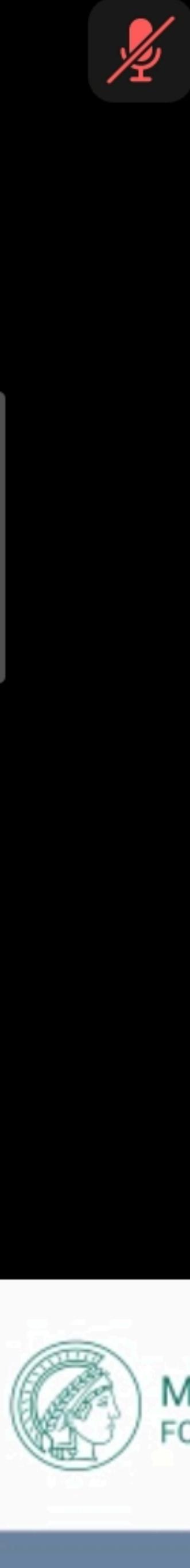
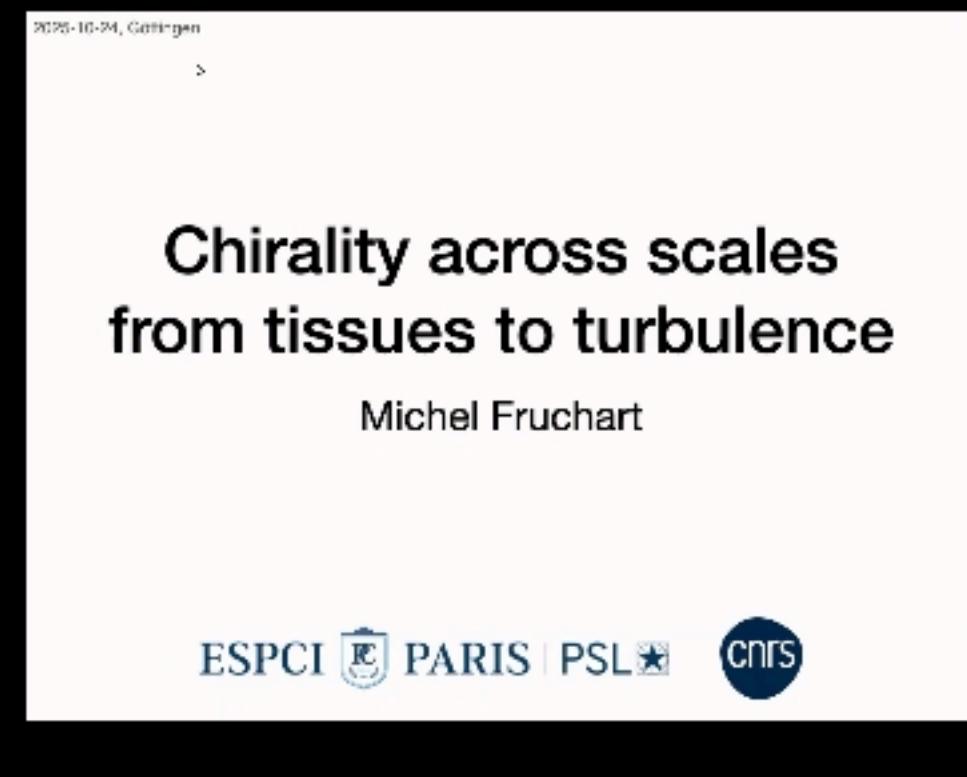
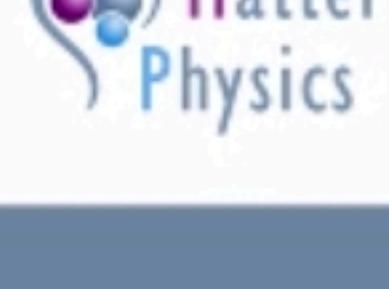


6:25

13.0 KB/s LTE 4G 16%



MAX PLANCK INSTITUTE
FOR DYNAMICS AND SELF-ORGANIZATION



LMP Seminar

The seminar will be starting soon...

MAX
PLANCK
SCHOOL

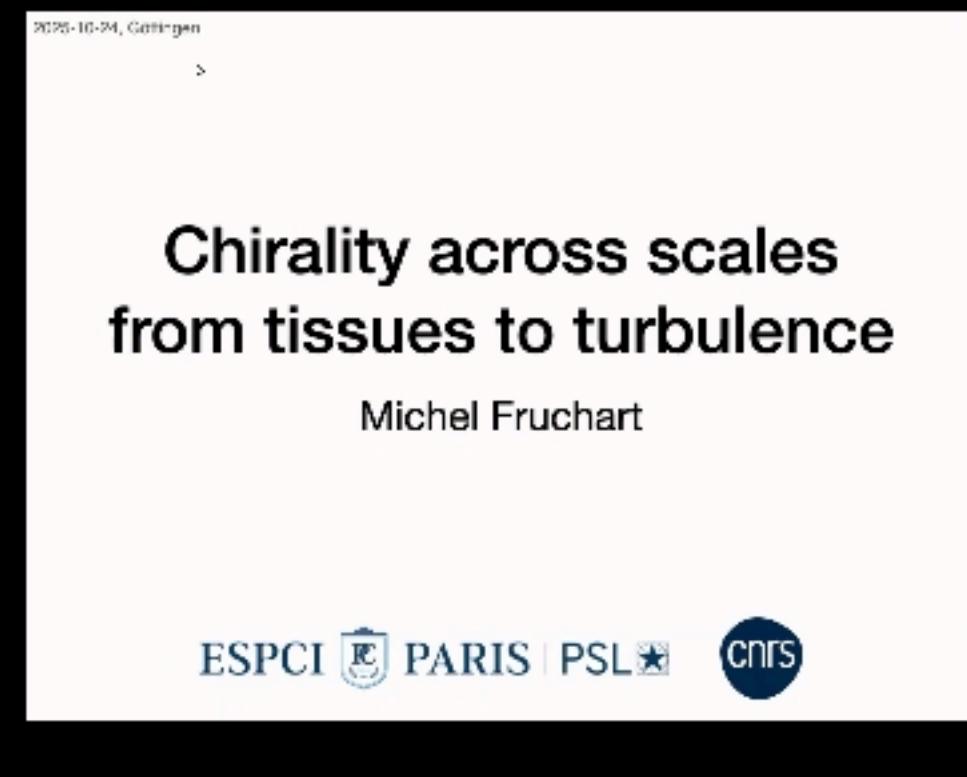
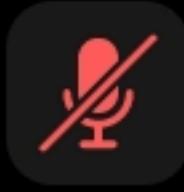
matter
to life



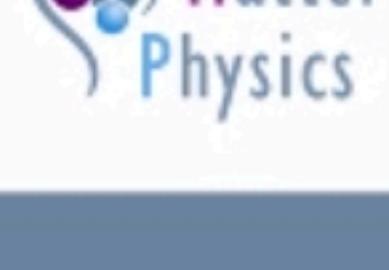
LMP Seminars

6:25

13.0 KB/s LTE 4G 16%



MAX PLANCK INSTITUTE
FOR DYNAMICS AND SELF-ORGANIZATION



LMP Seminar

The seminar will be starting soon...

MAX
PLANCK
SCHOOL

matter
to life



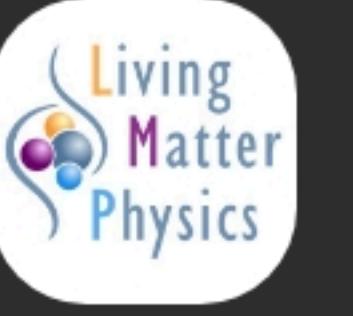
LMP Seminars



Participants (4)



Naman Dixit (me)



LMP Seminars (Host)



RB Robin Barta



Yuto Hosaka (保阪 悠人)

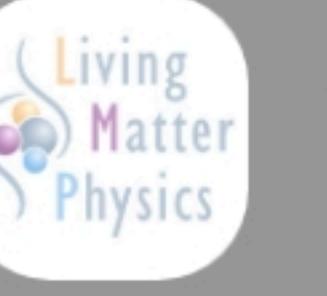


Invite

Participants (4)



Naman Dixit (me)



LMP Seminars (Host)



Robin Barta



Yuto Hosaka (保阪 悠人)



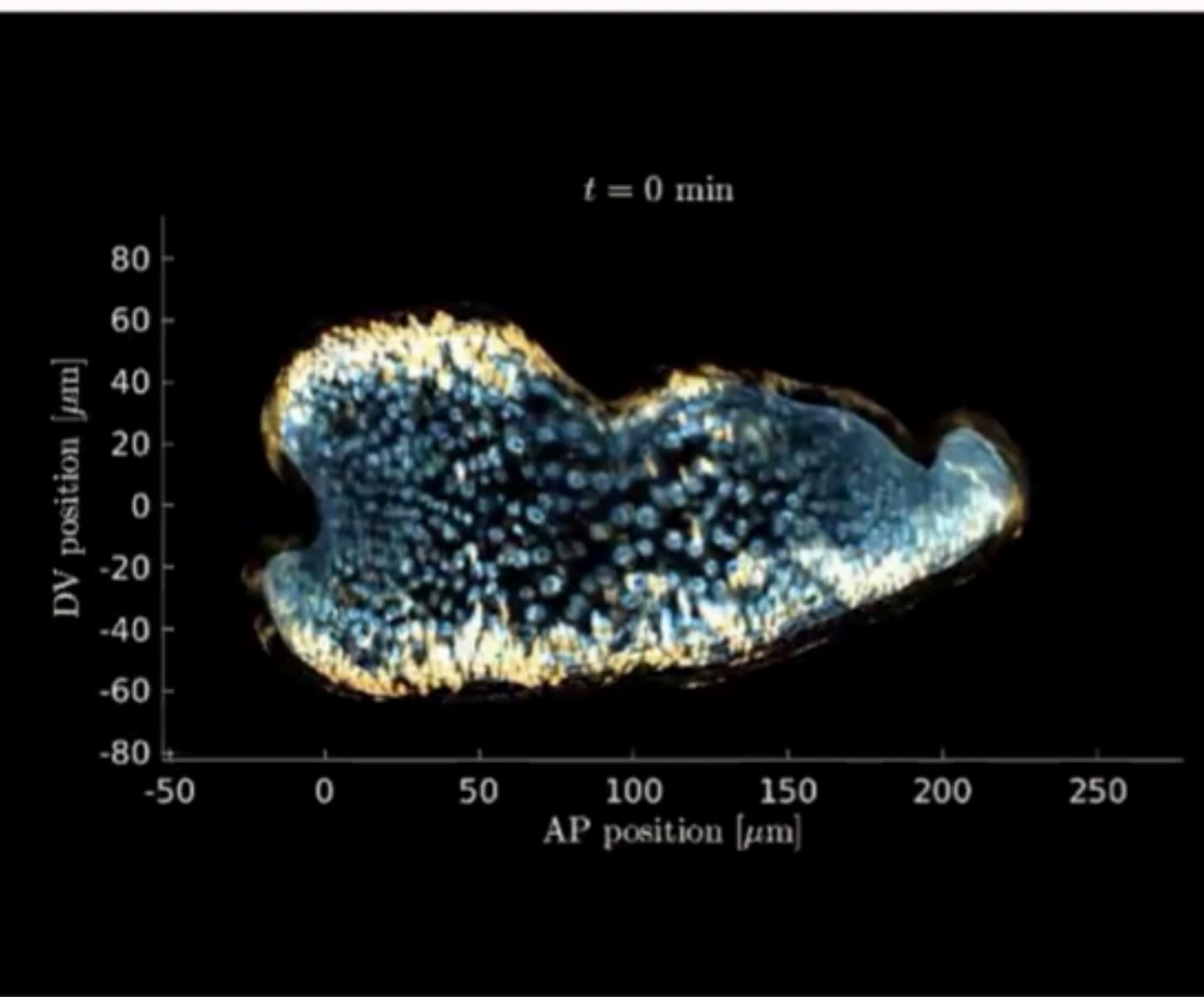
Invite

Inv



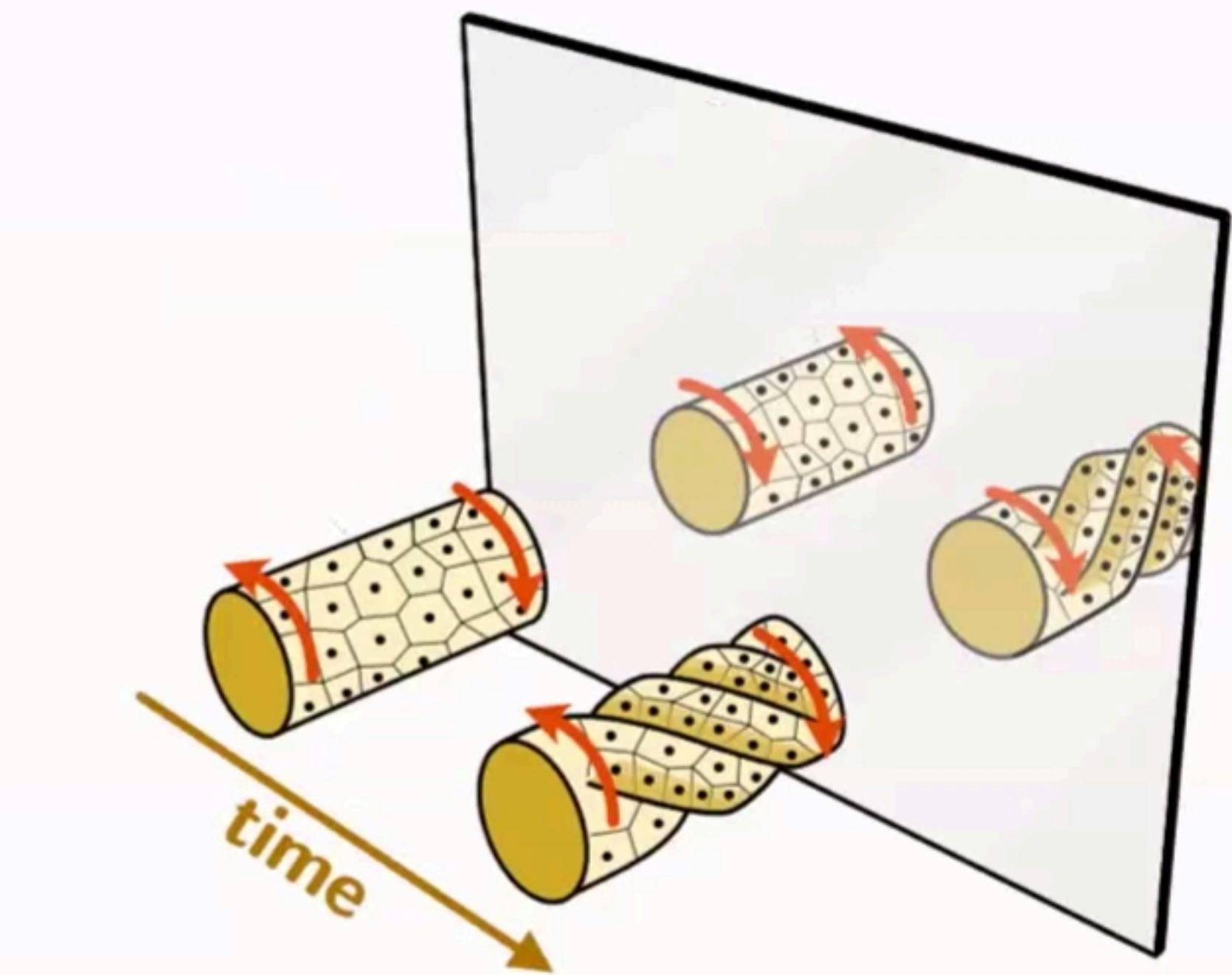
Chirality across scales in tissue dynamics

LMP Seminars



Drosophila midgut morphogenesis

Mitchell, Cislo, Shankar, Lin,
Shraiman, Streichan, [eLife \(2022\)](#)



¿how does chirality propagate across scales in biological tissues?



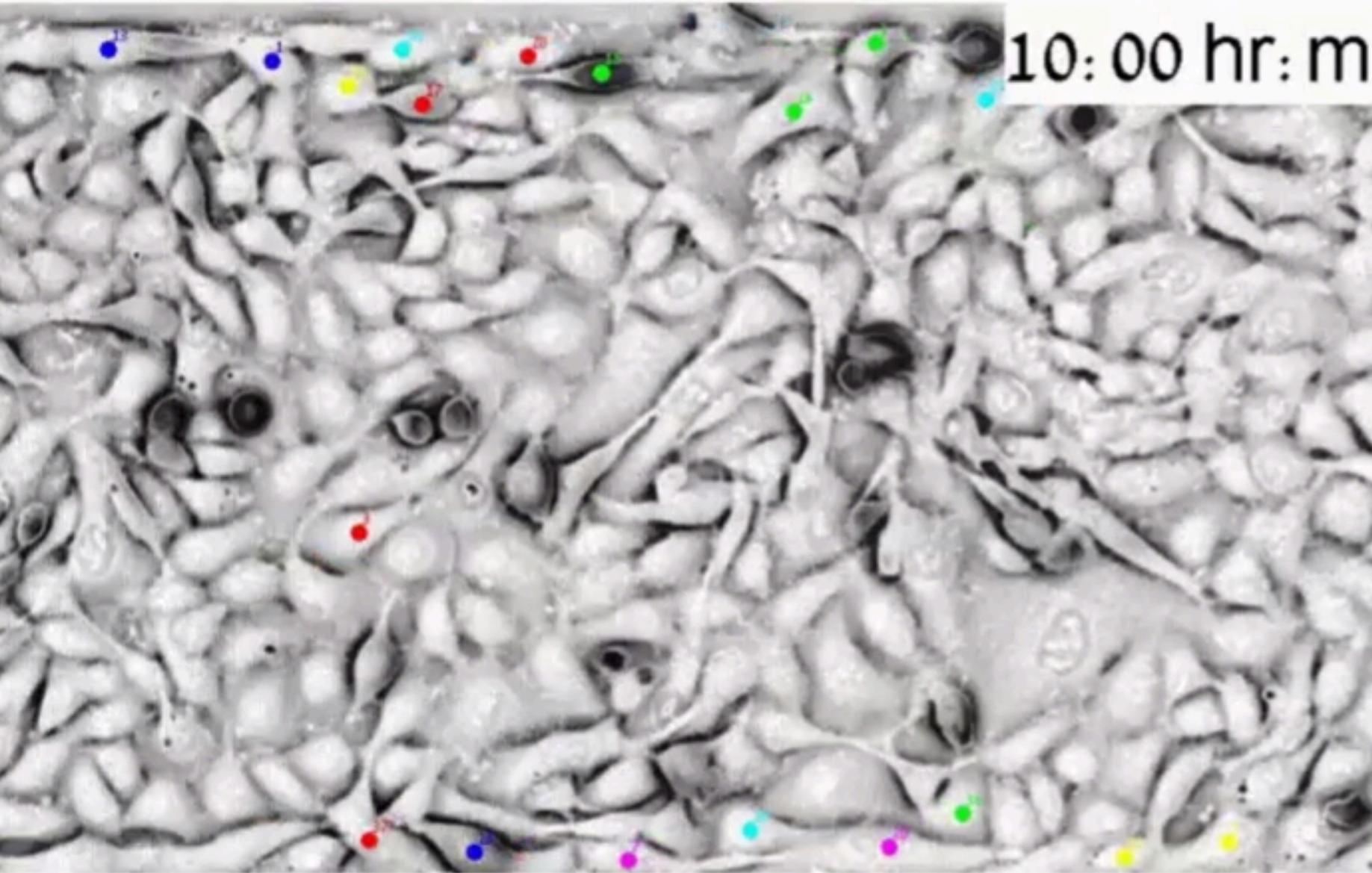
Chirality across scales in tissue dynamics



Pascal Silberzan

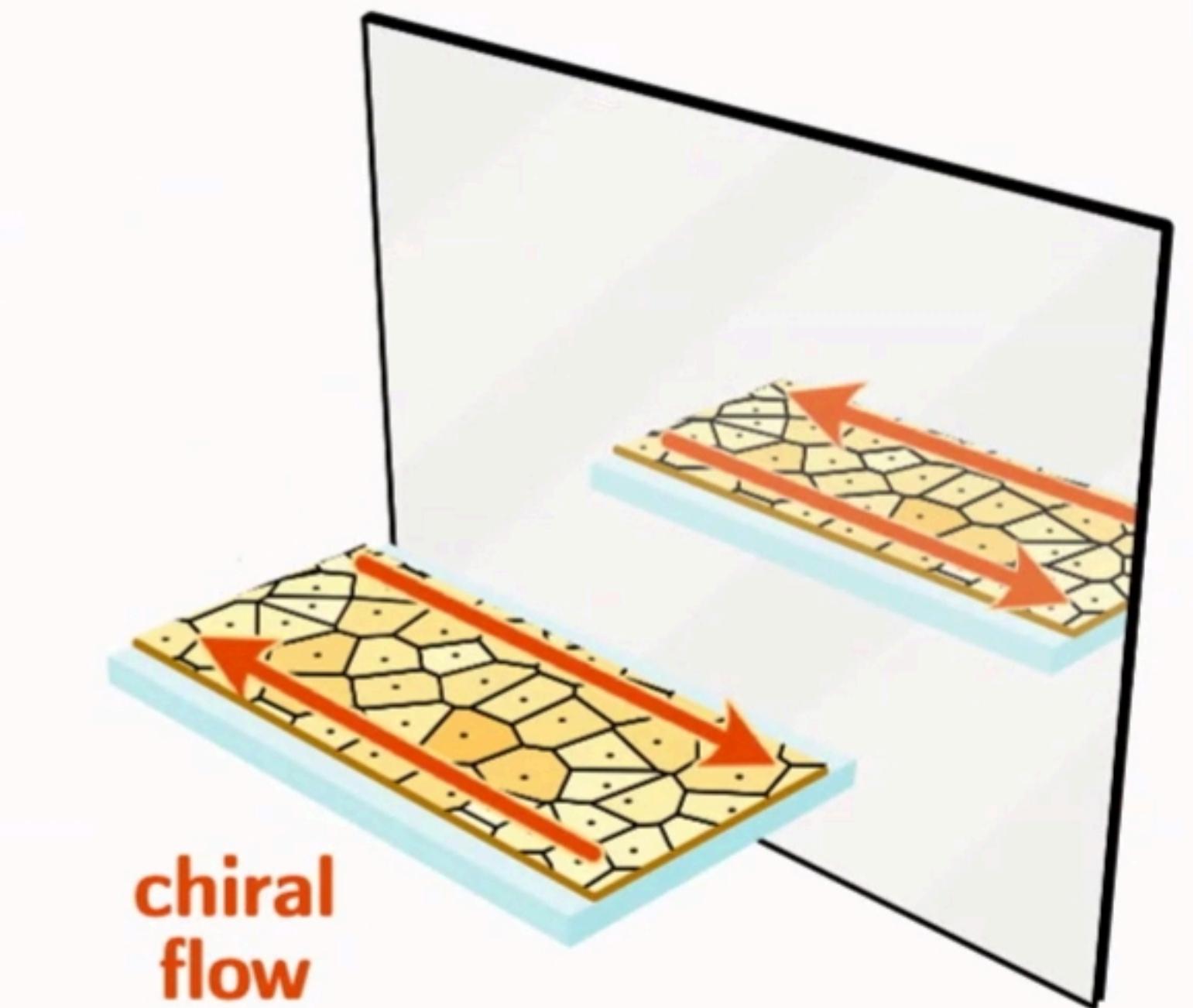


Victor Yashunsky

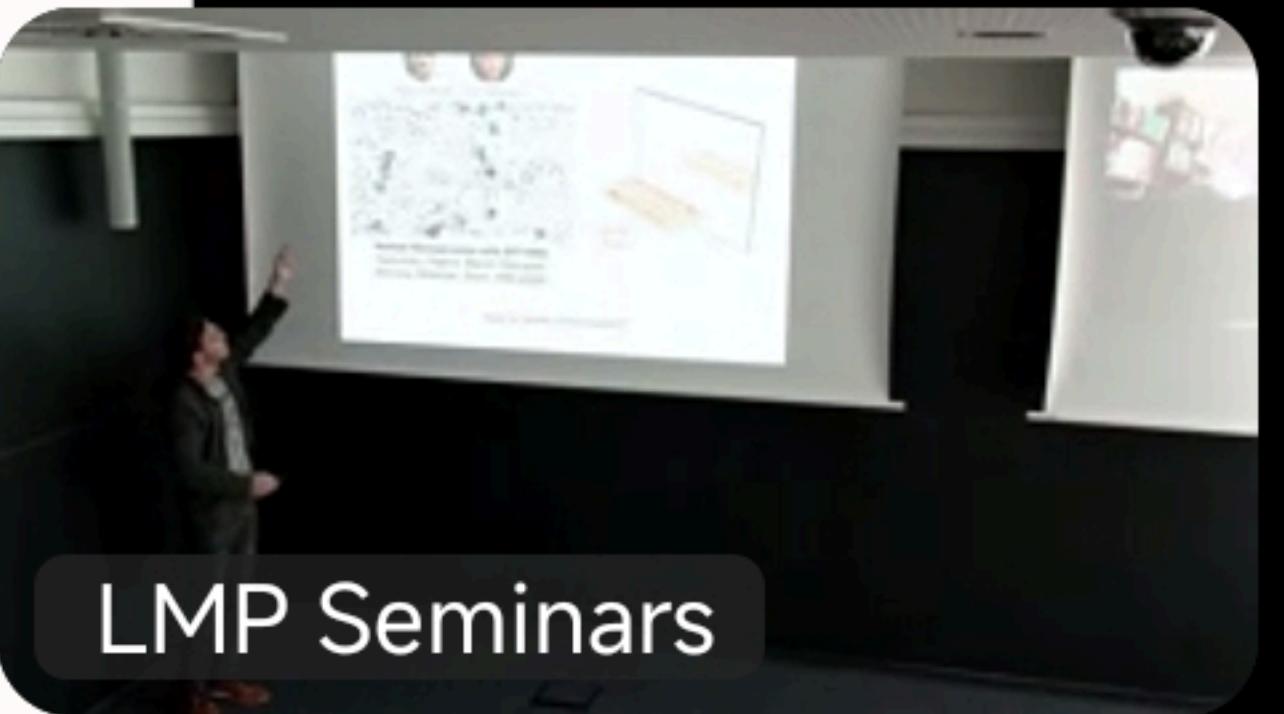


Human fibrosarcoma cells (HT1080)

Yashunsky, Pearce, Blanch-Mercader,
Ascione, Silberzan, Giomi, PRX (2022)



how to model chiral tissues?

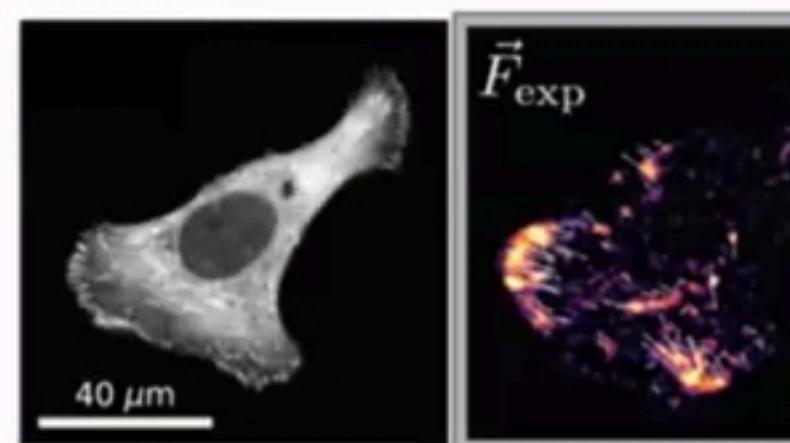




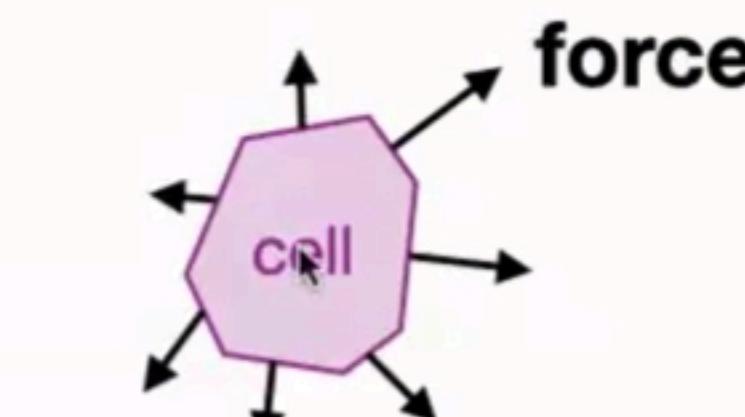
LMP Seminars's screen

How to model tissues?

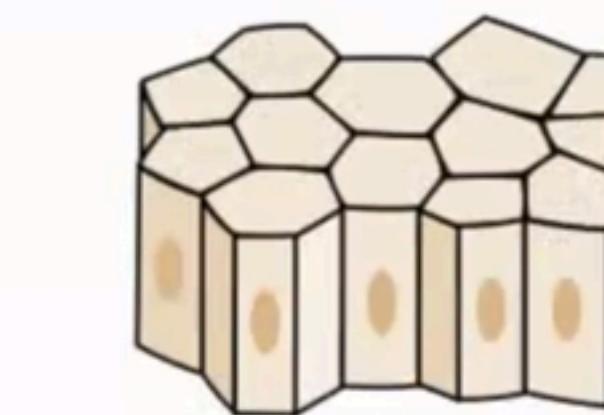
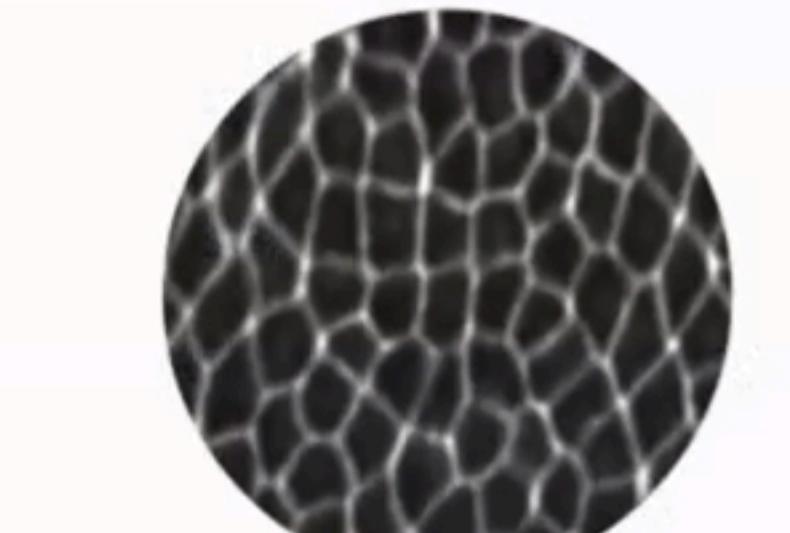
Human Osteosarcoma (U2OS) cells
Cell, 187(2), 481 (2024)



cell-level



mesoscopic



continuum



Drosophila, courtesy S. Streichan



vertex models are minimal descriptions of confluent biological tissues



see e.g.
CP Broedersz and FC MacKintosh [RMP \(2014\)](#)
DA Fletcher and RD Mullins [Nature \(2010\)](#)

vertex models

**complex
viscoelastic fluids/gels**

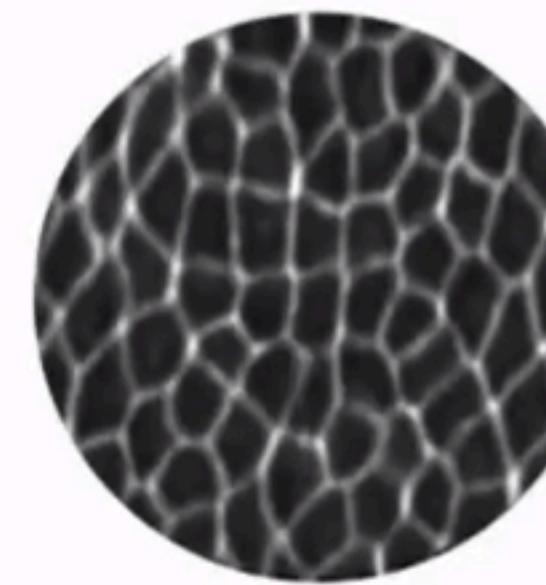
see e.g.
MC Marchetti et al., [RMP \(2013\)](#)
NI Petridou and CP Heisenberg, [EMBO \(2019\)](#)



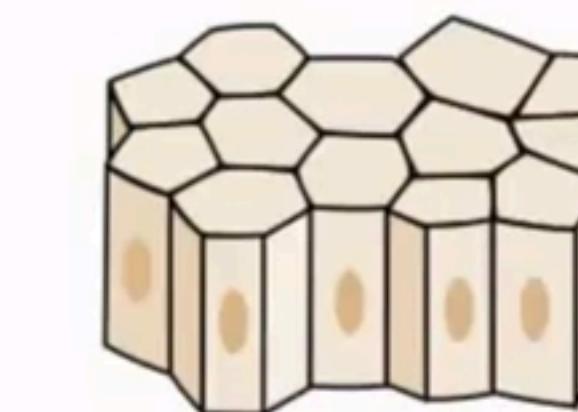
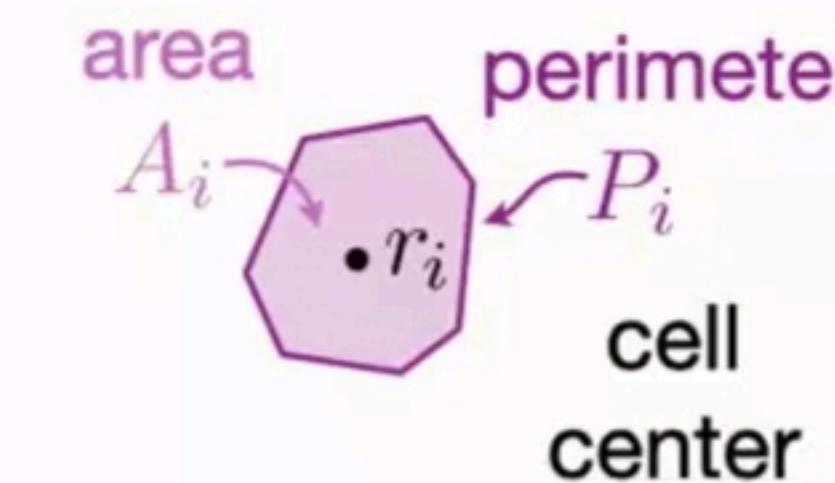
LMP Seminars



mesoscopic



Drosophila, courtesy S. Streichan



cells minimize an energy towards a preferred perimeter and area

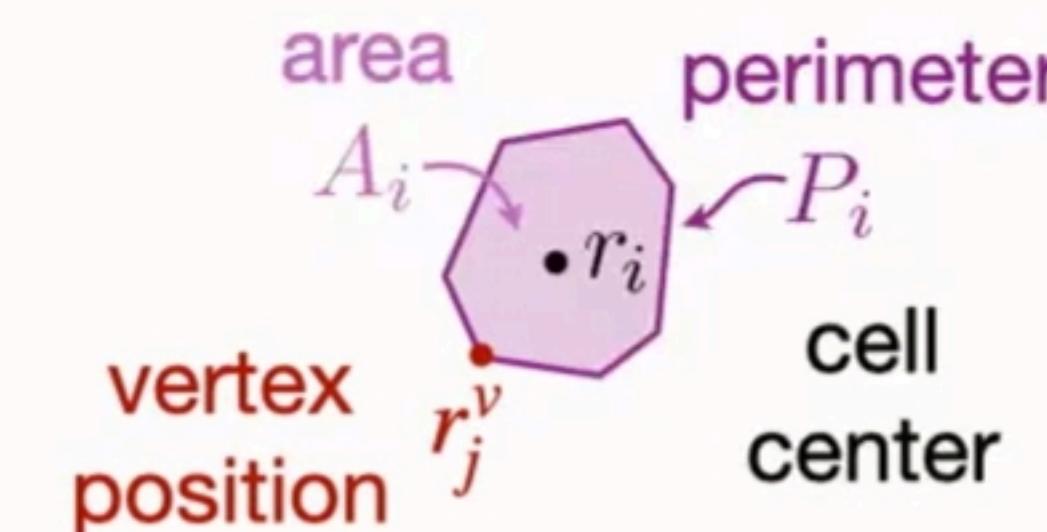




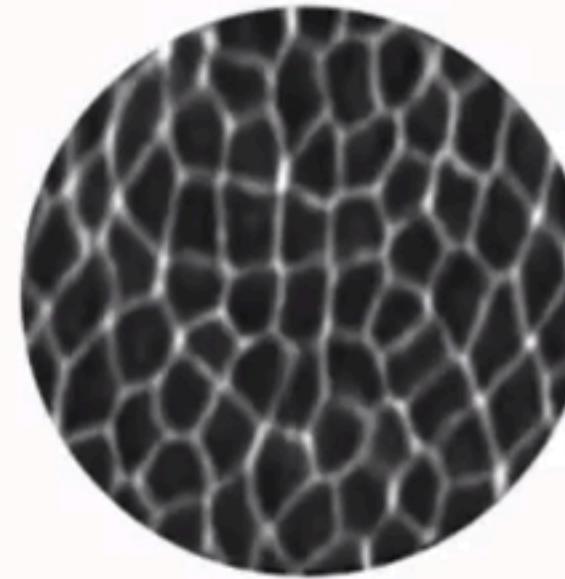
Vertex models

$$A_i = A_i(\{r_j^v\}_j)$$

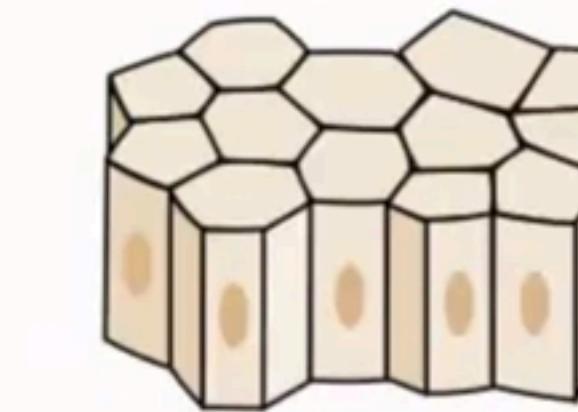
$$P_i = P_i(\{r_j^v\}_j)$$



mesoscopic



Drosophila, courtesy S. Streichan

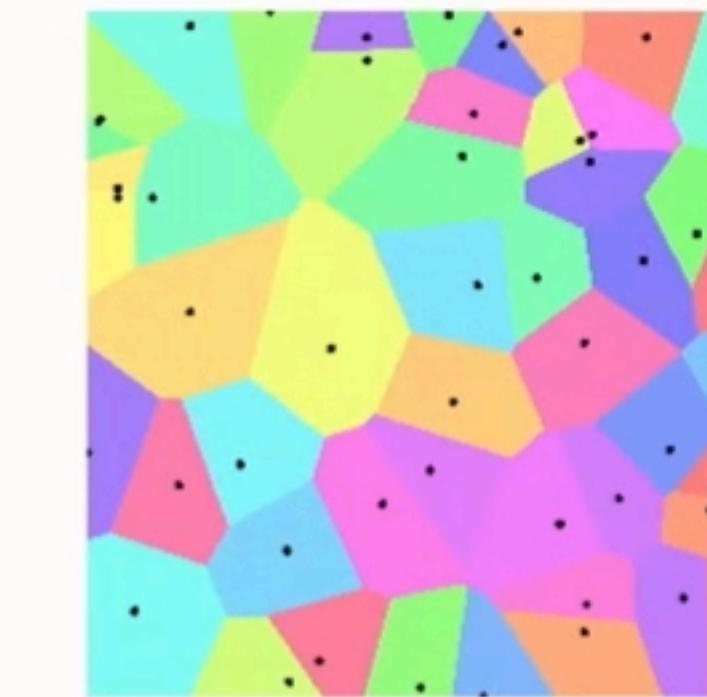


$$\gamma \partial_t \mathbf{r}_i = - \frac{\partial E}{\partial \mathbf{r}_j^v} \frac{\partial \mathbf{r}_j^v}{\partial \mathbf{r}_i} = - \frac{\partial E}{\partial \mathbf{r}_i}$$

$$E = \sum_i \left[\frac{K_P}{2} (P_i - P_0)^2 + \frac{K_A}{2} (A_i - A_0)^2 \right]$$

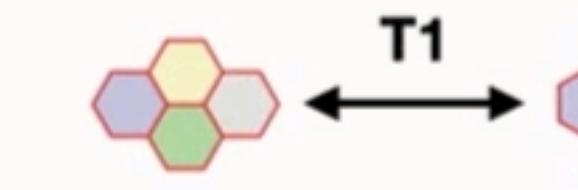
vertex positions can be computed through Voronoi tessellation

Voronoi tessellation



$$\mathbf{r}_j^v(\{r_i\}_i)$$

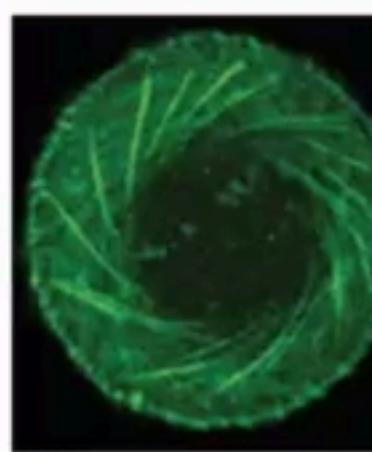
topological transitions



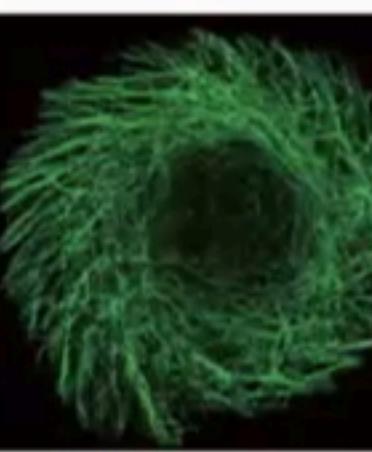
LMP Seminars



How to model chiral tissues?



