

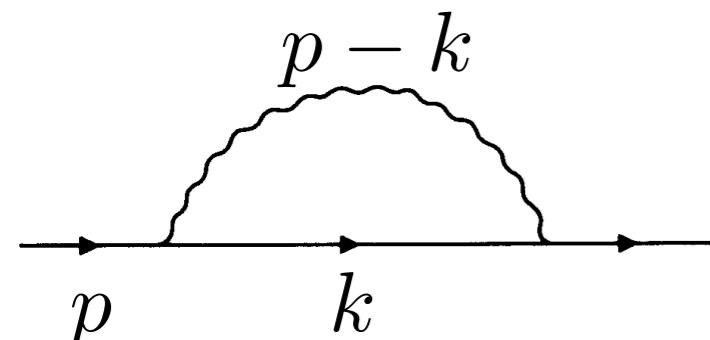
# Next-to-Leading Order with FeynRules

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Lund 2014

# Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

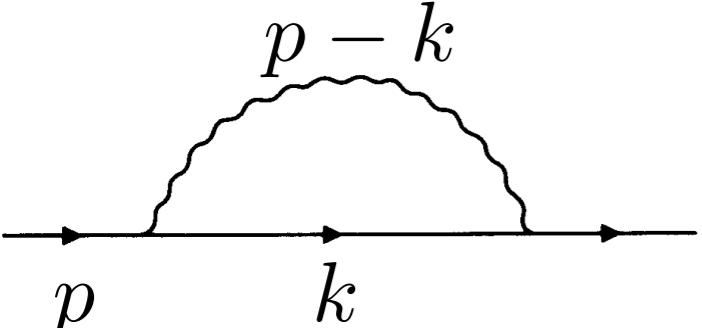
# Electron self-energy



A Feynman diagram illustrating the electron self-energy. It shows a horizontal incoming electron line labeled  $p$  entering from the left. This line splits into two parts: a lower straight line labeled  $k$  and an upper wavy line labeled  $p - k$ . The wavy line represents the virtual particle exchange between the electron and the external field.

$$= (-ie)^2 \int \frac{d^4 k}{(2\pi)^4} \gamma^\mu \frac{i(\cancel{k} + m)}{k^2 - m^2 + i\epsilon} \gamma^\nu \frac{-i\eta_{\mu\nu}}{(p - k)^2 + i\epsilon}$$

# Electron self-energy



Feynman diagram showing an incoming electron with momentum  $p$  (horizontal line) and an outgoing electron with momentum  $k$  (horizontal line). A wavy line representing the electron's interaction with its own field (self-energy) connects the two vertices.

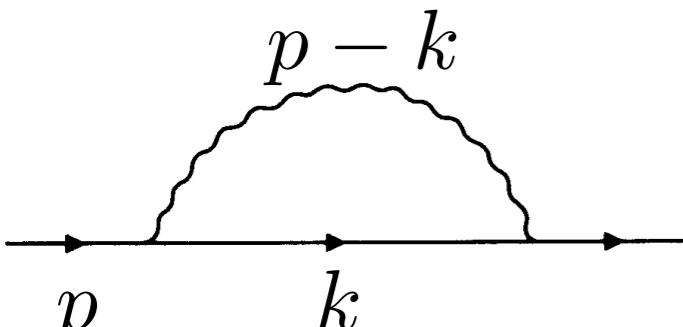
$$= (-ie)^2 \int \frac{d^4 k}{(2\pi)^4} \gamma^\mu \frac{i(\cancel{k} + m)}{k^2 - m^2 + i\epsilon} \gamma^\nu \frac{-i\eta_{\mu\nu}}{(p - k)^2 + i\epsilon}$$

Feynman parameter :

$$\frac{1}{AB} = \int_0^1 dx \frac{1}{(Ax + B(1 - x))}$$

$$= -e^2 \int_0^1 dx \int \frac{d^4 l}{(2\pi)^4} \frac{\gamma^\mu (\cancel{l} + x\cancel{p} + m) \gamma_\mu}{(l^2 - \Delta + i\epsilon)^2} \quad \leftarrow \begin{cases} l \equiv k - xp \\ \Delta = (1 - x)(m^2 - xp^2) \end{cases}$$

