

Lesson 4

Numerical Solution of PDE

4.1 BVP for 2nd Order Elliptic PDE

- *Finite Difference Method*
- *Finite Element Method*



Numerical Methods for PDE: 2nd Order Elliptic PDE



Now, lets try to solve the Poisson's equation with homogeneous boundary condition

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using the finite element method where Ω is a bounded domain.

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$$-\int_{\Omega} \nabla u \cdot \nabla v = \int_{\Omega} f v.$$

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that is

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We, therefore, see that the Galerkin approximation error is bounded by a constant multiple of the best approximation error for u by functions in V_h !

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- ***Finite Difference Method***
- ***Finite Element Method***
- ***Construction of FEM Approximation Spaces***



Numerical Methods for PDE: 2nd Order Elliptic PDE



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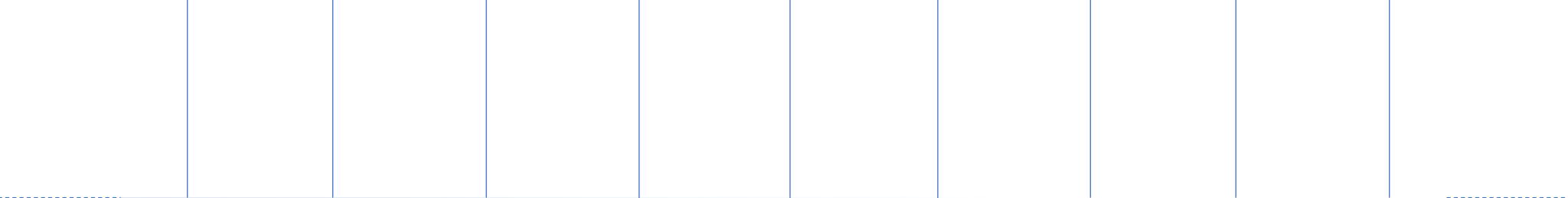


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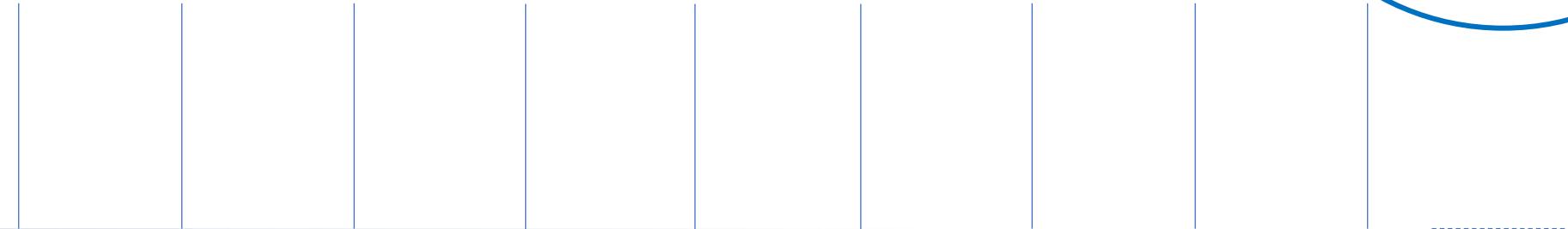


