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Gravitational waves from topological defects

Any short subtitle

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Contents

Notation

Constants and units. We use ‘natural units’ where $\hbar = c = 1$, where \hbar is the reduced Planck constant and c is the speed of light in vacuum. **Planck units? Set $k_B = G_N = 1$?** The Newtonian constant of gravitation G_N is referenced explicitly, and we use Planck units such as the Planck mass $M_{\text{Pl}} = (\hbar c / G_N)^{1/2} = G_N^{-1/2} \sim 10^{-8} \text{ kg}$.

Tensors. The metric signature $(-, +, +, +)$ is considered, i.e. $\det[g_{\mu\nu}] \equiv |g| < 0$. The Minkowski metric is denoted $\eta_{\mu\nu}$, whereas a general metric is denoted $g_{\mu\nu}$. A four-vector $p^\mu =$

$$[\eta_{\mu\nu}] = \text{diag}(-1, 1, 1, 1)$$

Christophel symbols:

$$\Gamma_{\mu\nu}^\rho = \dots \quad (1)$$

“Lambda tensor”:

$$\Lambda_{ij,kl} = \dots \quad (2)$$

Fourier transforms. We use the following convention for the Fourier transform of $f(x)$, $\tilde{f}(k)$, and its inverse, where x and k are Lorentz four-vectors:

$$\begin{aligned} f(x) &= \int \frac{d^4k}{(2\pi)^4} e^{-ik \cdot x} \tilde{f}(k) \\ \tilde{f}(k) &= \int d^4x e^{ik \cdot x} f(x) \end{aligned} \quad (3)$$

Here, $k \cdot x = k_\sigma x^\sigma = g_{\rho\sigma} k^\rho x^\sigma$.

Acronyms

CDM	<u>c</u> old <u>d</u> ark <u>m</u> atter
CMB	<u>c</u> osmic <u>m</u> icrowave <u>b</u> ackground (radiation)
DW	<u>d</u> omain <u>w</u> all
GR	<u>g</u> eneral <u>r</u> elativity
GW	<u>g</u> ravitational <u>w</u> ave
Λ CDM	Lambda (Greek Λ) <u>c</u> old <u>d</u> ark <u>m</u> atter model; Standard cosmological model

Nomenclature

In the table below is listed the most frequently used symbols in this paper, for reference.

Table 1: helo

Symbol	Referent	SI-value or definition
<i>Natural constants</i>		
G_N	Newtonian constant of gravitation	1.2 kg
k_B	Boltzmann's constant	1.2 K
<i>Fiducial quantities</i>		
h_0	Reduced Hubble constant	0.67
<i>Subscripts</i>		
Q_{gw}	Quantity Q related to gravitational wave	
<i>Functions and operators</i>		
$\Theta(\xi)$	Heaviside step function	$\begin{cases} 1 & \xi > 0 \\ 0 & \xi < 0 \end{cases}$
$\text{sgn}(\xi)$	Signum function	$2\Theta(\xi) - 1$
$\delta^{(n)}(\xi)$	Dirac-Delta function of $\xi \in \mathbb{R}^n$, $n \in \mathbb{N}$.	
$\delta^{\mu\nu}$	Kronecker delta.	

Chapter 1

Introduction

- GOALS:
 - Gather framework about GWs from DWs

$$\tilde{h}_{\otimes}'' + 2\mathcal{H}\tilde{h}_{\otimes}' + k^2\tilde{h}_{\otimes} = 16\pi G_N a^2 \tilde{\sigma}_{\otimes}; \quad \otimes = +, \times \quad (1.1)$$

$$\left(\tilde{h}^{\text{TT}}\right)_{ij}(\eta, \mathbf{k}) = \sum_{\otimes=+, \times} e_{ij}^{\otimes}(\hat{\mathbf{k}}) \tilde{h}_{\otimes}(\eta, \mathbf{k}) \quad (1.2)$$

$$\tilde{h}^{\text{TT}}_{ij}(\eta, \mathbf{k}) = \sum_{\otimes=+, \times} e_{ij}^{\otimes}(\hat{\mathbf{k}}) \tilde{h}_{\otimes}(\eta, \mathbf{k}) \quad (1.3)$$