# Electron self-energy



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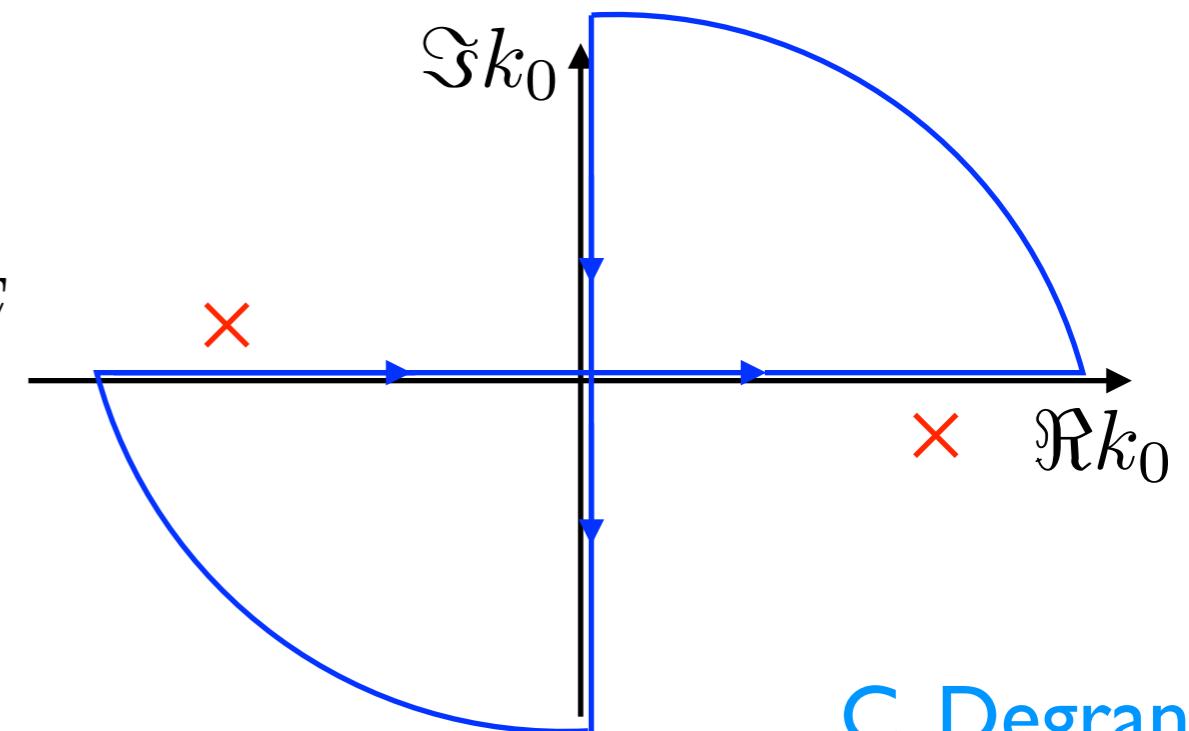
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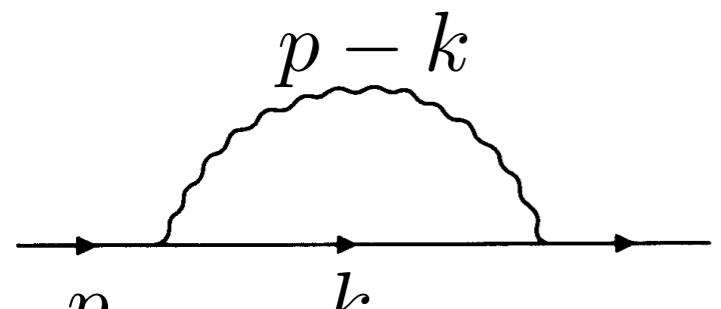
Minkowski to Euclidian space

$$\int_{-\infty}^{+\infty} dl_0 = - \int_{i\infty}^{-i\infty} dl_0 \quad dl_0 = i \int_{-\infty}^{+\infty} dl_0^E$$

$$l_0^E \equiv -il_0, \quad l^E \equiv 1$$



# Electron self-energy

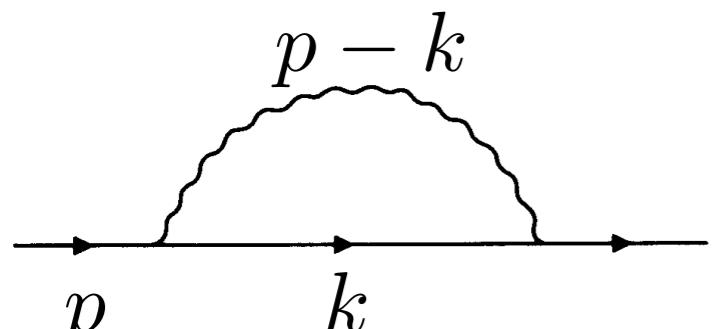


Feynman diagram illustrating the electron self-energy. A horizontal line with arrows at both ends represents an incoming electron with momentum  $p$ . A wavy line labeled  $p - k$  represents the virtual photon exchange. The outgoing electron has momentum  $k$ . The diagram is equated to the following expression:

$$= -ie^2 \int_0^1 dx \int \frac{d^4 l^E}{(2\pi)^4} \frac{\gamma^\mu (\not{k}^E + x \not{p} + m) \gamma_\mu}{((l^E)^2 + \Delta - i\epsilon)^2}$$

A blue arrow points from the term  $\not{k}^E$  to the text  $l^2 = -l_0^E - 1^E = -l^E$ , which is highlighted with a red box and an orange arrow.

# Electron self-energy



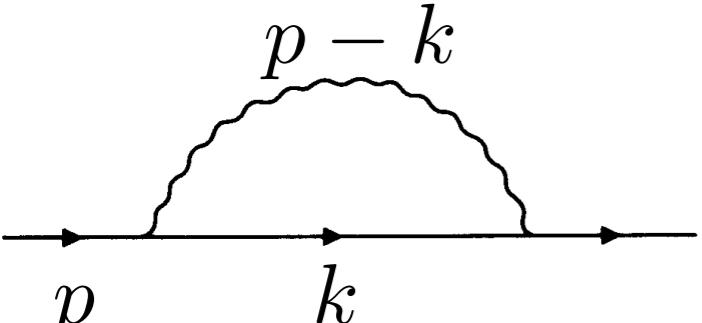
Feynman diagram showing an incoming electron with momentum  $p$  and an outgoing electron with momentum  $k$ . A wavy line representing a loop with momentum  $p - k$  connects the two. A blue arrow points from the loop to the outgoing momentum  $k$ , and another blue arrow points from the loop to the outgoing momentum  $0$ .

$$= -ie^2 \int_0^1 dx \int \frac{d^4 l^E}{(2\pi)^4} \frac{\gamma^\mu (\not{k}^E + x\not{p} + m) \gamma_\mu}{((l^E)^2 + \Delta - i\epsilon)^2}$$
$$l^2 = -l_0^{E2} - 1^{E2} = -l^{E2}$$

Dimensional regularization

$$= -ie^2 \int_0^1 dx \gamma^\mu (x\not{p} + m) \gamma_\mu \int \frac{d^d l^E}{(2\pi)^d} \frac{1}{((l^E)^2 + \Delta)^2}$$

# Electron self-energy

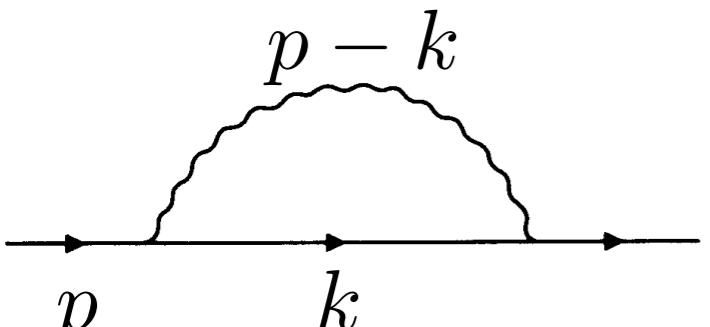

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$\uparrow$   
$$l^2 = -l_0^{E2} - l^{E2} = -l^{E2}$$

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## Dimensional regularization

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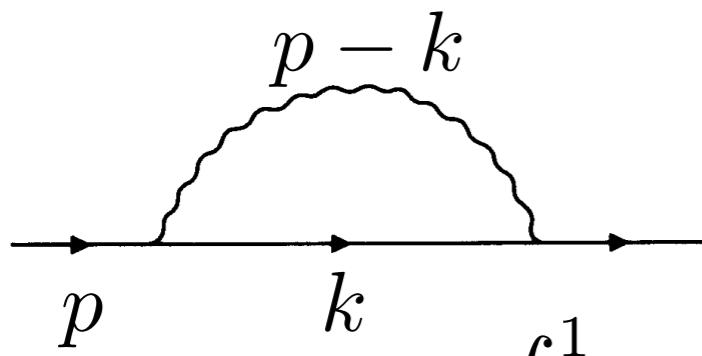
$$\boxed{\int \frac{d\Omega_d}{(2\pi)^d} \frac{1}{2} \int_0^\infty \frac{dl^{E^2} (l^{E^2})^{d/2-1}}{(l^{E^2} + \Delta)^2}} = \boxed{\frac{1}{(4\pi)^{d/2} \Gamma(d/2)}} \left(\frac{1}{\Delta}\right)^{2-d/2} \int_0^1 dz z^{1-d/2} (1-z)^{d/2-1}$$