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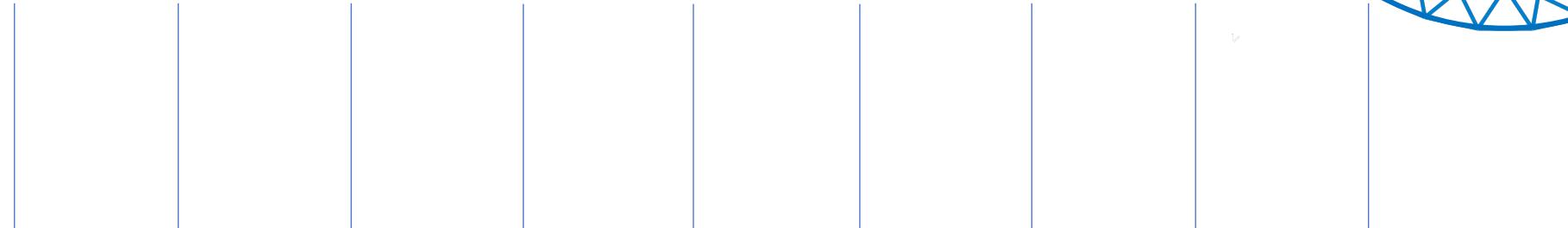
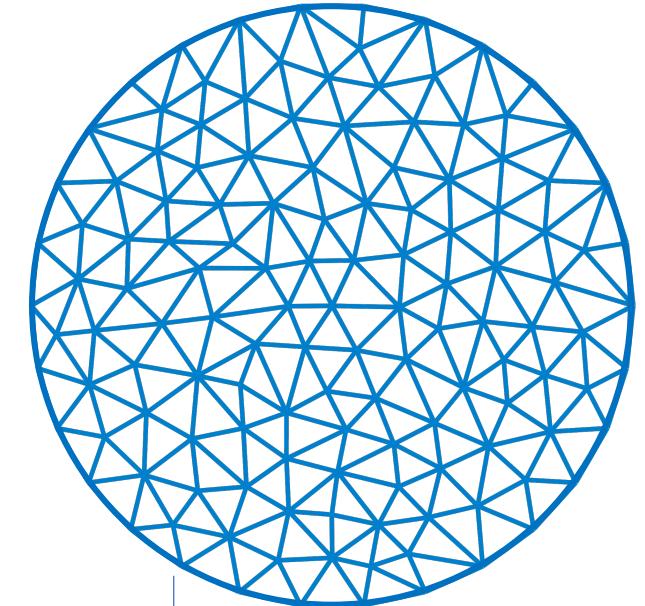


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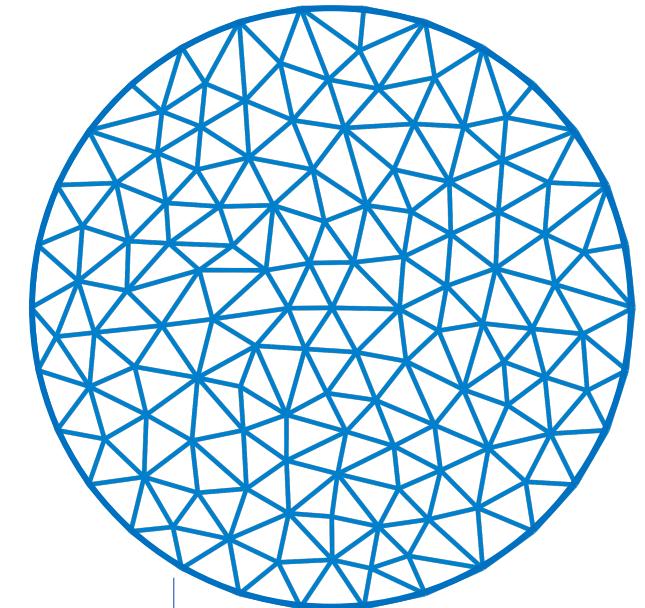