1.1 Preliminaries

- variational calculus/ varying action
- action
- pert. theory?
- line element
- gauge invariance
- FRW cosmology
- classical field theory

$$G_{\mu\nu} = 8\pi G_N T_{\mu\nu} \quad (1.4)$$

The action

$$S_{\text{ST}} = S_{\text{EH}} + S_{\phi} + S_{\text{m}} = \int d^4x \sqrt{-|g|} \left\{ \frac{M_{\text{Pl}}^2}{2} R - \frac{1}{2} \nabla_{\rho} \phi \nabla^{\rho} \phi - V(\phi) \right\} + S_{\text{m}} \quad (1.5)$$

GR as we know it is reconstructed when varying S_{E} with respect to the metric $g_{\mu\nu}(x)$

Part I

Background

Chapter 2

Classical Field Theory and Gravity

The action and blah blah

2.1 General Relativity

The Einstein–Hilbert action $\int d^4x$ in vacuum is $\int d^4x \mathcal{R}$ [check Planck mass def.](#)

$$S_{\text{EH}} = \frac{1}{2} M_{\text{Pl}}^2 \int d^4x \mathcal{R}, \quad (2.1)$$

where $\mathcal{R} = g^{\mu\nu} \mathcal{R}_{\mu\nu}$. By varying S_{EH} with respect to $g_{\mu\nu}$ one obtains the equation of motion

$$\mathcal{G}_{\mu\nu} \equiv \mathcal{R}_{\mu\nu} - \frac{1}{2} g_{\mu\nu} \mathcal{R} = 0. \quad (2.2)$$

Thus, we interpret GR *classical* field theory where the tensor field $g_{\mu\nu}$ is the gravitational field, $\int d^4x$ with the particle [realisation](#)? $\int d^4x$ named “graviton” $\int d^4x$.

- scalar field (ST theories)

2.2 Scalar-Tensor Theories?

2.3 Perturbation Theory

2.4 Classical Solitons

Chapter 3

Gravitational Waves

Chapter 4

Lattice- and N -body simulations

Part II

Project

Chapter 5

Calculating Gravitational Waves from Domain Walls

DRAFT

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Consider a planar domain wall in the xy -plane in a flat FRW universe, represented by a scalar field $\phi(\eta, \mathbf{x})$ and a potential $V(\phi)$. The background metric is

$$d\bar{s}^2 = \bar{g}_{\mu\nu} d\bar{x}^\mu d\bar{x}^\nu = -d\bar{t}^2 + a^2(t) \delta_{ij} d\bar{x}^i d\bar{x}^j = a^2(\eta) \{-d\eta^2 + d\mathbf{x}^2 + d\mathbf{y}^2 + d\mathbf{z}^2\}. \quad (5.1)$$

The solution to $\Box\phi = dV/d\phi$ is denoted $\bar{\phi}(\eta, z)$. We let indices $a, b, c = 1, 2$ and $i, j, k, l, \dots = 1, 2, 3$. Now we add a linear perturbation $\zeta(\eta, x^a)$ to the wall such that

$$\phi(\eta, \mathbf{x}) = \bar{\phi}(\eta, z; \zeta(\eta, x^a)) = \bar{\phi}(\eta, z; 0) + \zeta(\eta, x^a) \frac{\partial \bar{\phi}}{\partial z} \Big|_{z=0} + O(\zeta^2). \quad (5.2)$$