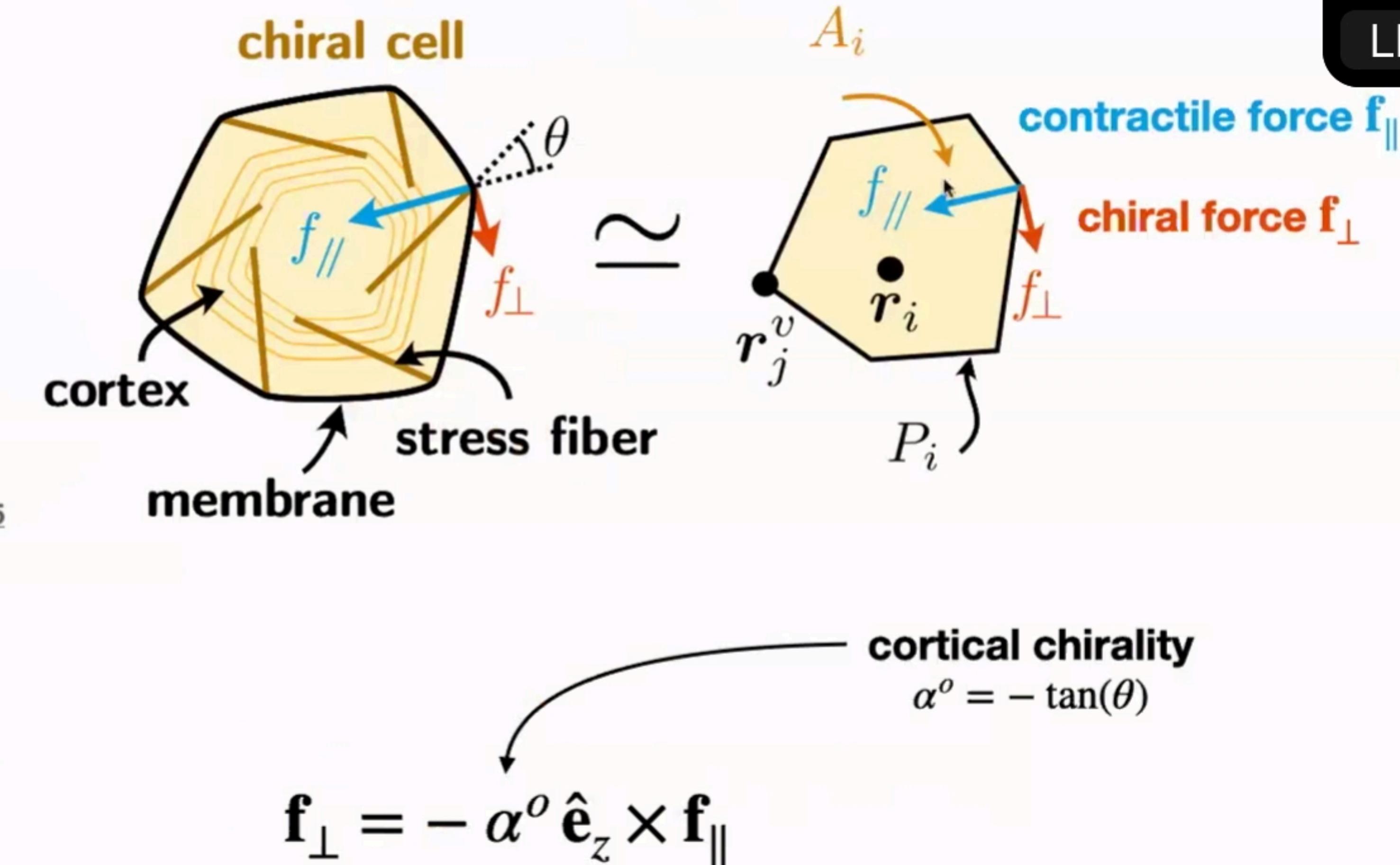
Fibroblasts
[Tee et al., 2015](#)



Carcinoma
[Yamamoto et al., 2025](#)

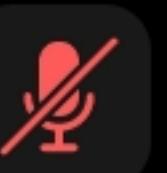


Sihan Chen



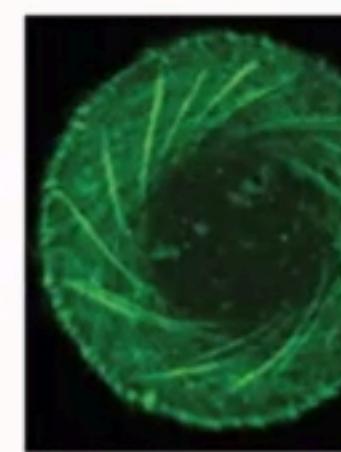
non-conservative force from chiral active forces



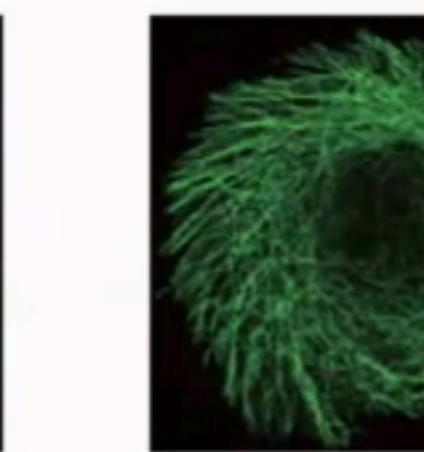


How to model chiral tissues?

cell-level



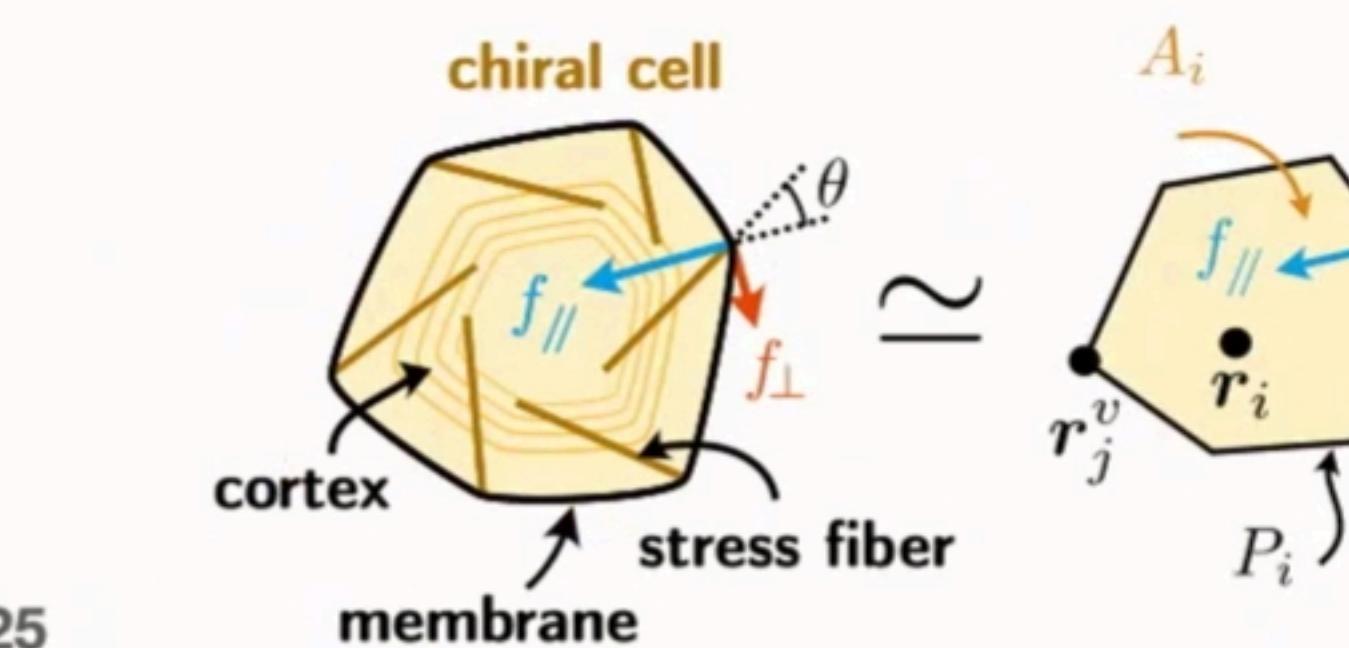
Fibroblasts
Tee et al., 2015



Carcinoma
Yamamoto et al., 2025

mesoscopic

$$E = \sum_i \left[\frac{K_P}{2} (P_i - P_0)^2 + \frac{K_A}{2} (A_i - A_0)^2 \right] = \sum_i \left[\frac{K_P P_i^2}{2} + \frac{K_P P_0 (P_0 - 2P_i)}{2} + \frac{K_A (A_i - A_0)^2}{2} \right]$$



chiral force
 $\mathbf{f}_\perp = -\alpha^o \hat{\mathbf{e}}_z \times \mathbf{r}_j^v$
 $\alpha^o = -\tan(\theta)$

LMP Seminars

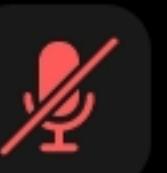
contractile force $\mathbf{f}_\parallel = -\frac{\partial E_C}{\partial \mathbf{r}_j^v}$

chiral force $\mathbf{f}_\perp = \alpha^o \hat{\mathbf{e}}_z \times \frac{\partial E_C}{\partial \mathbf{r}_j^v}$

$$\gamma \partial_t \mathbf{r}_i = \sum_j \left[-\frac{\partial E}{\partial \mathbf{r}_j^v} + \alpha^o \hat{\mathbf{e}}_z \times \frac{\partial E_C}{\partial \mathbf{r}_j^v} \right] \cdot \frac{\partial \mathbf{r}_j^v}{\partial \mathbf{r}_i}$$

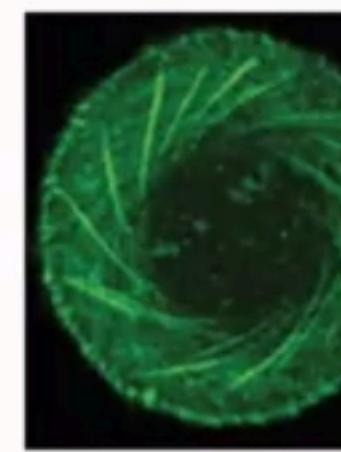
chain rule
does not save you here

chiral nonvariational extension of the vertex model

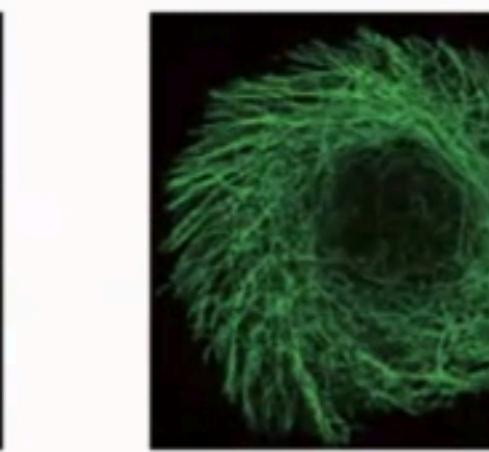


How to model chiral tissues?

cell-level



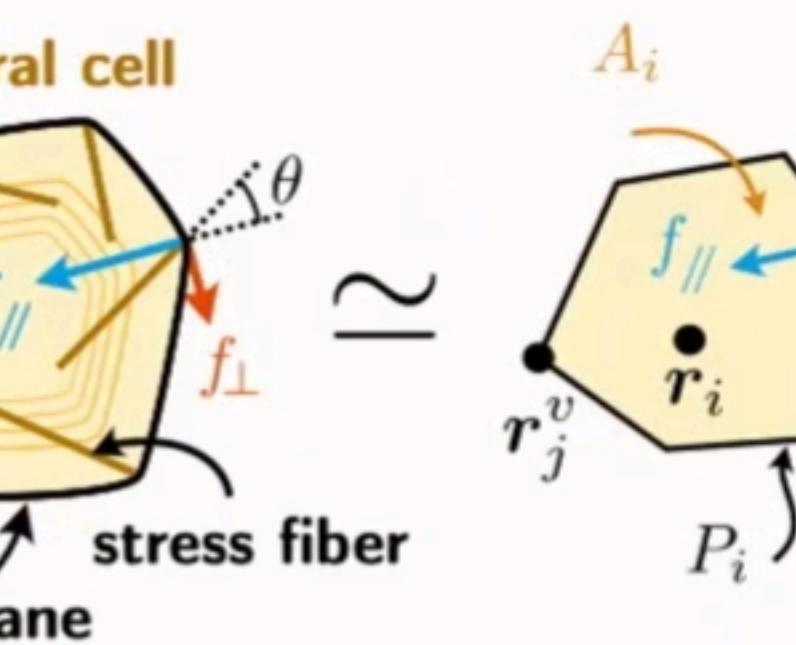
Fibroblasts
Tee et al., 2015



Carcinoma
Yamamoto et al., 2023

mesoscopic

$$E = \sum_i \left[\frac{K_P P_i^2}{2} + \frac{K_A}{2} (A_i - A_0)^2 \right]$$



chiral force
 $\mathbf{f}_\perp = -\alpha^o \hat{\mathbf{e}}_z \times \mathbf{r}_i^\nu$
 $\alpha^o = -\tan(\theta)$

LMP Seminars

continuum



$\mathbf{u}(\mathbf{r}, t)$ $\sigma(\mathbf{r}, t)$

coarse-graining

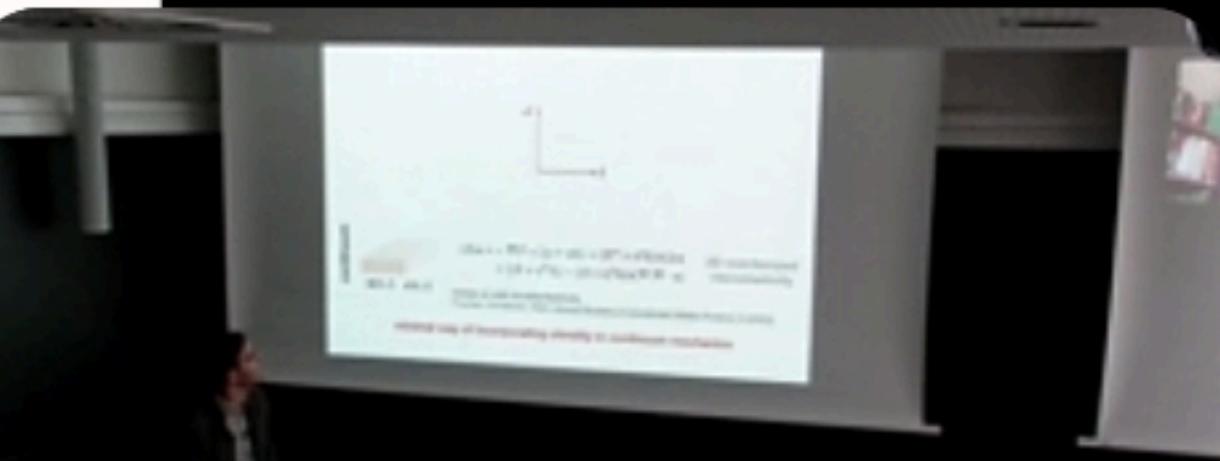
$$\begin{aligned} \zeta \partial_t \mathbf{u} = & -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \boldsymbol{\epsilon}] \Delta \mathbf{u} \\ & + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \boldsymbol{\epsilon}] \nabla (\nabla \cdot \mathbf{u}) \end{aligned}$$

chiral tissues as odd viscoelastic media in the continuum

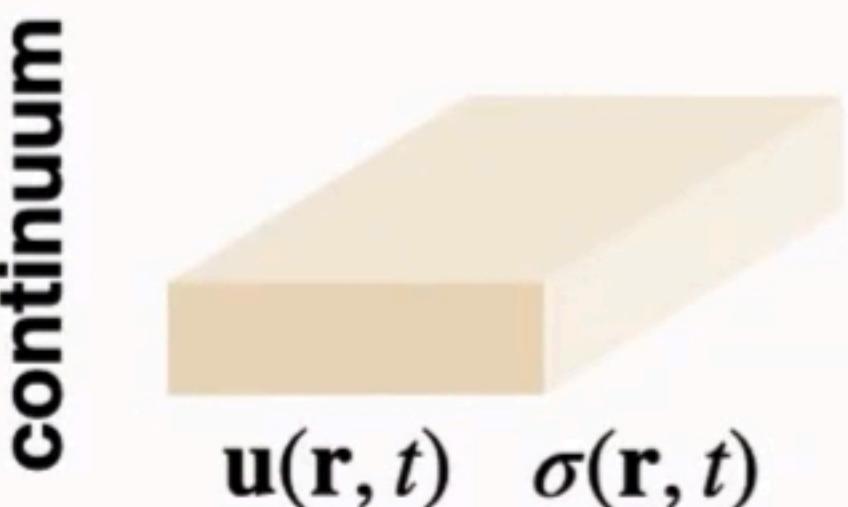
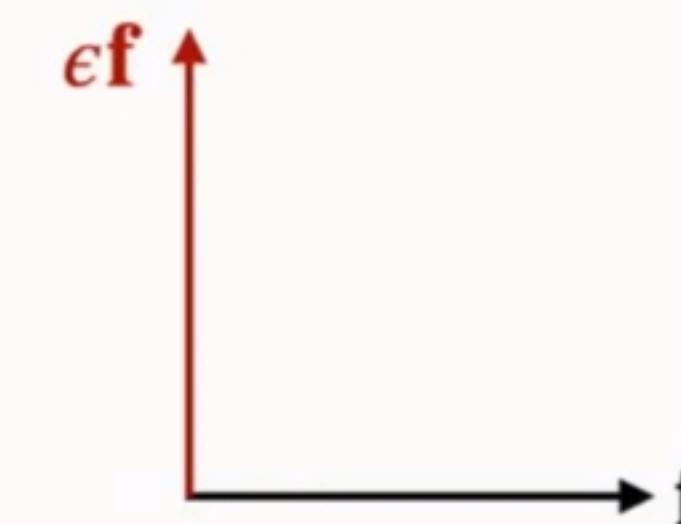




Odd viscosity and elasticity



LMP Seminars



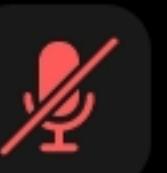
$$\begin{aligned}\zeta \partial_t \mathbf{u} = & -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t)\epsilon] \Delta \mathbf{u} \\ & + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t)\epsilon] \nabla (\nabla \cdot \mathbf{u})\end{aligned}$$

2D overdamped
viscoelasticity

review on odd viscosity/elasticity

Fruchart, Scheibner, Vitelli, [Annual Reviews of Condensed Matter Physics 14 \(2023\)](#)

minimal way of incorporating chirality in continuum mechanics



Odd viscosity and elasticity

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p \\ -\tau \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^h} + \begin{pmatrix} \zeta & \eta^B & 0 & 0 \\ \eta^A & \eta^R & 0 & 0 \\ 0 & 0 & \eta & \eta^o \\ 0 & 0 & -\eta^o & \eta \end{pmatrix}_{\eta_{\alpha\beta}} \begin{pmatrix} \bullet \\ \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{\dot{e}_\beta}$$

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p^{(\text{pre})} \\ -\tau^{(\text{pre})} \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^{(\text{pre})}} + \begin{pmatrix} B & \Lambda & 0 & 0 \\ A & \Gamma & 0 & 0 \\ 0 & 0 & \mu & K^o \\ 0 & 0 & -K^o & \mu \end{pmatrix}_{C_{\alpha\beta}} \begin{pmatrix} \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{e_\beta}$$

viscosity

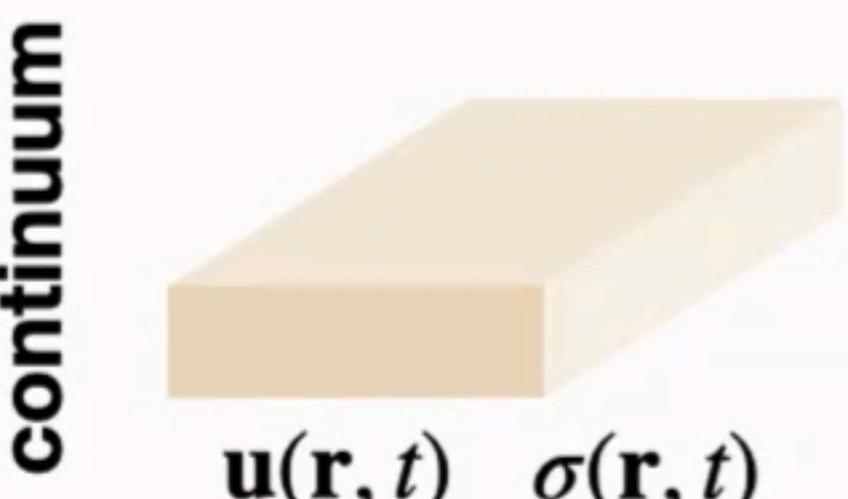
$$M_{ijkl}(\tau) = \eta_{ijkl}\delta(\tau)$$

elasticity

$$M_{ijkl}(\tau) = C_{ijkl}\Theta(\tau)$$

$$\sigma_{ij}(t) = \int_{-\infty}^{\infty} M_{ijkl}(\tau) \partial_\ell \dot{u}_k(t - \tau) d\tau$$

viscoelasticity



$$\begin{aligned} \zeta \partial_t \mathbf{u} = & -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \boldsymbol{\epsilon}] \Delta \mathbf{u} \\ & + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \boldsymbol{\epsilon}] \nabla (\nabla \cdot \mathbf{u}) = \text{Div}(\boldsymbol{\sigma}) \end{aligned}$$

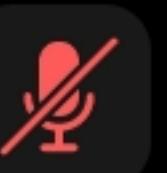
review on odd viscosity/elasticity

Fruchart, Scheibner, Vitelli, [Annual Reviews of Condensed Matter Physics 14 \(2023\)](#)

minimal way of incorporating chirality in continuum mechanics

[I'm not going to talk about convected derivatives because it does not matter here see the paper for references.]





Odd viscosity and elasticity

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p \\ -\tau \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^h} + \begin{pmatrix} \zeta & \eta^B & 0 & 0 \\ \eta^A & \eta^R & 0 & 0 \\ 0 & 0 & \eta & \eta^o \\ 0 & 0 & -\eta^o & \eta \end{pmatrix}_{\eta_{\alpha\beta}} \begin{pmatrix} \bullet \\ \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{\dot{e}_\beta}$$

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p^{(\text{pre})} \\ -\tau^{(\text{pre})} \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^{(\text{pre})}} + \begin{pmatrix} B & \Lambda & 0 & 0 \\ A & \Gamma & 0 & 0 \\ 0 & 0 & \mu & K^o \\ 0 & 0 & -K^o & \mu \end{pmatrix}_{C_{\alpha\beta}} \begin{pmatrix} \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{e_\beta}$$

viscosity

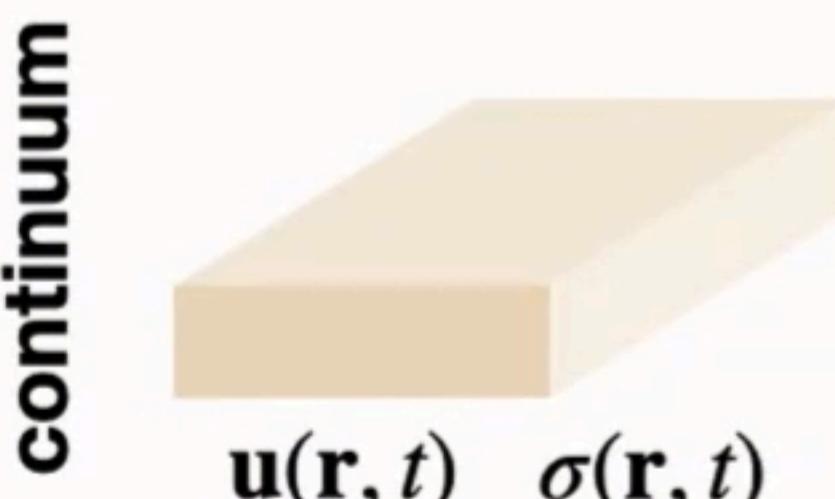
$$M_{ijkl}(\tau) = \eta_{ijkl}\delta(\tau)$$

elasticity

$$M_{ijkl}(\tau) = C_{ijkl}\Theta(\tau)$$

$$\sigma_{ij}(t) = \int_{-\infty}^{\infty} M_{ijkl}(\tau) \partial_\ell \dot{u}_k(t - \tau) d\tau$$

viscoelasticity



$$\begin{aligned} \zeta \partial_t \mathbf{u} = & -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \boldsymbol{\epsilon}] \Delta \mathbf{u} \\ & + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \boldsymbol{\epsilon}] \nabla (\nabla \cdot \mathbf{u}) = \text{Div}(\boldsymbol{\sigma}) \end{aligned}$$

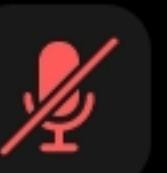
review on odd viscosity/elasticity

Fruchart, Scheibner, Vitelli, [Annual Reviews of Condensed Matter Physics 14 \(2023\)](#)

minimal way of incorporating chirality in continuum mechanics

[I'm not going to talk about convected derivatives because it does not matter here see the paper for references.]





Odd viscosity and elasticity

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p \\ -\tau \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^h} + \begin{pmatrix} \zeta & \eta^B & 0 & 0 \\ \eta^A & \eta^R & 0 & 0 \\ 0 & 0 & \eta & \eta^o \\ 0 & 0 & -\eta^o & \eta \end{pmatrix}_{\eta_{\alpha\beta}} \begin{pmatrix} \bullet \\ \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{\dot{e}_\beta}$$

$$\begin{pmatrix} \oplus \\ \circlearrowleft \\ \oplus \\ \times \end{pmatrix}_{\sigma_\alpha} = \begin{pmatrix} -p^{(\text{pre})} \\ -\tau^{(\text{pre})} \\ 0 \\ 0 \end{pmatrix}_{\sigma_\alpha^{(\text{pre})}} + \begin{pmatrix} B & \Lambda & 0 & 0 \\ A & \Gamma & 0 & 0 \\ 0 & 0 & \mu & K^o \\ 0 & 0 & -K^o & \mu \end{pmatrix}_{C_{\alpha\beta}} \begin{pmatrix} \square \\ \diamond \\ \square \\ \square \end{pmatrix}_{e_\beta}$$

viscosity

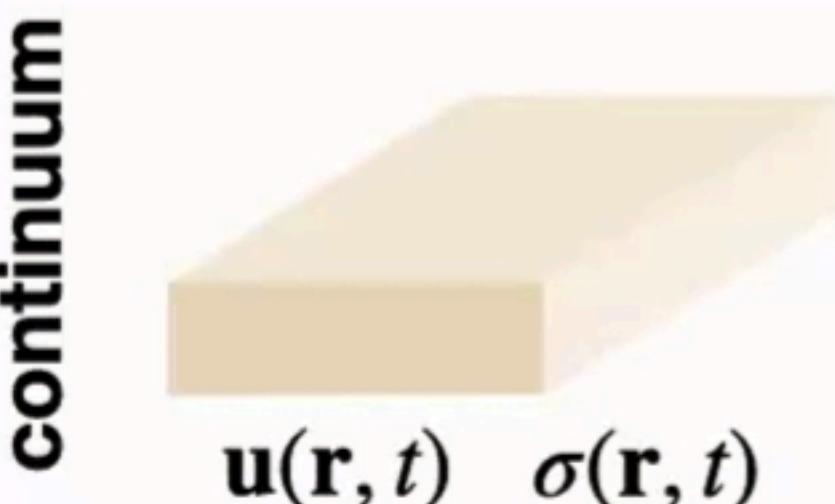
$$M_{ijkl}(\tau) = \eta_{ijkl}\delta(\tau)$$

elasticity

$$M_{ijkl}(\tau) = C_{ijkl}\Theta(\tau)$$

$$\sigma_{ij}(t) = \int_{-\infty}^{\infty} M_{ijkl}(\tau) \partial_\ell \dot{u}_k(t - \tau) d\tau$$

viscoelasticity



$$\begin{aligned} \zeta \partial_t \mathbf{u} = & -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t)\epsilon] \Delta \mathbf{u} \\ & + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t)\epsilon] \nabla (\nabla \cdot \mathbf{u}) = \text{Div}(\sigma) \end{aligned}$$

review on odd viscosity/elasticity

Fruchart, Scheibner, Vitelli, [Annual Reviews of Condensed Matter Physics 14 \(2023\)](#)

minimal way of incorporating chirality in continuum mechanics

[I'm not going to talk about convected derivatives because it does not matter here see the paper for references.]



6:46

29.8 KB/s 4G 9%



LMP Seminars



How to model chiral tissues?

cell-level

Fibroblasts Tee et al., 2015 Carcinoma Yamamoto et al., 2023

mesoscopic

chiral cell
cortex
stress fiber
membrane

$f_{\perp} = -\alpha^o \hat{\mathbf{e}}_z \times \mathbf{f}_{\parallel}$

$\alpha^o = -\tan(\theta)$

$$E = \sum_i \left[\frac{K_P}{2} P_i^2 + \frac{K_P}{2} P_0(P_0 - 2P_i) + \frac{K_A}{2} (A_i - A_0)^2 \right]$$

$$\gamma \partial_t \mathbf{r}_i = \sum_j \left[-\frac{\partial E}{\partial \mathbf{r}_j^v} + \alpha^o \hat{\mathbf{e}}_z \times \frac{\partial E_C}{\partial \mathbf{r}_j^v} \right] \cdot \frac{\partial \mathbf{r}_j^v}{\partial t}$$

continuum

$\zeta \partial_t \mathbf{u} = -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \epsilon] \Delta \mathbf{u} + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \epsilon] \nabla (\nabla \cdot \mathbf{u}).$

$\mathbf{u}(\mathbf{r}, t) \quad \sigma(\mathbf{r}, t)$

how to measure parameters from noisy experimental data of living tissues?

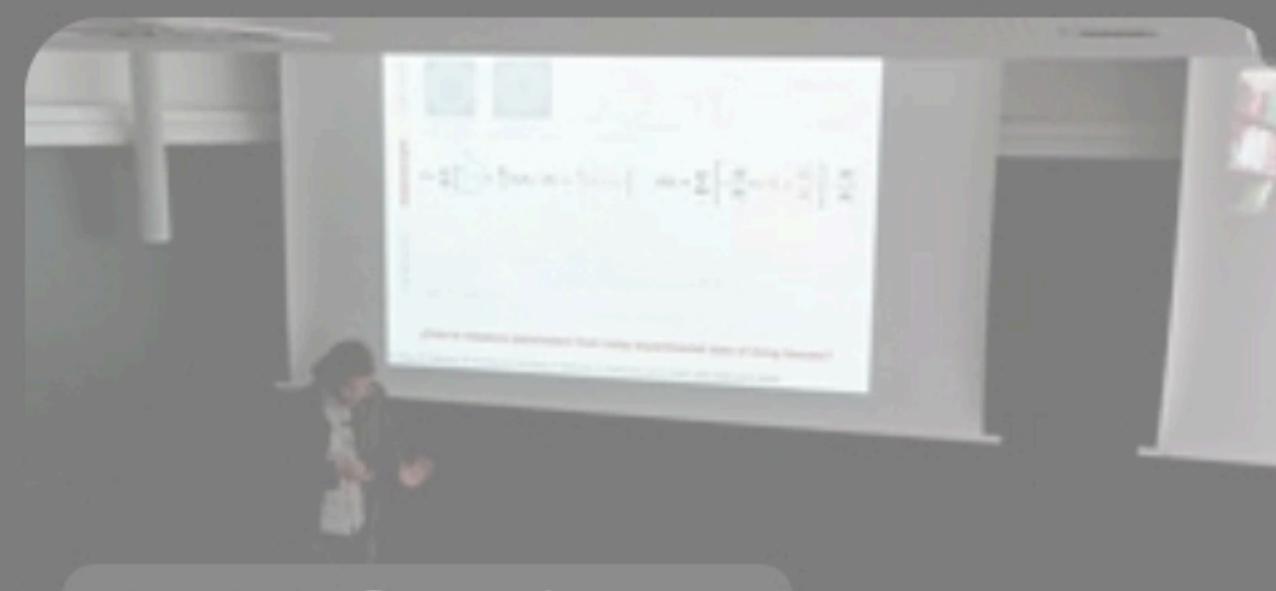
S. Chen, E. Gökmen, M. Fruchart, M. Krumbein, P. Silberzan, V. Yashunsky and V. Vitelli, arXiv 2506.12276 (2025)



LMP Seminars's screen

6:46

29.8 KB/s 4G 9%



LMP Seminars



How to model chiral tissues?

cell-level

Fibroblasts
Tee et al., 2015

Carcinoma
Yamamoto et al., 2023

mesoscopic

$$E = \sum_i \left[\frac{K_P}{2} P_i^2 + \frac{K_P}{2} P_0(P_0 - 2P_i) + \frac{K_A}{2} (A_i - A_0)^2 \right]$$
 $\gamma \partial_t \mathbf{r}_i = \sum_j \left[-\frac{\partial E}{\partial \mathbf{r}_j^\nu} + \alpha^\circ \hat{\mathbf{e}}_z \times \frac{\partial E_C}{\partial \mathbf{r}_j^\nu} \right] \cdot \frac{\partial \mathbf{r}_j^\nu}{\partial \mathbf{r}_i}$

continuum

$$\zeta \partial_t u = -\nabla P + [(\mu + \eta \partial_t) + (K^\circ + \eta^\circ \partial_t)\epsilon] \Delta u + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t)\epsilon] \nabla (\nabla \cdot u).$$
 $u(\mathbf{r}, t) \quad \sigma(\mathbf{r}, t)$

chiral force

$$\mathbf{f}_\perp = -\alpha^\circ \hat{\mathbf{e}}_z \times \mathbf{f}_\parallel$$

$$\alpha^\circ = -\tan(\theta)$$

Diagram: A schematic of a chiral cell showing internal structures: cortex, membrane, and stress fiber. A coordinate system is defined with axes \mathbf{f}_\parallel and \mathbf{f}_\perp . The angle θ is shown between the stress fiber and the \mathbf{f}_\parallel axis.