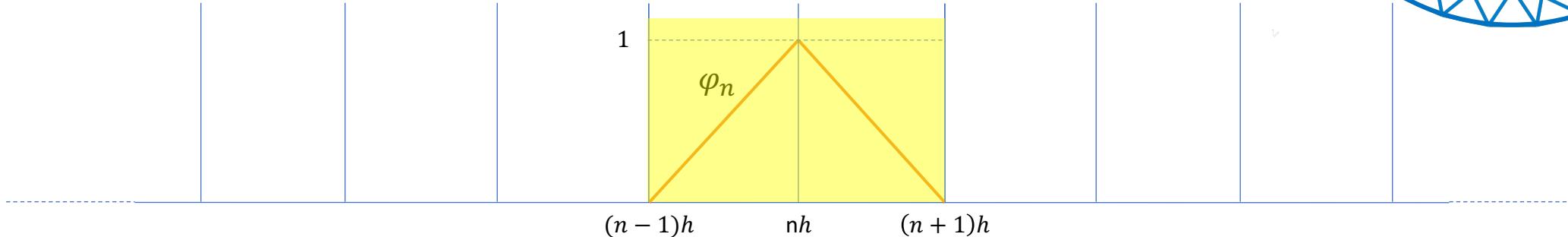
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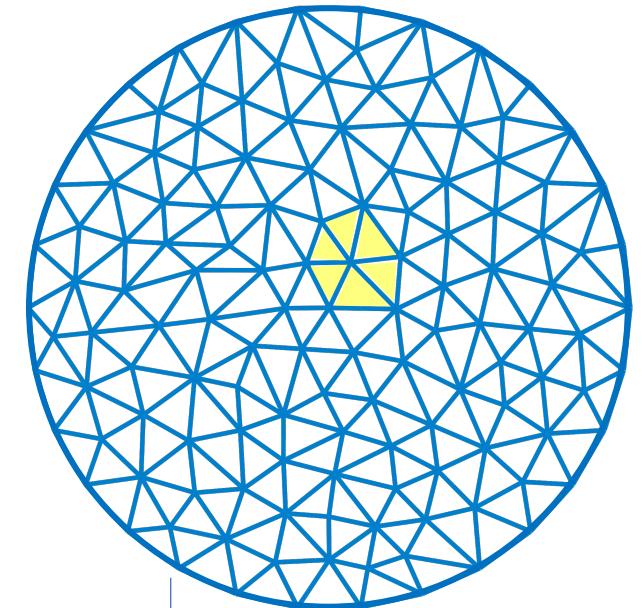
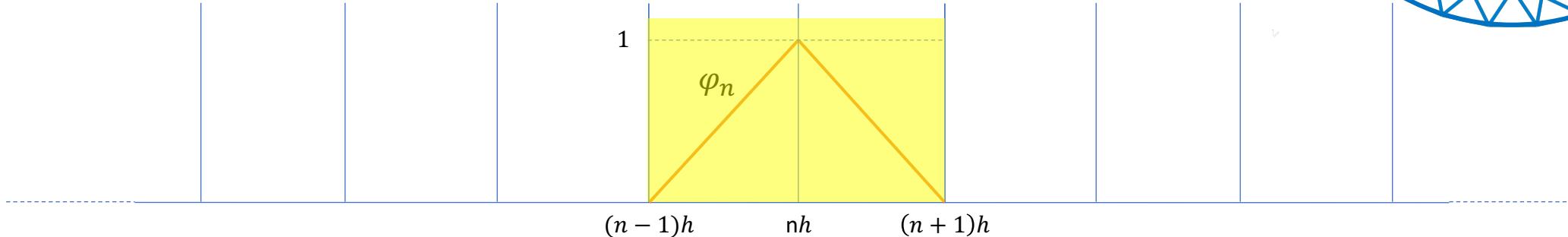


