

# 1 Introduction

## 2 Preliminaries

### 2.1 Quaternion representation of color images

A quaternion,  $q$ , was defined by Hamilton [9] as a generalization of the complex numbers:

$$q = a + b\mathbf{i} + c\mathbf{j} + d\mathbf{k}$$

The real number  $a, b, c$  and  $d$  are called the components of  $q$ , and the imaginary units  $i, j$  and  $k$  are defined according to the following rules:

$$\begin{aligned} \mathbf{i}^2 = \mathbf{j}^2 = \mathbf{k}^2 = \mathbf{ijk} = -1, \\ \mathbf{ij} = -\mathbf{ji} = \mathbf{k}, \mathbf{jk} = -\mathbf{kj} = \mathbf{i}, \mathbf{ki} = -\mathbf{ik} = \mathbf{j}. \end{aligned}$$

Therefore, the set of quaternions  $\mathbb{H}$  is an *algebra*, where a quaternion is called pure quaternion when  $a = 0$ . The conjugate and modulus of a quaternion are respectively defined by

$$\begin{aligned} q^* &= a - b\mathbf{i} - c\mathbf{j} - d\mathbf{k}, \\ |q| &= \sqrt{a^2 + b^2 + c^2 + d^2}. \end{aligned}$$

Ell and Sangwine [10] utilized quaternions to represent a color image,  $f : \mathbb{R}^2 \rightarrow \mathbb{R}^3$ , as follows:

$$f(x, y) = \mathbf{i}f_R(x, y) + \mathbf{j}f_G(x, y) + \mathbf{k}f_B(x, y),$$

where functions  $f_R, f_G, f_B : \mathbb{R}^2 \rightarrow \mathbb{R}$  represent the red, green and blue components of the  $(x, y)$  pixel, respectively.

### 2.2 Quaternion Zernike moments

Quaternion Zernike moments, just as regular Zernike moments, are defined over the complex unit disk

$$\mathbb{D} = \left\{ z = re^{i\theta} \in \mathbb{C} : r \in [0, 1], \theta \in [0, 2\pi) \right\},$$

and should be expressed using polar coordinates as in [11]: for an  $f : (r, \theta) \mapsto \mathbb{R}^3$ , the right-side QZMs of order  $n$  with repetition  $m$  are

$$Z_{n,m}^R(f) = \frac{n+1}{\pi} \int_0^1 \int_0^{2\pi} f(r, \theta) R_{n,m}(r) e^{-\mu m \theta} r \, d\theta dr, \quad (1)$$

where  $|m| \leq n$  and  $n - |m|$  being even,  $\boldsymbol{\mu}$  is an arbitrary unit pure quaternion, regularly chosen as  $\boldsymbol{\mu} = (\mathbf{i} + \mathbf{j} + \mathbf{k})/\sqrt{3}$  in the literature and so in this paper too, and  $R_{n,m}(r)$  is the real-valued radial polynomial given by

$$R_{n,m}(r) = \sum_{k=0}^{\frac{n-|m|}{2}} \frac{(-1)^k (n-k)!}{k! \left(\frac{n+|m|}{2} - k\right)! \left(\frac{n-|m|}{2} - k\right)!} r^{n-2k}.$$

It is important to note that these radial polynomials are symmetric in terms of  $m$  and satisfy an orthogonality relation over  $[0, 1]$  w.r.t the weight  $r$  (see e.g. [12]), i.e.

$$\int_0^1 R_n^{[m]} R_{n'}^{[m]} = \frac{1}{2n+1} \delta_{n,n'},$$

and can be expressed in terms of the shifted Jacobi polynomials in the following way:

$$R_{n,m}(r) = r^{|m|} P_{\frac{n-|m|}{2}}^{(0,|m|)}(2r^2 - 1),$$

where we denoted the  $k$ -th degree classical Jacobi polynomials by  $P_k^{(\alpha,\beta)}$  (see e.g. [13]).

Consequently, the quaternion generalizations of classical Zernike functions

$$\phi_{n,m}(r, \theta) = R_{n,m}(r) e^{-\boldsymbol{\mu} m \theta}$$

satisfy the orthogonality relation

$$\frac{n+1}{\pi} \int_0^1 \int_0^{2\pi} \phi_{n,m}(r, \theta) \phi_{n',m'}^*(r, \theta) r \, d\theta dr = \delta_{n,n'} \delta_{m,m'}, \quad (2)$$

and we have

$$f(r, \theta) = \sum_{n=0}^{\infty} \sum_{m=-\infty}^{\infty} Z_{n,m}^R(f) R_{n,m}(r) e^{\boldsymbol{\mu} m \theta} = \sum_{n=0}^{\infty} \sum_{m=-\infty}^{\infty} Z_{n,m}^R(f) \phi_{n,m}^*(r, \theta)$$

w.r.t. norm convergence in the Hilbert space of square integrable functions.

