

Currents induced by magnetic impurities in superconductors with spin-orbit coupling

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Skyrmions are nice

PACS numbers:

I. INTRODUCTION

General context of skyrmions and topological excitations: memory, manipulation, local creation via SP STM. Extension of skyrmion discussion to the case of hybrid structures: SC and Skyrmion. What are the consequences of bringing topological exchange field into SC. Question we address is the possible local spectroscopic signatures of SC quasiparticles in SC due to skyrmion field. We know from the past discussion that there are impurity bound states in SC near magnetic impurities. We have now the framework to address formation of bound states. Talk about local single impurity limit (YSR) and show the cartoon of the local and extended skyrmion and spectra. There are two effects: local scattering and Zeeman field hence the DOS will be split etc. Draw similarities and differences with single imp. In parallel with skyrmion discovery the local imaging using magnetic probes like MFM and SP-STM allowed one to image the matter at atomic resolution while also resolving spin content of electron carriers in the substrate. Here we prove the existence of the new type of localized excitation on the skyrmion core we call Sc-YSR state (alternative is skyrmion bound state (sbs)). Show the main results upfront in the introduction. Both LDOS and SP-LDOS. Main section:

Introduce T matrix and results for analytic solution. Introduce the numerical approach and present the results as a function of position and as a function of energy. Kind of same figs as in Shos talk. Discuss the results and what it means, how big the signal is etc. Unfortunately we do not see any topological state at zero energy and as such these results represent a new kind of magnetic texture induced states that exhibit intragap states.

II. SKYRMIONS IN FERROMAGNETIC FILMS

Define skyrmions. Write topological number. Define magnetic moments of the skyrmion (anapole, monopole etc). Discuss different types of skyrmions and that they are equivalent for electrons.

III. T-MATRIX ANALYSIS

Superconductor-ferromagnet heterostructures were recently proposed as a viable platform for realizing topological superconductivity (TS) [1–3], which can host Majorana fermion quasiparticles at vortex cores and boundaries [4–6]. Majorana fermions obey non-Abelian statistics and may be utilized for topological quantum computation [7–9]. The key ingredients driving these systems in the topologically non-trivial regime are the spin-orbit coupling (SOC) and magnetism. Recently, the search for experimental realizations of TS has also led to engineering the impurity bands of the Yu-Shiba-Rusinov (YSR) states [10–12], induced by magnetic atoms on the surface of a superconductor [13–24]. Following this recipe, zero-energy peaks in the tunneling spectrum were recently measured at the ends of a one-dimensional (1D) chain of magnetic atoms [25]. Such a tunneling spectrum could be the evidence of Majorana edge states, although alternative explanations are also possible [26].

The interplay of SOC and magnetism has another remarkable consequence. Consider a two-dimensional (2D) surface of a 3D superconductor. The effective Hamiltonian of the surface $h(\mathbf{p}) = \frac{\mathbf{p}^2}{2m} + \lambda(\boldsymbol{\sigma} \times \mathbf{p})_z$ contains Rashba SOC due to the absence of inversion symmetry at the surface. Then, the velocity operator $\mathbf{v} = \frac{dh(\mathbf{p})}{d\mathbf{p}} = \frac{\mathbf{p}}{m} + \lambda \hat{\mathbf{z}} \times \boldsymbol{\sigma}$ contains a spin-dependent term that gives an extra contribution to the current

$$\mathbf{j}_{\text{extra}} = \lambda \hat{\mathbf{z}} \times \langle \boldsymbol{\sigma} \rangle. \quad (1)$$

A ferromagnet proximity-coupled to the superconductor would render a finite spin polarization $\langle \boldsymbol{\sigma} \rangle \neq 0$ and thus generate a current as schematically shown in Fig. The phenomenon of driving a current with magnetism is known as the magnetoelectric effect. This effect may vanish in metals due to dissipation but survives in superconductors lacking inversion symmetry [27–31]. The magnetoelectric effect was also recently discussed in a pure 1D model of TS [32].

t-matrix is good too.