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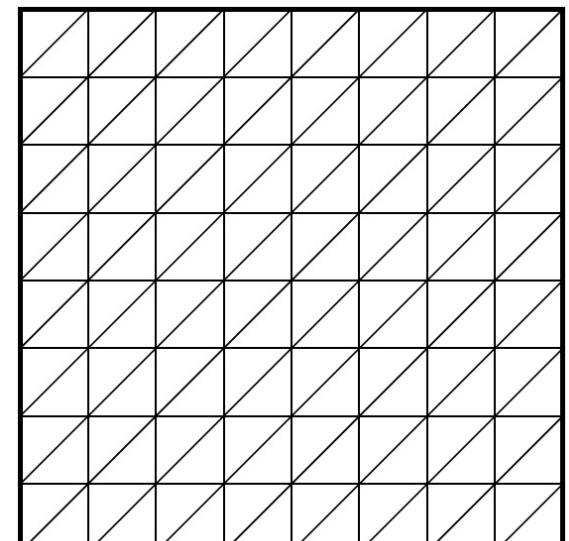
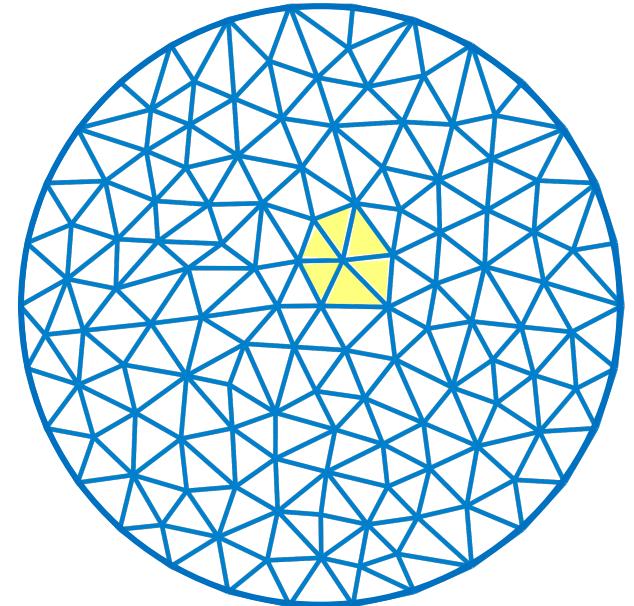
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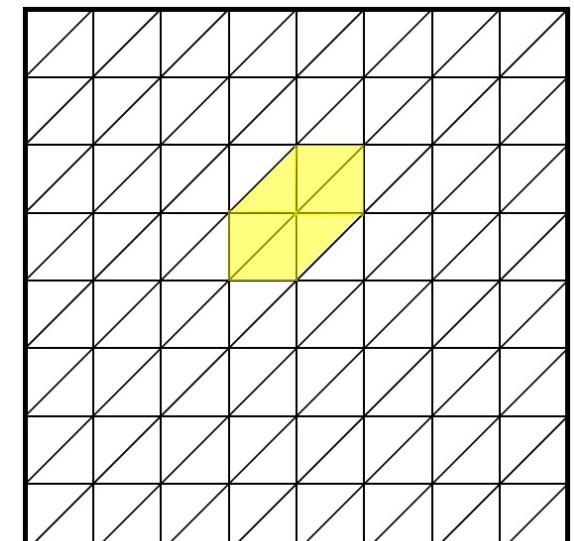
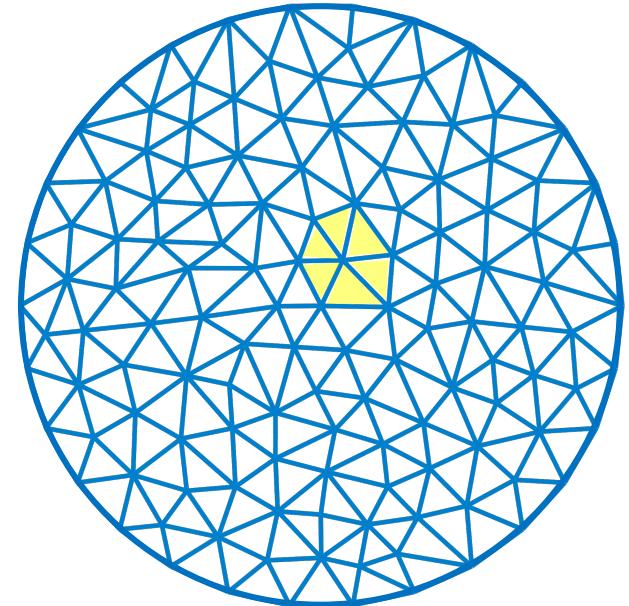
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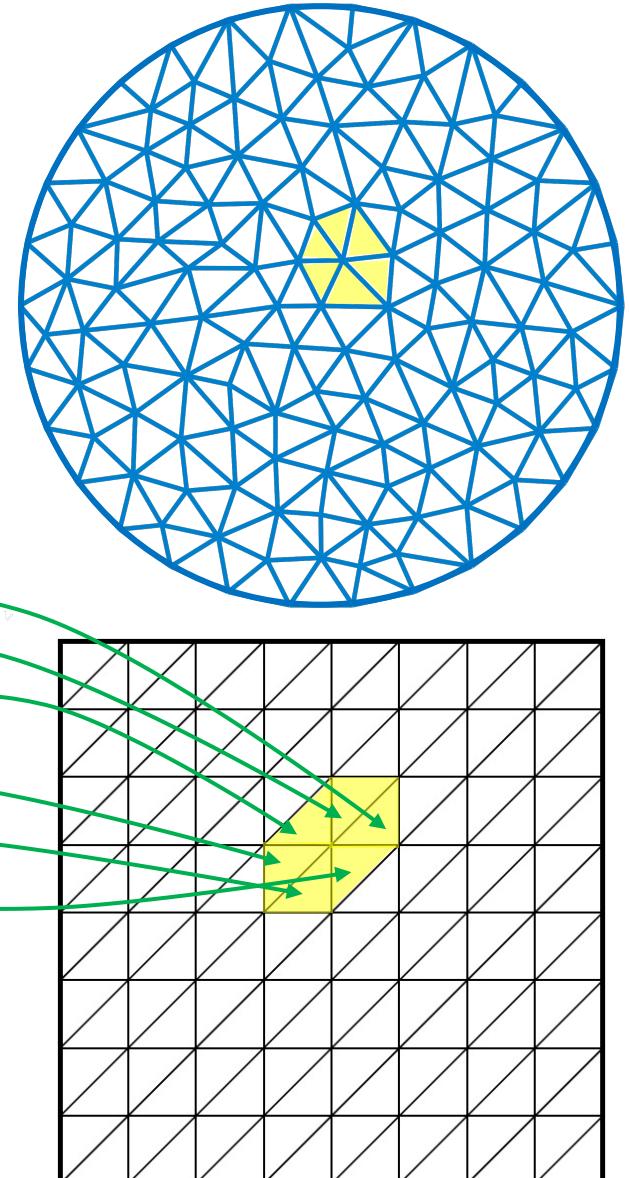