↳ how to measure parameters from noisy experimental data of living tissues?

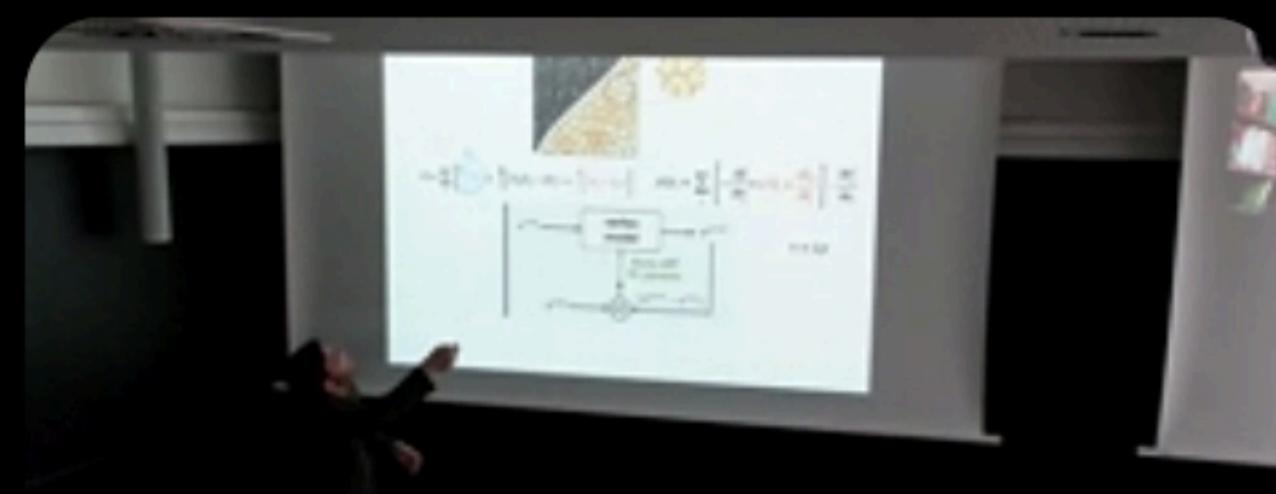
S. Chen, E. Gökmen, M. Fruchart, M. Krumbein, P. Silberzan, V. Yashunsky and V. Vitelli, arXiv 2506.12276 (2025)



LMP Seminars's screen

6:46

41.6 KB/s 4G 9%



LMP Seminars



Let us fit the parameters!

$$E = \sum_i \left[\frac{K_P}{2} P_i^2 + \frac{K_P}{2} P_0(P_0 - 2P_i) + \frac{K_A}{2} (A_i - A_0)^2 \right] \quad \gamma \partial_t \mathbf{r}_i = \sum_j \left[-\frac{\partial E}{\partial \mathbf{r}_j^\nu} + \alpha^o \hat{\mathbf{e}}_z \times \frac{\partial E_C}{\partial \mathbf{r}_j^\nu} \right] \cdot \frac{\partial \mathbf{r}_j^\nu}{\partial \mathbf{r}_i}$$

r^{exp} → **vertex model** → v^{pred}

v^{exp} → \oplus → $\|v^{\text{pred}} - v^{\text{exp}}\|$

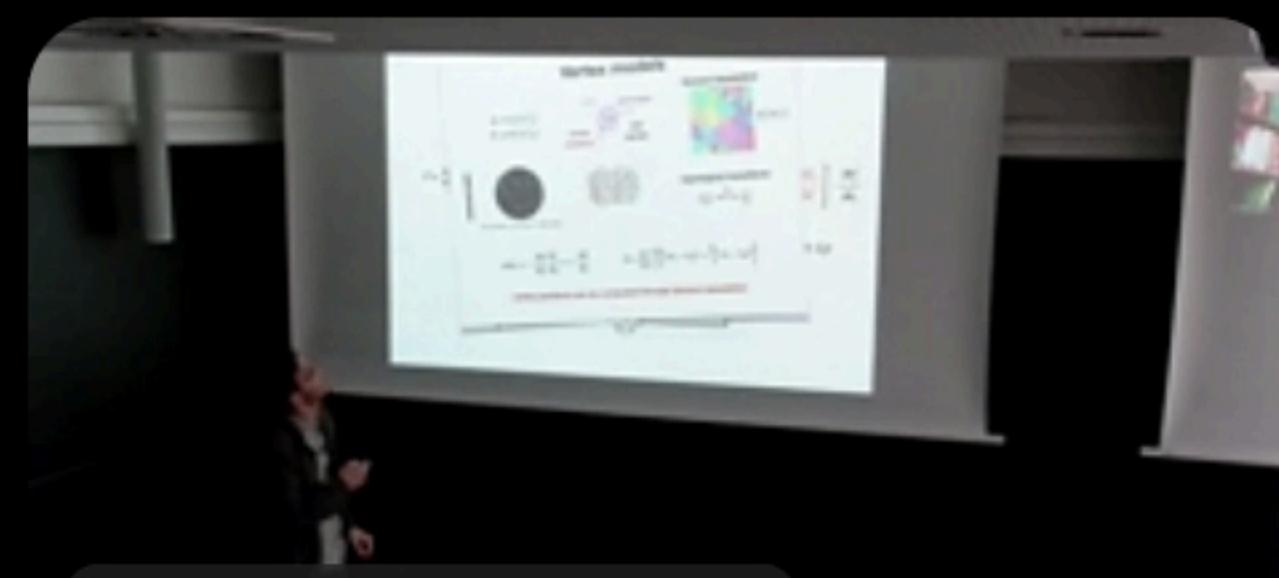
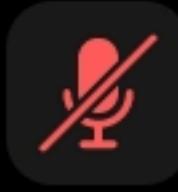
Auto-diff fit params

$\mathbf{v} = \partial_t \mathbf{r}$



6:47

36.5 KB/s 4G 9%



LMP Seminars



Let us fit the parameters!

Vertex models

$A_i = A_i(\{r_j^v\}_j)$ $P_i = P_i(\{r_j^v\}_j)$

area perimeter
vertex position r_j^v cell center

Voronoi tessellation

$\mathbf{r}_j^v(\{\mathbf{r}_i\}_i)$

Drosophila, courtesy S. Streicher

$E = \sum_i$

mesoscopic

$\gamma \partial_i \mathbf{r}_i = - \frac{\partial E}{\partial \mathbf{r}_i} \frac{\partial r_j^v}{\partial \mathbf{r}_i} = - \frac{\partial E}{\partial \mathbf{r}_i}$

$E = \sum_i \left[\frac{K_p}{2} (P_i - P_0)^2 + \frac{K_A}{2} (A_i - A_0)^2 \right]$

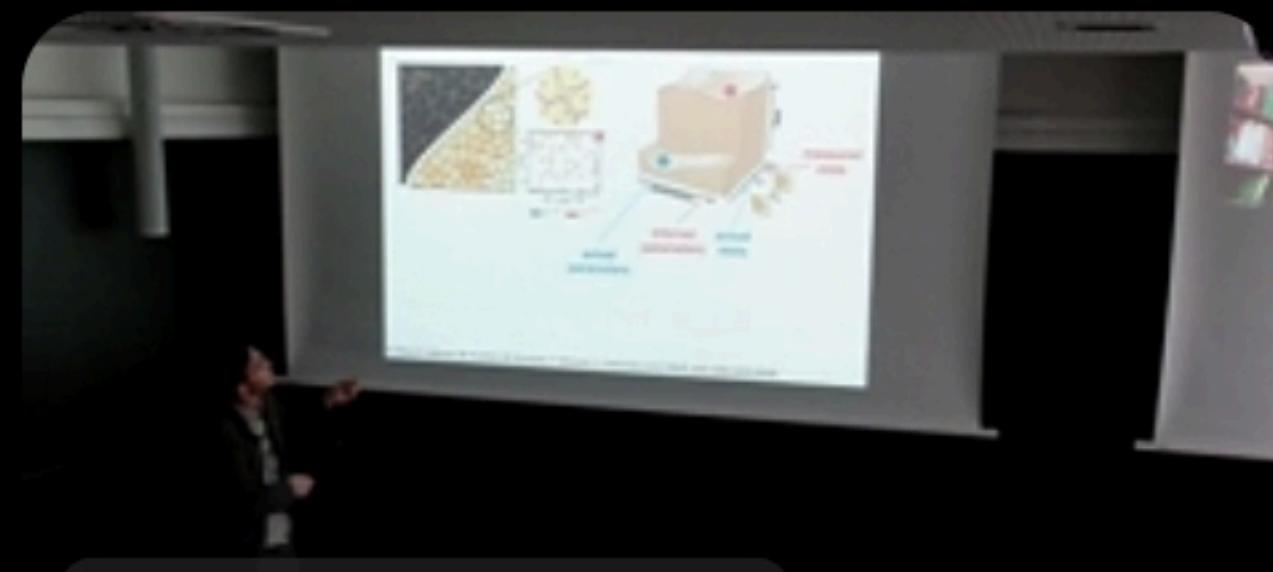
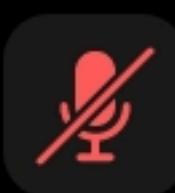
topological transitions

$\frac{\partial E_C}{\partial \mathbf{r}_j^v} \cdot \frac{\partial \mathbf{r}_j^v}{\partial \mathbf{r}_i}$

$= \partial_i \mathbf{r}$

vertex positions can be computed through Voronoi tessellation

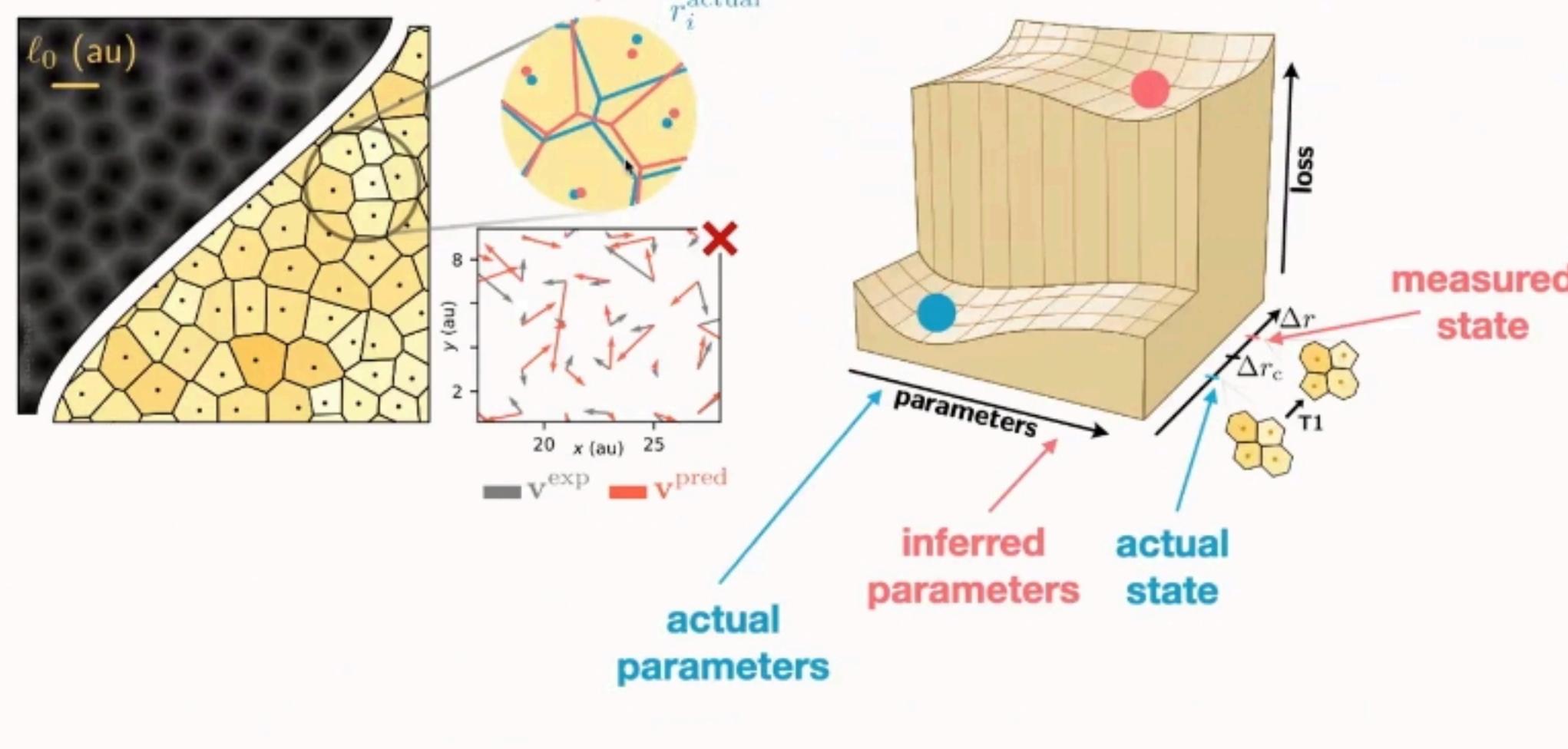




LMP Seminars



On the dark side of the topological transition

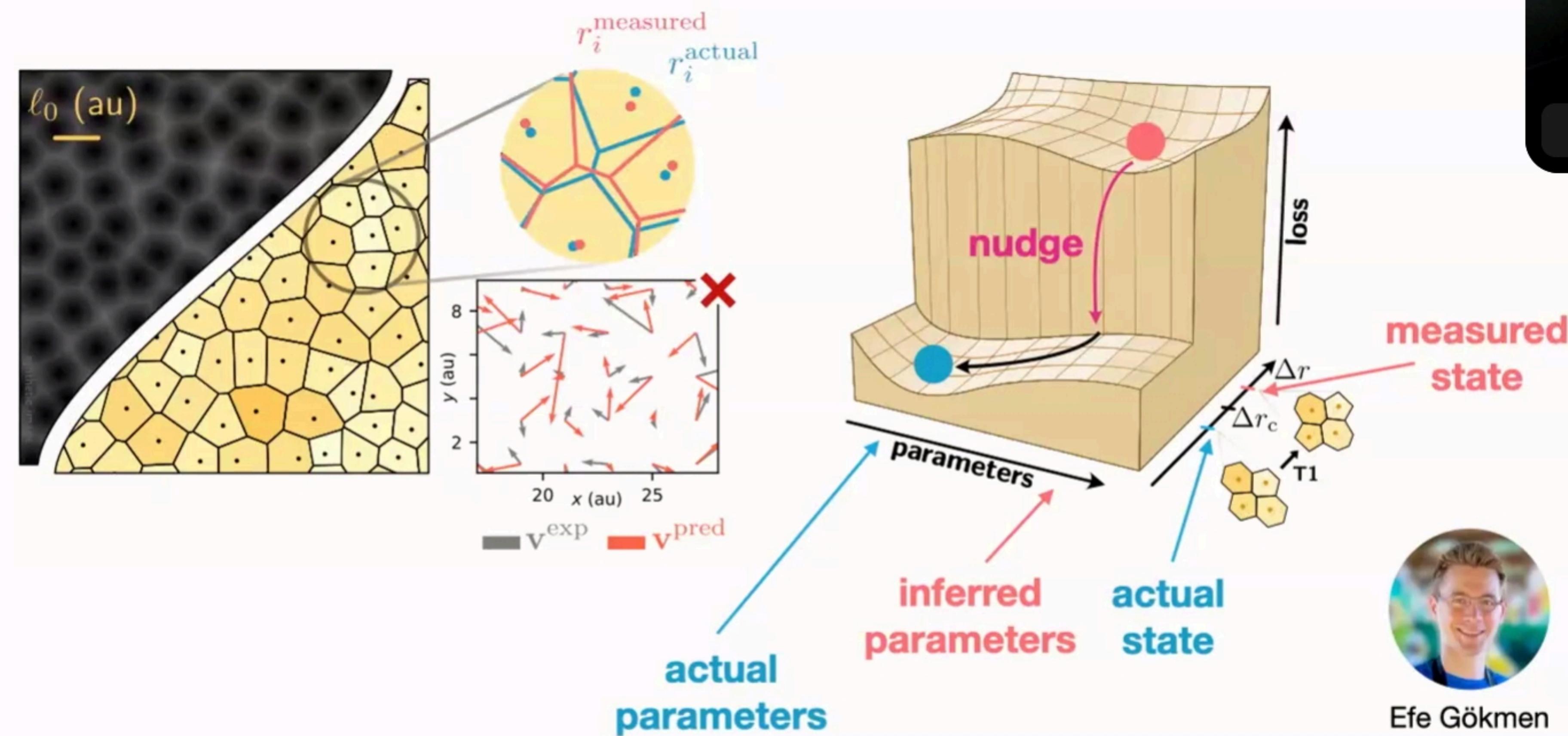


S. Chen, E. Gökm en, M. Fruchart, M. Krumbein, P. Silberzan, V. Yashunsky and V. Vitelli, arXiv 2506.12276 (2025)





On the dark side of the topological transition



LMP Seminars

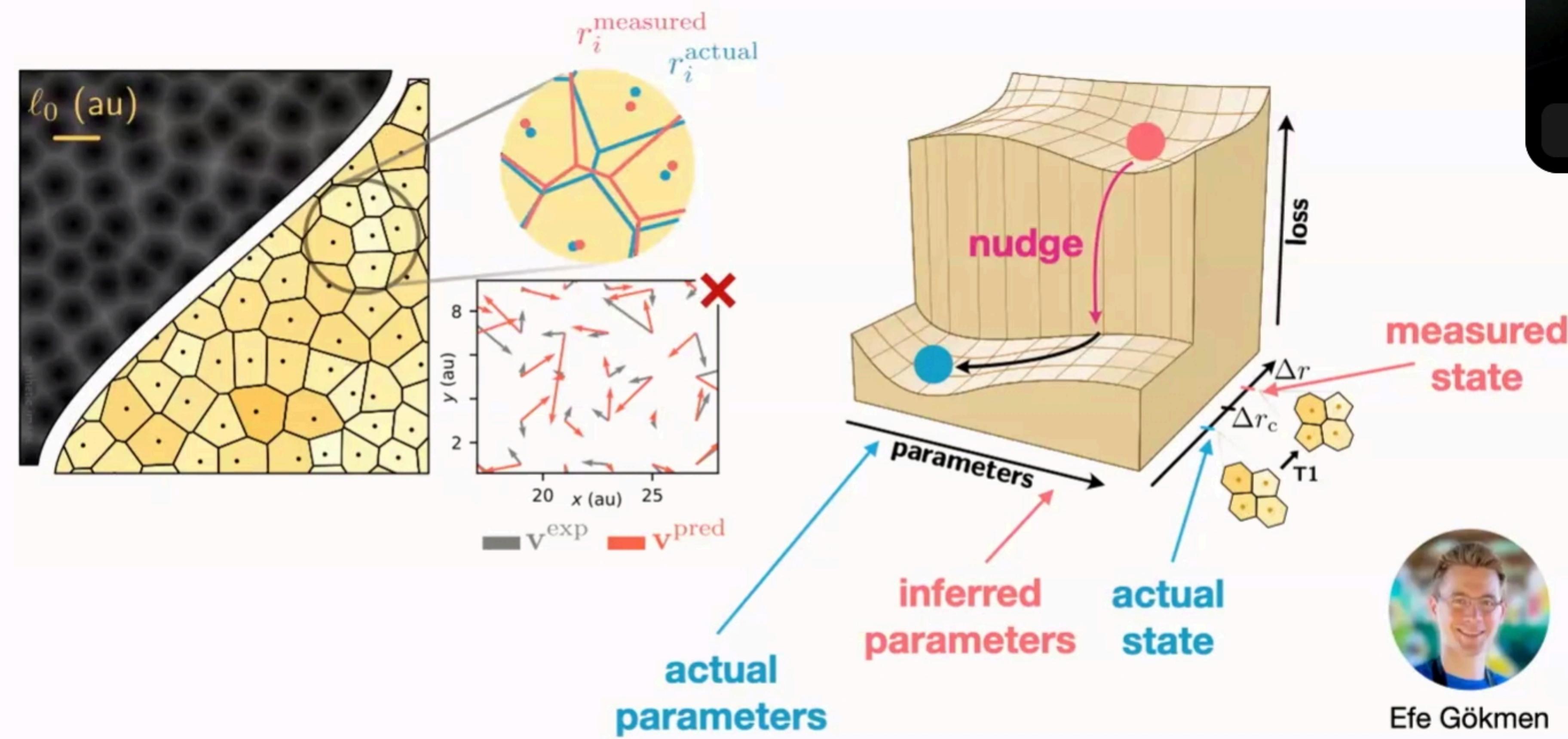


Efe Gökmen

fit data to model to fit model to data



On the dark side of the topological transition



LMP Seminars

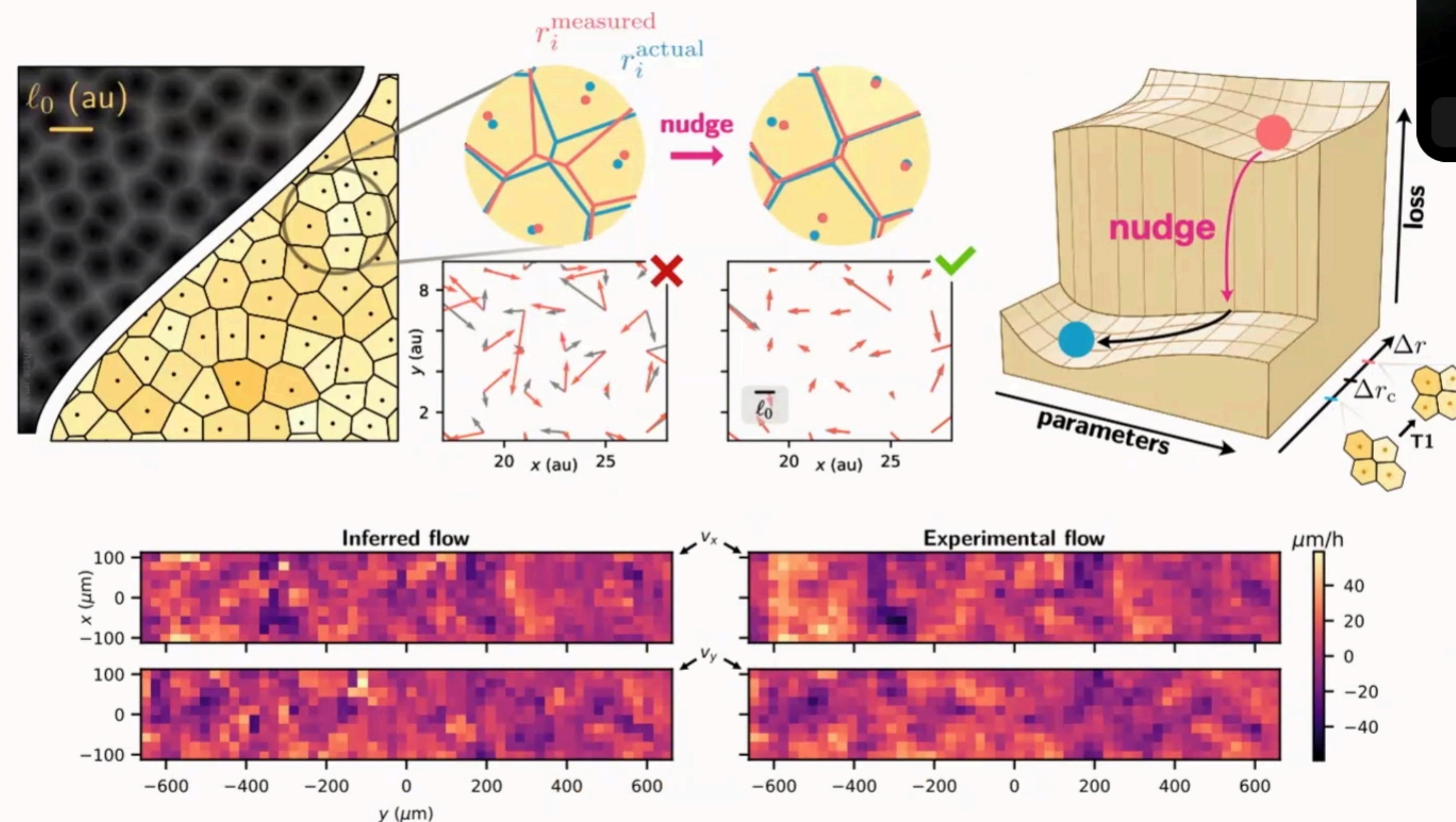


Efe Gökmen

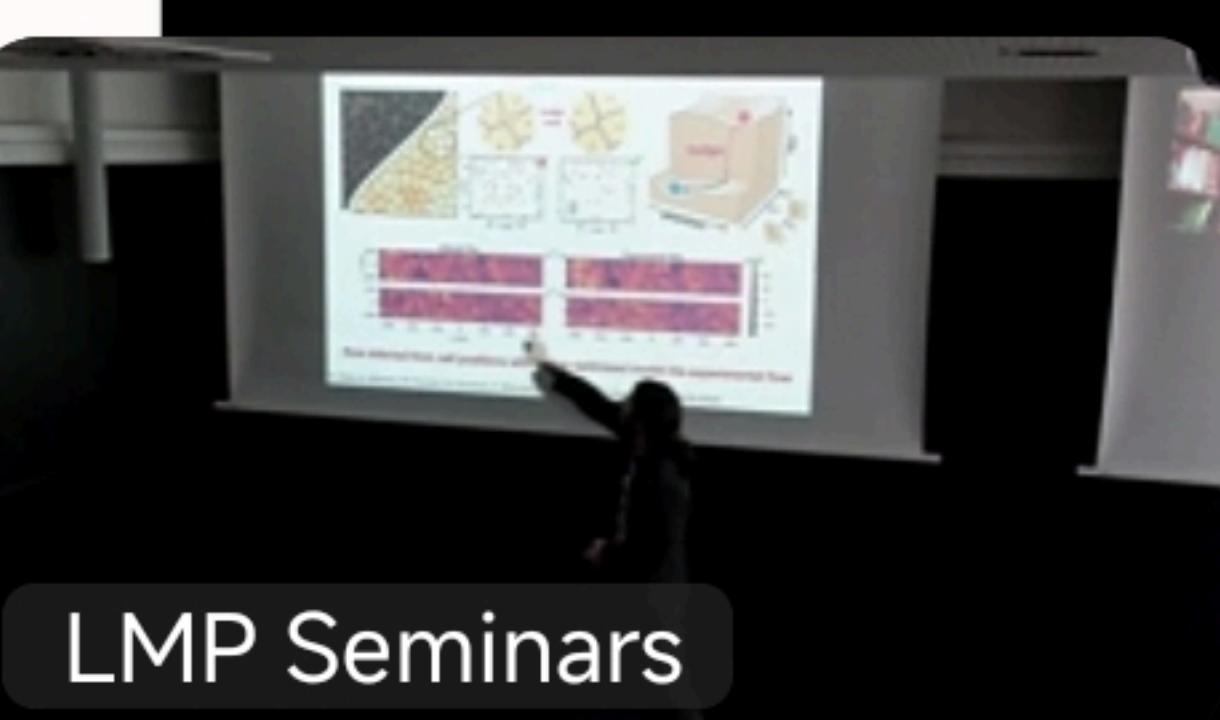
fit data to model to fit model to data



On the dark side of the topological transition

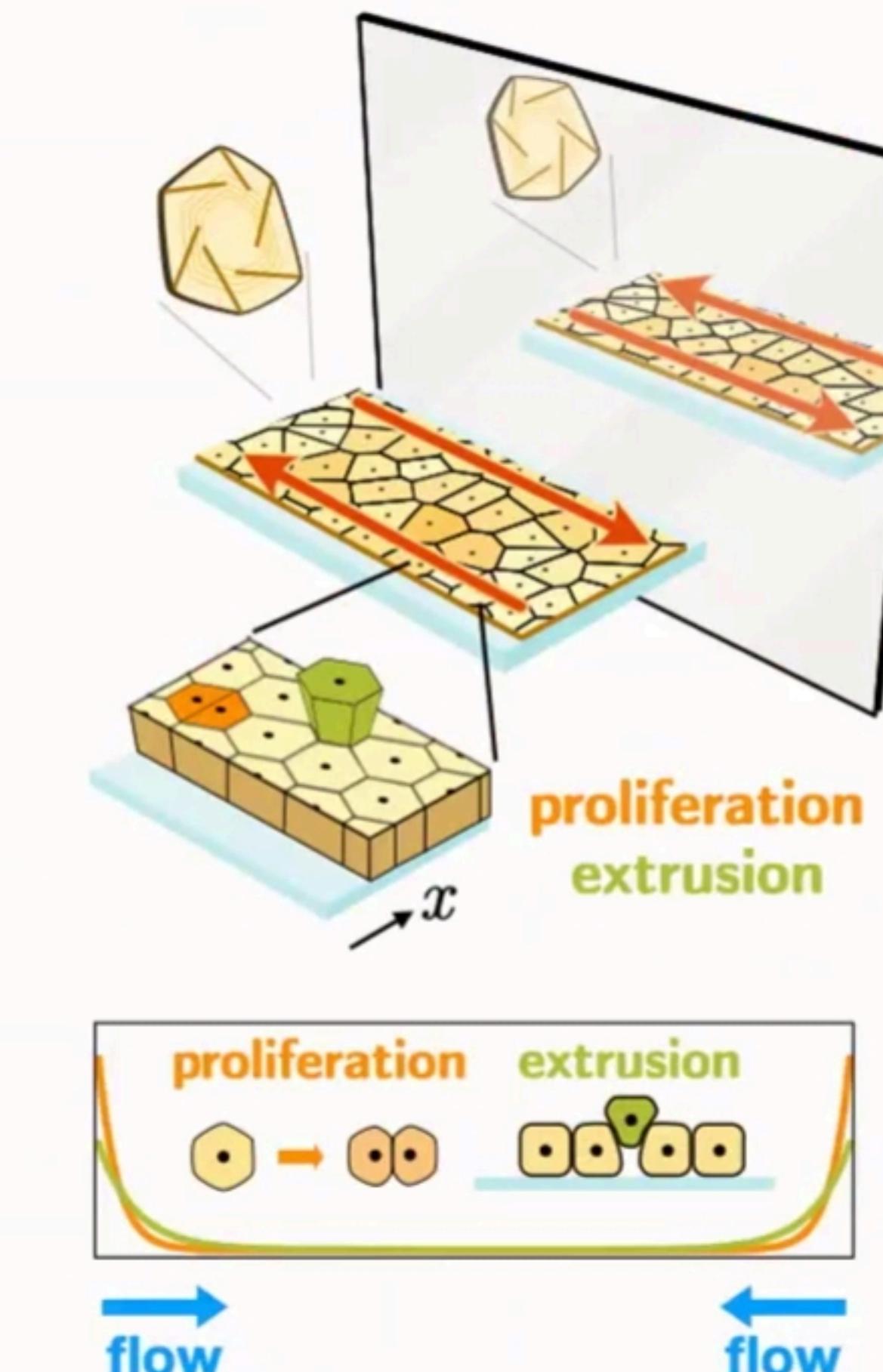
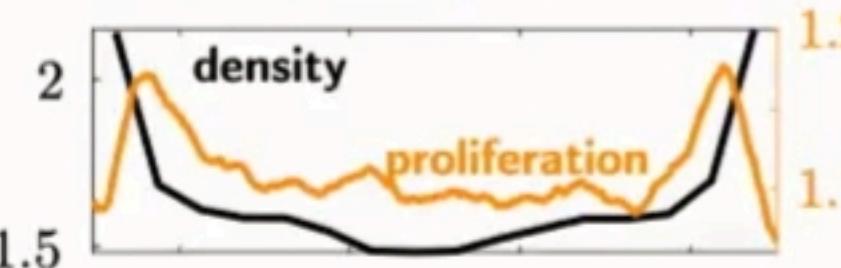
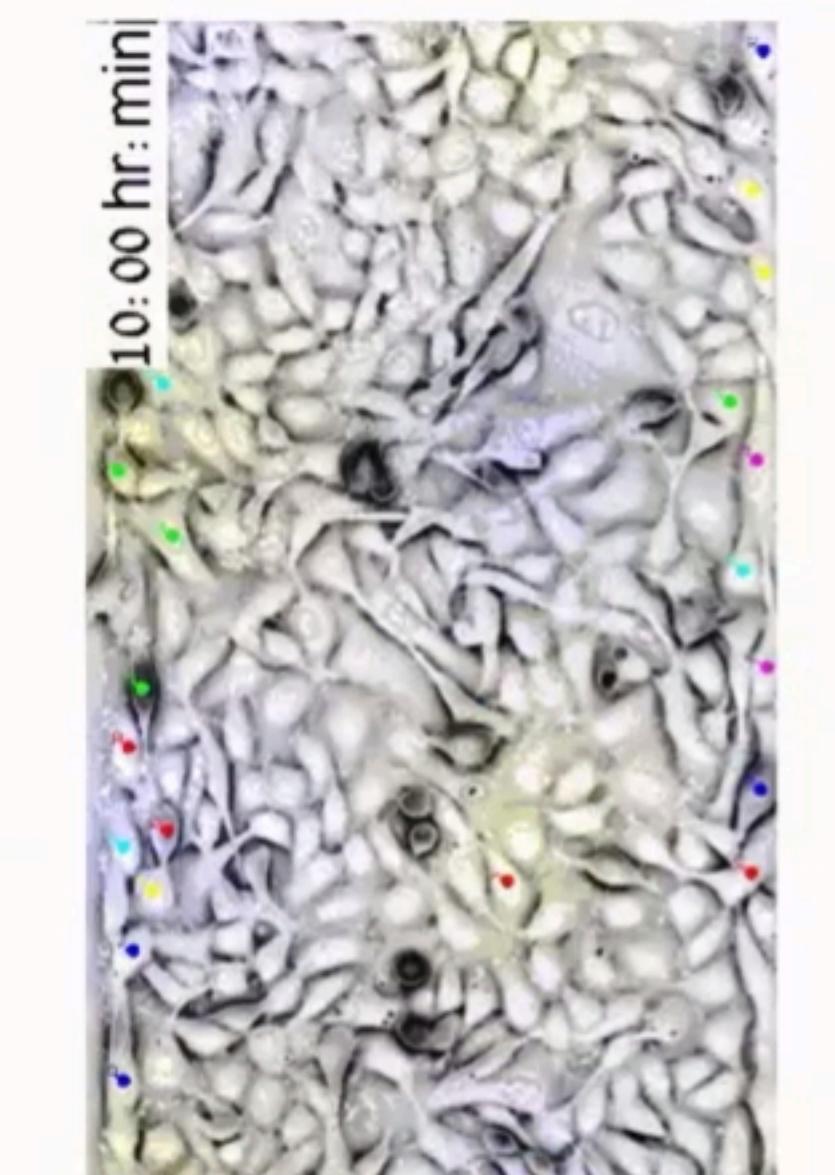


flow inferred from cell positions with nudge-optimized model fits experimental flow



