

Nikhil Patten  
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Dr. Moe  
ASTR5420

94%

Please show all work. If you collaborate with other students, write their names at the top of your homework. Please also write the total number of hours you spent on this. If you use any code or plotting routines, print those out and attach them to your hand-written solutions.

I collaborated with Alex on this homework.

This took me approximately 8 hours to complete.

### 1. (20%) Luminosity Relations

- (a) The first-order homologous relation for the mass conservation equation  $dM(r)/dr = 4\pi r^2 \rho(r)$  becomes  $M/R \propto R^2 \rho$  or  $\rho \propto M/R^3$ . Write first-order homologous relations for hydrostatic equilibrium ( $P$  in terms of  $M$  and  $R$ ) and radiative energy transport ( $T$  in terms of  $L$ ,  $M$ , and  $R$ ; you can assume opacity  $k_R$  is a constant due to Thompson electron scattering).

$$\begin{aligned}\frac{dP}{dR} &= -\rho g \\ \frac{dP}{dR} &\propto \frac{M}{R^3} \frac{M}{R^2} \\ dP &\propto \frac{M^2}{R^5} dR\end{aligned}$$

$$P \propto \frac{M^2}{R^4}$$

$$\begin{aligned}\frac{dT}{dR} &= -\frac{3k_R \rho}{64\pi R^2 \sigma T^3} L \\ \frac{dT}{dR} &\propto \frac{M}{R^3} \frac{1}{R^2 T^3} L \\ T^3 dT &\propto M L \frac{dR}{R^5} \\ T^4 &\propto \frac{M L}{R^4}\end{aligned}$$

$$T \propto (ML)^{1/4} \frac{1}{R}$$

- (b) Solar-type main-sequence stars are pressure supported by nearly an ideal gas, i.e.,  $P \propto \rho T$ . Now solve for the mass-luminosity  $L \propto M^x$  relation. Recall that  $x = 3.5$  is the actual exponent for solar-type MS stars. Is your approximation close to the true value?

$$\begin{aligned}\frac{M^2}{R^4} &\propto (ML)^{1/4} \frac{1}{R} \frac{M}{R^3} \\ M^2 &\propto L^{1/4} M^{1.25} \\ M^{0.75} &\propto L^{1/4}\end{aligned}$$

$$L \propto M^3$$

This approximation is reasonable close to the actual luminosity-mass relations for MS stars. We found a value of 3 for the exponent using the first order homologous approximations when the actual value is 3.5.

index  $n = 1$ ). More massive stars develop radiative cores while the rocky cores of less massive planets strongly affect their equation of state. For  $n = 1$  polytropes, solve for the radius  $R$  in terms of  $K$ ,  $\gamma$ , and other constants, independent of mass  $M$ . This question should demonstrate why Jovian planets, brown dwarfs, and late-M dwarfs all have roughly the same radius of  $R = 0.1R_{\odot} = 1R_J$  despite spanning 3 orders of magnitude in mass  $M = 0.3 - 300M_J$ .

$$P \propto K \rho^{\gamma} \quad (1)$$

$$\frac{M^2}{R^4} \propto K \left( \frac{M}{R^3} \right)^2 \quad (2)$$

$$\frac{M^2}{R^4} \propto K \frac{M^2}{R^6} \quad (3)$$

$$R^2 \propto K \quad (4)$$

$$\boxed{R \propto K^{0.5}}$$

Wanted you to keep constants through.  
 $R = \sqrt{\frac{K\pi}{2G}} \Rightarrow$  but right idea

As shown above, radius is independent of mass for fully-convective cores (Jovian planets, brown dwarfs, late-M dwarfs).

#### 4. (20%) White Dwarfs (WDs)

(a) Low-mass WDs are supported by degeneracy pressure of non-relativistic electrons, which have an equation of state of  $P \propto \rho^{\gamma}$  with  $\gamma = 5/3$  (polytropic index  $n = 1.5$ ). Show that such WDs shrink with increasing mass according to a  $R \propto M^{-1/3}$  mass-radius relation.

$$P \propto \rho^{1.5}$$

$$\frac{M^2}{R^4} \propto \left( \frac{M}{R^3} \right)^{3/2}$$

$$\frac{M^2}{R^4} \propto \frac{M^{5/3}}{R^5}$$

$$R \propto \frac{M^{5/3}}{M^2}$$

$$\boxed{R \propto M^{-1/3}}$$

(b) The centers of massive WDs are supported by degeneracy pressure of relativistic electrons, which have an EOS of  $P = K \rho^{\gamma}$  with  $\gamma = 4/3$  ( $n = 3$ ) and  $K = (3/\pi)^{1/3} hc / [8(\mu_e m_H)^{4/3}]$ . Show that the radius dependence of a  $n = 3$  polytrope disappears, and solve for the mass (in  $M_{\odot}$ ) of a fully relativistic WD composed of fully ionized C, O, and Ne. What is the significance of this result?