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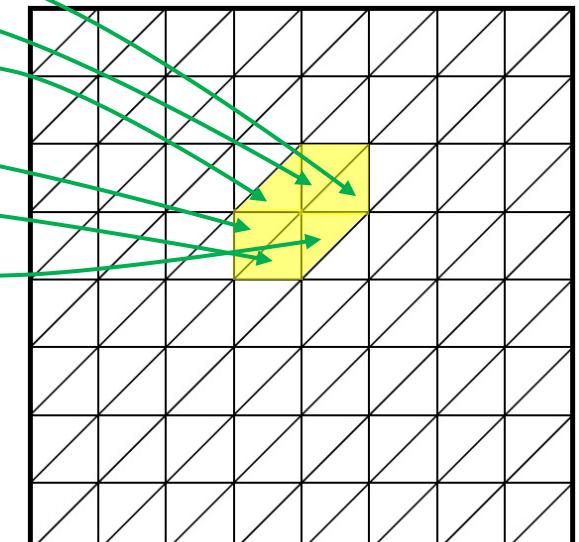
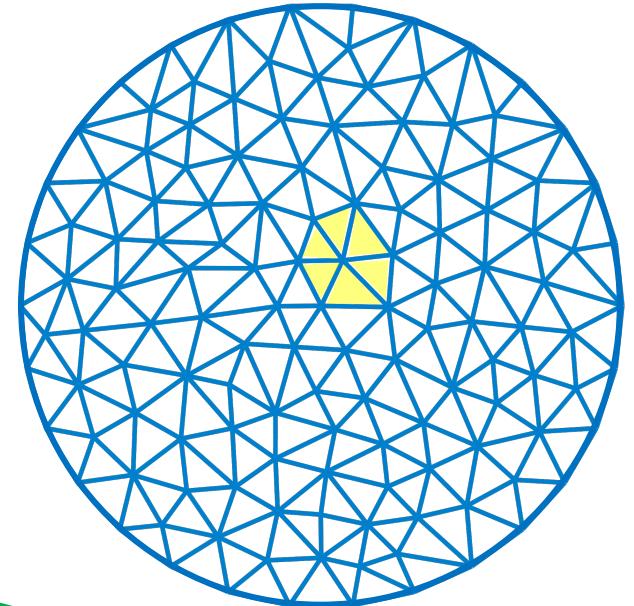
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$$\varphi_{mn}(x_1, x_2) = \begin{cases} 1 - (x_1 - mh)/h & \\ 1 - (x_2 - nh)/h & \\ 1 + (x_1 - mh)/h - (x_2 - nh)/h & \\ 1 + (x_1 - mh)/h & \\ 1 + (x_2 - nh)/h & \\ 1 - (x_1 - mh)/h + (x_2 - nh)/h & \end{cases}$$