$\uparrow$

$$z \equiv \frac{\Delta}{l^{E^2} + \Delta}$$

# Electron self-energy

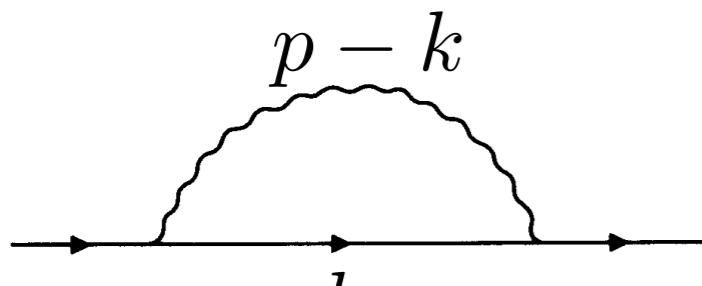
$$B(a, b) \equiv \frac{\Gamma(a)\Gamma(b)}{\Gamma(a+b)} \equiv \int dz z^{a-1} (1-z)^{b-1}$$



$$= -ie^2 \int_0^1 dx \gamma^\mu (x \not{p} + m) \gamma_\mu \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(2-d/2)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^{2-d/2}$$

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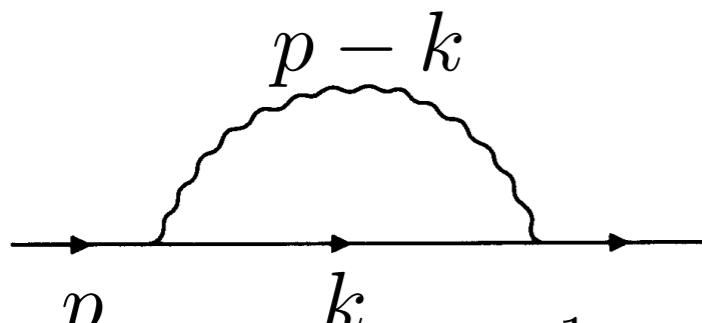
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$$\frac{1}{(4\pi)^2} \left( \frac{1}{\epsilon} - \gamma + \log(4\pi) - \log \Delta \right)$$

$\downarrow d \equiv 4 - 2\epsilon$

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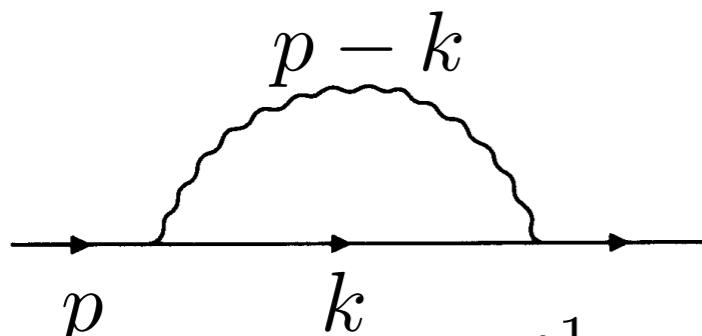
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$$\frac{1}{\bar{\epsilon}}$$

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$d \equiv 4 - 2\epsilon$

$$\frac{1}{\bar{\epsilon}}$$

$$= -i \frac{e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} (-\not{p} + 4m) + \text{finite}$$

# Renormalization

$$x_0 \rightarrow x + \delta x,$$

$$\phi_0 \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi + \sum_{\chi} \frac{1}{2}\delta Z_{\phi\chi}\chi.$$

$$\mathcal{L}_{Dirac} \rightarrow \bar{\psi} (\cancel{Q} - m) (1 + \delta Z_{\psi\psi}) \psi - \delta m \bar{\psi} \psi$$