IV. NUMERICAL ANALYSIS

V. CONCLUSION

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- [1] R. M. Lutchyn, J. D. Sau, and S. Das Sarma, “Majorana fermions and a topological phase transition in semiconductor-superconductor heterostructures,” *Phys. Rev. Lett.* **105**, 077001 (2010).
- [2] Y. Oreg, G. Refael, and F. von Oppen, “Helical liquids and Majorana bound states in quantum wires,” *Phys. Rev. Lett.* **105**, 177002 (2010).
- [3] J. D. Sau, R. M. Lutchyn, S. Tewari, and S. Das Sarma, “Generic new platform for topological quantum computation using semiconductor heterostructures,” *Phys. Rev. Lett.* **104**, 040502 (2010).
- [4] A. Yu. Kitaev, “Unpaired Majorana fermions in quantum wires,” *Phys. Usp.* **44**, 131 (2001).
- [5] J. Alicea, “New directions in the pursuit of Majorana fermions in solid state systems,” *Rep. Prog. Phys.* **75**, 076501 (2012).
- [6] C. W. J. Beenakker, “Search for Majorana fermions in superconductors,” *Annu. Rev. Condens. Matter Phys.* **4**, 113 (2013).
- [7] N. Read and D. Green, “Paired states of fermions in two dimensions with breaking of parity and time-reversal symmetries and the fractional quantum Hall effect,” *Phys. Rev. B* **61**, 10267 (2000).
- [8] D. A. Ivanov, “Non-Abelian Statistics of Half-Quantum Vortices in p-Wave Superconductors,” *Phys. Rev. Lett.* **86**, 268 (2001).
- [9] C. Nayak, S. H. Simon, A. Stern, M. Freedman, and S. Das Sarma, “Non-Abelian anyons and topological quantum computation,” *Rev. Mod. Phys.* **80**, 1083–1159 (2008).
- [10] L. Yu, “Bound state in superconductors with paramagnetic impurities,” *Acta Phys. Sin.* **21**, 75 (1965).
- [11] H. Shiba, “Classical Spins in Superconductors,” *Prog. Theor. Phys.* **40**, 435 (1968).
- [12] A. I. Rusinov, *JETP Lett.* **9**, 85 (1969).
- [13] T. P. Choy, J. M. Edge, A. R. Akhmerov, and C. W. J. Beenakker, “Majorana fermions emerging from magnetic nanoparticles on a superconductor without spin-orbit coupling,” *Phys. Rev. B* **84**, 1 (2011).
- [14] S. Nadj-Perge, I. K. Drozdov, B. A. Bernevig, and A. Yazdani, “Proposal for realizing Majorana fermions in chains of magnetic atoms on a superconductor,” *Phys. Rev. B* **88**, 1 (2013).
- [15] J. Klinovaja, P. Stano, A. Yazdani, and D. Loss, “Topological Superconductivity and Majorana Fermions in RKKY Systems,” *Phys. Rev. Lett.* **111**, 186805 (2013).
- [16] M. M. Vazifeh and M. Franz, “Self-Organized Topological State with Majorana Fermions,” *Phys. Rev. Lett.* **111**, 206802 (2013).
- [17] B. Braunecker and P. Simon, “Interplay between Classical Magnetic Moments and Superconductivity in Quantum One-Dimensional Conductors: Toward a Self-Sustained Topological Majorana Phase,” *Phys. Rev. Lett.* **111**, 147202 (2013).
- [18] F. Pientka, L. I. Glazman, and F. von Oppen, “Topological superconducting phase in helical Shiba chains,” *Phys. Rev. B* **88**, 1 (2013).
- [19] S. Nakosai, Y. Tanaka, and N. Nagaosa, “Two-dimensional p-wave superconducting states with magnetic moments on a conventional s-wave superconductor,” *Phys. Rev. B* **88**, 180503 (2013).
- [20] K. Pöyhönen, A. Westström, J. Röntynen, and T. Ojanen, “Majorana states in helical Shiba chains and ladders,” *Phys. Rev. B* **89**, 115109 (2014).
- [21] I. Reis, D. J. J. Marchand, and M. Franz, “Self-organized topological state in a magnetic chain on the surface of a superconductor,” *Phys. Rev. B* **90**, 085124 (2014).
- [22] P. M. R. Brydon, S. Das Sarma, H.-Y. Hui, and J. D. Sau, “Topological Yu-Shiba-Rusinov chain from spin-orbit coupling,” *Phys. Rev. B* **91**, 064505 (2015).
- [23] J. Röntynen and T. Ojanen, “Topological superconductivity and high Chern numbers in 2D ferromagnetic Shiba lattices,” *arXiv:1412.5834*.
- [24] J. Li, T. Neupert, Z. J. Wang, A. H. MacDonald, A. Yazdani, and B. A. Bernevig, “A novel platform for two-dimensional chiral topological superconductivity,” *arXiv:1501.00999v1*.
- [25] S. Nadj-Perge, I. K. Drozdov, J. Li, H. Chen, S. Jeon, J. Seo, A. H. MacDonald, B. A. Bernevig, and A. Yazdani, “Observation of Majorana fermions in ferromagnetic atomic chains on a superconductor,” *Science* **346**, 602 (2014).
- [26] J. D. Sau and P. M. R. Brydon, “Bound states of a ferromagnetic wire in a superconductor,” *arXiv:1501.03149*.
- [27] L.S. Levitov, Yu.V. Nazarov, and G.M. Eliashberg, “Magnetostatics of superconductors without an inversion center,” *JETP Lett.* **41**, 445 (1985).
- [28] V. M. Edelstein, “Characteristics of the Cooper pairing in two-dimensional noncentrosymmetric electron systems,” *Sov. Phys. JETP* **68**, 1244 (1989).
- [29] V. M. Edelstein, “Magnetoelectric effect in polar superconductors,” *Phys. Rev. Lett.* **75**, 2004 (1995).
- [30] S. K. Yip, “Two-dimensional superconductivity with strong spin-orbit interaction,” *Phys. Rev. B* **65**, 144508 (2001).
- [31] E. Bauer and M. Sigrist, eds., *Non-Centrosymmetric Superconductors*, Vol. 847 (Springer Berlin Heidelberg, 2012).
- [32] T. Ojanen, “Magnetoelectric Effects in Superconducting Nanowires with Rashba Spin-Orbit Coupling,” *Phys. Rev. Lett.* **109**, 226804 (2012).