Since multiplication in  $\mathbb{H}$  is not commutative, a similar set of left-side quaternion Zernike moments can also be defined as

$$Z_{n,m}^L(f) = \frac{n+1}{\pi} \int_0^1 \int_0^{2\pi} \phi_{n,m}(r, \theta) f(r, \theta) r \, d\theta dr, \quad (3)$$

having

$$f(r, \theta) = \sum_{n=0}^{\infty} \sum_{m=-\infty}^{\infty} \phi_{n,m}^*(r, \theta) Z_{n,m}^L(f).$$

For computation purposes, it is important to recall [11] that the left- and right-side moments can be obtained by computing the conventional complex moments for the individual color channel functions  $f_R, f_G$  and  $f_B$  respectively:

$$Z_{n,m}^R = A_{n,m}^R + \mathbf{i}B_{n,m}^R + \mathbf{j}C_{n,m}^R + \mathbf{k}Z_{n,m}^R,$$

where

$$\begin{aligned} A_{n,m}^R &= -\frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_R)) + \text{Im}(Z_{n,m}(f_G)) + \text{Im}(Z_{n,m}(f_B))], \\ B_{n,m}^R &= \text{Re}(Z_{n,m}(f_R)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_G)) - \text{Im}(Z_{n,m}(f_B))], \\ C_{n,m}^R &= \text{Re}(Z_{n,m}(f_G)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_B)) - \text{Im}(Z_{n,m}(f_R))], \\ D_{n,m}^R &= \text{Re}(Z_{n,m}(f_B)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_R)) - \text{Im}(Z_{n,m}(f_G))], \end{aligned}$$

and similarly for the left-side ones

$$Z_{n,m}^L = A_{n,m}^L + \mathbf{i}B_{n,m}^L + \mathbf{j}C_{n,m}^L + \mathbf{k}Z_{n,m}^L,$$

where

$$\begin{aligned} A_{n,m}^L &= -\frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_R)) + \text{Im}(Z_{n,m}(f_G)) + \text{Im}(Z_{n,m}(f_B))], \\ B_{n,m}^L &= \text{Re}(Z_{n,m}(f_R)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_B)) - \text{Im}(Z_{n,m}(f_G))], \\ C_{n,m}^L &= \text{Re}(Z_{n,m}(f_G)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_R)) - \text{Im}(Z_{n,m}(f_B))], \\ D_{n,m}^L &= \text{Re}(Z_{n,m}(f_B)) + \frac{1}{\sqrt{3}} [\text{Im}(Z_{n,m}(f_G)) - \text{Im}(Z_{n,m}(f_R))]. \end{aligned}$$

## 2.3 Invariance properties

The previously defined QZMs should be used to obtain certain invariants for similarity transforms. The proofs of these results are found in [11].

For a function rotated around the origin, i.e. for  $f'(r, \theta) = f(r, \theta - \alpha)$  with some rotation angle  $\alpha$ , we have

$$\begin{aligned} Z_{n,m}^R(f') &= Z_{n,m}^R(f)e^{-\mu m \theta}, \\ Z_{n,m}^L(f') &= e^{-\mu m \theta} Z_{n,m}^L(f). \end{aligned}$$

It is easy to spot that the moduli of moments are left unchanged, and so they could be used as rotational invariants, just as the combined values

$$\Phi_{n,k}^m = Z_{n,m}^R(f) Z_{k,-m}^L(f) = -Z_{n,m}^R(f) (Z_{k,m}^R(f))^*,$$

for any  $|m| \leq n, |m| \leq k$  and  $n - |m|, k - |m|$  even. These will be referred to as quaternion Zernike moment rotation invariant (QZMRI). Note that the latter are quaternion-valued invariants, which could be identified with 4 real values.

For non-negative integers  $m$  and  $l$ , the quaternions defined by

$$c_{m,l}^{t,k} = (-1)^{l-k} \frac{(m+2l+1)(m+k+l)!}{(l-k)!(k-t)!(m+k+t+1)!},$$

$$L_{m+2l,m}^R(f) = \sum_{t=0}^l \sum_{k=t}^l \left( \sqrt{|Z_{0,0}^R(f)|} \right)^{-(m+2k+2)} c_{m,l}^{t,k} Z_{m+2t,m}^R(f)$$

are invariant under scaling of the function, i.e. they are the same for  $f$  and  $f''(r, \theta) = f(r/\lambda, \theta)$  with some scaling factor  $\lambda > 0$ .