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real constants

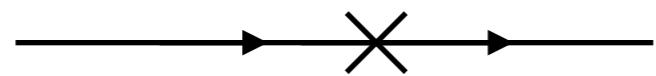
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**real constants**

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$$i (\cancel{p} - m) \delta Z_{\psi\psi} - i \delta m$$

# Renormalization

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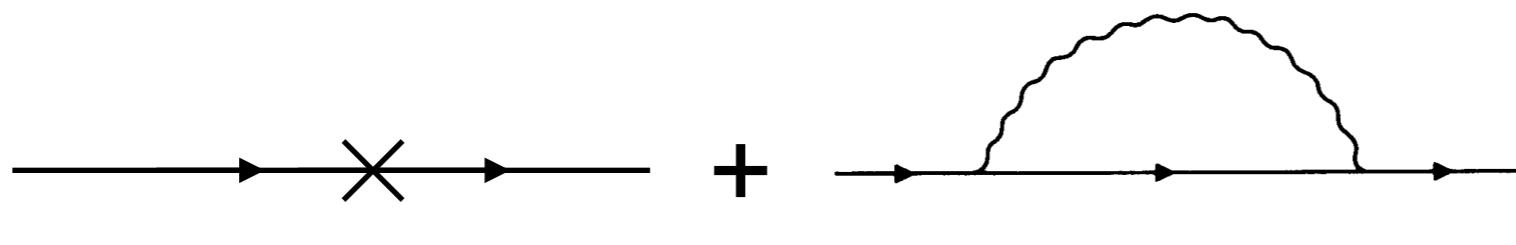
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$$= 0 \quad \frac{1}{\bar{\epsilon}} \quad + \text{finite}$$

$$\delta Z_{\psi\psi} = -\frac{e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} + \dots$$

$$\delta m = -\frac{3e^2}{(4\pi)^2} \frac{1}{\bar{\epsilon}} + \dots$$

# Renormalization

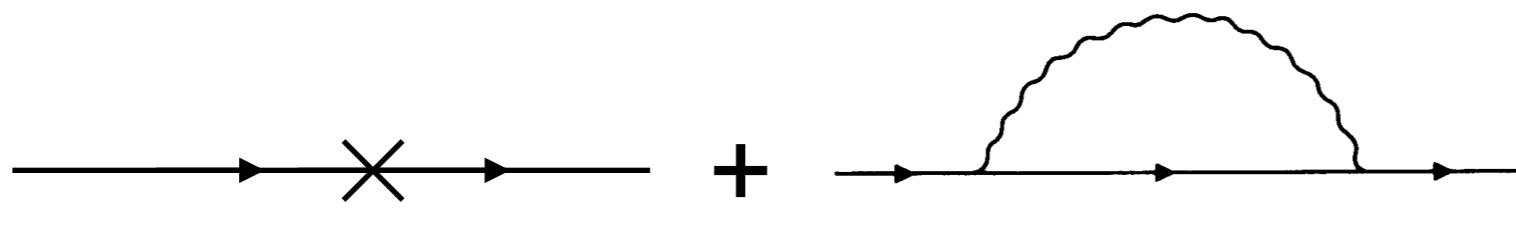
$$x_0 \rightarrow x + \boxed{\delta x}, \quad \text{real constants}$$

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$$\mathcal{L}_{Dirac} \rightarrow \bar{\psi} (\cancel{D} - m) (1 + \delta Z_{\psi\psi}) \psi - \delta m \bar{\psi} \psi$$



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**Arbitrary:  
Renormalization  
scheme**  
C. Degrande

# Operator dimension

- Fermion fields : 3/2
- Boson fields : 1
- derivatives : 1

Dimension 4

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi \partial^\mu \phi - m^2 \phi^2) \quad \mathcal{L} = \bar{\psi}(x) [i\cancel{D} - m] \psi(x)$$

$$D_\mu = \partial_\mu - ieQ A_\mu$$

# Operator dimension

- Fermion fields : 3/2
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Dimension 4

The diagram illustrates the concept of operator dimension. A central teal arrow points to the text "Dimension 4". Two other teal arrows point from this central label to two separate Lagrangian terms. The first term is the Lagrangian for a scalar field,  $\mathcal{L} = \frac{1}{2} (\partial_\mu \phi \partial^\mu \phi - m^2 \phi^2)$ . The second term is the Lagrangian for a fermion field,  $\mathcal{L} = \bar{\psi}(x) [i\cancel{D} - m] \psi(x)$ .