The functions φ_{mn} form a basis for subspace of piecewise linear functions with respect to the given partition/triangulation.



Numerical Methods for PDE: 2nd Order Elliptic PDE

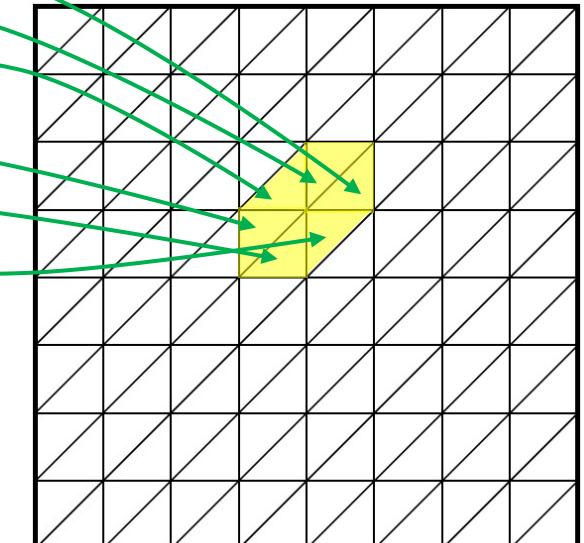


Note that (*exercise*)

$$\int_{\Omega} \nabla \varphi_{mn} \cdot \nabla \varphi_{kl} = \begin{cases} 4, & m = k, n = l, \\ -1, & m = k \pm 1, n = l \text{ or } m = k, n = l \pm 1, \\ 0, & \text{otherwise.} \end{cases}$$

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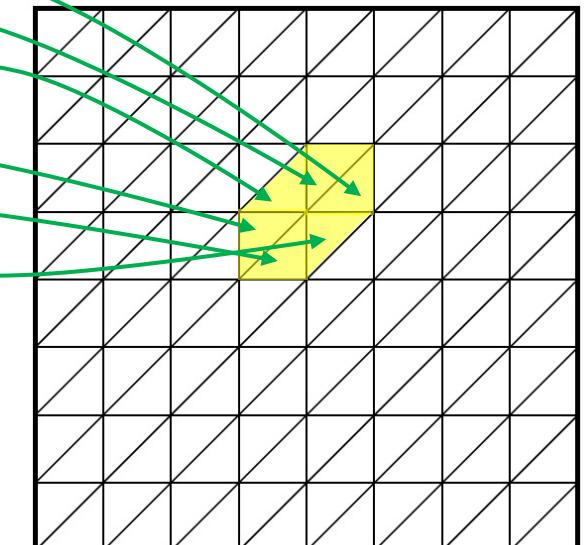
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Thus, for $u_h = \sum u_{mn} \varphi_{mn}$, the linear system reads

$$\frac{u_{m-1,n} + u_{m+1,n} + u_{m,n-1} + u_{m,n+1} - 4u_{mn}}{h^2} = \frac{1}{h^2} \int_{\Omega} f \varphi_{mn} = \tilde{f}_{mn}, \quad 1 \leq m, n \leq N.$$

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Numerical Methods for PDE: 2nd Order Elliptic PDE



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We see that matrix on the left hand side of the linear matrix (called stiffness matrix) for the piecewise linear finite elements for the Laplace operator on the unit square using a uniform mesh is exactly the matrix of the 5-point Laplacian.

Numerical Methods for PDE: 2nd Order Elliptic PDE



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Numerical Methods for PDE: 2nd Order Elliptic PDE



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Note that (*exercise*)

$$\int_{\Omega} \varphi_{mn} = h^2$$

and if $f \in C^2$, then (*exercise*)

$$\tilde{f}_{mn} = \frac{1}{h^2} \int_{\Omega} \left(f(mh, nh) + f_{x_1}(mh, nh)(x_1 - mh) + f_{x_2}(mh, nh)(x_2 - nh) + O(h^2) \right) \varphi_{mn} = f_{mn} + O(h^2)$$

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Therefore, at vertices, we see that the finite element method converges with order 2 (*why?*).

Numerical Methods for PDE: 2nd Order Elliptic PDE



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Recall that, from the general error analysis, we have

$$\|u_h - u\|_{H^1} \leq (1 + C/\gamma) \inf_{w \in V_h} \|u - w\|_{H^1}.$$

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Exercise: Find C and γ for the piecewise linear finite elements.

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From the approximation theory, we have the following result on the best approximation error.

