin\left(n\pi \frac{x}{L}\right)$

Idea: $\langle \theta_0(x) | \theta_n(x, 0) \rangle = \langle \theta(x, 0) | \theta_n(x, 0) \rangle$

Project initial conditions over every base vector of the solution space

Because the solution space is orthogonal:

$$d_n = \frac{2}{L} \int_0^L \theta_0(x) \sin\left(n\pi \frac{x}{L}\right) dx$$

Coefficients: Development details.

$$\begin{aligned}\langle \theta(x, 0) | \theta_n(x, 0) \rangle &= \int_0^L \sum_{m=1}^{+\infty} d_m \sin\left(m\pi \frac{x}{L}\right) \sin\left(n\pi \frac{x}{L}\right) dx \\ &= \sum_{m=1}^{+\infty} d_m \int_0^L \sin\left(m\pi \frac{x}{L}\right) \sin\left(n\pi \frac{x}{L}\right) dx = d_n \int_0^L \sin^2\left(n\pi \frac{x}{L}\right) dx\end{aligned}$$

$$\int_0^L \sin\left(m\pi \frac{x}{L}\right) \sin\left(n\pi \frac{x}{L}\right) dx = \begin{cases} 0 \Leftrightarrow m \neq n \\ \sin^2\left(n\pi \frac{x}{L}\right) \Leftrightarrow m = n \end{cases}$$

$$\langle \theta(x, 0) | \theta_n(x, 0) \rangle = d_n \frac{L}{2} = \langle \theta_0(x) | \theta_n(x, 0) \rangle$$

$$d_n = \frac{2}{L} \int_0^L \theta_0(x) \sin\left(n\pi \frac{x}{L}\right) dx$$

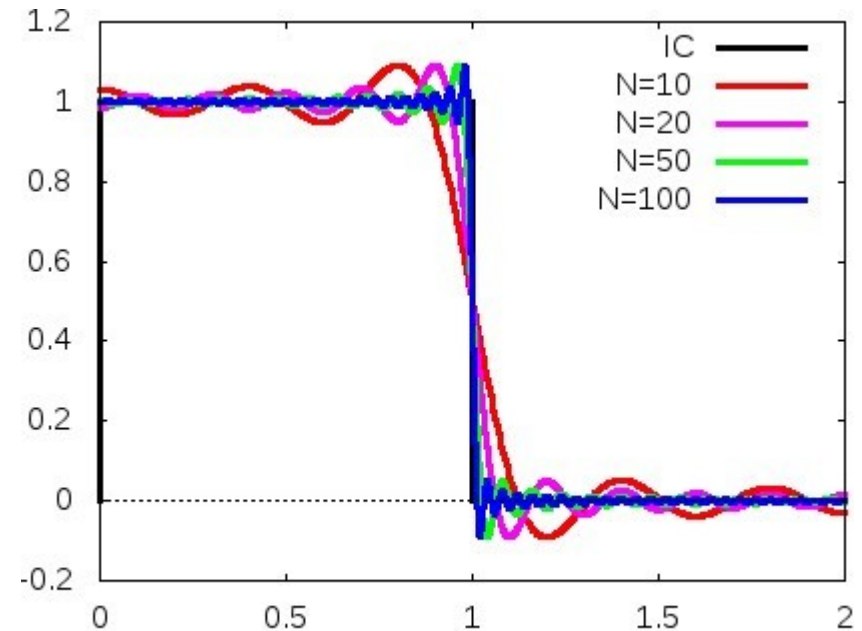
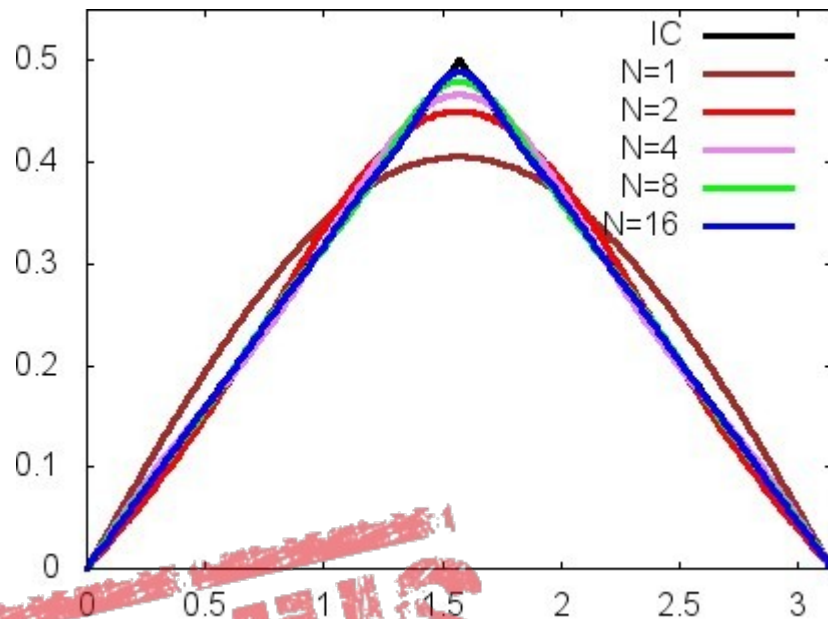
Finally, the solution will look like:

$$\theta(x, t) = \sum_{n=1}^{+\infty} d_n e^{-\left(\frac{n\pi}{L}\right)^2 t} \sin\left(n\pi \frac{x}{L}\right) \quad \forall n \in \mathbb{N} > 0$$
$$d_n = \frac{2}{L} \int_0^L \theta_0(x) \sin\left(n\pi \frac{x}{L}\right) dx$$