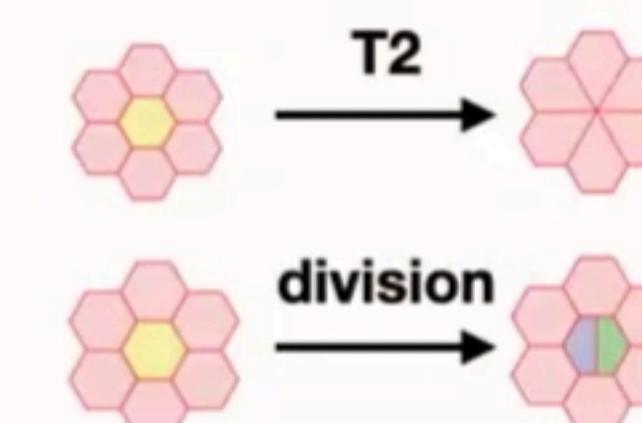


Cells divide



inhomogeneous proliferation and extrusion near walls

mesoscopic



continuum

$$\partial_t \rho + \nabla \cdot [\rho v] = [k_{\text{prolif}}(\mathbf{r}) - k_{\text{extr}}(\mathbf{r})]\rho$$

see e.g.
Hallatschek et al., *Nat Rev Phys* (2023)

LMP Seminars





Participants (5)



Naman Dixit (me)



LMP Seminars (Host)



PB Philip Bittihn (Co-host)



GL Giorgio Lovato



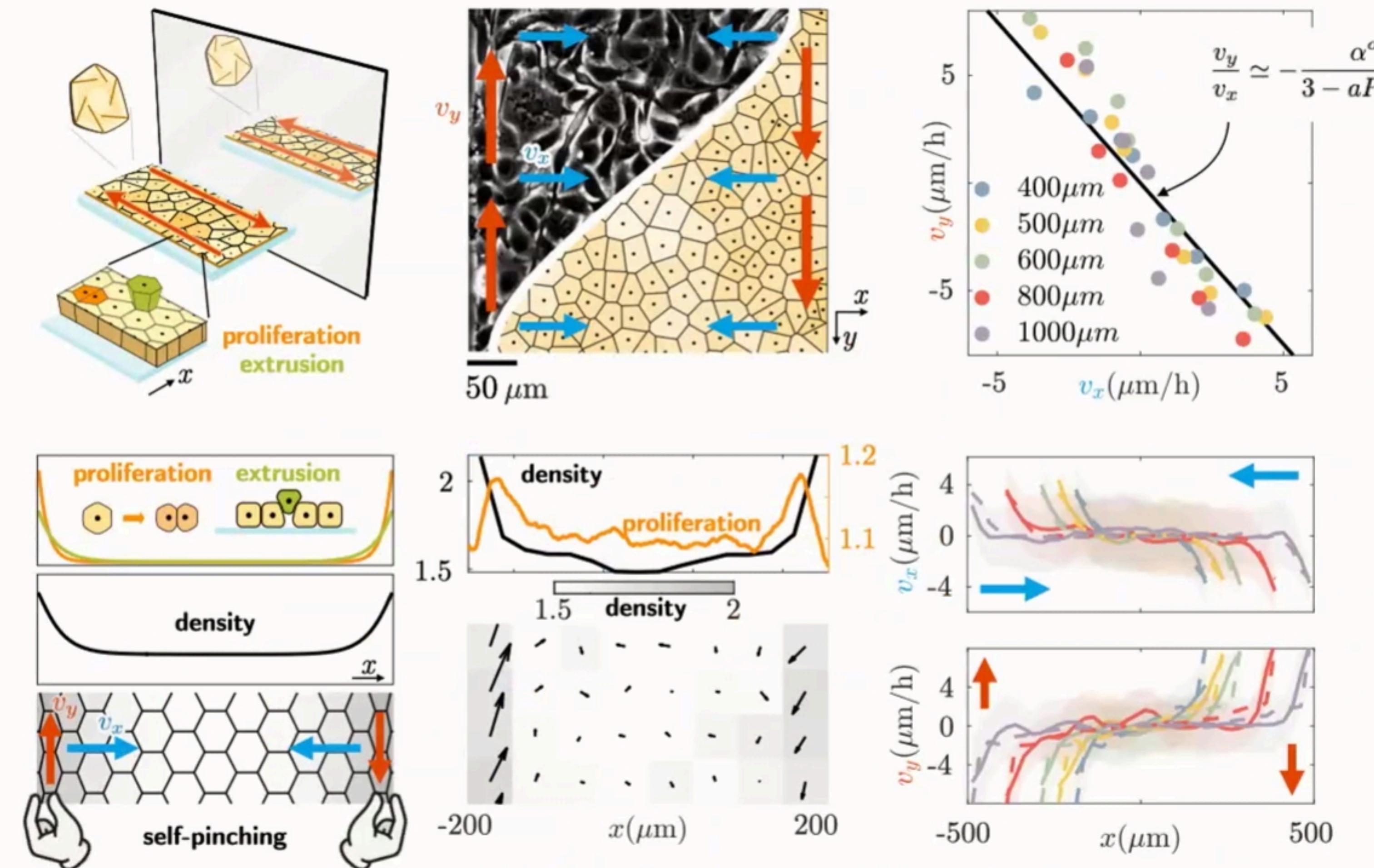
Yuto Hosaka (保阪 悠人)



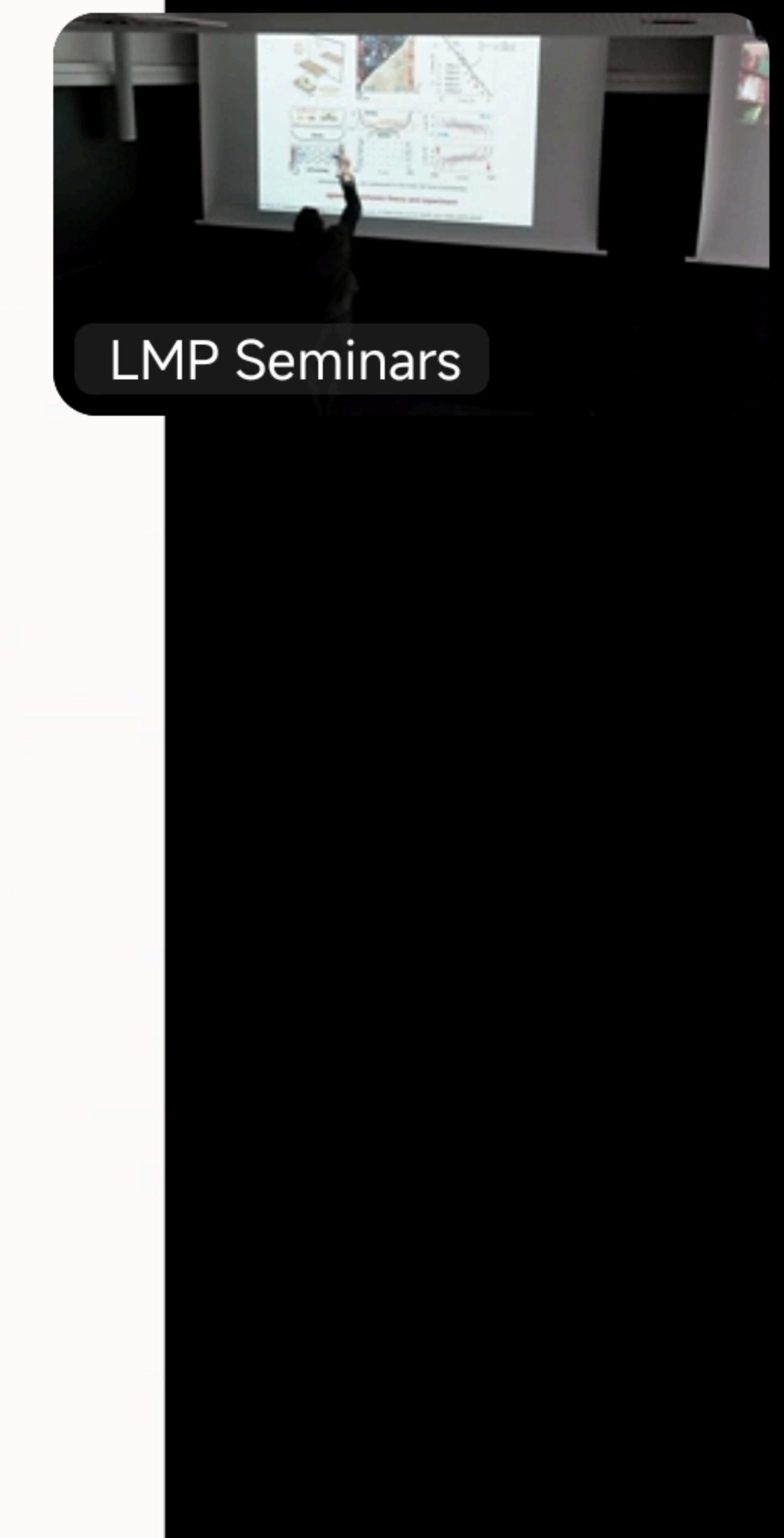
Invite



From self-pinching to chiral edge flow



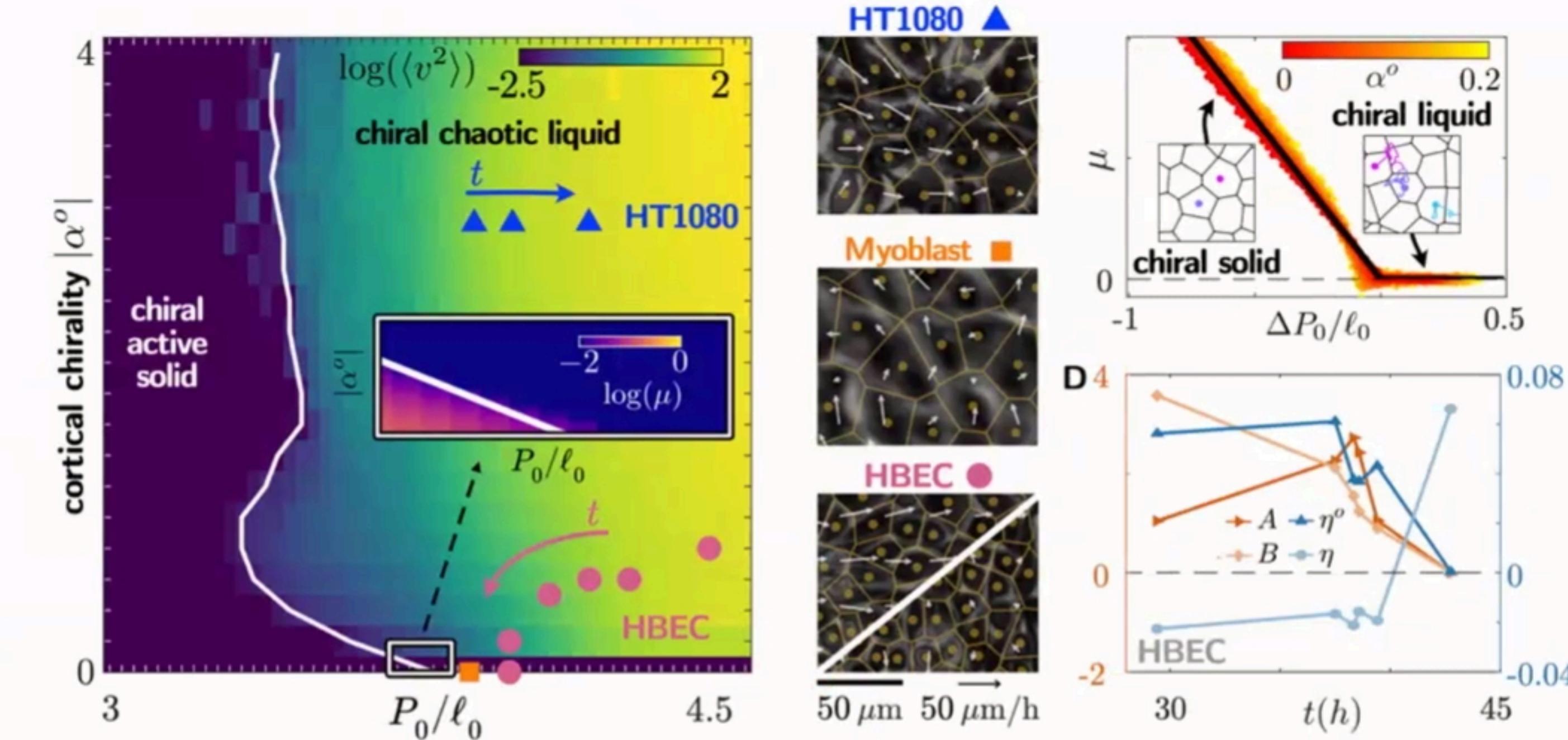
agreement between theory and experiment





Measuring viscoelastic moduli

continuum
 $\mathbf{u}(\mathbf{r}, t)$ $\sigma(\mathbf{r}, t)$
 $\rho(\mathbf{r}, t)$



$$\zeta \partial_t \mathbf{u} = -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \epsilon] \Delta \mathbf{u} + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \epsilon] \nabla (\nabla \cdot \mathbf{u}).$$

$$\partial_t \rho + \nabla \cdot [\rho \mathbf{v}] = [k_{\text{prolif}}(\mathbf{r}) - k_{\text{extr}}(\mathbf{r})] \rho$$

odd
viscoelasticity

inhomogeneous
proliferation

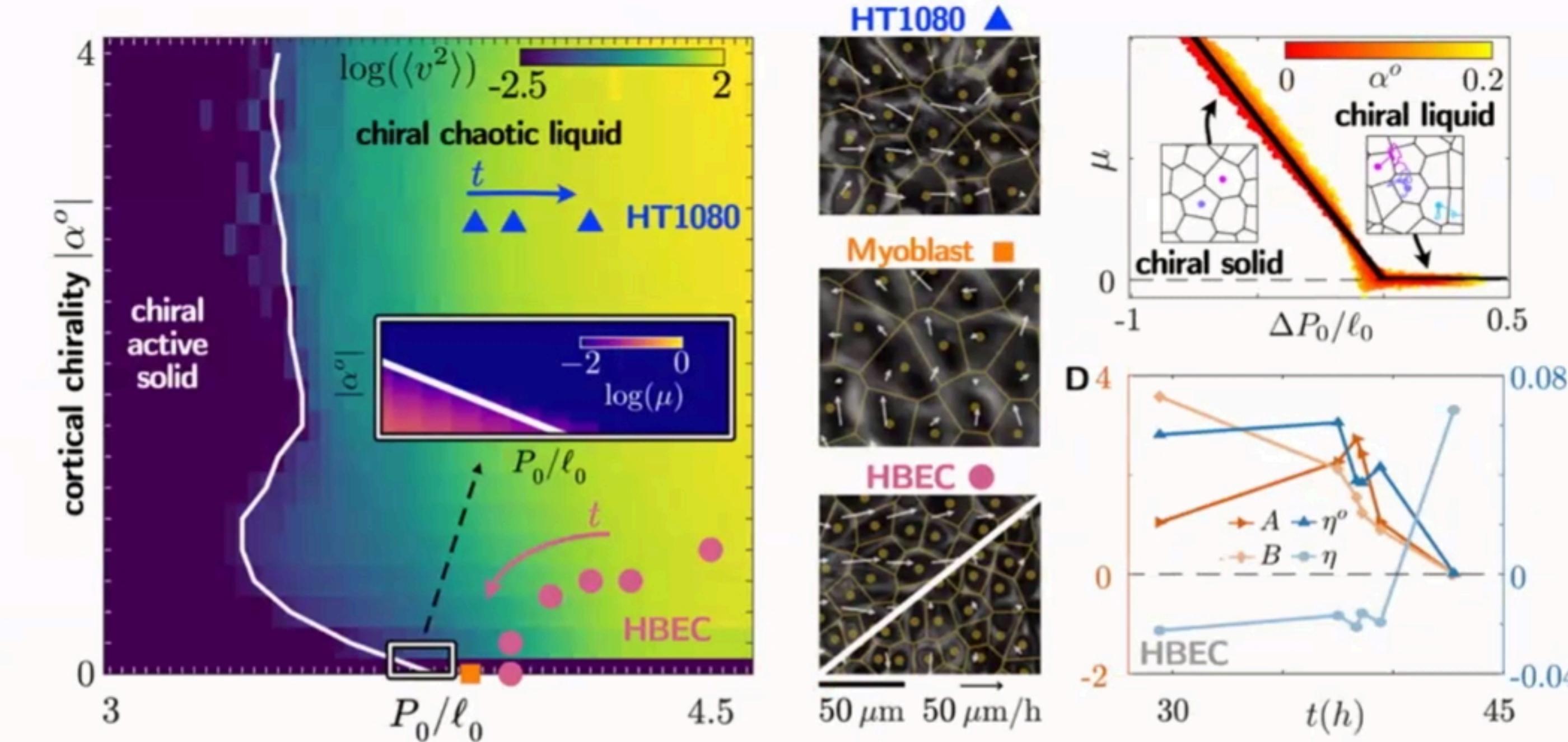
moduli can be time-dependent





Measuring viscoelastic moduli

continuum
 $\mathbf{u}(\mathbf{r}, t)$ $\sigma(\mathbf{r}, t)$
 $\rho(\mathbf{r}, t)$



$$\zeta \partial_t \mathbf{u} = -\nabla P + [(\mu + \eta \partial_t) + (K^o + \eta^o \partial_t) \epsilon] \Delta \mathbf{u} + [(B + \eta^B \partial_t) - (A + \eta^A \partial_t) \epsilon] \nabla (\nabla \cdot \mathbf{u}).$$

$$\partial_t \rho + \nabla \cdot [\rho \mathbf{v}] = [k_{\text{prolif}}(\mathbf{r}) - k_{\text{extr}}(\mathbf{r})] \rho$$

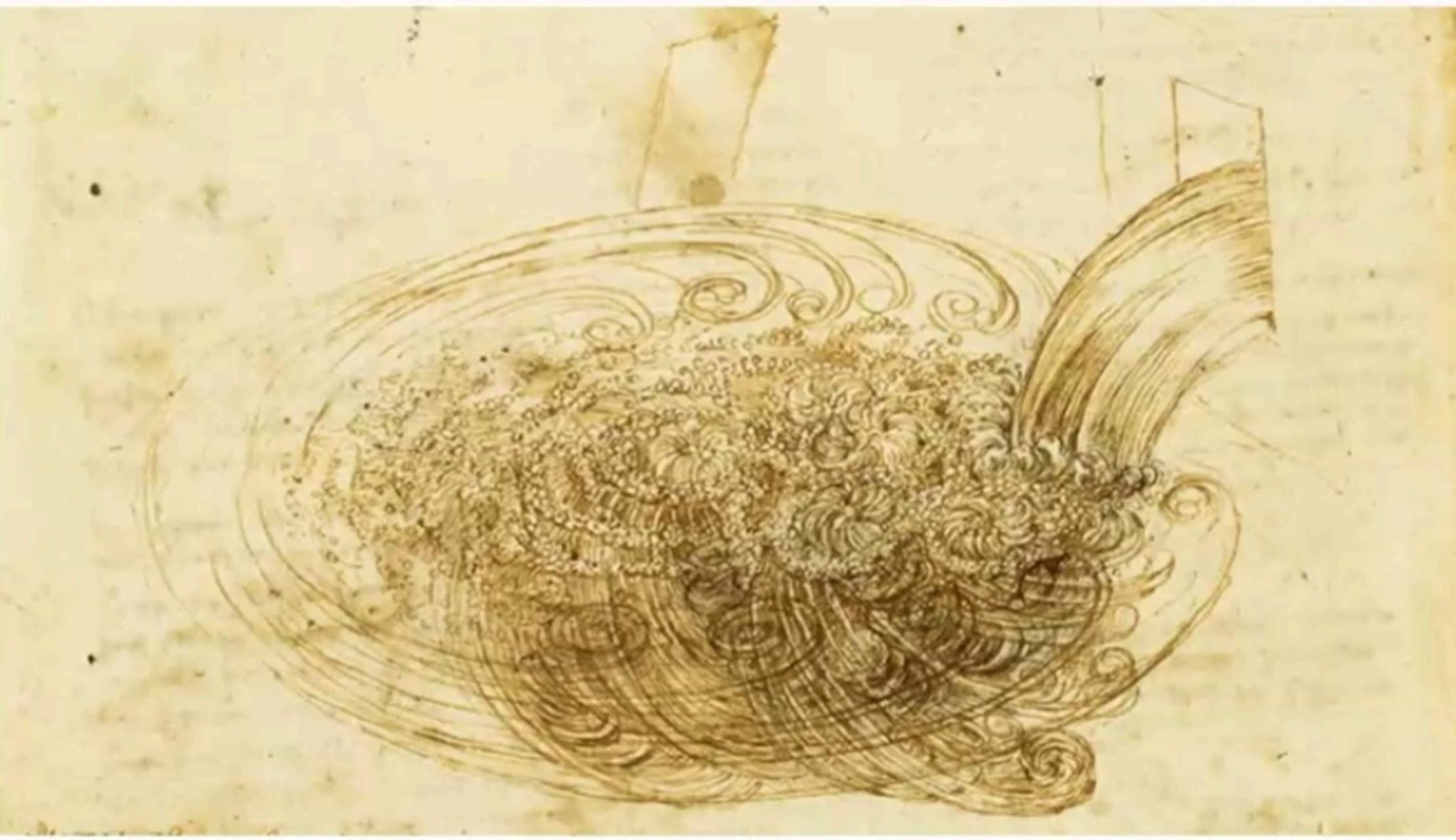
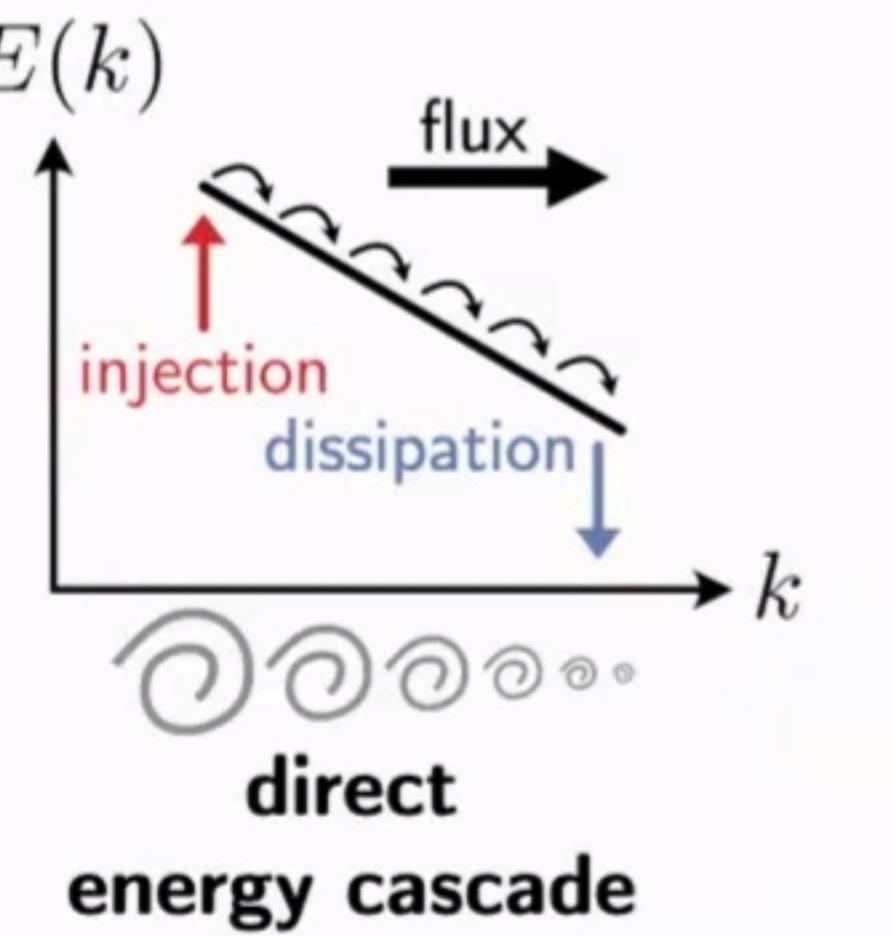
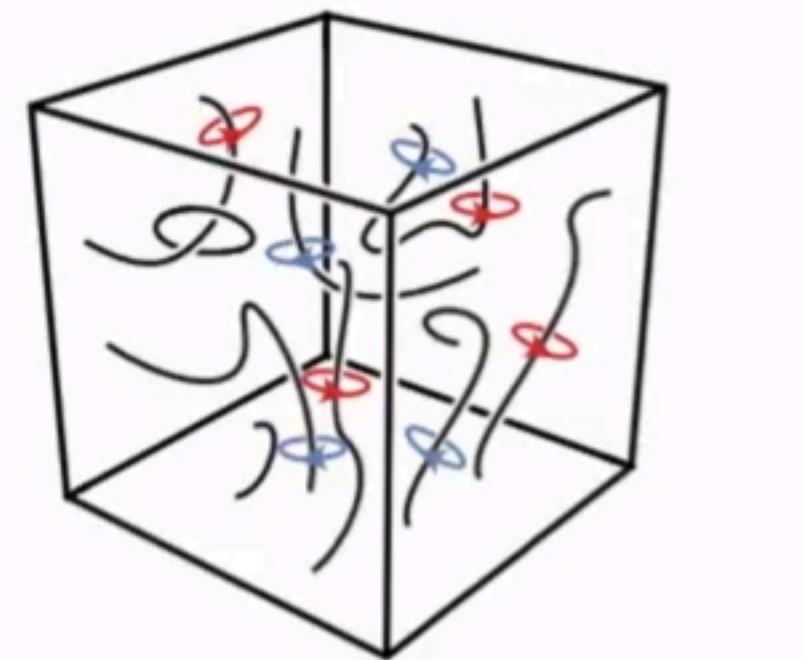
odd
viscoelasticity

inhomogeneous
proliferation

moduli can be time-dependent



What is turbulence?



L. da Vinci (1508)

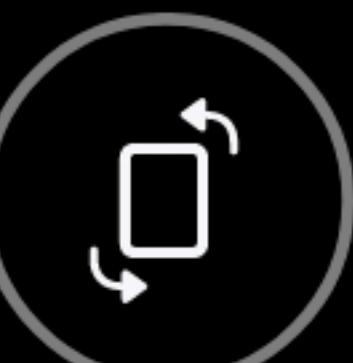
LMP Seminar

The seminar will be starting soon...

LMP Seminars



LMP Seminars's screen

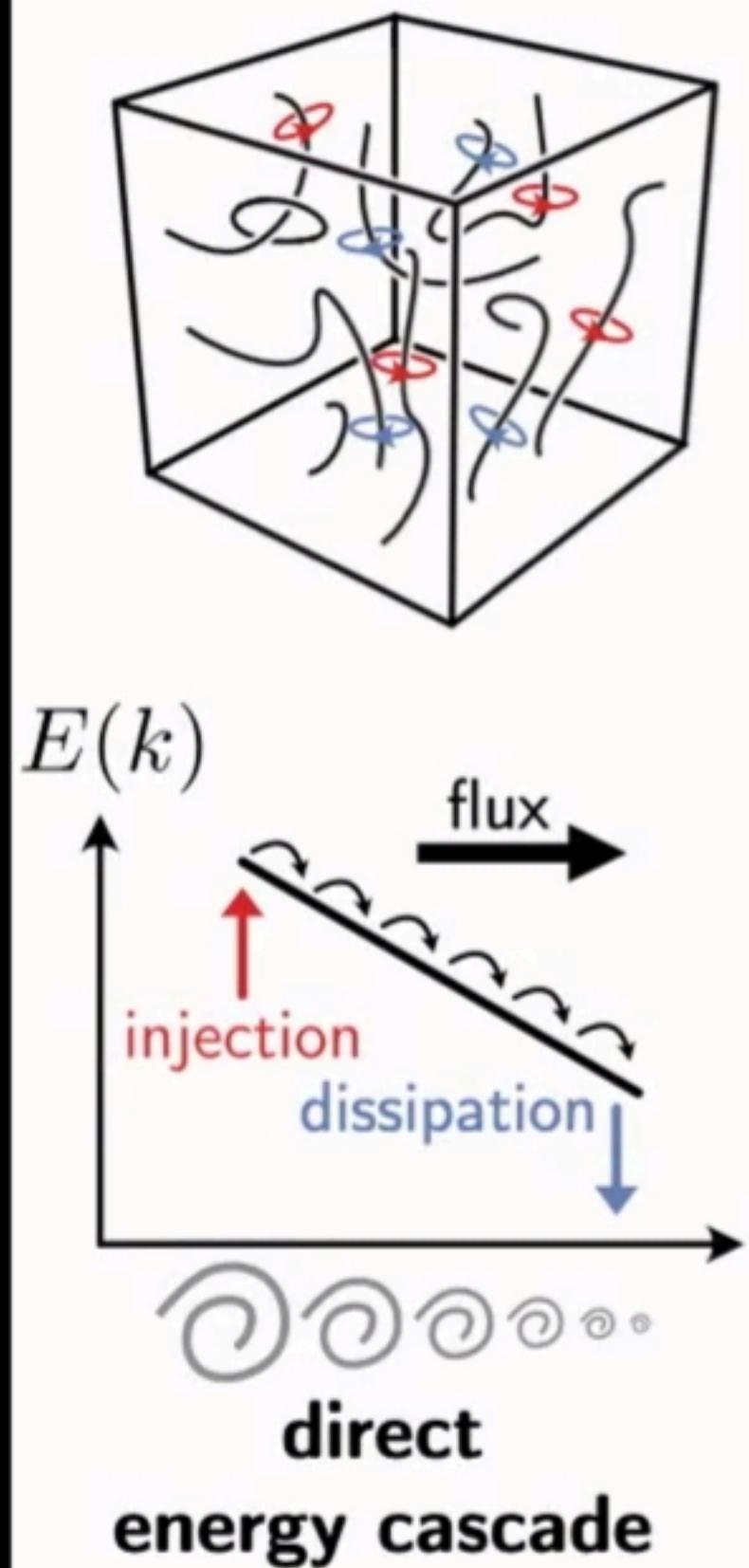


What is turbulence?

LMP Seminar

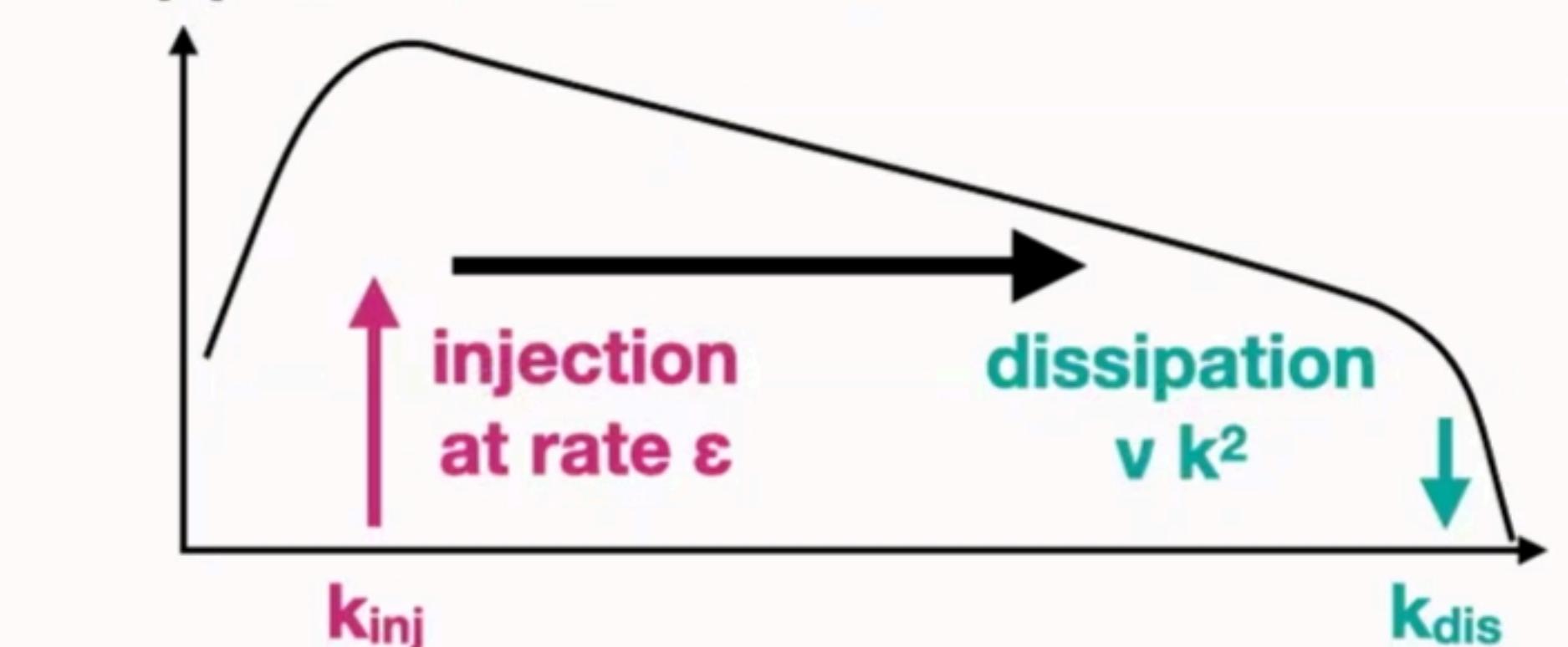
The seminar will be starting soon...