Theorem

Let there be given a family of triangulations $\{\mathcal{T}_h\}$ of a polygonal domain Ω and let $h = \max_{T \in \mathcal{T}_h} \text{diam}(T)$. Let r be a positive integer. For each h let $P_h: C(\Omega) \rightarrow M_0^r(\mathcal{T}_h)$ denote the nodal interpolant, where $M_0^r(\mathcal{T}_h)$ is the space of continuous functions which restrict to polynomials of degree at most r when restricted to any triangle $T \in \mathcal{T}_h$. Then, there is a constant c such that

$$\begin{aligned} \|u - P_h u\|_{L^\infty(\Omega)} &\leq ch^{r+1} \|u^{(r+1)}\|_{L^\infty(\Omega)}, & u \in C^{r+1}(\overline{\Omega}), \\ \|u - P_h u\|_{L^2(\Omega)} &\leq ch^{r+1} \|u^{(r+1)}\|_{L^2(\Omega)}, & u \in H^{r+1}(\Omega). \end{aligned}$$

Moreover, if the family of triangulations are shape regular (the minimal angle of each triangulation is bounded below uniformly), then there is a constant C such that

$$\begin{aligned} \|\nabla(u - P_h u)\|_{L^\infty(\Omega)} &\leq Ch^r \|u^{(r+1)}\|_{L^\infty(\Omega)}, & u \in C^{r+1}(\overline{\Omega}), \\ \|\nabla(u - P_h u)\|_{L^2(\Omega)} &\leq Ch^r \|u^{(r+1)}\|_{L^2(\Omega)}, & u \in H^{r+1}(\Omega). \end{aligned}$$

Lesson 4

Numerical Solution of PDE

4.2 Parabolic PDE



Numerical Methods for PDE: Parabolic PDE



As a model problem, consider the heat equation on a spatial domain Ω for a time interval $[0, T]$. The solution u satisfies

$$\frac{\partial u}{\partial t} = c \Delta u + f, \quad x \in \Omega, t \in [0, T], \quad c > 0.$$

Numerical Methods for PDE: Parabolic PDE



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To obtain a well-posed problem, we need to give boundary conditions

$$u = 0, \quad x \in \Gamma, t \in [0, T]$$

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Lesson 4

Numerical Solution of PDE

4.2 Parabolic PDE

- Semi-discretization



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Let's assume, for simplicity, that $\Omega = (0,1) \times (0,1)$. Then, we have the 5-point Laplacian Δ_h as the discretization of the Laplacian Δ . Thus, we have the solution u_h satisfying

$$\begin{aligned} \frac{\partial u_h}{\partial t} &= c\Delta_h u_h + f, & x \in \Omega_h, t \in [0, T] \\ u_h &= 0, & x \in \Gamma_h, t \in [0, T] \\ u_h &= u_0, & x \in \Omega_h, t = 0. \end{aligned}$$

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We may view this as an initial value problem for a system of $(N - 1)^2$ ODEs.

Numerical Methods for PDE: Parabolic PDE



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The process of reducing the evolutionary PDE to a system of ODEs by using a finite difference approximation of the spatial operator is called **semi-discretization** or the **method of lines**.

Numerical Methods for PDE: Parabolic PDE



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The process of reducing the evolutionary PDE to a system of ODEs by using a finite difference approximation of the spatial operator is called **semi-discretization** or the **method of lines**.

This is not a full discretization as we still have to choose a numerical method to solve the ODEs.

Lesson 4

Numerical Solution of PDE

4.2 Parabolic PDE

- *Semi-discretization*
- ***Finite difference discretization***



Numerical Methods for PDE: Parabolic PDE



To further simplify the presentation, let's drop down to one space dimension, that is, take $\Omega = (0,1)$.

Numerical Methods for PDE: Parabolic PDE



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Numerical Methods for PDE: Parabolic PDE



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To solve the ODEs, we first employ the Euler method.