Similarly to using QZMs to define rotation invariants, the previously defined scaling invariants can also be used to construct quaternions

$$\Psi_{n,k}^m = L_{n,m}^R(f) (L_{k,m}^R(f))^*,$$

which are in fact invariant to *both* rotation and scaling, and is called quaternion Zernike moment invariant (QZMI).

In order to ensure certain translation invariance, Suk and Flusser [14] defined the common centroid  $(x_c, y_c)$  of all three channels as follows:

$$m_{0,0} = m_{0,0}(f_R) + m_{0,0}(f_G) + m_{0,0}(f_B),$$

$$x_c = (m_{1,0}(f_R) + m_{1,0}(f_G) + m_{1,0}(f_B)) / m_{0,0},$$

$$y_c = (m_{0,1}(f_R) + m_{0,1}(f_G) + m_{0,1}(f_B)) / m_{0,0},$$

where the  $m_{i,j}$ -s are the respective values of the zero-order and first-order geometric moments. Translation invariance is then obtained by translating every  $f$  such that the point  $(x_c, y_c)$  is mapped to the origin  $(0, 0)$ .

### 3 Discretization

In the applications of such theoretical results belonging to continuous domains, an actual image of size  $N \times N$  pixels is generally considered as a

sample, obtained by the equidistant sampling of a continuous function along both axis. In order to make use of QZMs and their invariants, a linear transformation is required to map the pixels to a suitable domain inside the complex unit circle  $\mathbb{D}$ .

Afterwards, using the polar form of the transformed coordinates, the continuous integrals of (1) and (3) are replaced by discrete approximations of form

$$Z_{n,m}^R(f) \approx \lambda_{n,m} \sum_{k=1}^{N_1} \sum_{j=1}^{N_2} f(r_k, \theta_j) \phi(r_k, \theta_j) w(r_k, \theta_j), \quad (4)$$

with some discretization weight function  $w$  and normalization constants  $\lambda_{n,m}$ . E.g. in [11] this is done as

$$Z_{n,m}^R(f) \approx \frac{2n+2}{\pi(N-1)^2} \sum_{x=1}^N \sum_{y=1}^N f(r_{x,y}, \theta_{x,y}) \phi(r_{x,y}, \theta_{x,y}),$$

which is actually the conventional Cartesian coordinate-based calculation method, coupled with the mapping proposed by Chong et al. [15]

$$\begin{aligned} r_{x,y} &= \sqrt{(c_1 x + c_2)^2 + (c_1 y + c_2)^2}, \\ \theta_{x,y} &= \tan^{-1} \left( \frac{c_1 y + c_2}{c_1 x + c_2} \right), \end{aligned}$$

with  $c_1 = \sqrt{2}/(N-1)$  and  $c_2 = -1/\sqrt{2}$ .

A similar approach was utilized numerous times in the literature, substituting the radial polynomial system  $R_{n,m}$  of Zernike functions with other well-known orthogonal systems over the radial interval  $[0, 1]$ . For color image analysis, further examples of this include the QFMMs [1], the QPETs, QPCTs and QPSTs [3], the QCMs [4], the QBFMs [5], the QRHFM [6], the QG-CHFM and QG-PJFM [7] and the QSBFM [8].

This natural method was proven susceptible to inaccuracies of both geometric and numeric nature when applied to greyscale images (see e.g. [16], [17]), and these problems are carried over to the quaternion generalizations as well. Later in [18], a more accurate computational technique was proposed, where  $\mathbb{D}$  is partitioned into polar sectors of approximately equal areas, and the original pixels are transformed to the centroids of these via cubic interpolation, and the weights are computed by performing the integration of functions  $\phi_{n,m}$  over the respective polar sectors.

The same idea, adapted to quaternion-type moments for any aforementioned radial system, could substantially improve performances for tasks like image reconstruction, as well as recognition after rotation, scaling and translation (RST) transformations and addition of different kinds of noise; and the reason for this is essentially the improved accuracy gained for the computation of respective moments and invariants. Examples of this technique applied for color images include the papers of Hosny et al. for Legendre [19] and Chebyshev [20] radial systems.