LMP Seminars



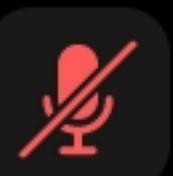
energy spectrum

$$E(k) \sim \varepsilon^{2/3} k^{-5/3}$$

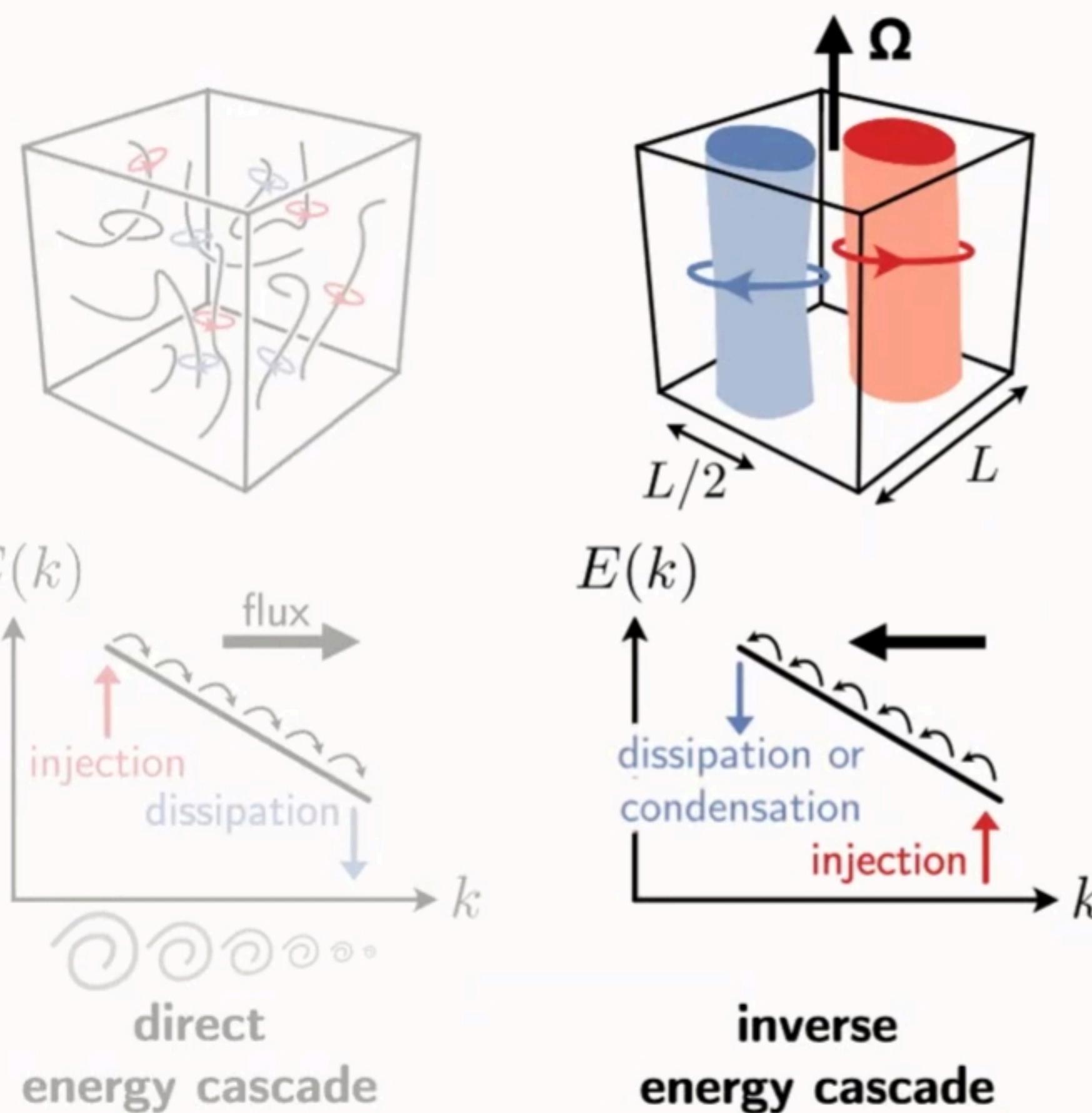


large vortices break down into smaller vortices

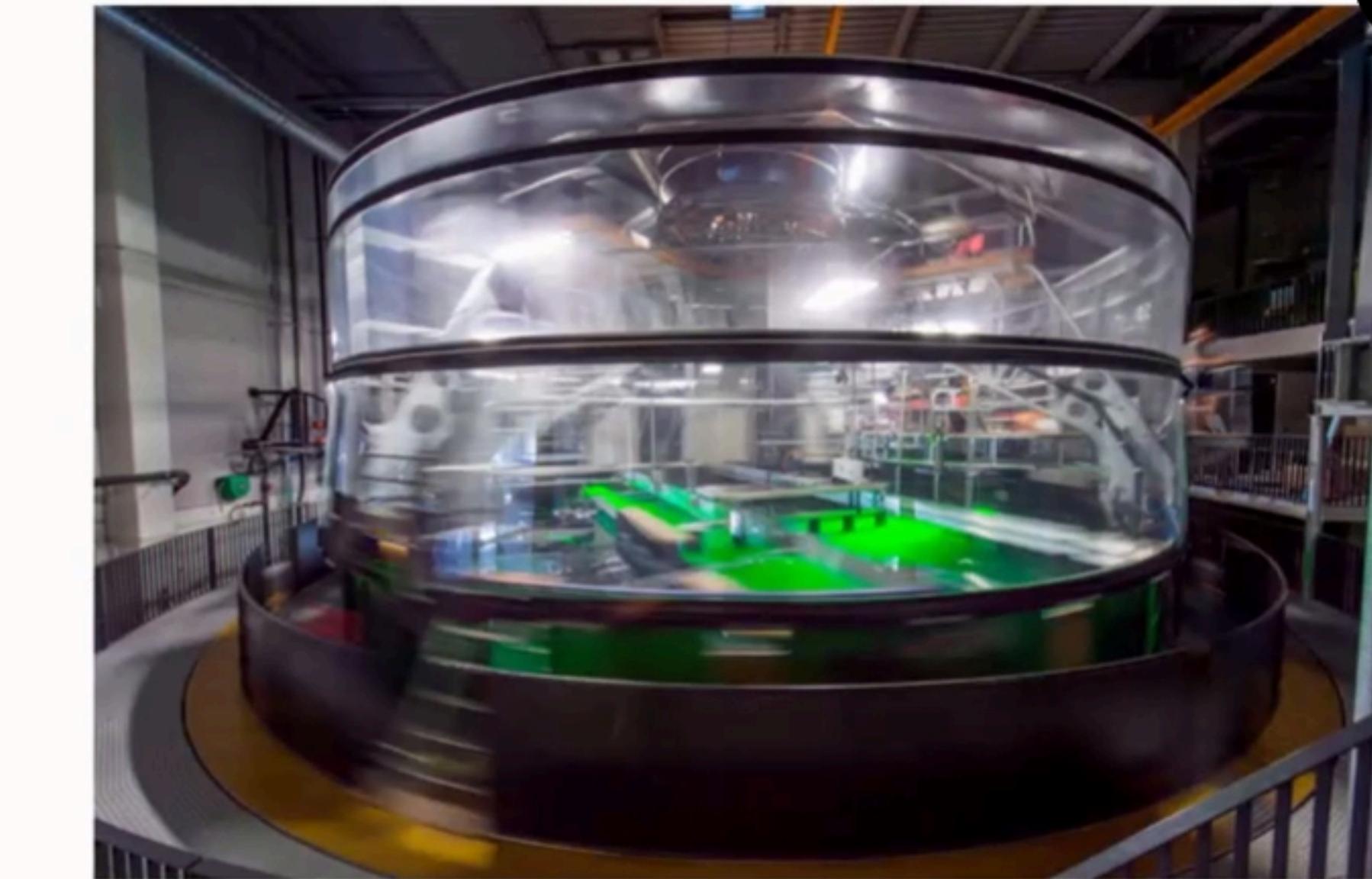
turbulent cascades are scale invariant, apparently featureless, chaotic states



What is turbulence?



rotation leads to an inverse cascade



Coriolis @ LEGI

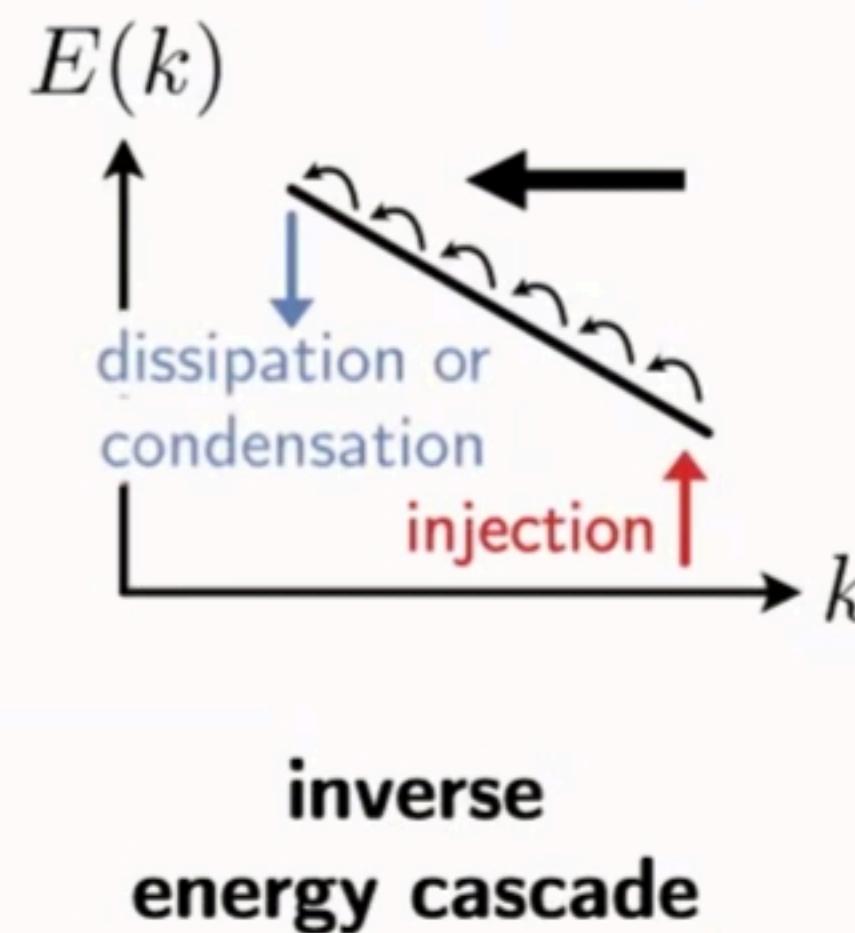
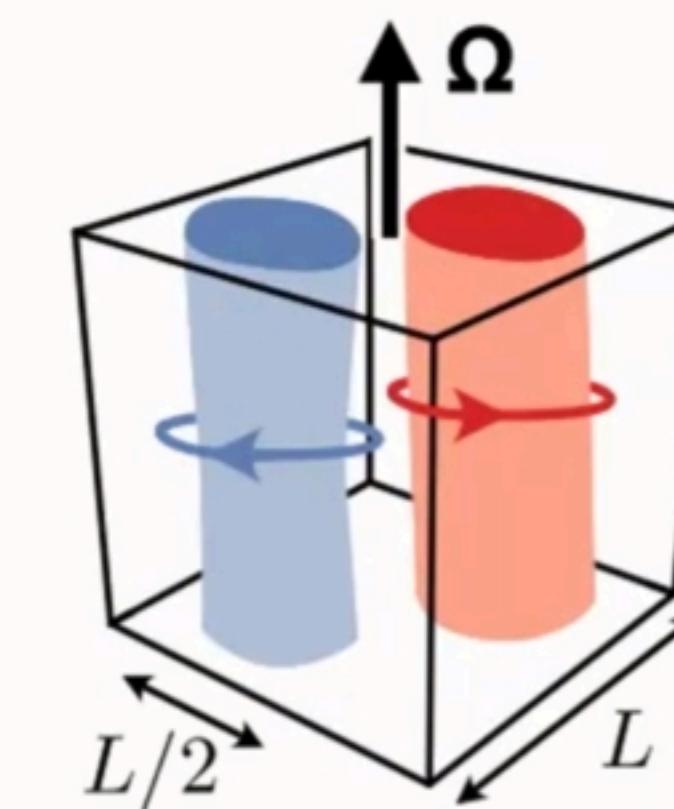
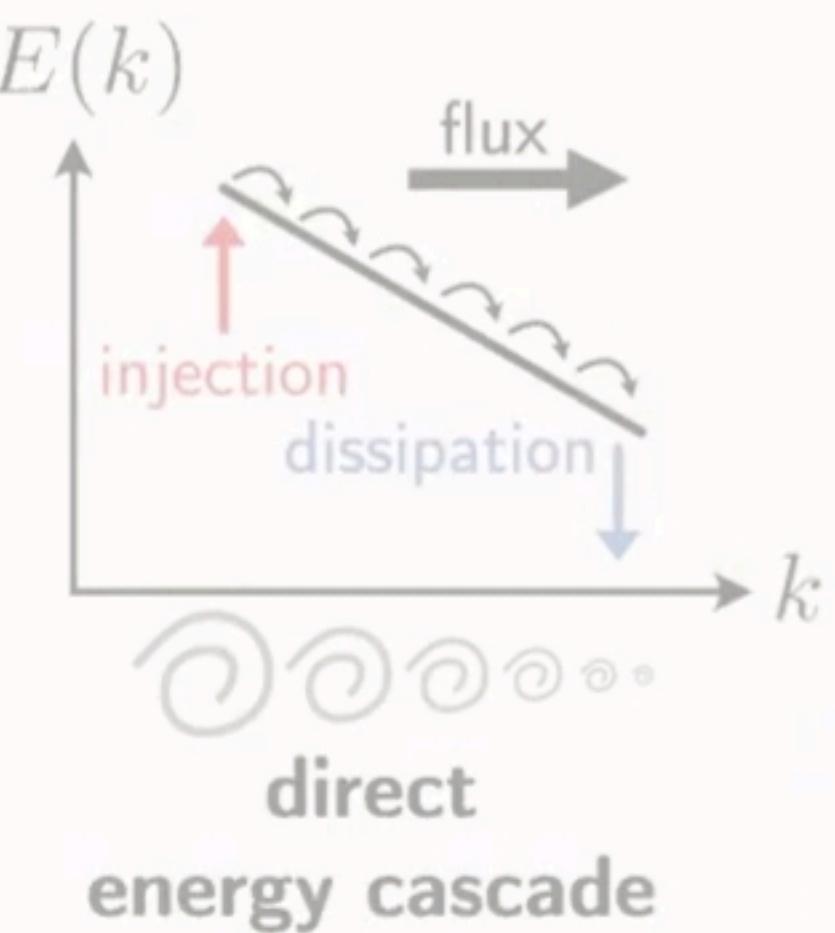
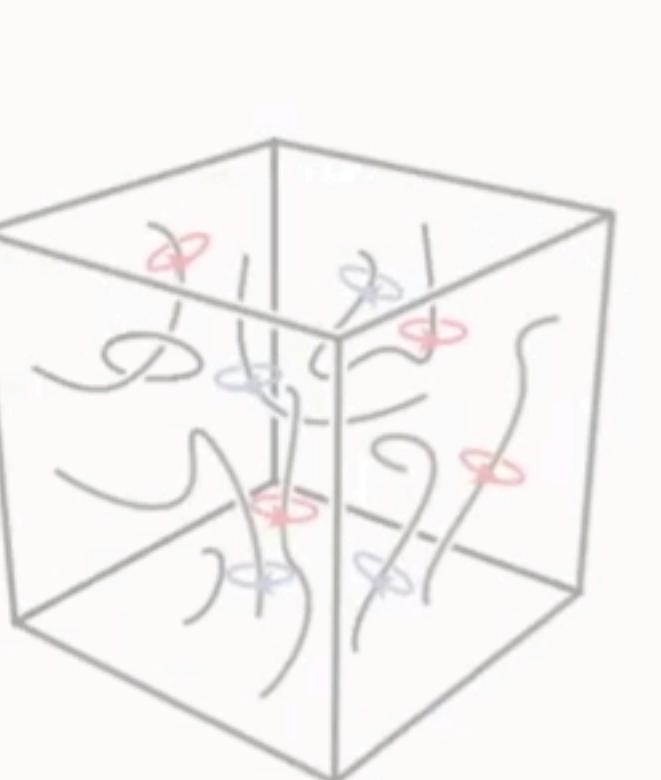
tank : 13 m diameter ; 130 tons of water



LMP Seminars



What is turbulence?

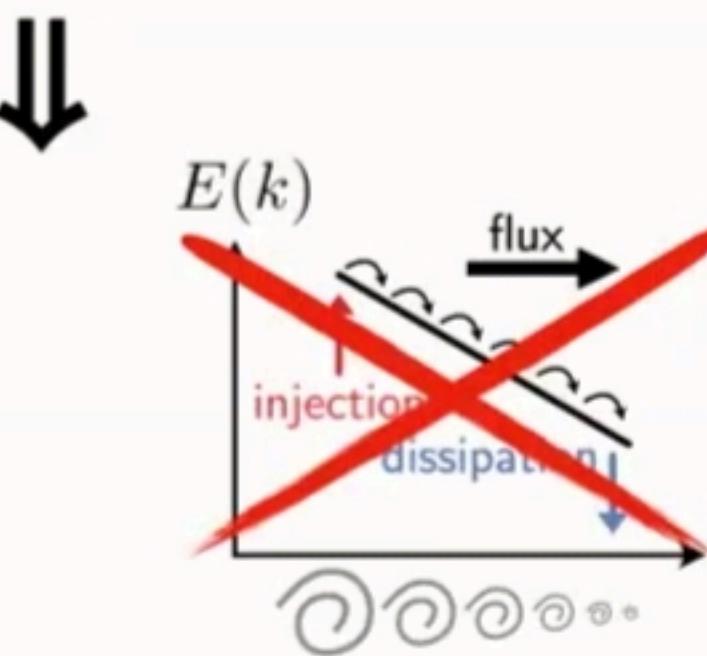


**flow quasi-2D
(Taylor-Proudman)**

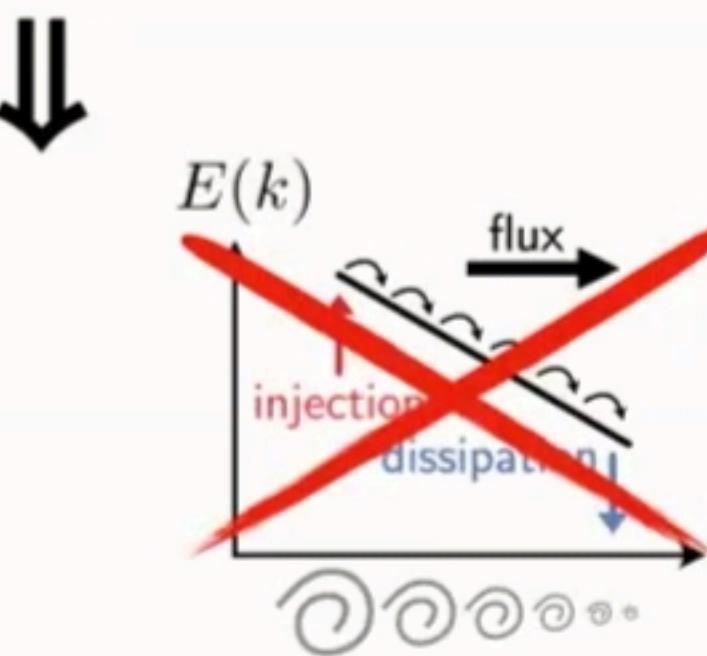
$$\partial_z v \simeq 0$$



**vortices
do not split**

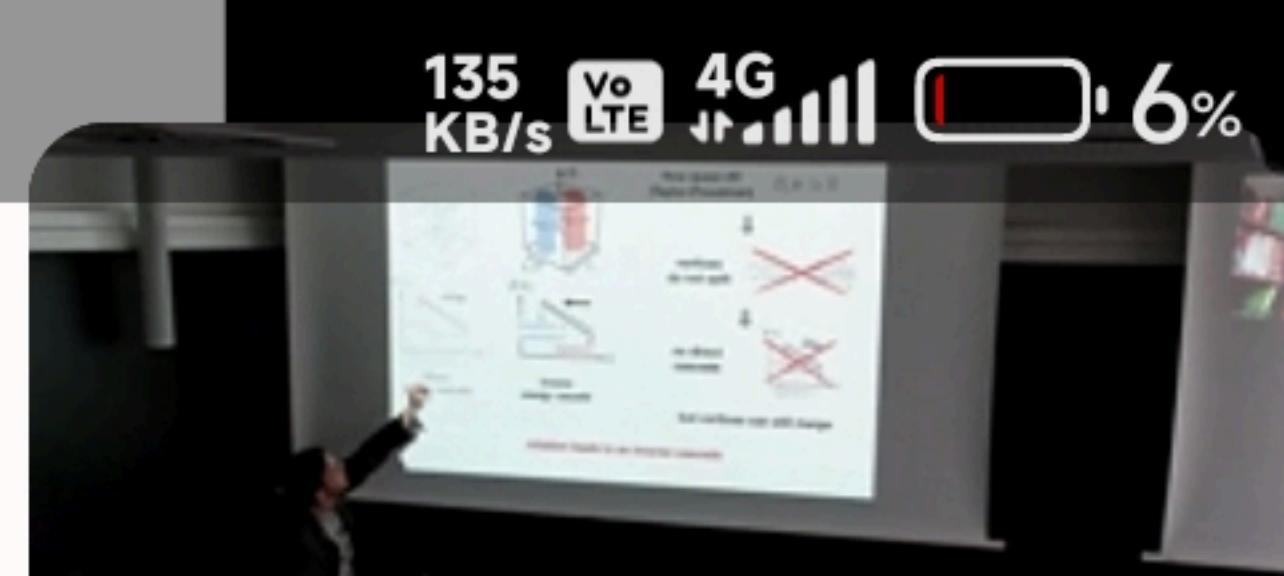


**no direct
cascade**



but vortices can still merge

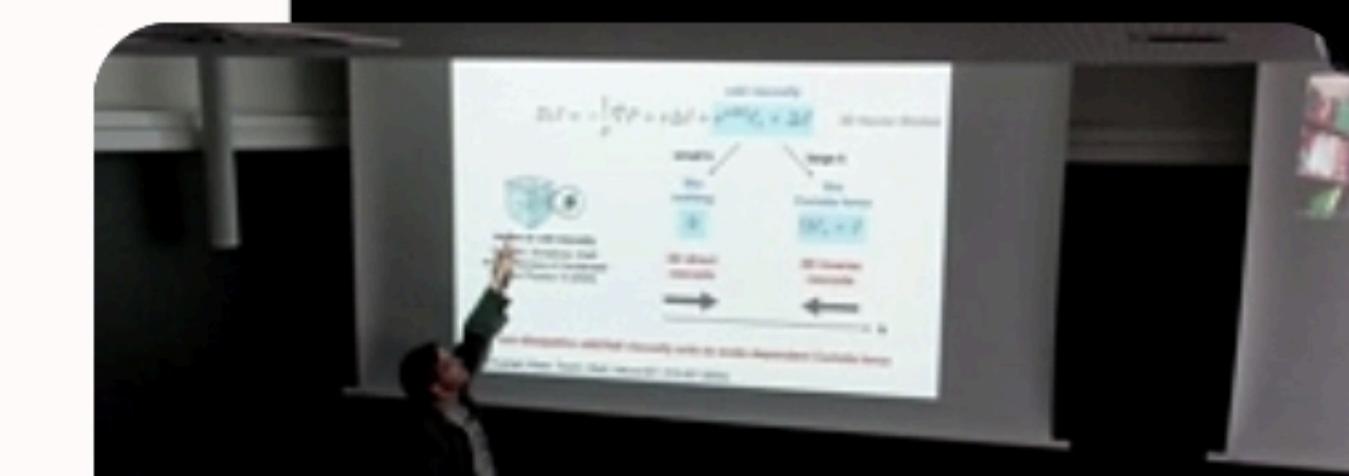
rotation leads to an inverse cascade



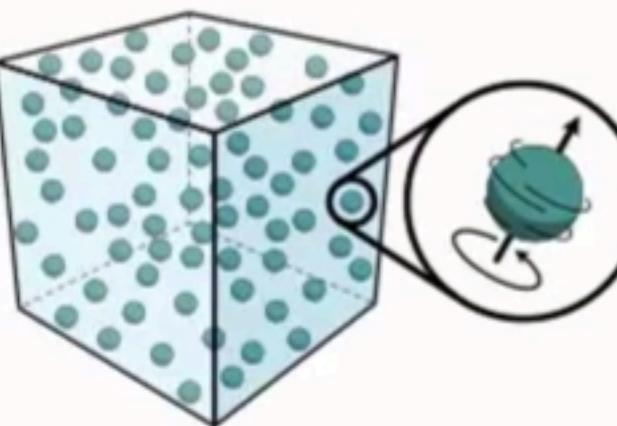
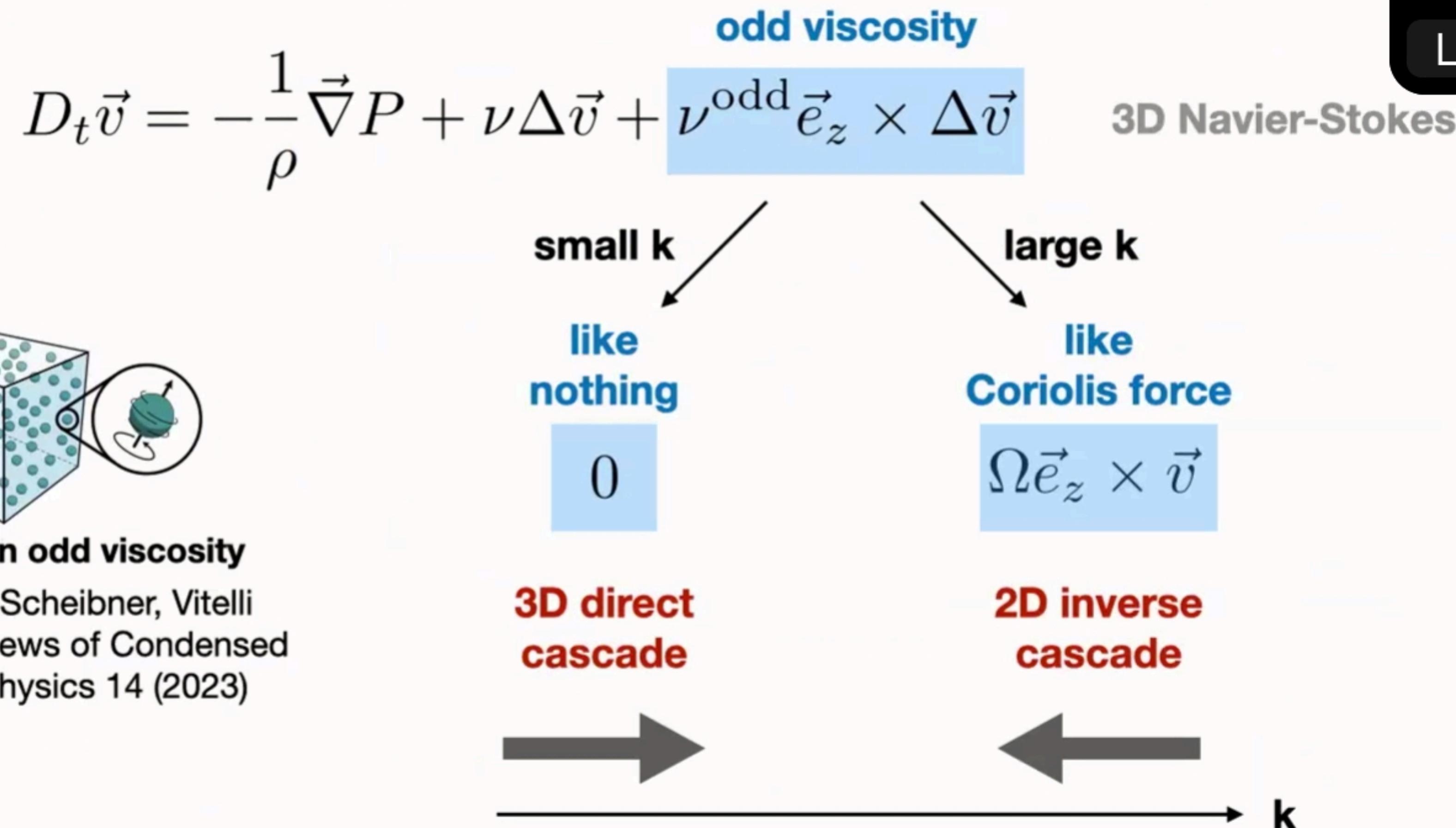
LMP Seminars



Odd viscosity



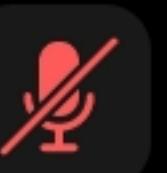
LMP Seminars



review on odd viscosity

Fruchart, Scheibner, Vitelli
Annual Reviews of Condensed
Matter Physics 14 (2023)

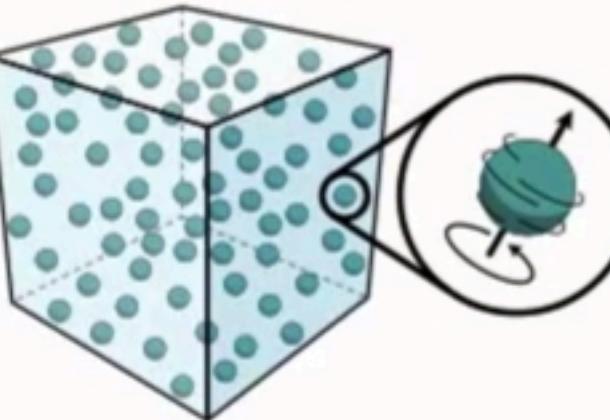
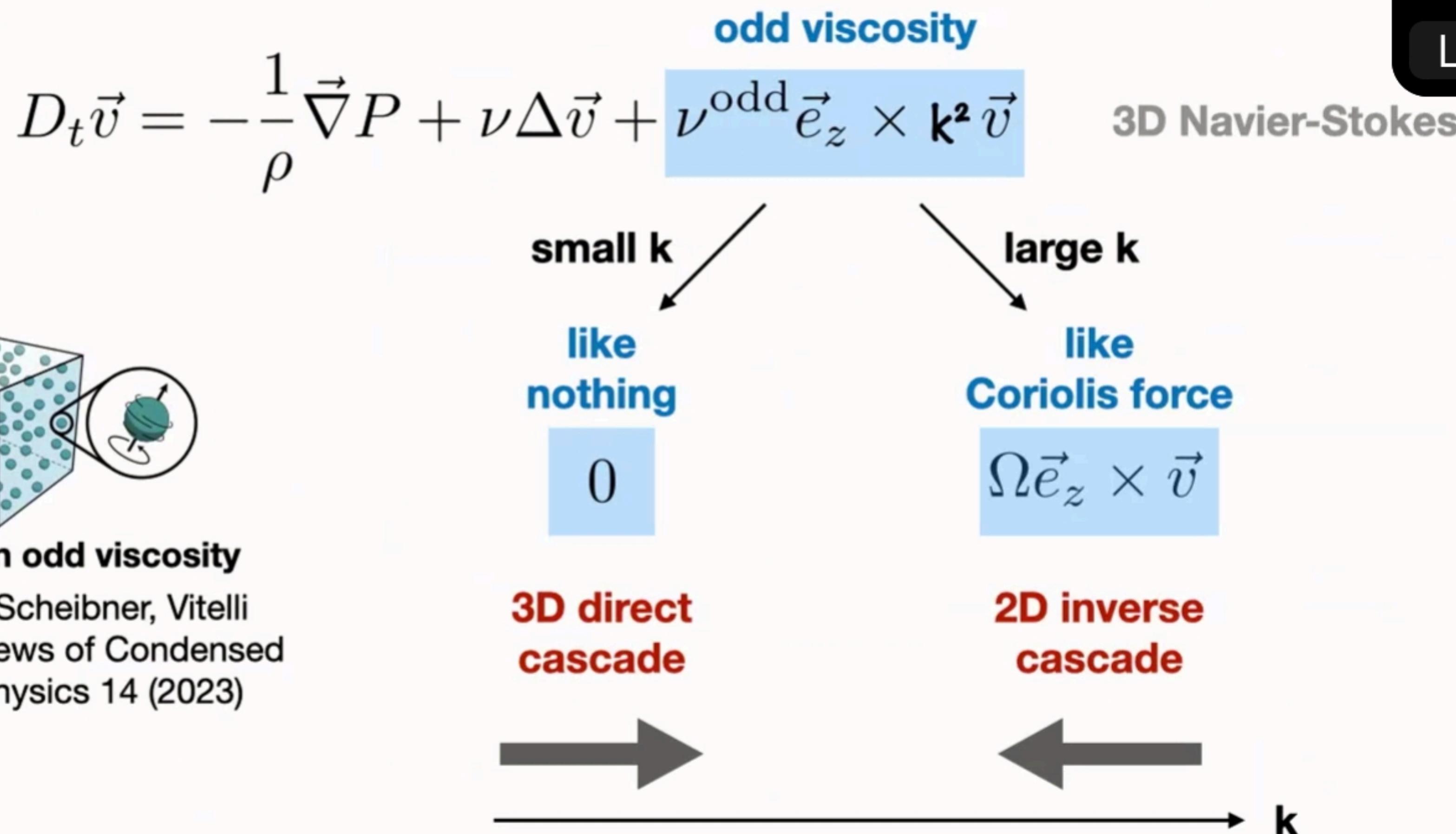
non-dissipative odd/Hall viscosity acts as scale-dependent Coriolis force



Odd viscosity



LMP Seminars



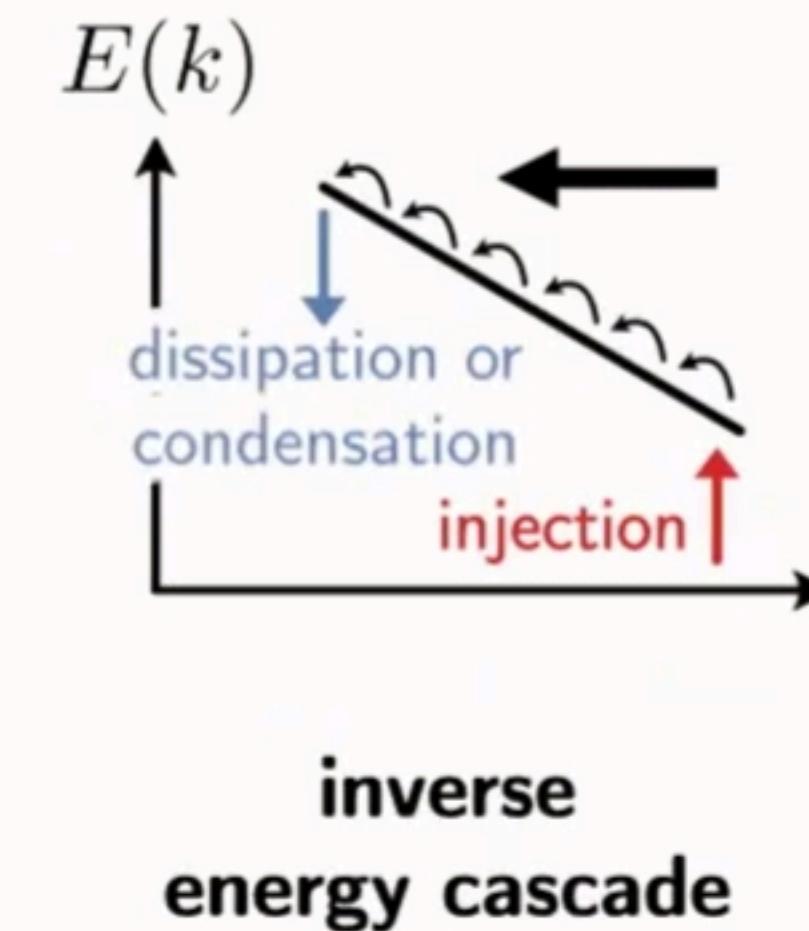
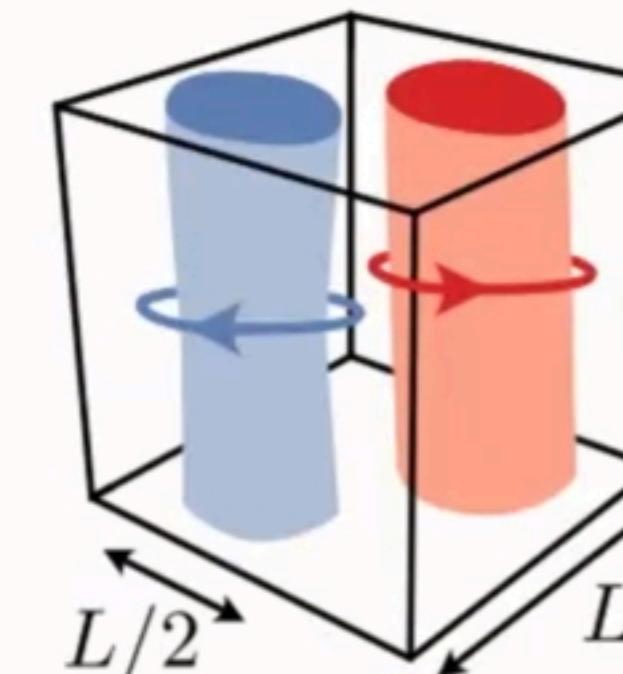
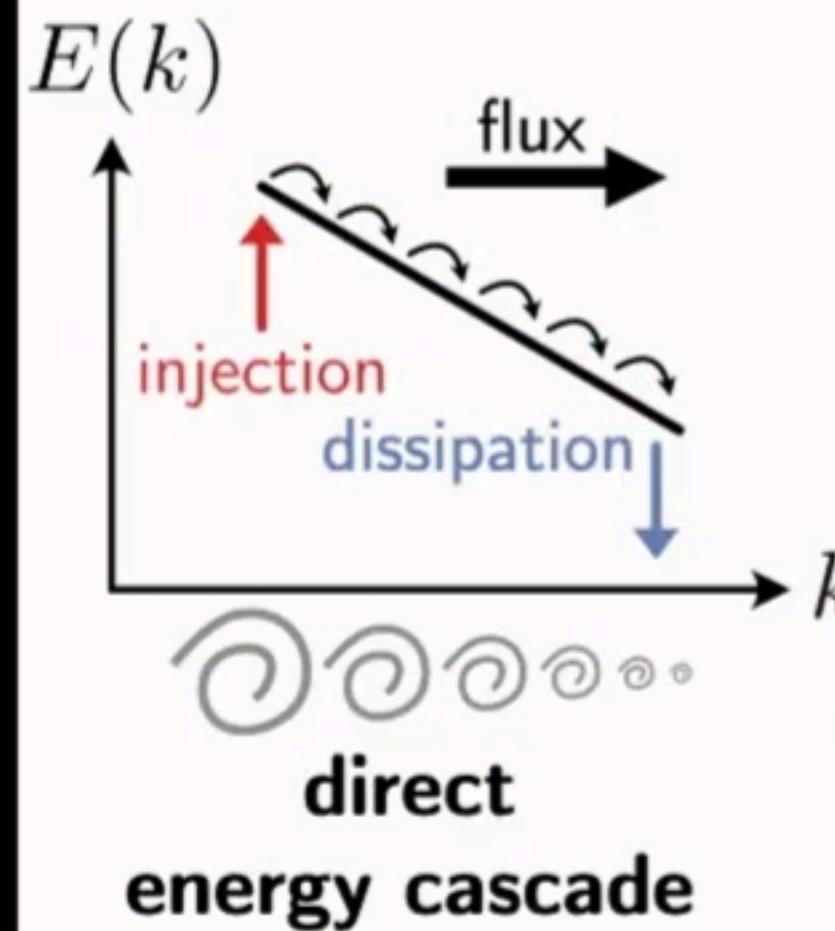
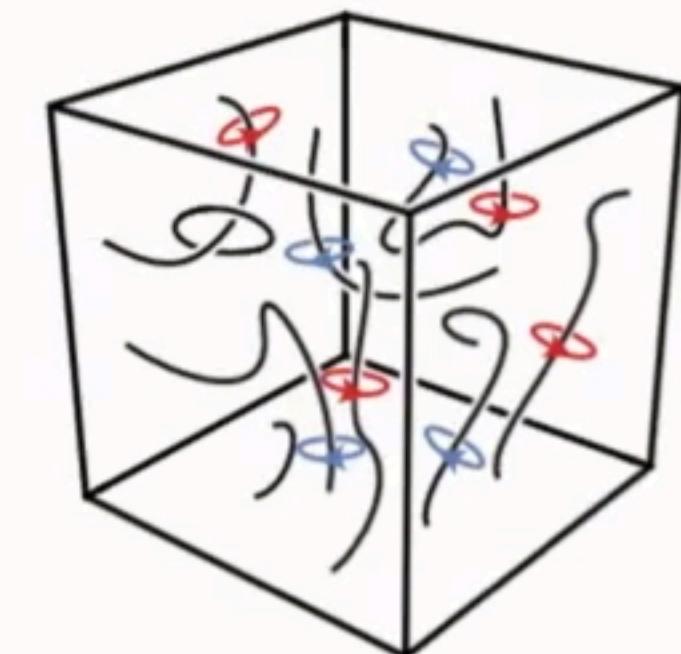
review on odd viscosity