$$M_* = -4\pi \alpha^3 \left[ \frac{(n+1)K}{4\pi G \alpha^2} \right]^{n/n-1} \xi_0^2 \left( \frac{d\theta}{d\xi} \right)_{\xi_0}$$

$$\alpha = R_*/\xi_0$$

$$\xi_0^2 \left( \frac{d\theta}{d\xi} \right)_{\xi_0} = -2.02$$

$$M_* = -4\pi \frac{R_*^3}{\xi_0^3} \frac{(3+1)^{3/2} K^{3/2}}{(4\pi G)^{3/2} \left( \frac{R_*}{\xi_0} \right)^3} (-2.02)$$

$$M_* = (2.02) \frac{4\pi (8)}{(4\pi)^{3/2} G^{3/2}} \left[ \left( \frac{3}{\pi} \right)^{1/3} \frac{hc}{8(\mu_e m_H)^{4/3}} \right]^{3/2}$$

$$M_* = \frac{(2.02) 8}{\sqrt{4\pi} G^{3/2}} \left( \frac{3}{\pi} \right)^{1/2} \frac{(hc)^{3/2}}{8^{3/2} (\mu_e m_H)^2}$$

```
In [148]: opa = pd.read_csv('/d/users/nikhil/Downloads/OPAL_SolarComposition.txt')
# NP Reading in raw text file
goodopa = opa.to_numpy()
# NP Converting to numpy array
betteropa = [np.fromstring(goodopa[i][0], dtype = float, sep = ' ').tolist()\
              for i in range(len(goodopa))]
# NP Creating a table for opacities and temps., skipping first row
bestopa = [betteropa[i+1][1:] for i in range(len(betteropa) - 1)]
# NP Creating a list of only opacities
row_lengths = []
for row in bestopa:
    row_lengths.append(len(row))
max_length = max(row_lengths)
for row in bestopa:
    while len(row) < max_length:
        row.append(0)
pristineopa = np.array(bestopa)
# NP Creating numpy array of opacities to better parse through array
logTs = np.array([betteropa[i+1][0] for i in range(len(betteropa)-1)])
# NP Creating array of temperatures
logRs = np.array(np.fromstring(goodopa[0][6:], dtype = float, sep = ' '))
# NP Creating array of R's
Tgrid = np.array([[t for r in range(len(logRs))] for t in logTs])
# NP Creating 2D numpy array of temperatures
Rgrid = np.array([[r for r in logRs] for t in range(len(logTs))])
# NP Creating 2D numpy array of temperatures
rhos = np.array([np.array([np.round(r + 3 * (t - 6) , 3) for r in logRs])\
                      for t in logTs])
# NP Creating density grip for each opacity

/tmp/ipykernel_3361101/2143050322.py:5: DeprecationWarning: string or file could
ched data; this will raise a ValueError in the future.
    betteropa = [np.fromstring(goodopa[i][0], dtype = float, sep = ' ').tolist()\
```

Interpolating over opacity table.

Code:

```
In [149]: fillinopa = interp2d(logRs, logTs, pristineopa, kind = 'cubic')
# NP Interpolating over Rs and Ts to fill in all values for opacity
```

Reading in Solar Model.

Code:

```
In [150]: r, m_r, l_r, T_r, rho_r, logP_r = np.loadtxt('/d/users/nikhil/'
              'Downloads/SolarModel.csv', unpack = True, skiprows = 1,\
              usecols = (0, 1, 2, 3, 4, 5), delimiter = ',')
# NP Reading in table quantities
```

Defining temperature gradient. Note, the convective temperature gradient changes in different layers of the Sun corresponding to the different states of matter in different layers.

Code:

```
In [197]: n_rad = np.array([(3e-2 * fillinopa(usefulRs[i], usefulTs[i])\
              * 10 ** (logP_r[i] * l_r[i] * L_o) / (64 * np.pi\
              * G * m_r[i] * M_o * 0 * (10 ** (usefulTs[i])) ** 4)\
              for i in range(len(r))])
# NP Calculating temperature gradient for the Sun for all layers
n_conv = [0.4 for ra in r if ra < .7] + [0.0216 for ra in r\
              if ra >= .7]
# NP Calculating convective temperature gradient for the whole
# NP Sun
```

Plotting  $\nabla_{rad}$  as a function of radius.

Code:

calculate how many times more massive a star would have to be to achieve a fully convective core.

$$L \propto M^{3.5}$$

$$\frac{L}{L_{\odot}} = \left( \frac{M}{M_{\odot}} \right)^{3.5}$$

$$\frac{L}{L_{\odot}} = 9.68$$

$$\left( \frac{M}{M_{\odot}} \right)^{3.5} = 9.68$$

$$M = (9.68)^{1/3.5} M_{\odot}$$

$$M = 1.91 M_{\odot}$$

Code:

```
In [213]: print('Luminosity difference: ' + format(.4/n_rad\
[r == 0.09][0][0], '.2E') + ' times.')
# NP Calculating luminosity difference to make core convective
print('Mass difference: ' + format((.4 /n_rad[r ==\
0.09][0][0]) ** (1 /3.5), '.2E') + ' times.')
# NP Calculating mass difference to make core convective

Luminosity difference: 9.68E+00 times.
Mass difference: 1.91E+00 times.
```