Furthermore, Fourier transforming [[←show this!](#)]_■ the spatial components gives

$$\phi(\eta, \mathbf{k}) = \int d^3x e^{ik_i x^i} \phi(\eta, \mathbf{x}) = \left[(2\pi)^2 \delta^{(2)}(k_a) - ik_3 \zeta(\eta, k_a) \right] \bar{\phi}(\eta, k_3; 0) + O(\zeta^2). \quad (5.3)$$

The TT-part of the energy-momentum tensor is [[←refer to some section](#)]_■

$$T_{ij}^{\text{TT}}(\eta, \mathbf{k}) = \Lambda_{ij,kl}(\hat{\mathbf{k}}) \int \frac{d^3p}{(2\pi)^3} p_k p_l \phi(\eta, \mathbf{p}) \phi(\eta, \mathbf{k} - \mathbf{p}). \quad (5.4)$$

We define a quantity t_{kl} by

$$T_{ij}^{\text{TT}}(\eta, \mathbf{k}) = \Lambda_{ij,kl}(\hat{\mathbf{k}}) \left(\frac{1}{2\pi} \cdot t_{kl}(\eta, \mathbf{k}) + O(\zeta^2) \right), \quad (5.5)$$

and the additional function

$$\mathfrak{I}_n(\eta, q_0) = \int_{\mathbb{R}} dq q^n \bar{\phi}(\eta, q; 0) \bar{\phi}(\eta, q_0 - q; 0). \quad (5.6)$$

After some manipulation [[←show this!](#)]_■, we get the following:

$$t_{ab}(\eta, \mathbf{k}) = k_a k_b [-i\zeta(\eta, k_c)] \mathfrak{I}_1(\eta, k_3) \quad (5.7a)$$

$$t_{a3}(\eta, \mathbf{k}) = k_a [-i\zeta(\eta, k_c)] \mathfrak{I}_2(\eta, k_3) \quad (5.7b)$$

$$t_{33}(\eta, \mathbf{k}) = k_3 [-i\zeta(\eta, k_c)] \mathfrak{I}_2(\eta, k_3) + (2\pi)^2 \delta^{(2)}(k_a) \mathfrak{I}_2(\eta, k_3) \quad (5.7c)$$

┐ There are some *small* constraint on the perturbation from this. Need to be commented! ┘

Gravitational waves sourced by this field is – to first order in ζ – given by

$$\begin{aligned} ah_{ij}(\eta, \mathbf{k}) &= \frac{16\pi G_N}{k} \int_{\eta_i}^{\eta} d\eta' \sin(k[\eta - \eta']) a(\eta') T_{ij}^{\text{TT}}(\eta', \mathbf{k}) \\ &= \frac{8G_N}{k} \Lambda_{ij,kl}(\hat{\mathbf{k}}) \int_{\eta_i}^{\eta} d\eta' \sin(k[\eta - \eta']) a(\eta') t_{kl}(\eta', \mathbf{k}) + O(\zeta^2). \end{aligned} \quad (5.8)$$

Remaining is the $\Lambda_{ij,kl}t_{kl}$ -elements, which in total are [6]_? terms per ij , due to symmetry in t_{kl} :

$$\begin{aligned}\Lambda_{ij,kl}(\hat{\mathbf{k}})t_{kl}(\eta, \mathbf{k}) &= 2 \left\{ \Lambda_{ij,12}t_{12} + \Lambda_{ij,13}t_{13} + \Lambda_{ij,23}t_{23} \right\} (\eta, k\hat{\mathbf{k}}) \\ &\quad + \left\{ \Lambda_{ij,11}t_{11} + \Lambda_{ij,22}t_{22} + \Lambda_{ij,33}t_{33} \right\} (\eta, k\hat{\mathbf{k}})\end{aligned}\tag{5.9}$$

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5.1 General Formalism

Chapter 6

Simulating Gravitational Waves from Domain Walls

Chapter 7

Studying Gravitational Waves from Domain Walls

Chapter 8

Discussion

Chapter 9

Conclusion and Outlook

Bibliography

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Bibliography

Appendix A

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