Fruchart, Scheibner, Vitelli
Annual Reviews of Condensed
Matter Physics 14 (2023)

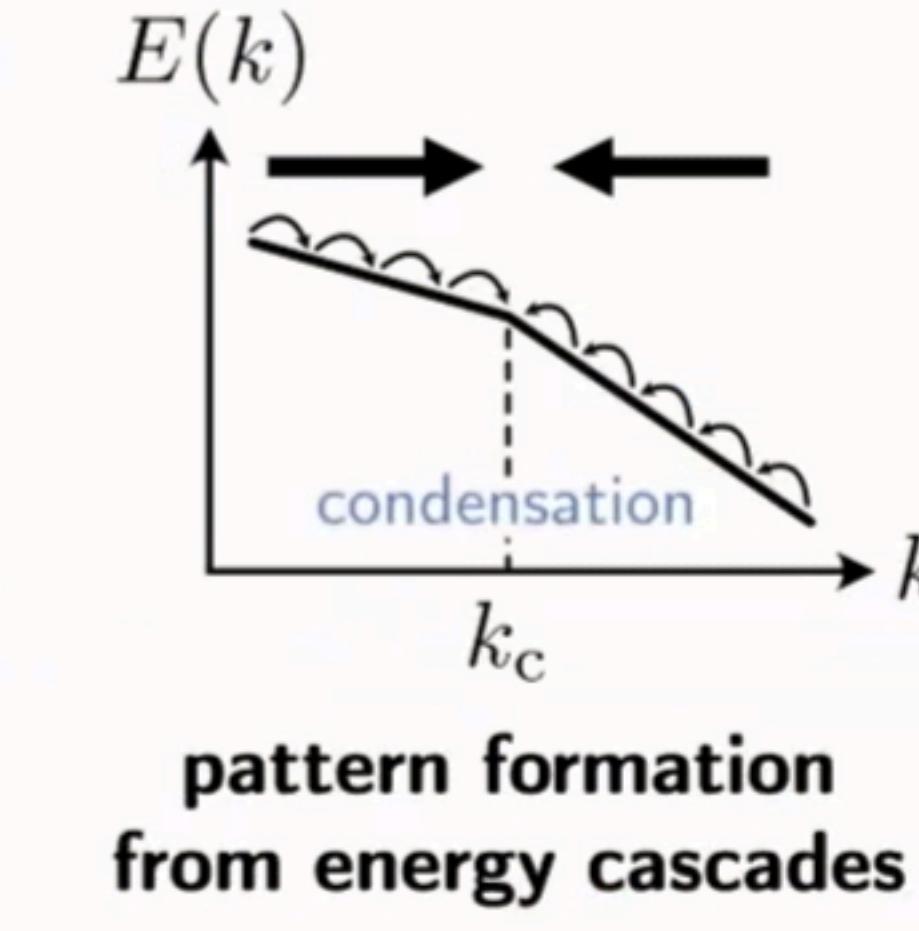
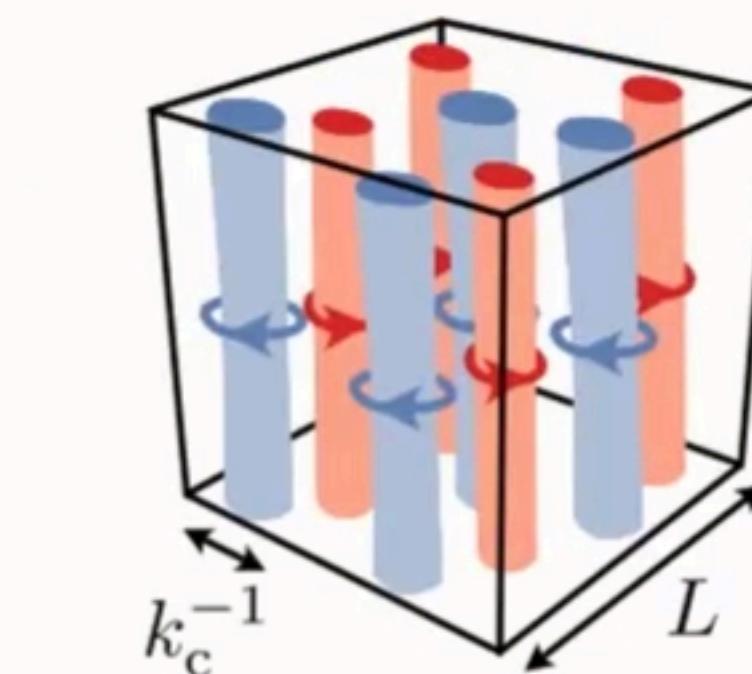
non-dissipative odd/Hall viscosity acts as scale-dependent Coriolis force



Pattern formation by turbulent cascades



condensation



Xander De



LMP Seminars



Tali Khain



Federico Toschi



Vincenzo Vitelli

large vortices split while small vortices merge \Rightarrow wavelength selection



Odd viscosity

fluid with odd viscosity

$$D_t \vec{v} = -\frac{1}{\rho} \vec{\nabla} P + \begin{pmatrix} \nu & \nu^{\text{odd}} & 0 \\ -\nu^{\text{odd}} & \nu & 0 \\ 0 & 0 & \nu \end{pmatrix} \Delta \vec{v}$$

LMP Seminars



odd waves

$$\omega(\mathbf{k}) = \pm \nu_{\text{odd}} k_z |\mathbf{k}|$$



odd viscosity induces waves in the fluid



Odd viscosity

fluid with odd viscosity

$$D_t \vec{v} = -\frac{1}{\rho} \vec{\nabla} P + \begin{pmatrix} \nu & \nu^{\text{odd}} & 0 \\ -\nu^{\text{odd}} & \nu & 0 \\ 0 & 0 & \nu \end{pmatrix} \Delta \vec{v}$$

LMP Seminars

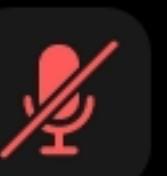


odd waves

$$\omega(\mathbf{k}) = \pm \nu_{\text{odd}} k_z |\mathbf{k}|$$

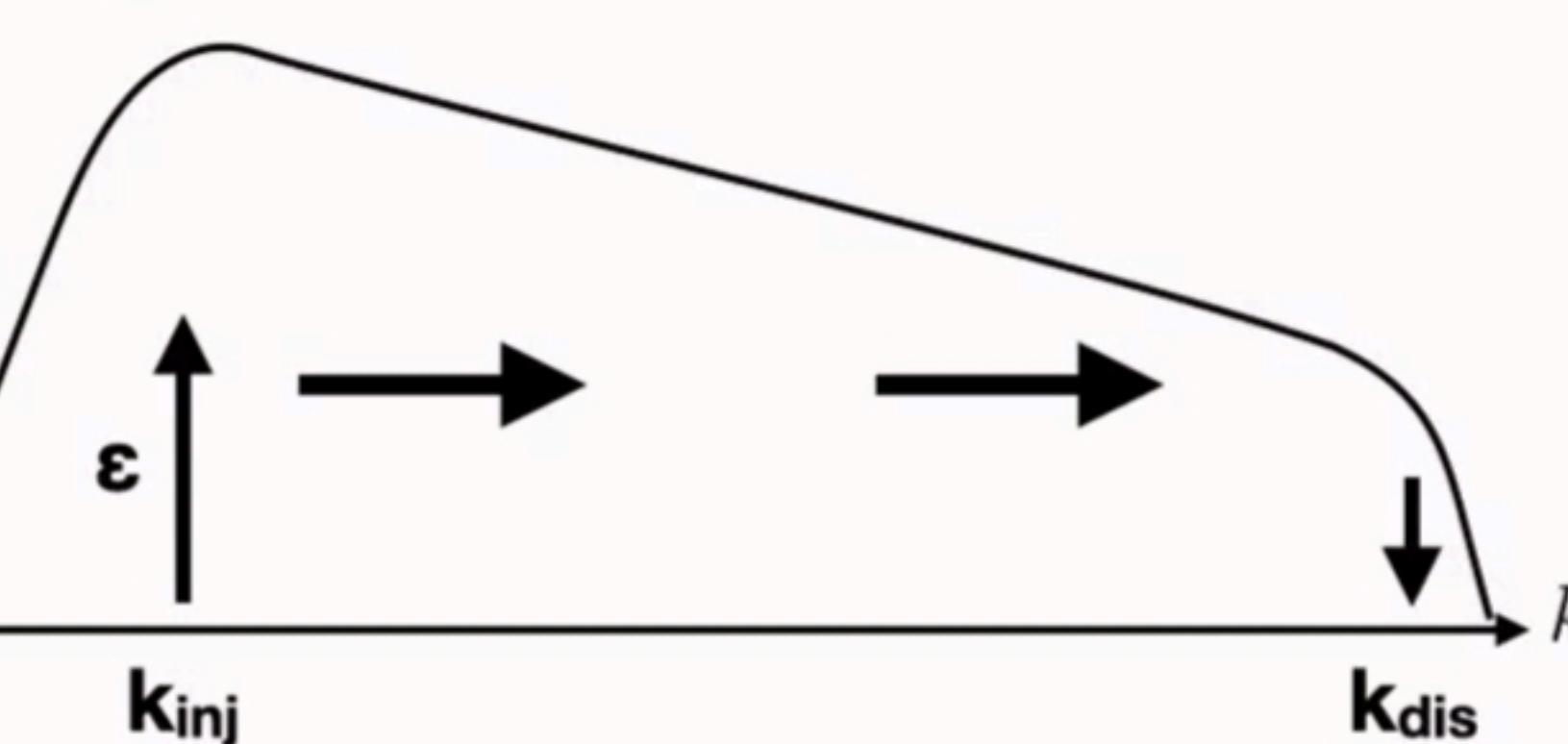


odd viscosity induces waves in the fluid



Waves vs Vortices

$$E(k) \sim \varepsilon^{2/3} k^{-5/3}$$

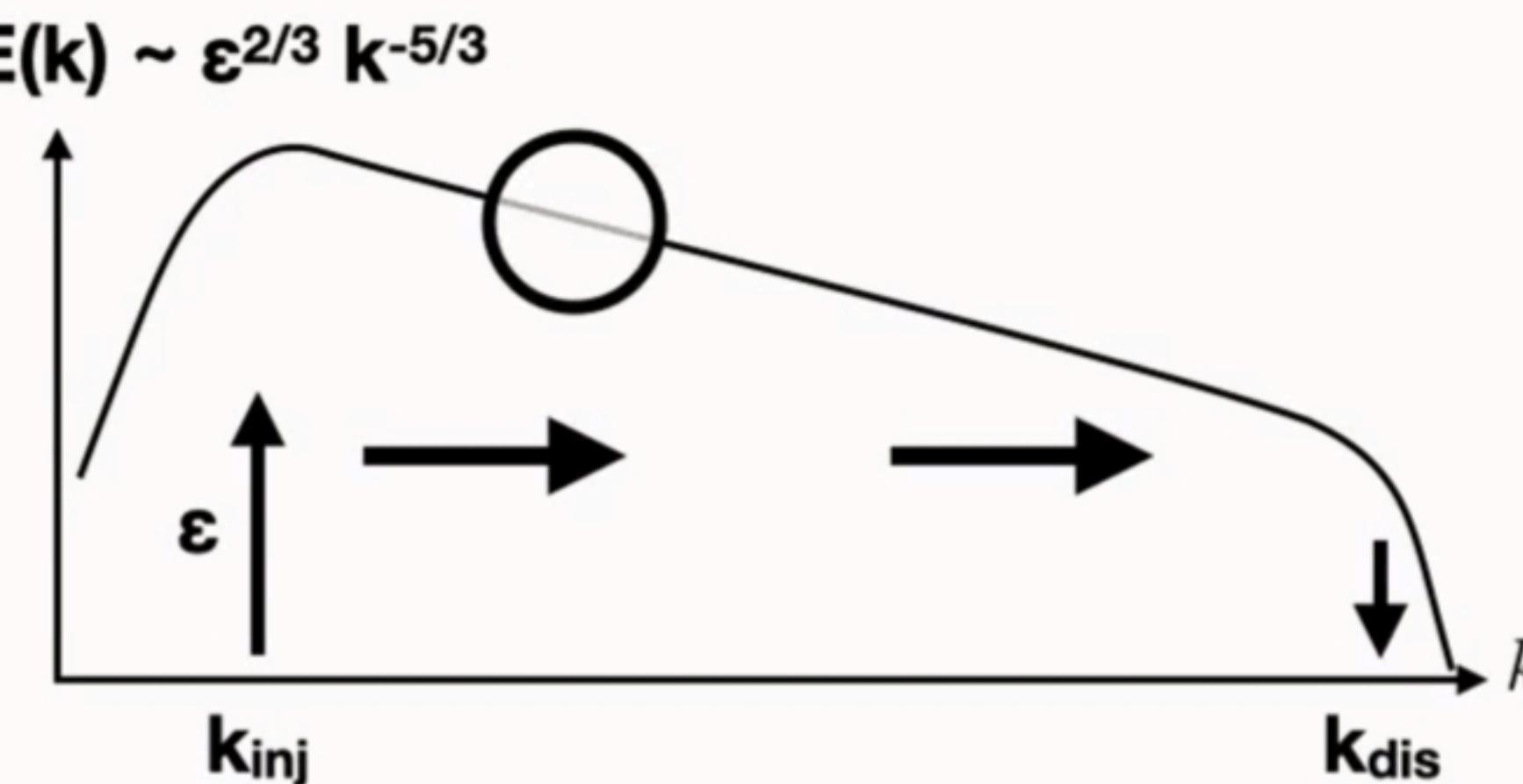


$$k_{\text{dis}} \sim (v^3/\varepsilon)^{1/4}$$

LMP Seminars



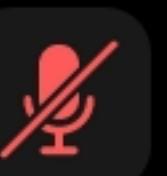
Waves vs Vortices



🌀 → ⚡ characteristic time
(eddy turnover time) $\tau_{\text{nl}}(k)$

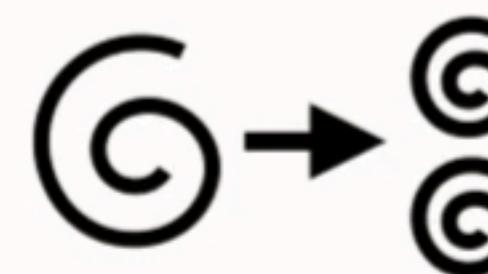
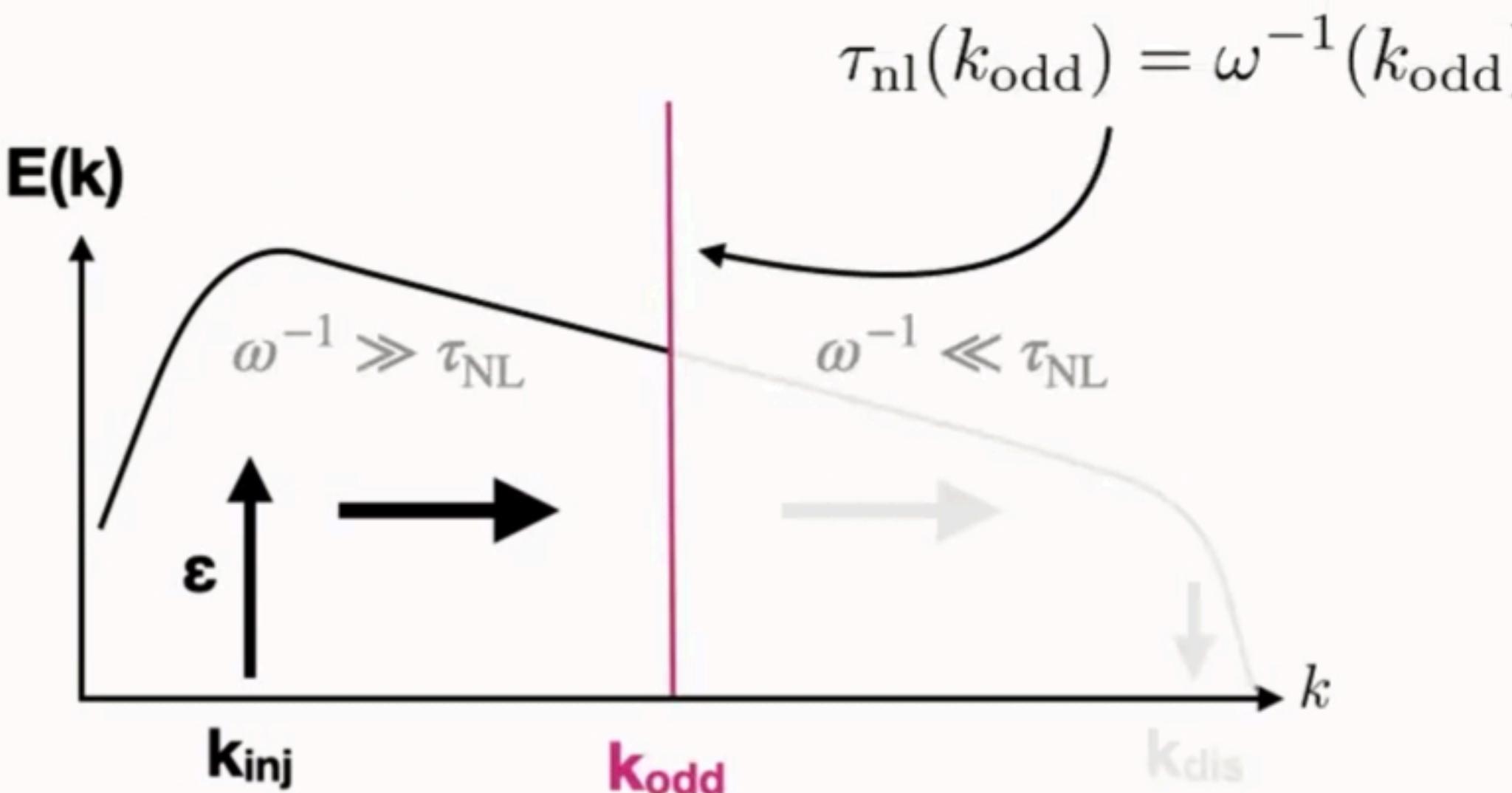
$$k_{\text{dis}} \sim (v^3/\varepsilon)^{1/4}$$



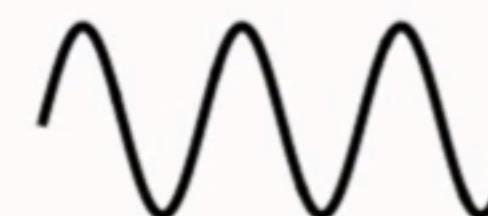


Waves vs Vortices

$$k_{\text{odd}} \equiv \epsilon^{\frac{1}{4}} \nu_{\text{odd}}^{-\frac{3}{4}}$$

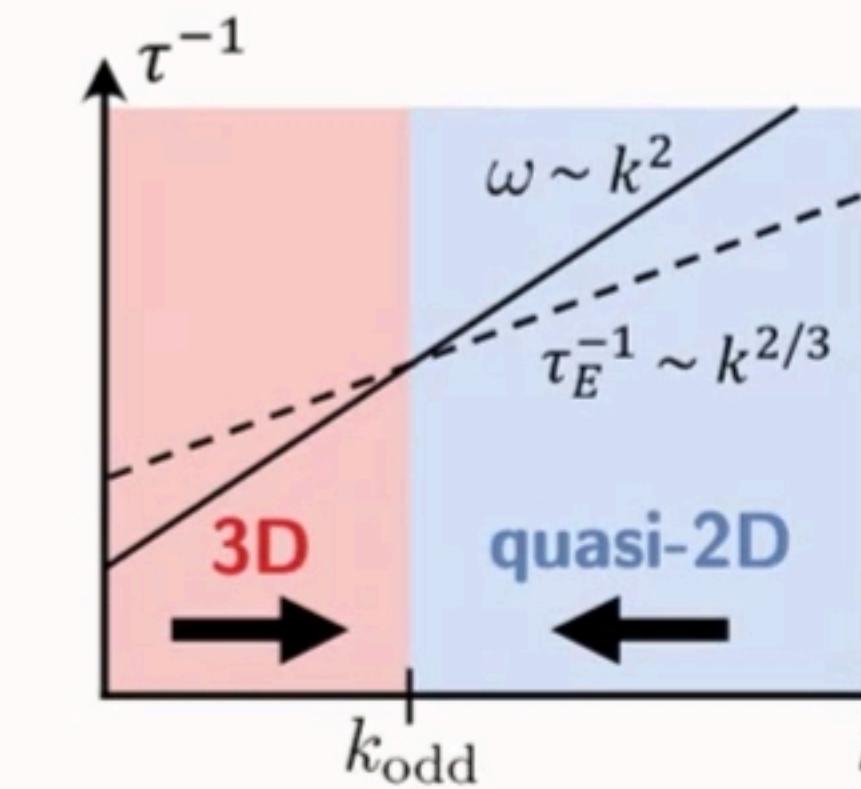


**characteristic time
(eddy turnover time)**



**dispersion of
odd waves**

$$\tau_{\text{nl}}(k) \quad \omega(k)$$



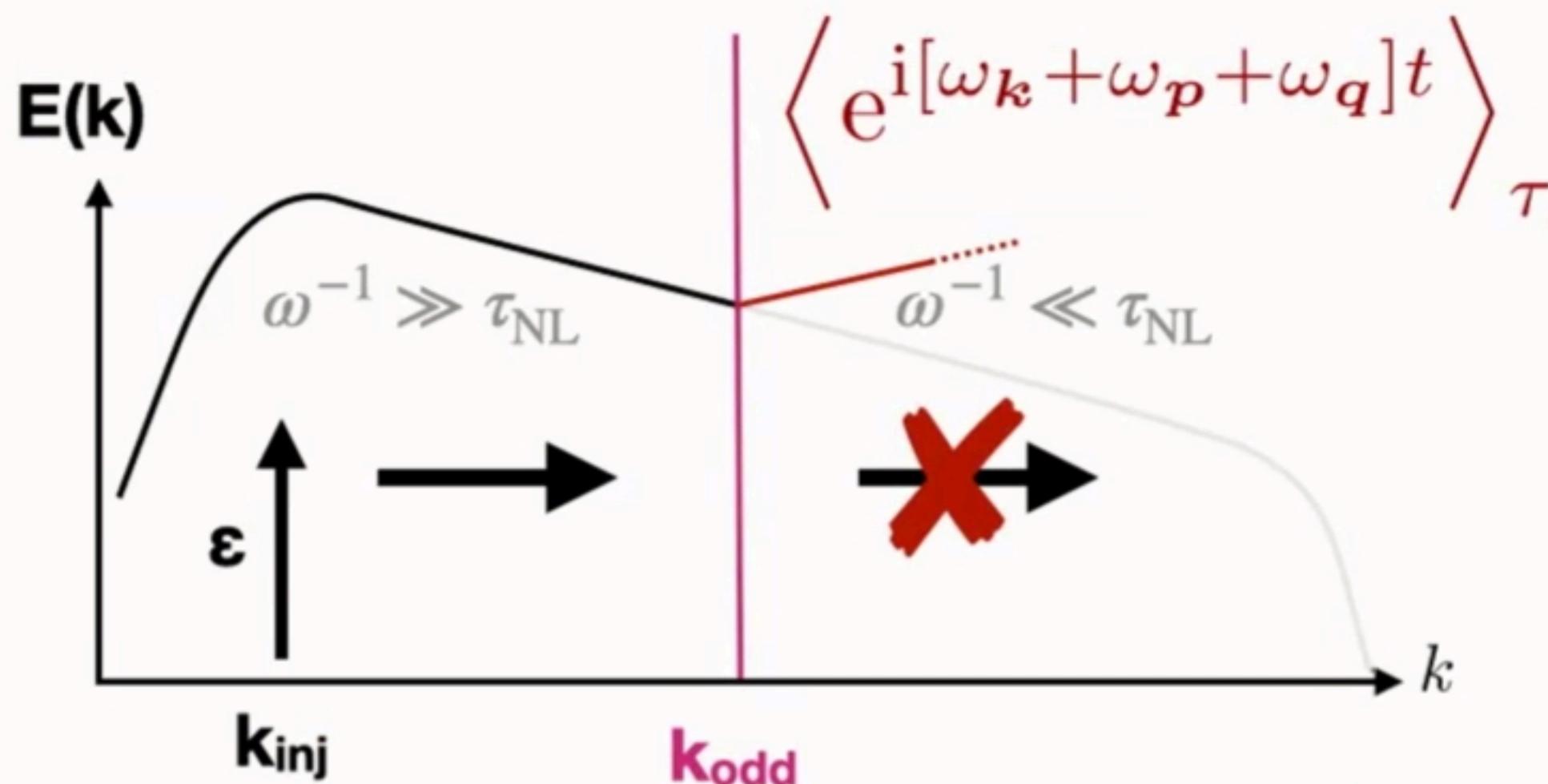
odd waves affect eddy turbulence

LMP Seminars



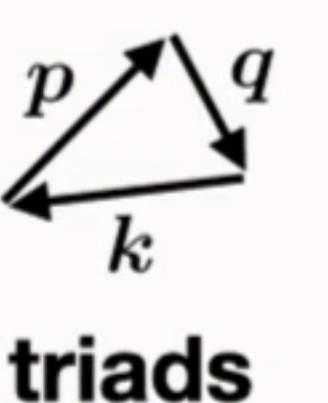
Waves vs Vortices

$$k_{\text{odd}} \equiv \epsilon^{\frac{1}{4}} \nu_{\text{odd}}^{-\frac{3}{4}}$$

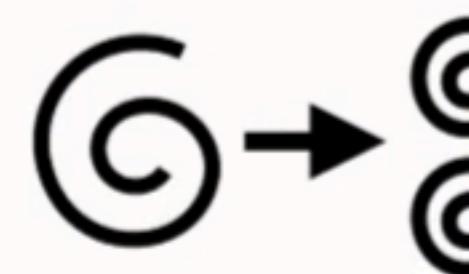


cascade is arrested
so energy piles up

LMP Seminars

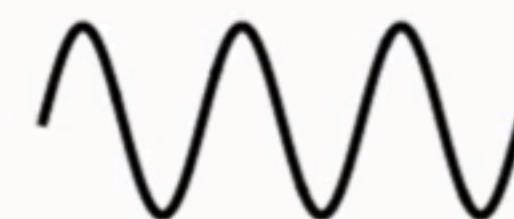


triads



characteristic time
(eddy turnover time)

$$\tau_{\text{nl}}(k)$$



dispersion of
odd waves

$$\omega(k)$$

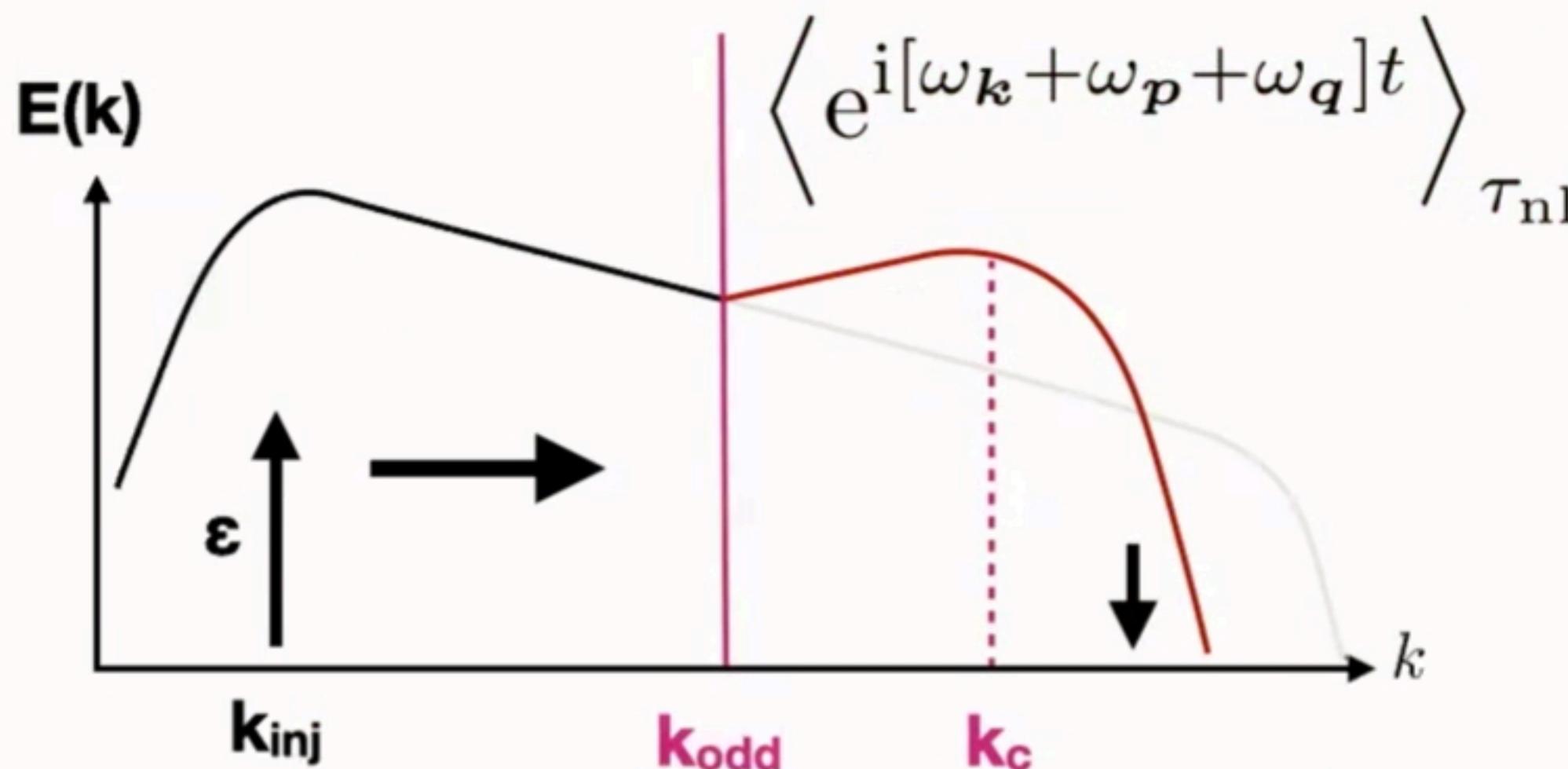
odd waves affect eddy turbulence



Waves vs Vortices

$$k_{\text{odd}} \equiv \epsilon^{\frac{1}{4}} \nu_{\text{odd}}^{-\frac{3}{4}}$$

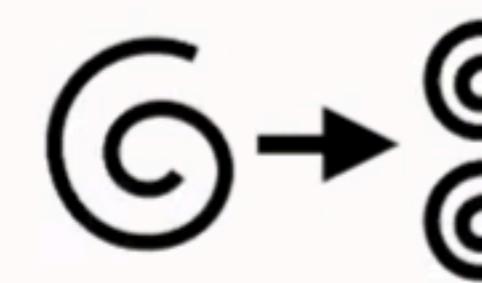
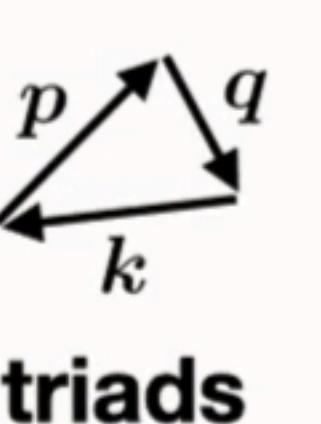
$$k_c \propto \epsilon^{\frac{1}{4}} \nu^{-\frac{1}{2}} \nu_{\text{odd}}^{-\frac{1}{4}}$$



cascade is arrested
so energy piles up

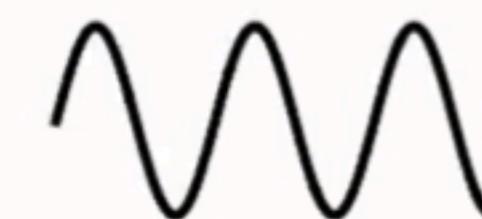
LMP Seminars

... until viscous
dissipation $\sim \nu k^2$
takes over



characteristic time
(eddy turnover time)

$$\tau_{\text{nl}}(k)$$

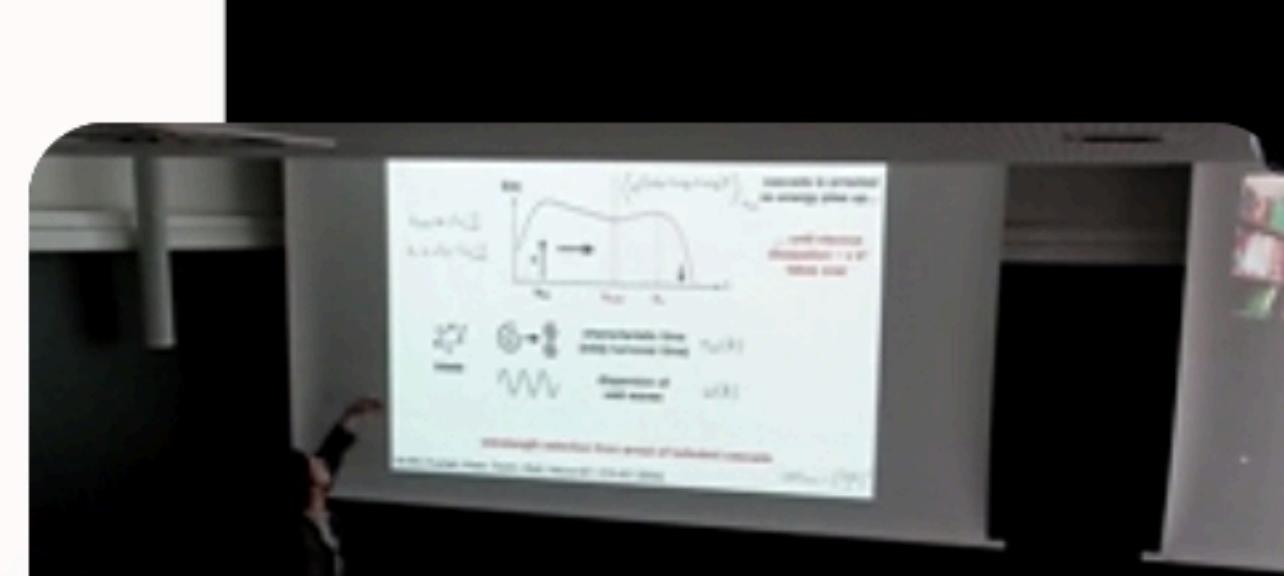


dispersion of
odd waves

$$\omega(k)$$

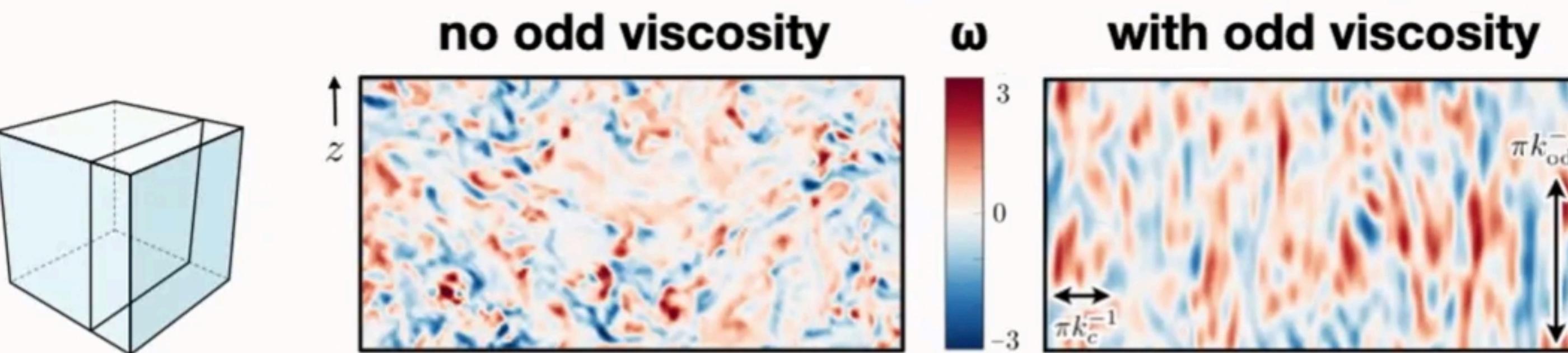
wavelength selection from arrest of turbulent cascade

$$(\Delta E)_{\text{max}} \sim \left(\frac{\nu_{\text{odd}}}{\nu} \right)^{\frac{1}{3}}$$

