

3 Rayleigh-Benard Convection

Chandrasekhar⁴ has carefully studied linear stability of Rayleigh-Benard convection using the Boussinesq approximation, given in terms of the Rayleigh (Ra) and Prandtl (Pr) numbers as

$$\frac{\partial \mathbf{u}}{\partial t} + \mathbf{u} \cdot \nabla \mathbf{u} = -\nabla p + Pr \nabla^2 \mathbf{u} - Ra Pr T, \quad \frac{\partial T}{\partial t} + \mathbf{u} \cdot \nabla T = \nabla^2 T,$$

along with continuity, $\nabla \cdot \mathbf{u} = 0$. Above a certain critical Rayleigh number Ra_c , the conduction due to adverse temperature gradient becomes linearly unstable to small perturbations and convection rolls develop (Fig. 10 top). The table below shows predictions of Ra_c at the most unstable wavenumber k_c for $\Omega = [0 : 2\pi/k_c] \times [0 : 1]$ having periodic conditions in x . Dirichlet conditions $T = 1 - y$ are specified for temperature on the horizontal boundaries, and three different conditions are considered for velocity: both walls (Dirichlet-Dirichlet), both stress-free (Neumann-Neumann) and a mix (Dirichlet-Neumann).

| Critical Rayleigh number for 3 types of boundaries | | | | | | | |
|--|-----------|-----------|----------|---------|--------|--------|--------|
| BC | Ra_{12} | Ra_{23} | Ra_c^4 | k_c^4 | Ra_1 | Ra_2 | Ra_3 |
| D-D | 1707.75 | 1707.74 | 1707.76 | 3.117 | 1760 | 1740 | 1725 |
| D-N | 1100.71 | 1100.64 | 1100.65 | 2.682 | 1144 | 1122 | 1111 |
| N-N | 657.639 | 657.566 | 657.511 | 2.2214 | 690 | 680 | 670 |

According to dynamical systems theory, the saturation amplitude (U) and kinetic energy (\bar{E}_k) grow, respectively, as $\sqrt{\epsilon}$ and ϵ , for $\epsilon := (Ra - Ra_c)/Ra_c \ll 1$. Thus Ra_c can be determined from a linear fit of (volume-averaged steady state) \bar{E}_k versus Ra for two or more values of Ra as shown in Fig. 10 (left). In the table above, the estimates Ra_{12} and Ra_{23} are determined from solution pairs at (Ra_1, Ra_2) and (Ra_2, Ra_3) , respectively. For each case, E_k is computed in a single unsteady run (Fig. 10 right) by varying Ra after time marching to a steady state such that $|dE_k/dt| \leq \sigma \times \bar{E}_k$ with $\sigma = 10^{-4}$. Larger σ or lower polynomial orders $N < 7$ lead to a drop in accuracy for the estimates of Ra_c . Iteration tolerances were controlled by setting $TOLREL = 10^{-5}$. Larger values (10^{-3}) did not yield a clear initial linear stage with exponential growth but the Ra_c estimate was not affected.

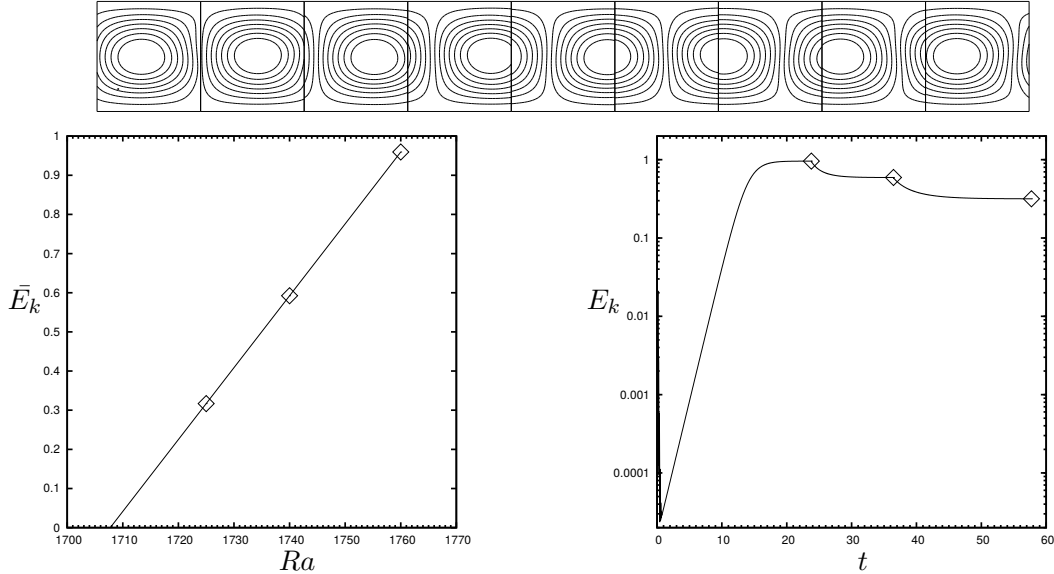


Figure 2: Rayleigh-Benard convection with walls: (top) streamlines for $E=9$ elements and $N=7$; (left) kinetic energy for $E=3$ and $N=7$ versus Rayleigh number and (right) time.

⁴S. Chandrasekhar, “Hydrodynamic and Hydromagnetic Stability,” Oxford University Press (1961)