A drawback of this approach is that the cubic interpolation applied for coordinate transformation is irreversible, so another transformation with further introduced error is required if one wishes to reconstruct the original square/rectangle shaped image. Instead, in the previous works [16]-[20], reconstruction error was measured for the transformed image and its reconstructions on disc  $\mathbb{D}$  only. Another, although minor issue is the natural smoothing provided by the cubic interpolation, which unintentionally improves the obtained results for experiments on noisy images by filtering some of the noise, a sometimes unwanted side effect for one who wants to compare the performances of the moments themselves. To ensure fair comparisons with the first transformation group, linear interpolation should be used, ensuring invertible coordinate transformations, and also excluding unwanted side effects.

Finally, the results of Wang et al. [21] and Liu et al. [22] for radial harmonic Fourier moments are some of the recent advances in the study of quaternion moments, introducing a novel coordinate transformation, which seems to improve accuracy for this specific function system while maintaining fast computation times, without utilizing cubic interpolation at all. These improvements may seem surprising at first, but hopefully our results will clarify the theoretical reason behind them.

Now we shall describe our approach to the discretization of continuous quaternion Zernike moments. Our aim is to ensure simple and reversible computability, but achieve considerable reduction in potential computational inaccuracies at the same time. For this, we define a system of sampling points  $(r_k, \theta_j)$ , over which the integral discretization (4) maintains advantageous theoretical properties. The idea behind our approach is motivated by the work [12].

Fix a positive integer  $N$ . Let us denote by  $\rho_{k,N}$ ,  $k = 1, \dots, N$  the roots of the  $N$ -th order Legendre polynomial [13]. Notating the fundamental polynomials of Lagrange interpolation w.r.t. roots  $\rho_{k,N}$  as  $\ell_{k,N}$ , we define the

constants

$$\mathcal{A}_{k,N} = \int_{-1}^1 \ell_{k,N}(x) dx, \quad (k = 1, \dots, N).$$

Then, the system of sampling points  $X_N$  over the unit disk  $\mathbb{D}$  is defined in polar form by

$$X_N \ni (r_k, \theta_j) = \left( \sqrt{\frac{1 + \rho_{k,N}}{2}}, \frac{2\pi j}{4N} \right), \quad (k = 1, \dots, N, j = 1 \dots, 4N),$$

and the respective weight values and constants are

$$w(r_k, \theta_j) = \frac{\mathcal{A}_{k,N}}{8N}, \quad \lambda_{n,m} = 2n + 2,$$

so the generated integral approximation is

$$Z_{n,m}^R(f) \approx (2n + 2) \sum_{k=1}^N \sum_{j=1}^{4N} f(r_k, \theta_j) \phi(r_k, \theta_j) \frac{\mathcal{A}_{k,N}}{8N},$$

and in fact any integral over  $\mathbb{D}$  can be approximated by the discretization

$$\int_{\mathbb{D}} f \approx \int_{X_N} f = \sum_{k=1}^N \sum_{j=1}^{4N} f(r_k, \theta_j) \frac{\mathcal{A}_{k,N}}{8N}.$$

We remark that the values  $\mathcal{A}_{k,N}$  are somewhat close to  $\frac{2}{N}$ , so the weighting is close to being uniform.

What is appealing in this choice of  $X_N$  is the fact that the orthogonality of quaternion Zernike functions (2) is preserved under changing to discrete integration over this set of points, i.e.

**Theorem 1.** *Suppose that for  $n, n'$  naturals and  $m, m'$  integers we have  $n - |m| + n' + |m'| < 4N$  and  $n + |m| + n' - |m'| < 4N$ . Then*

$$(2n + 2) \int_{X_N} \phi_{n,m} \phi_{n',m'}^* = \delta_{n,n'} \delta_{m,m'}.$$

**Proof.**

□

Theoretically, the discrete orthogonality relation of Theorem 1 means that the exact values of right and left moments can be computed, provided

we can measure the function values over  $X_N$ , for a sufficiently large  $N$ : let us consider an arbitrary linear combination of right moments (1) in the form

$$f \approx f_N(r, \theta) = \sum_{n < \frac{4N}{3}} \sum_m Z_{n,m}^R(f) \phi_{n,m}^*(r, \theta),$$

then the exact value of the moment  $Z_{n,m}^R(f)$ , using the previous result, is

$$Z_{n,m}^R(f) = \int_{X_N} f_N \phi_{n,m},$$

and the same can be said for the left moments. This proves that for system  $X_N$ , the discretization error is zero, a property that no other previously used quaternion moment method possessed. Besides this, only numerical roundoffs and the moment order threshold  $f \approx f_N$  generates computational inaccuracies.

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