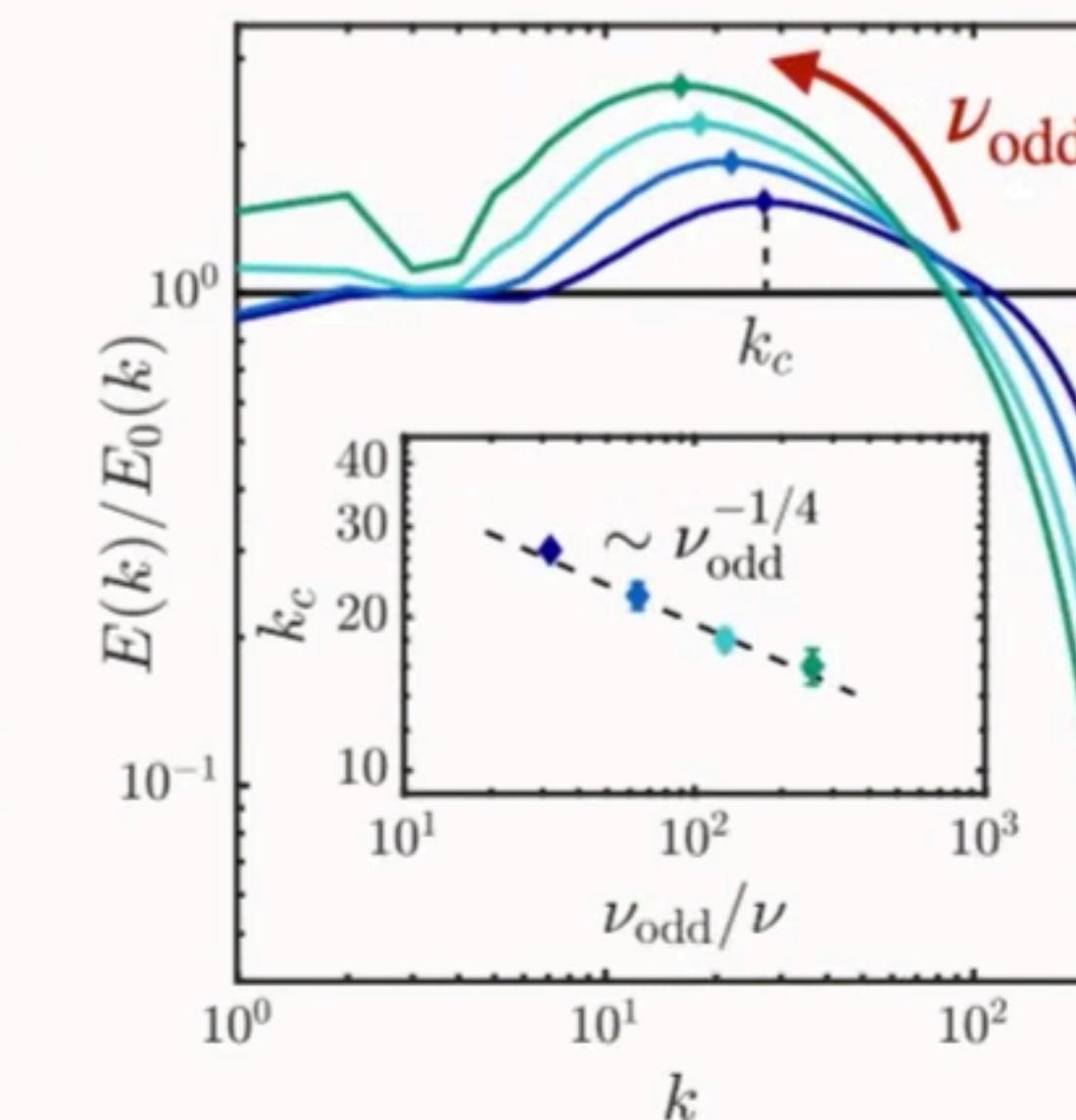




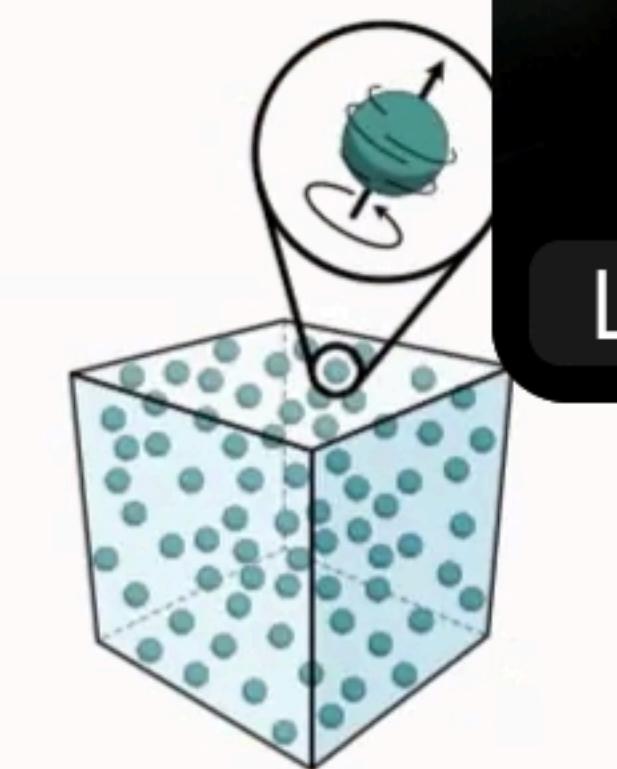
Direct numerical simulations



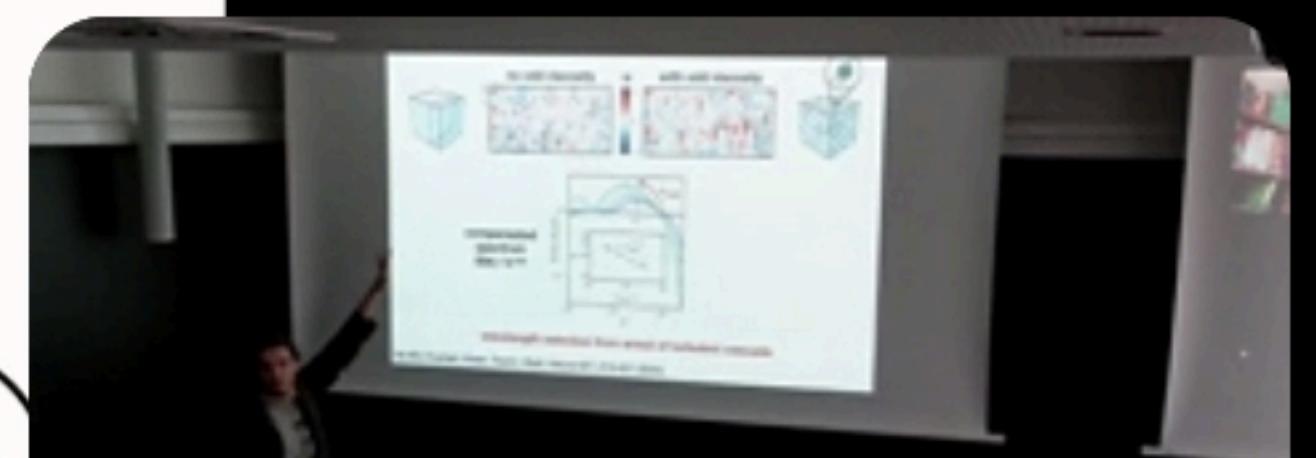
**compensated
spectrum
 $E(k) / k^{-5/3}$**

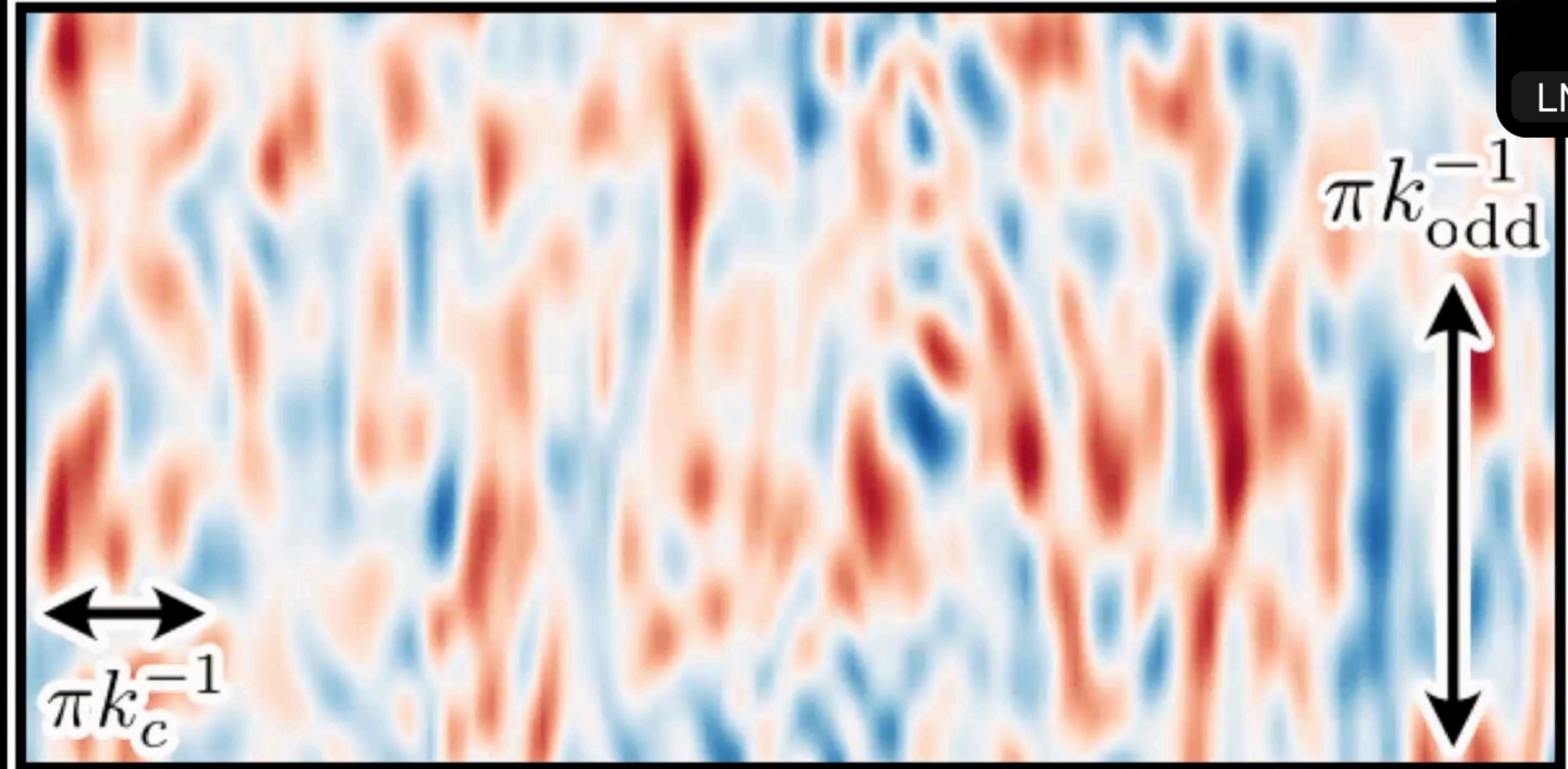


wavelength selection from arrest of turbulent cascade

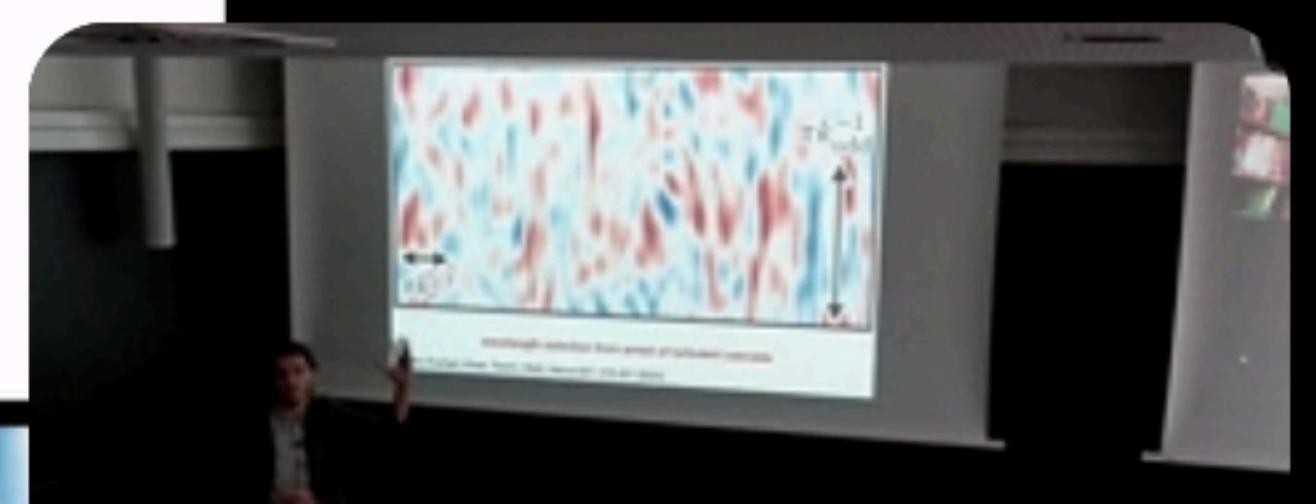


LMP Seminars



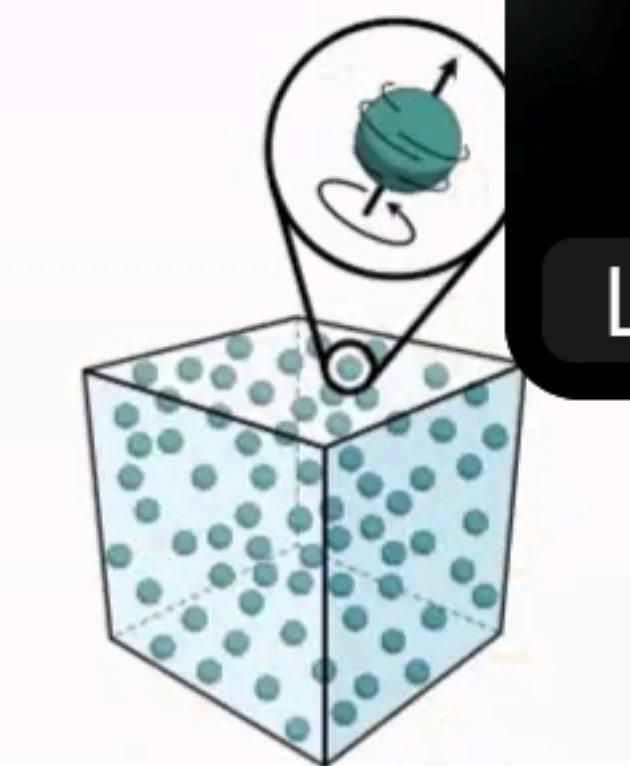
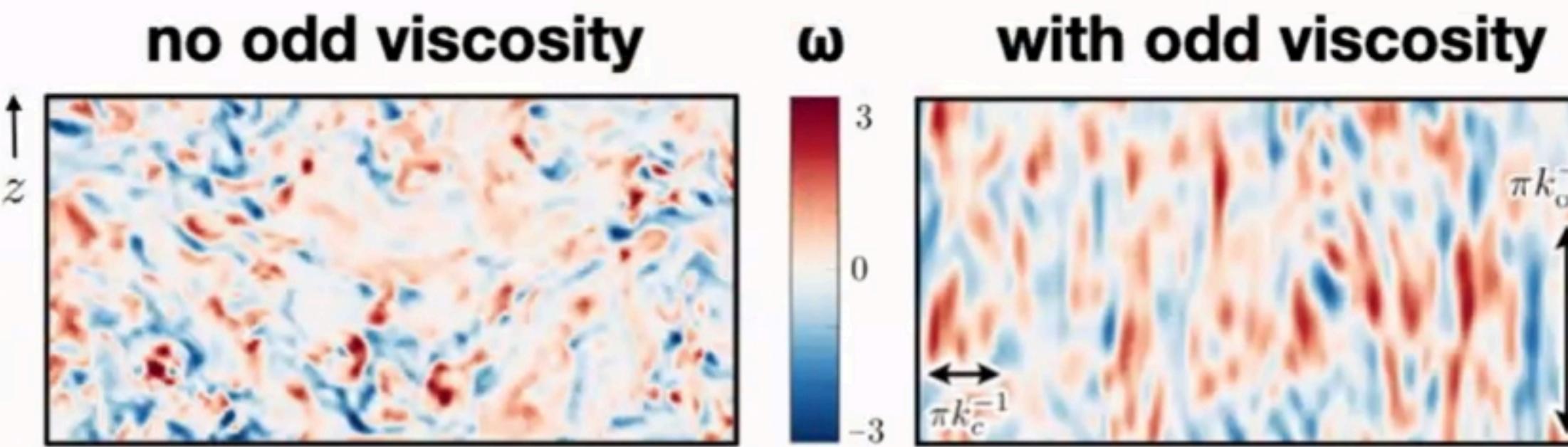
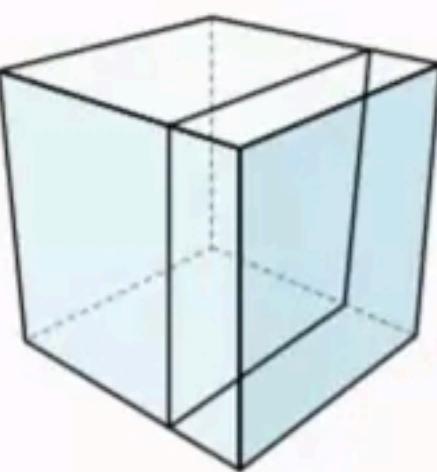


wavelength selection from arrest of turbulent cascade



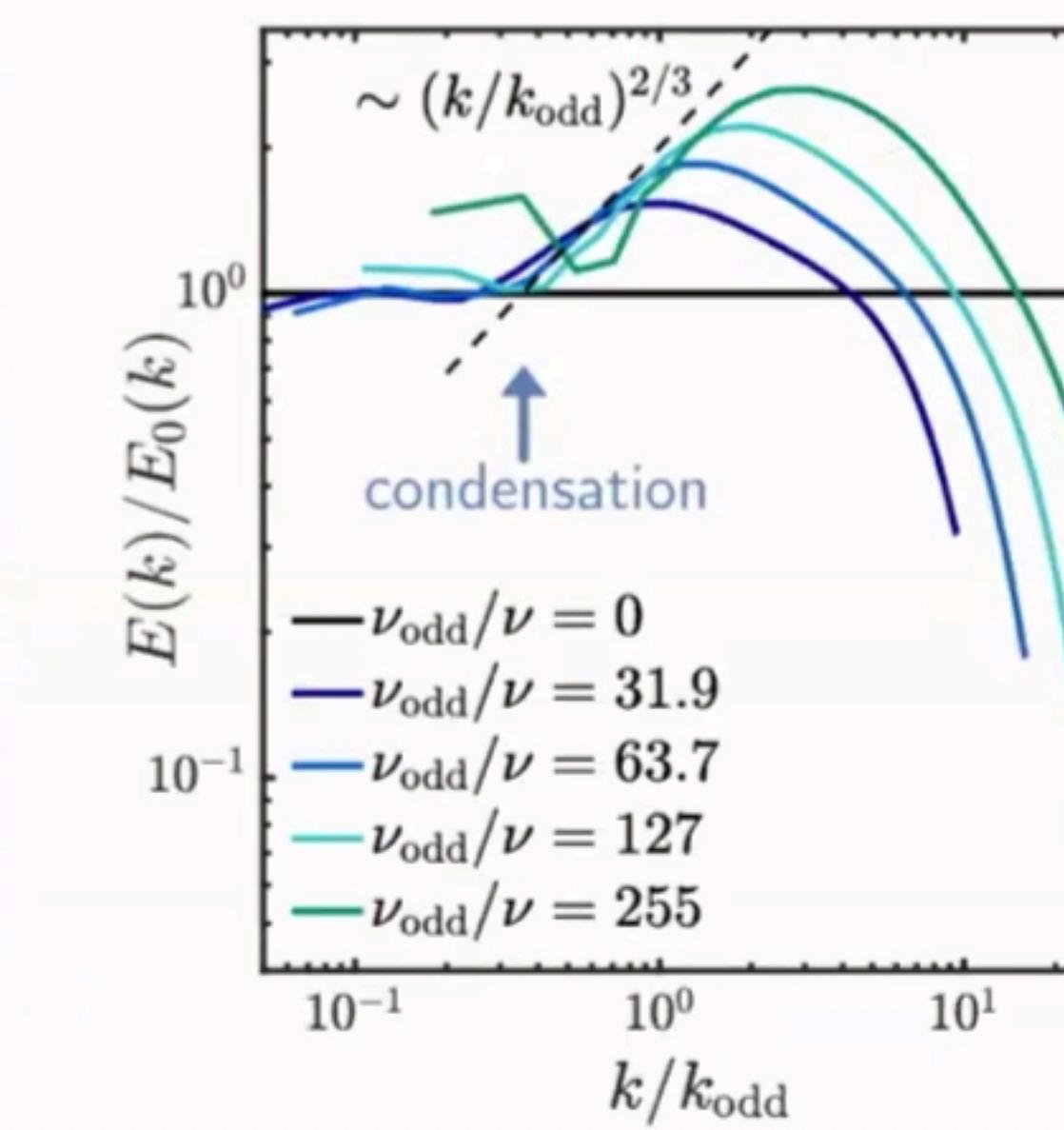
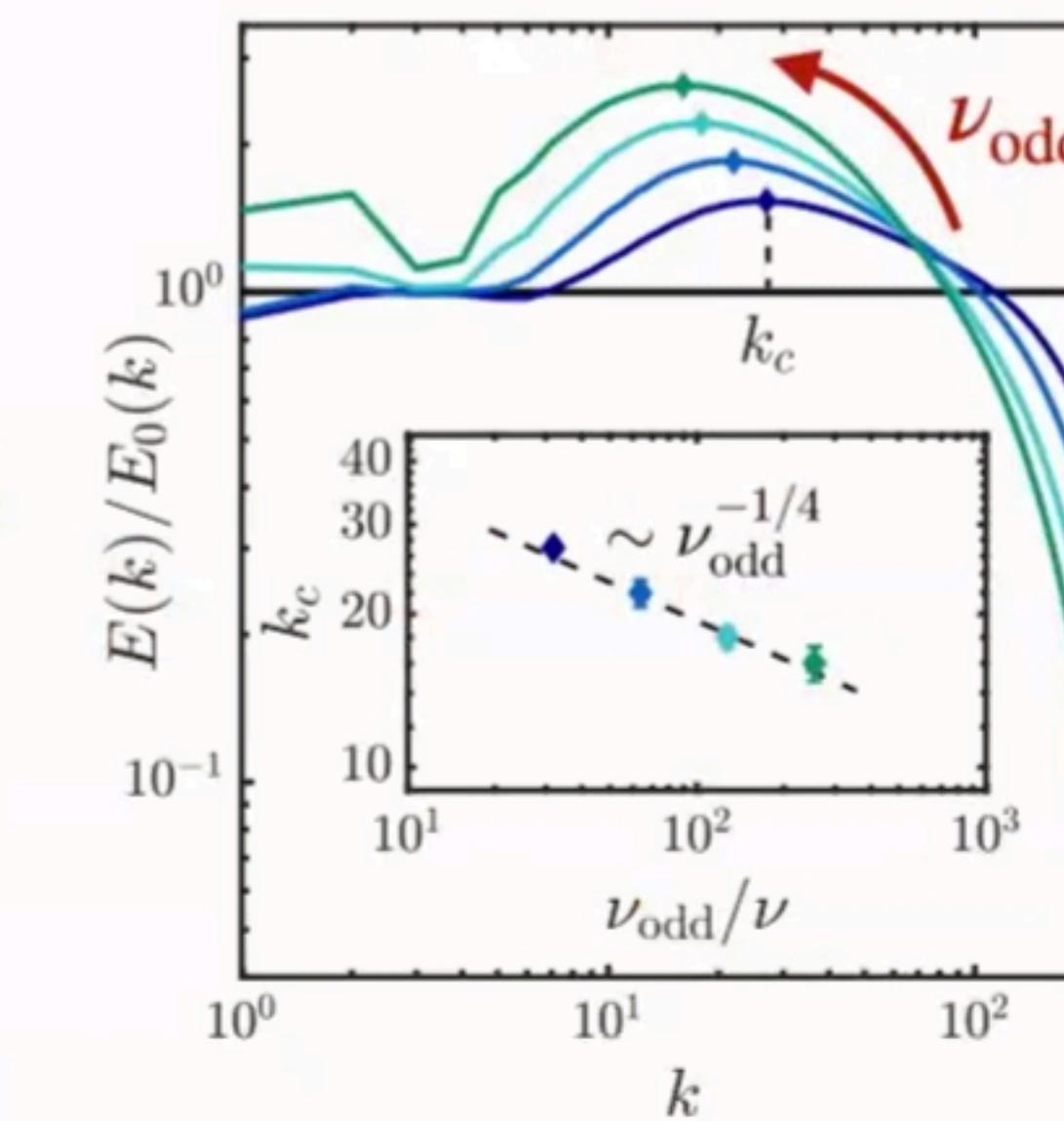


Direct numerical simulations

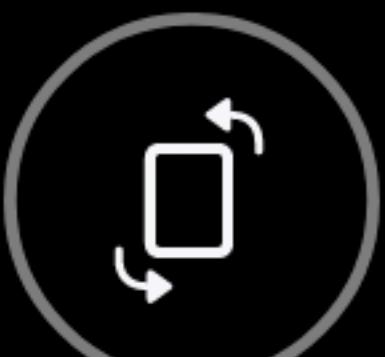


LMP Seminars

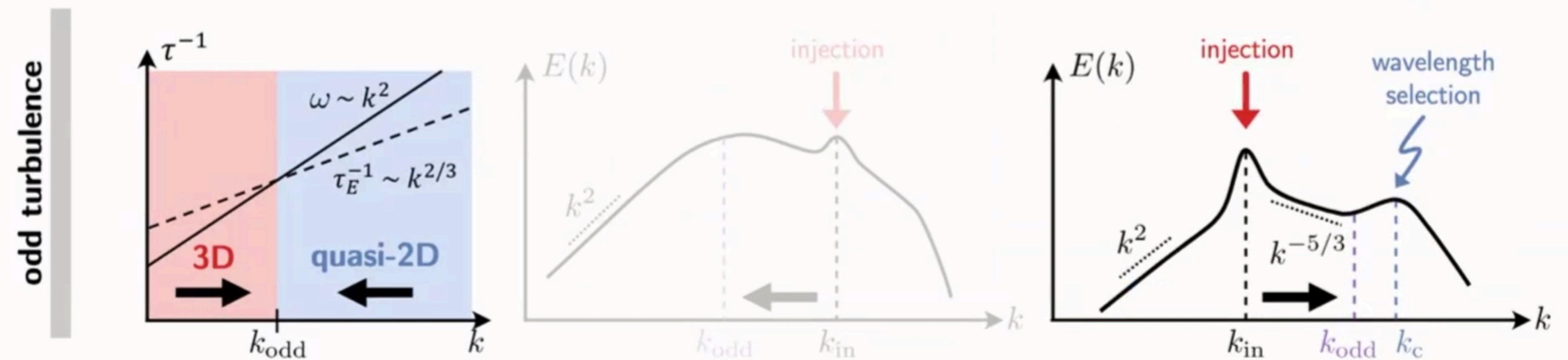
compensated spectrum
 $E(k) / k^{-5/3}$

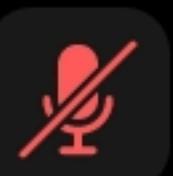


wavelength selection from arrest of turbulent cascade

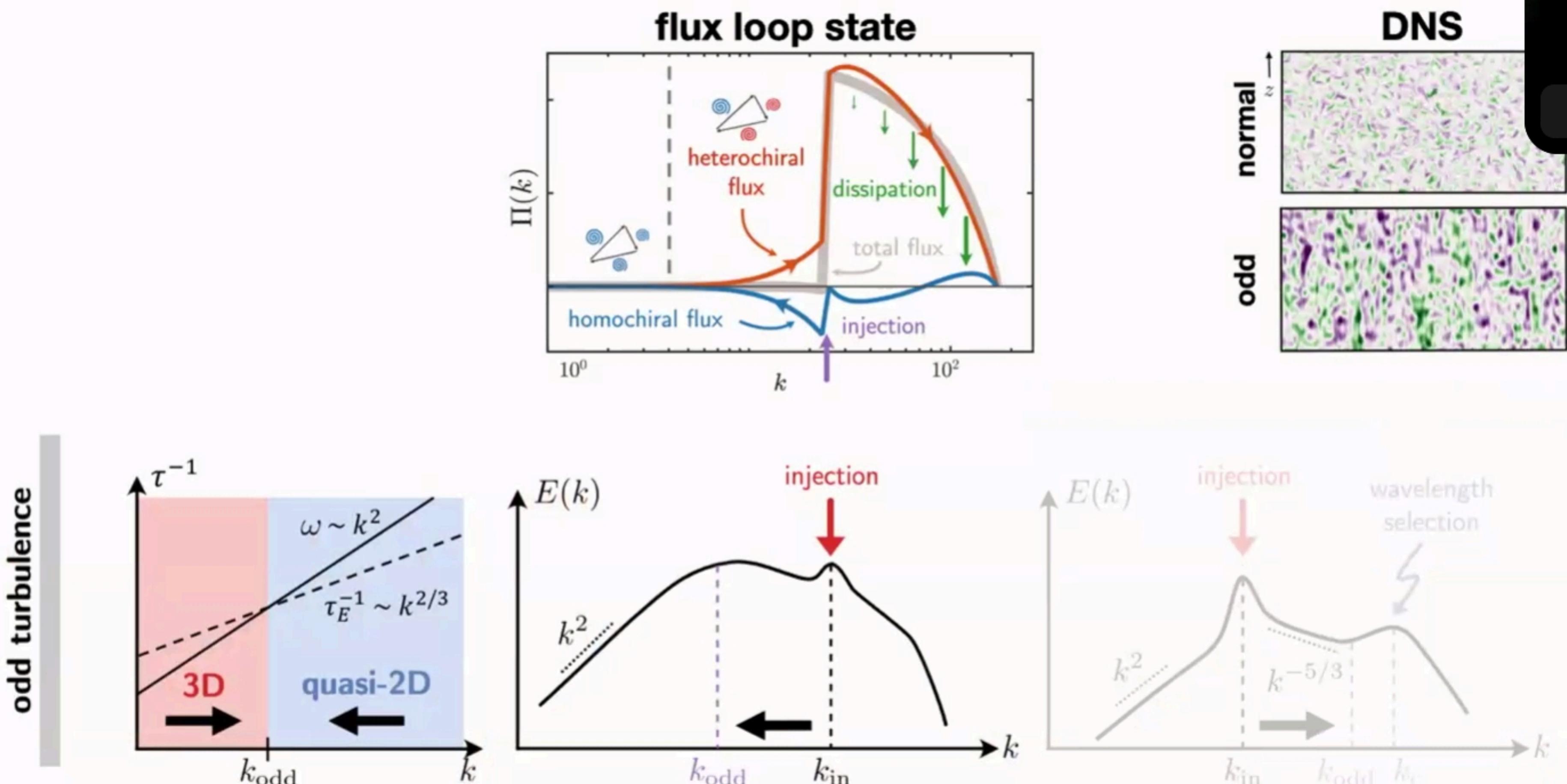


Odd turbulence





Odd turbulence



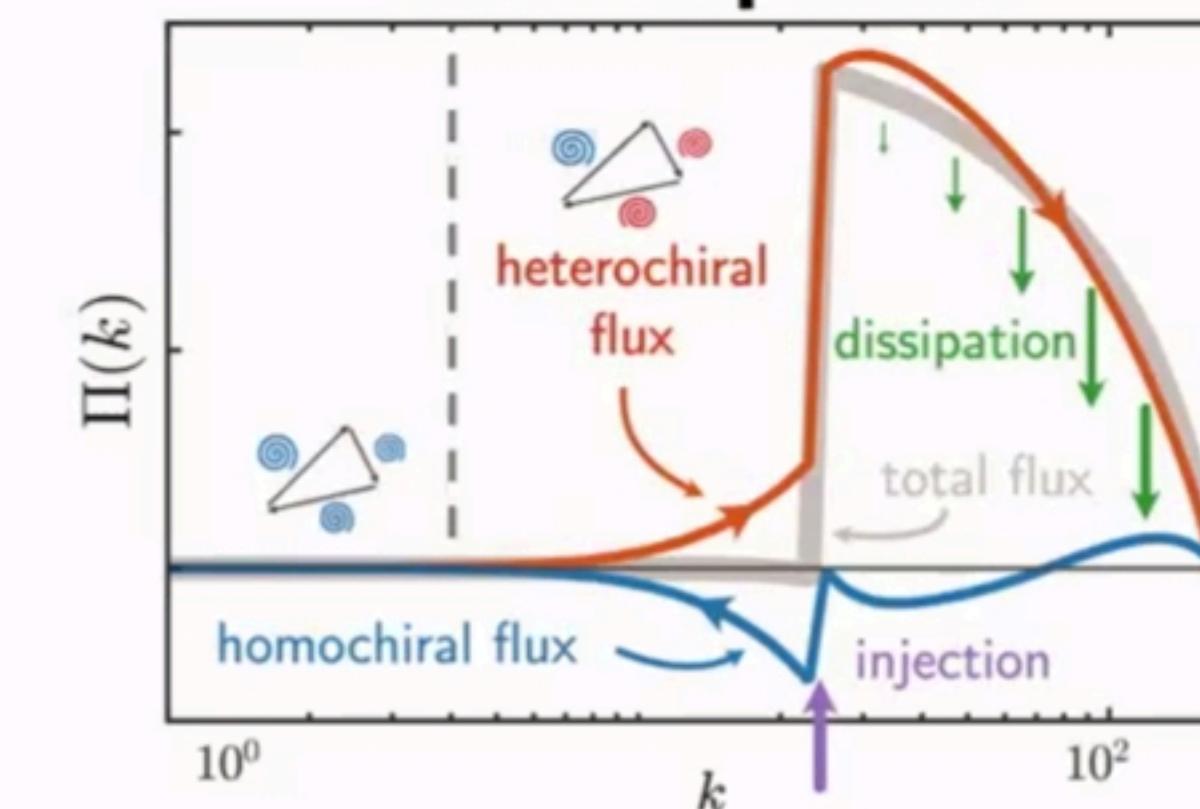
the inverse cascade can also be arrested