Writing $u_n^j = u_h(nh, jk)$, the fully discrete system reads

$$\frac{u_n^{j+1} - u_n^j}{k} = c \frac{u_{n+1}^j - 2u_n^j + u_{n-1}^j}{h^2} + f_n^j, \quad 0 < n < N, j = 0, 1, \dots, M-1,$$

$$u_0^j = u_N^j = 0, \quad j = 0, 1, \dots, M-1,$$

$$u_n^0 = u_0(nh), \quad 0 < n < N.$$

Numerical Methods for PDE: Parabolic PDE



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We call this the **forward-centered difference method** for the heat equation. Since the Euler's method is explicit, we don't need to solve a linear system:

$$u_n^{j+1} = (1 - 2\lambda)u_n^j + \lambda u_{n+1}^j + \lambda u_{n-1}^j + kf_n^j, \quad j = 0, 1, \dots, M-1,$$

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To analyze the method, consider the local (truncation) error

$$\ell_n^j = \frac{u(nh, (j+1)k) - u(nh, jk)}{k} - c \frac{u((n+1)h, jk) - 2u(nh, jk) + u((n-1)h, jk)}{h^2} - f_n^j.$$

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By Taylor's theorem,

$$\ell_n^j = \frac{k}{2} \frac{\partial^2 u}{\partial t^2} - c \frac{h^2}{12} \frac{\partial^4 u}{\partial x^4}$$

where derivatives are evaluated at appropriate points.

Numerical Methods for PDE: Parabolic PDE



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Let $e_n^j = u_n^j - u(nh, jk)$. Then,

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If $\lambda \leq 1/2$, and $\ell = \max_{n,j} |\ell_n^j|$, we get

$$E^{j+1} \leq E^j + k\ell.$$

Numerical Methods for PDE: Parabolic PDE



To analyze the method, consider the local (truncation) error

$$\ell_n^j = \frac{u(nh, (j+1)k) - u(nh, jk)}{k} - c \frac{u((n+1)h, jk) - 2u(nh, jk) + u((n-1)h, jk)}{h^2} - f_n^j.$$

By Taylor's theorem,

$$\ell_n^j = \frac{k}{2} \frac{\partial^2 u}{\partial t^2} - c \frac{h^2}{12} \frac{\partial^4 u}{\partial x^4}$$

where derivatives are evaluated at appropriate points.

Let $e_n^j = u_n^j - u(nh, jk)$. Then,

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$$\max_{n,j} |e_n^j| \leq T\ell \leq C(k + h^2) \leq C(1 + 1/(2c))h^2.$$

Lesson 4

Numerical Solution of PDE

4.2 Parabolic PDE

- *Semi-discretization*
- *Full finite difference discretization*
- ***Fourier Analysis***



Numerical Methods for PDE: Parabolic PDE



Another useful way to analyze is to use Fourier analysis.

Numerical Methods for PDE: Parabolic PDE



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Recall that, on $L(I_h)$, we define the inner product

$$\langle u, v \rangle_h = h \sum_{k=1}^{N-1} u(kh)v(kh)$$

with the corresponding norm $\|v\|_h$ and $\varphi_m(x) = \sin \pi mx, m = 1, \dots, N - 1$, form an orthogonal basis.

Also,

$$D_h^2 \varphi_m = -\lambda_m \varphi_m, \quad \lambda_m = \frac{2}{h^2} (\cos \pi mh - 1) = \frac{4}{h^2} \sin^2 \frac{\pi mh}{2},$$

where the eigenvalues satisfy

$$8 < \lambda_1 < \lambda_2 < \dots < \lambda_{N-1} < \frac{4}{h^2}.$$

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$$\frac{\partial u_h}{\partial t} = \sum_{m=1}^{N-1} \frac{da_m^h}{dt} \varphi_m, \quad D_h^2 u_h = - \sum_{m=1}^{N-1} a_m^h \lambda_m \varphi_m.$$

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Numerical Methods for PDE: Parabolic PDE



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The difference equation then gives

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If we assume $ck/h^2 \leq 1/2$, then $ck\lambda_m \leq ck(4/h^2) \leq 2$ and hence $|1 - ck\lambda_m| \leq 1$ for all m and the solution remains bounded. On the other hand, if $|1 - ck\lambda_m| > 1$ for some m , the initial data will increase exponentially.

Numerical Methods for PDE: Parabolic PDE



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where $(I + ckD_h^2)$ is a symmetric operator on $L(I_h)$. Thus,

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The stability result then yields

$$\|e^j\|_h \leq \left(\max_m |1 - ck\lambda_m| \right)^j \|e^0\|_h + Mk \max_j \|\ell^j\|_h \leq T \max_j \|\ell^j\|_h.$$