Important facts:

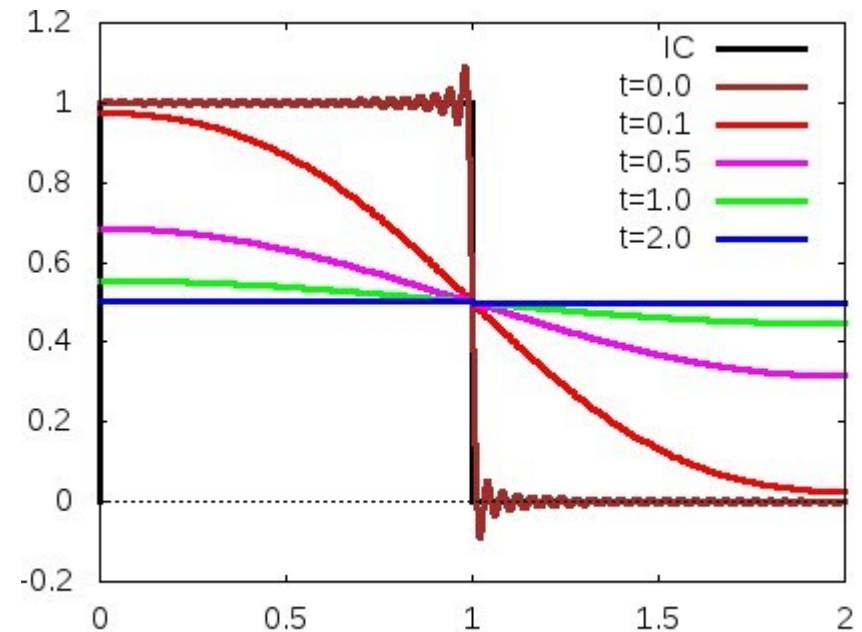
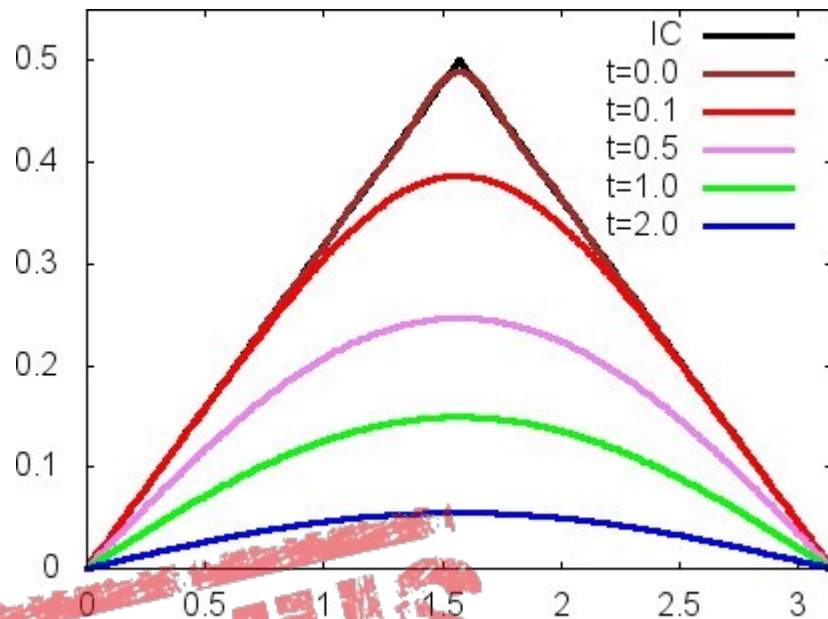
- Eigenvalues determine the **shape of the solution**.
- Eigenvalues are still infinite but **discrete**.
- Eigenvalues are determined by the **equation**
(i.e., equality, operators and boundary conditions)
- The solution base we are working with is **just a convenient base** of the solution space. We could use others!

Coefficients: Graphical example.



Note that this is valid for **continuous functions** only! Otherwise it appears **Gibbs phenomenon**.

Final solution evolution: Graphical example.



Note that this is valid for **continuous functions** only! Otherwise it appears **Gibbs phenomenon**.

Summary:

- Operators can be diagonalized in a **convenient base**.
- Solution space is going to be **infinite dimensional**.
- The basis is **guaranteed to be orthogonal** by the Sturm-Liouville theorem
- B.C. provide information about the solution structure.
- I.C. provide information about the scaling of such structure.

Coming soon...

What kind of information can we obtain from the following equations?

$$\frac{\partial \theta}{\partial t} - \nabla^2 \theta = 0$$

$$\frac{\partial \theta}{\partial t} + u \cdot \nabla \theta = 0$$

$$\nabla^2 \theta = 0$$

What can eigenvalues and eigenfunctions tell us?

See you in next session!