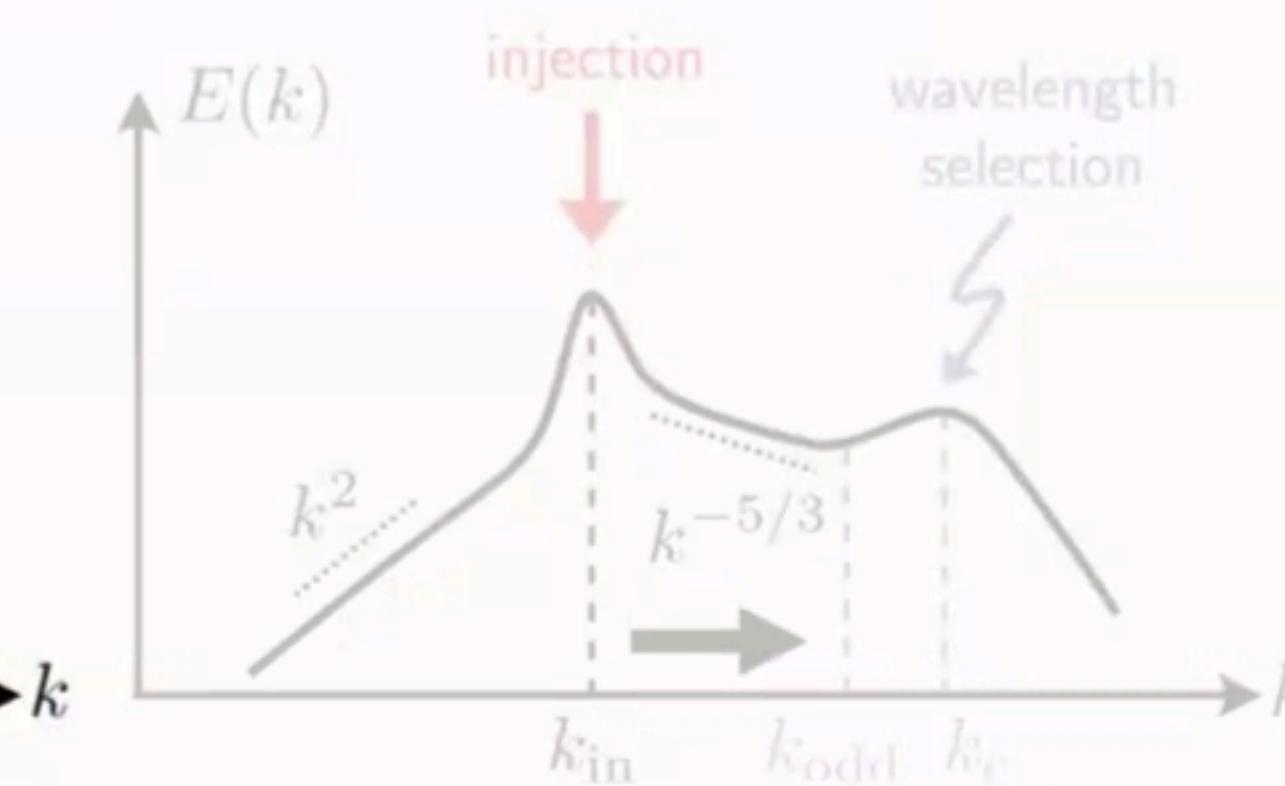
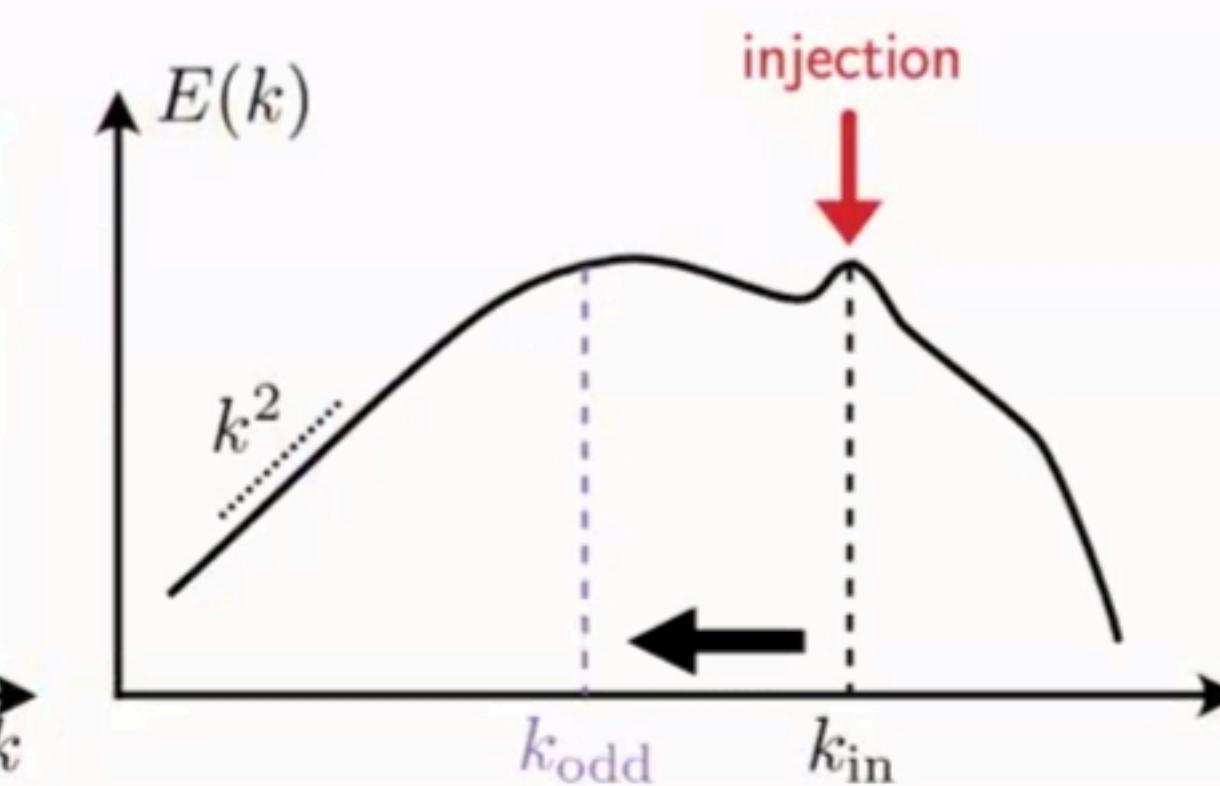
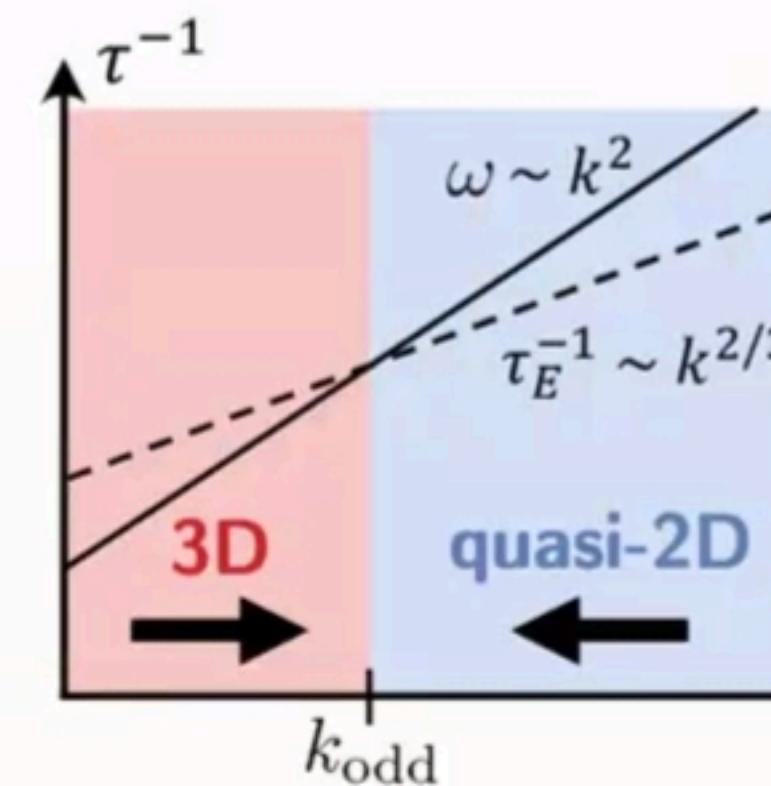


Odd turbulence

flux loop state

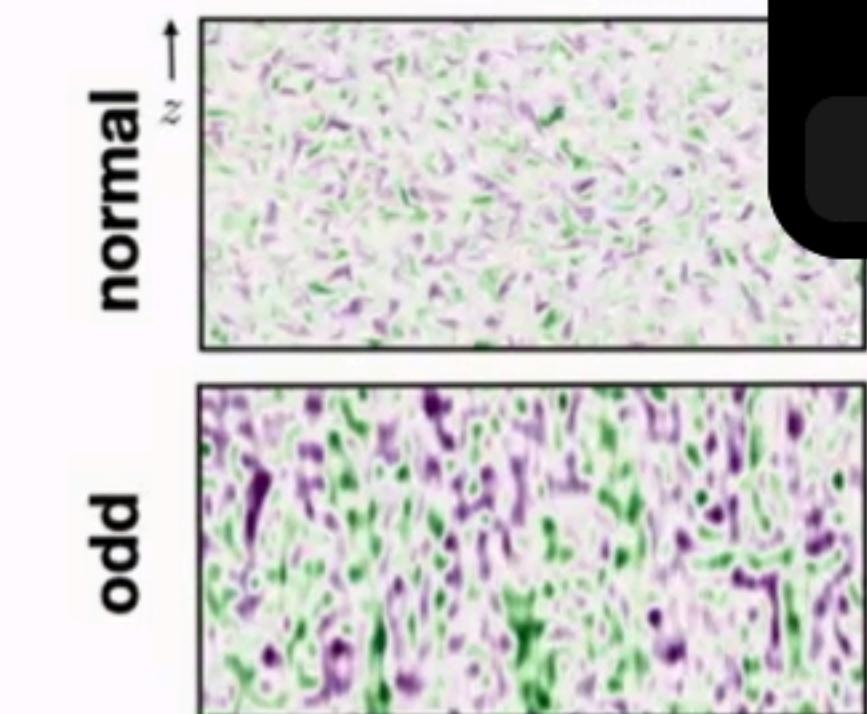


odd turbulence

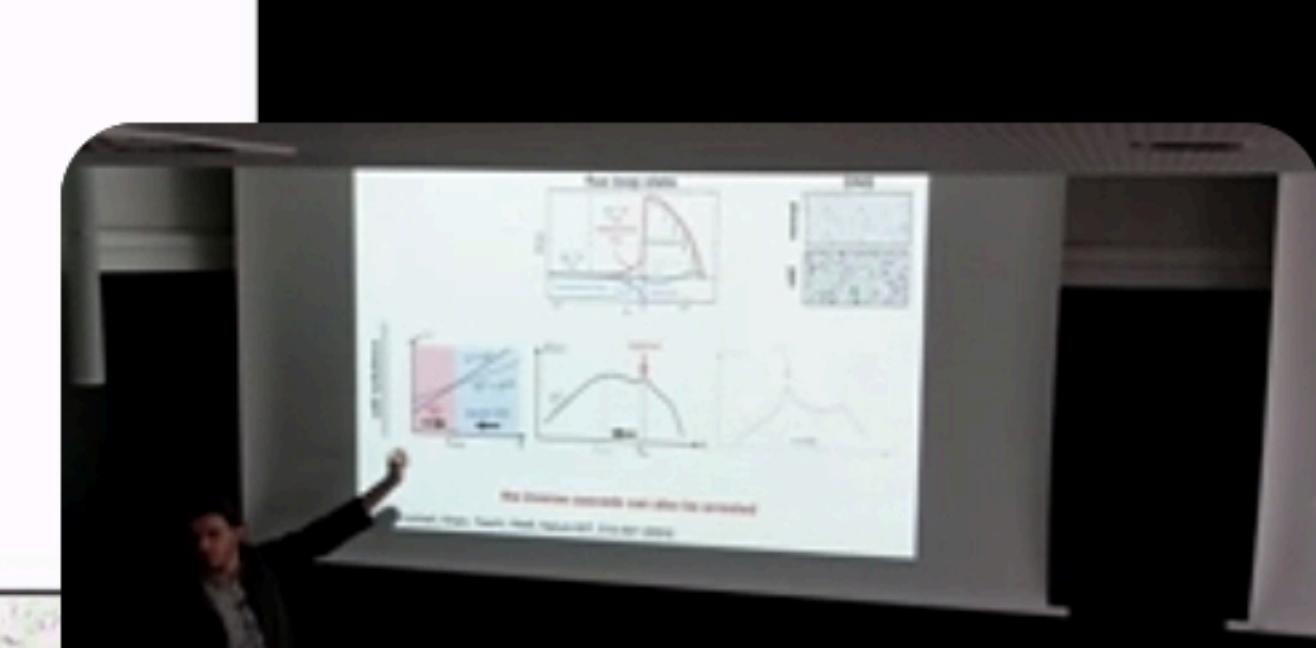


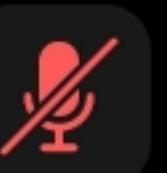
the inverse cascade can also be arrested

DNS



LMP Seminars





LMP Seminars

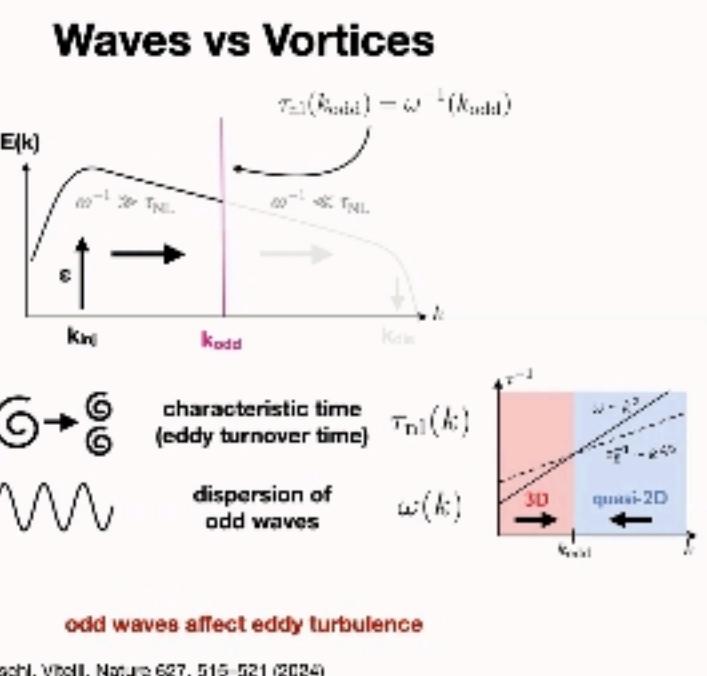
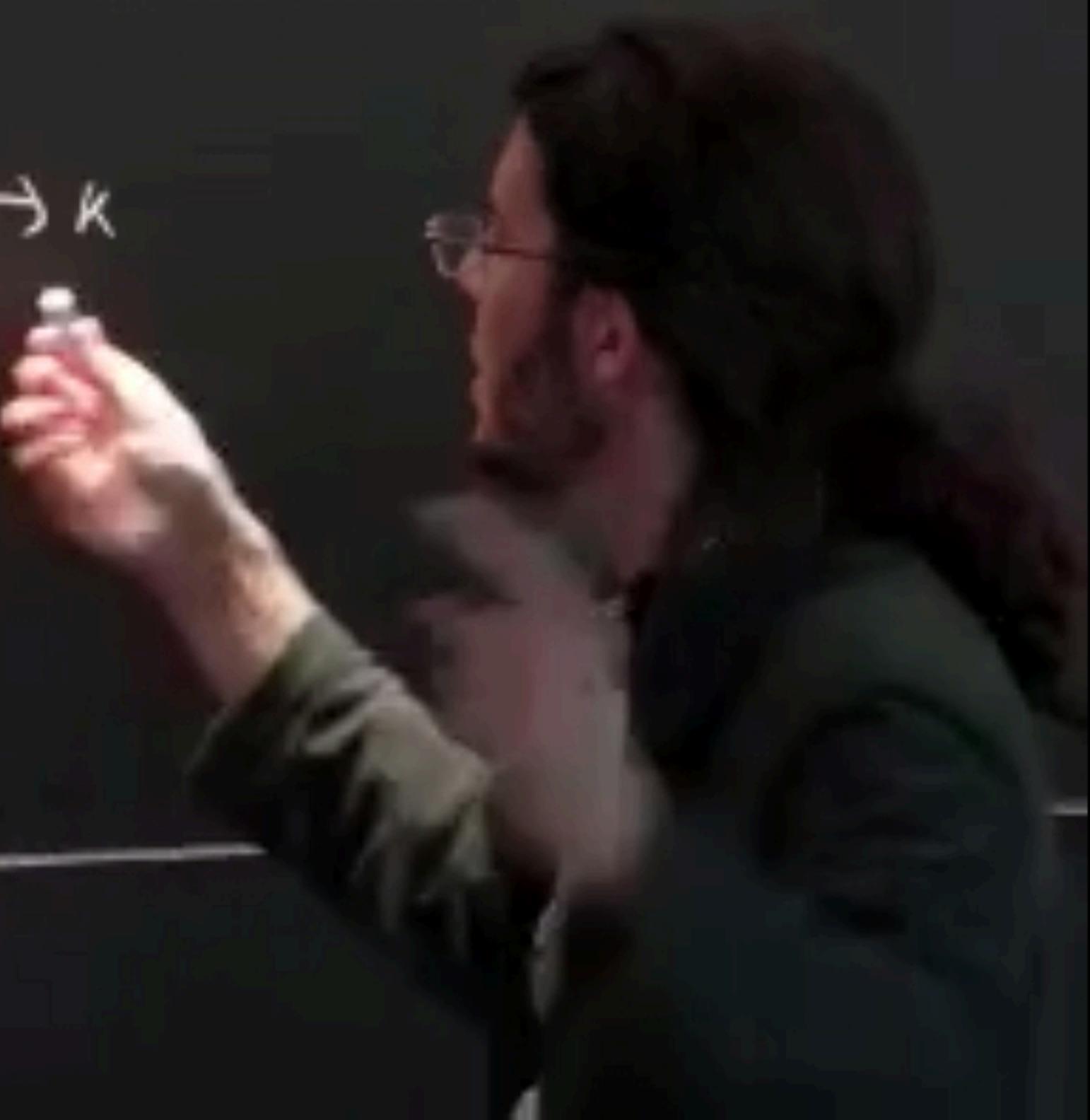
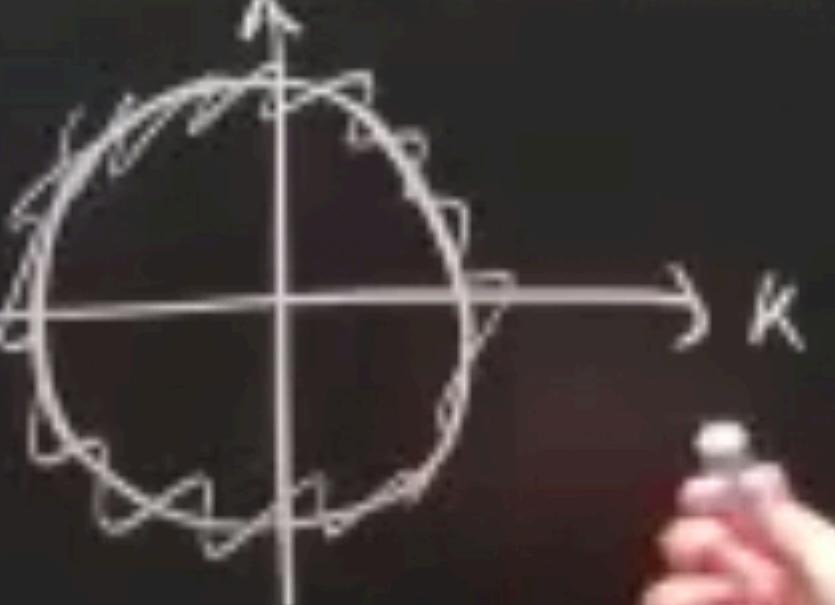


$$\rho D_t v = -\nabla P + \eta \Delta v + \eta \xi_2 \Delta v + \xi(t, x)$$

$$\downarrow$$

$$\partial_t v + v \cdot \nabla v$$

$$\langle \hat{\xi}(t, k) \hat{\xi}(t, -k) \rangle$$

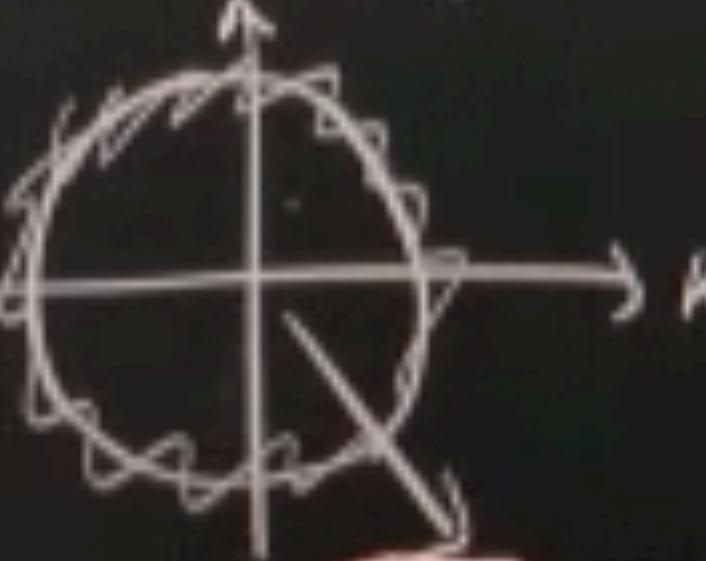




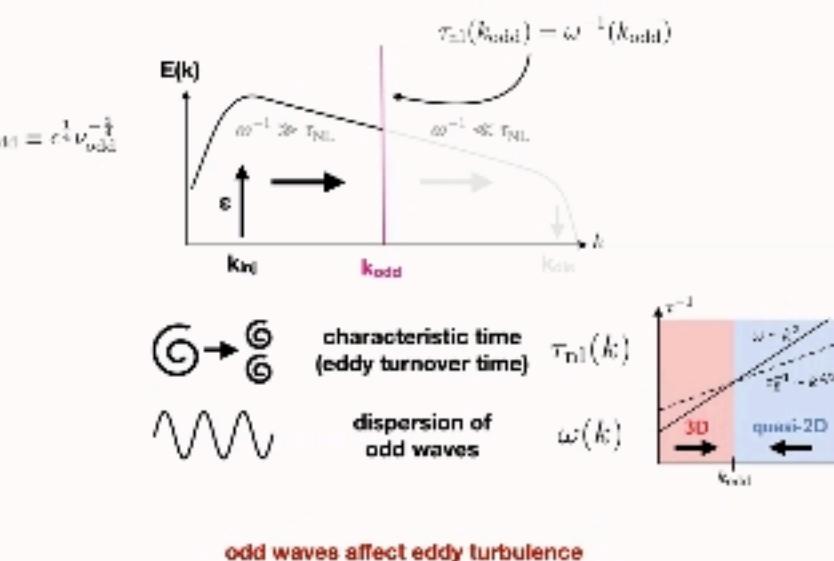
$$\rho D_t v = -\nabla P + \eta \Delta v + \cancel{\eta \frac{\partial}{\partial z} v} + \xi(t, z)$$

$$\downarrow$$

$$\partial_t v + v \cdot \nabla v$$

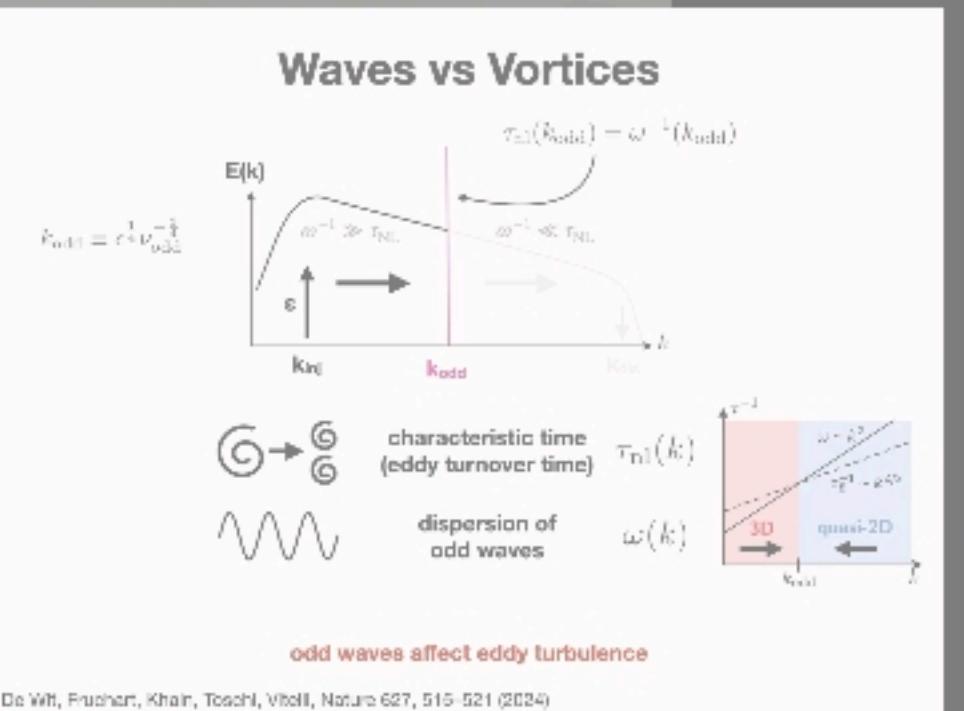
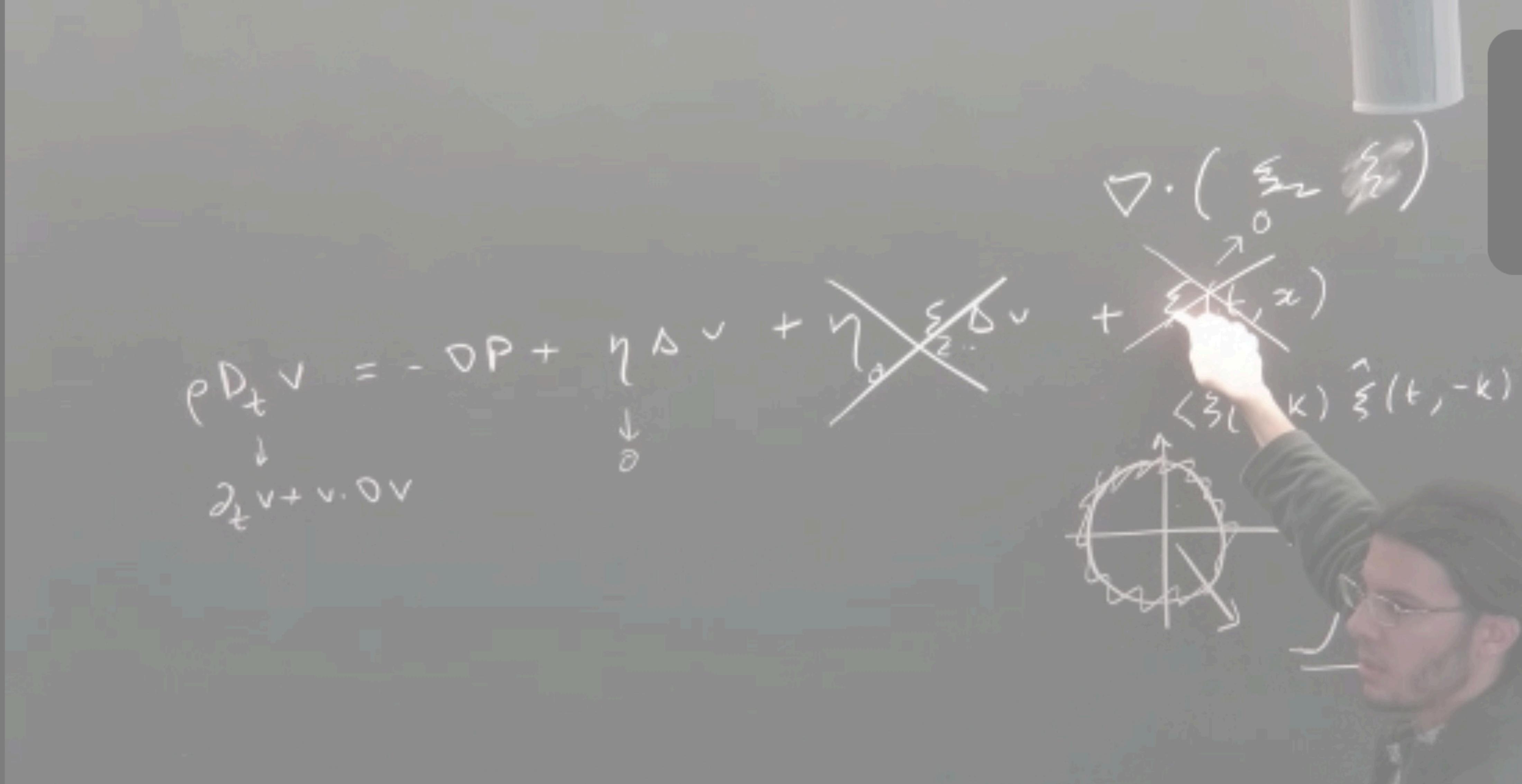


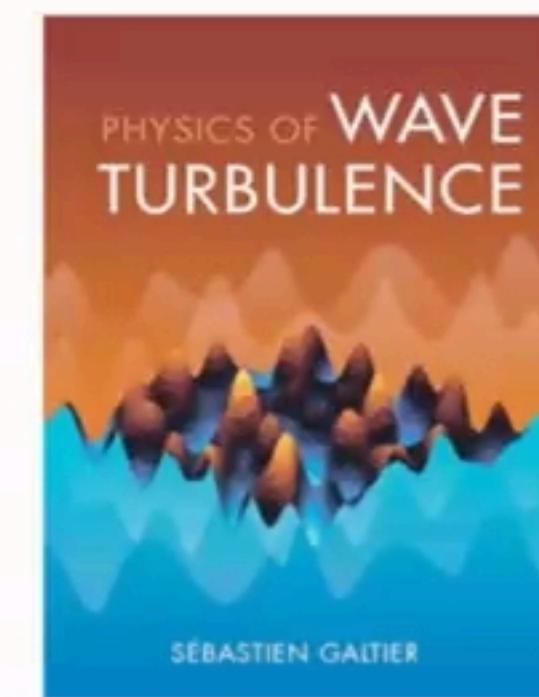
Waves vs Vortices



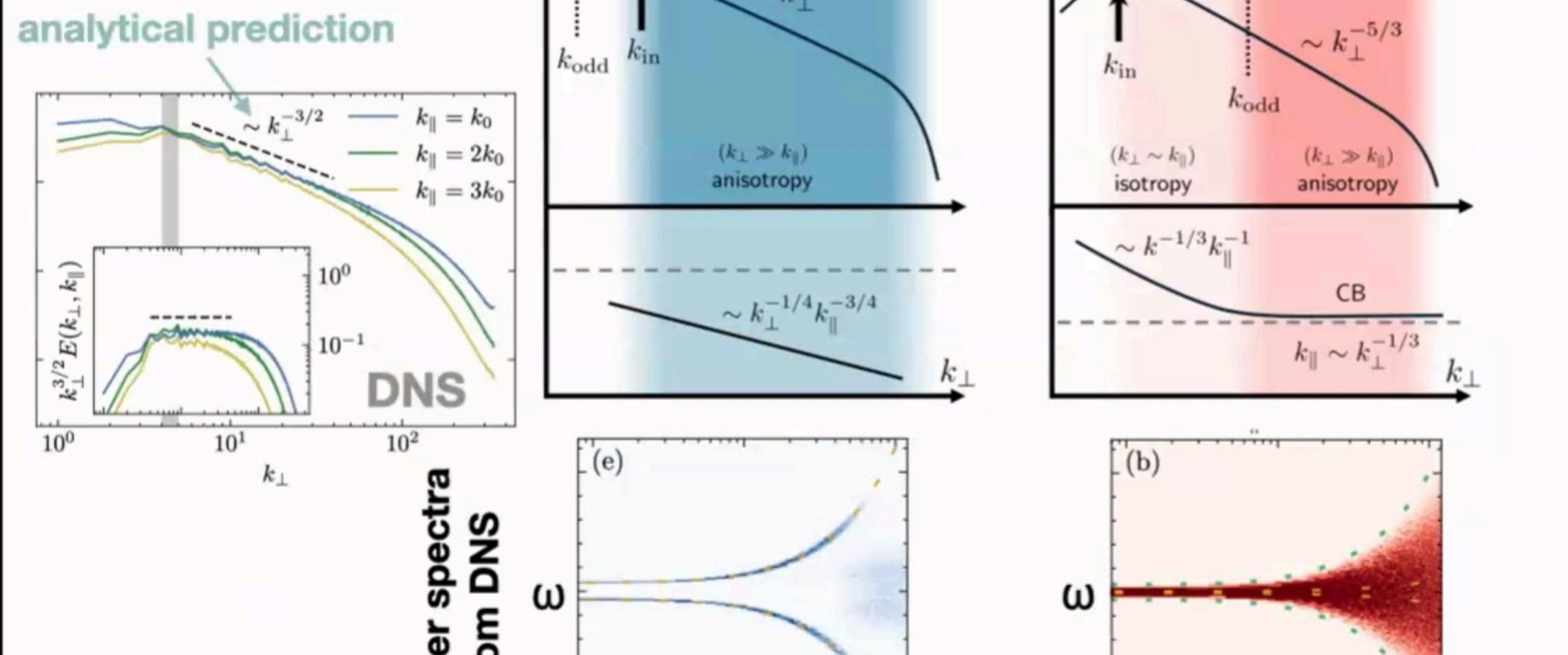
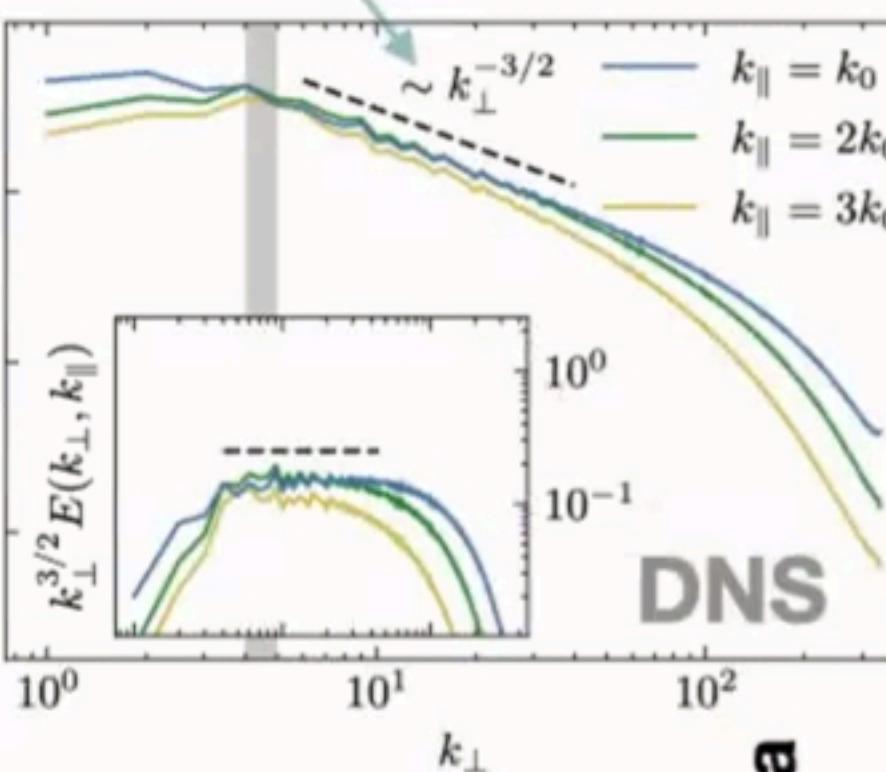
De Wit, Puchart, Khan, Toschi, Voth, Nature 627, 515-521 (2024)







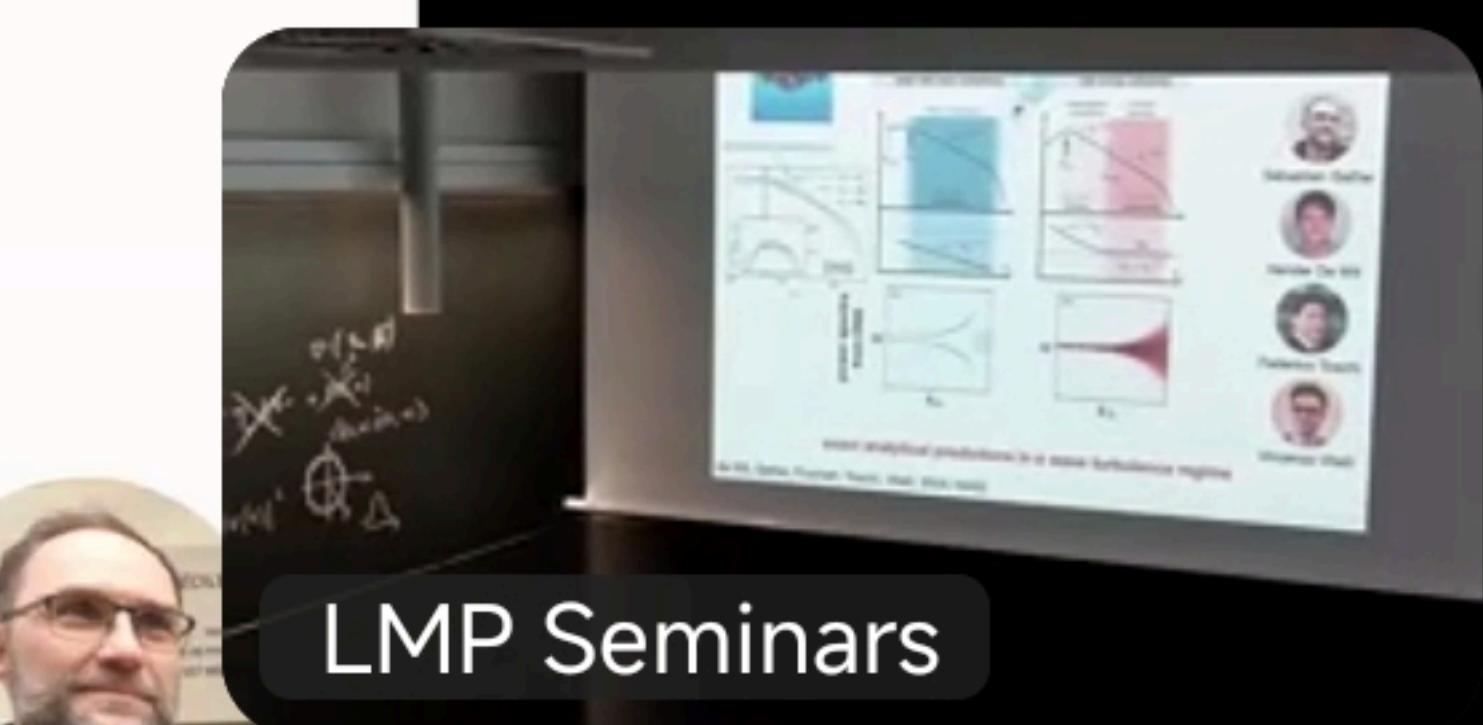
analytical prediction



power spectra
from DNS

exact analytical predictions in a wave turbulence regime

Odd wave turbulence



Sébastien Galtier



Xander De Wit



Federico Toschi



Vincenzo Vitelli