$$\mathcal{L} = \frac{1}{2} (\partial_\mu \phi \partial^\mu \phi - m^2 \phi^2)$$
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**Definition** : dimension of the operator is the sum of the dimensions of its fields and derivatives

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$d=2$

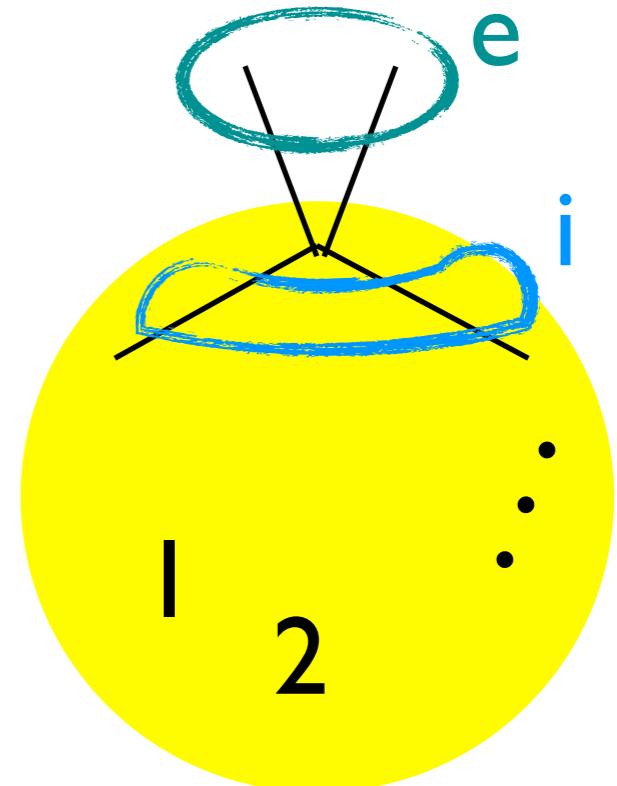
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**Definition** : dimension of the operator is the sum of the dimensions of its fields and derivatives

Exclude the dimension of the coefficient!

# Divergent amplitudes

Each vertex:  $p^{d-e-i}$



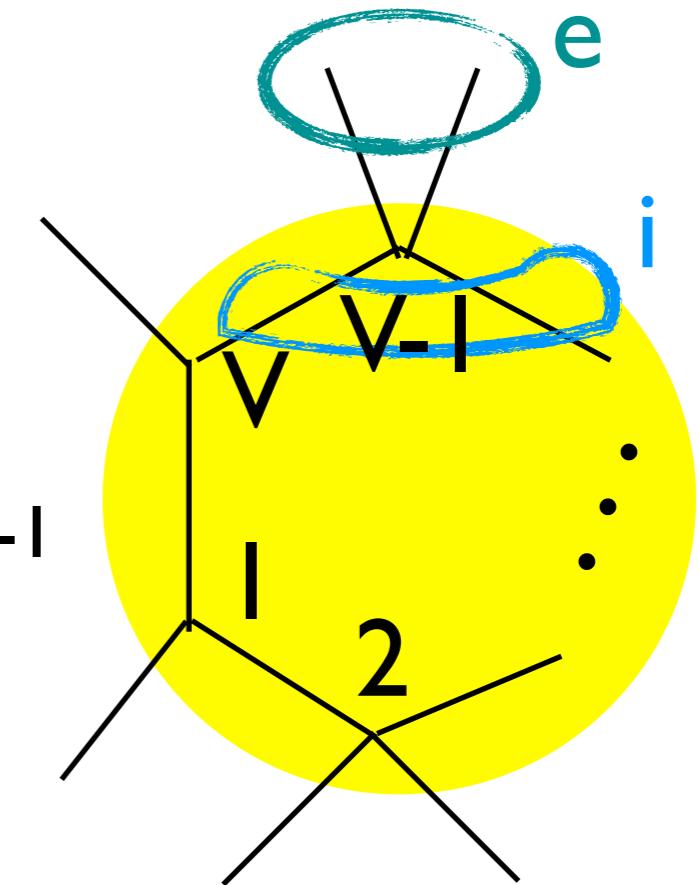
# Divergent amplitudes

Each vertex:  $p^{d-e-i}$

One loop:  $V$  propagators

Each internal fermion propagator:  $p^{-1}$

Each internal boson propagator:  $p^{-2}$



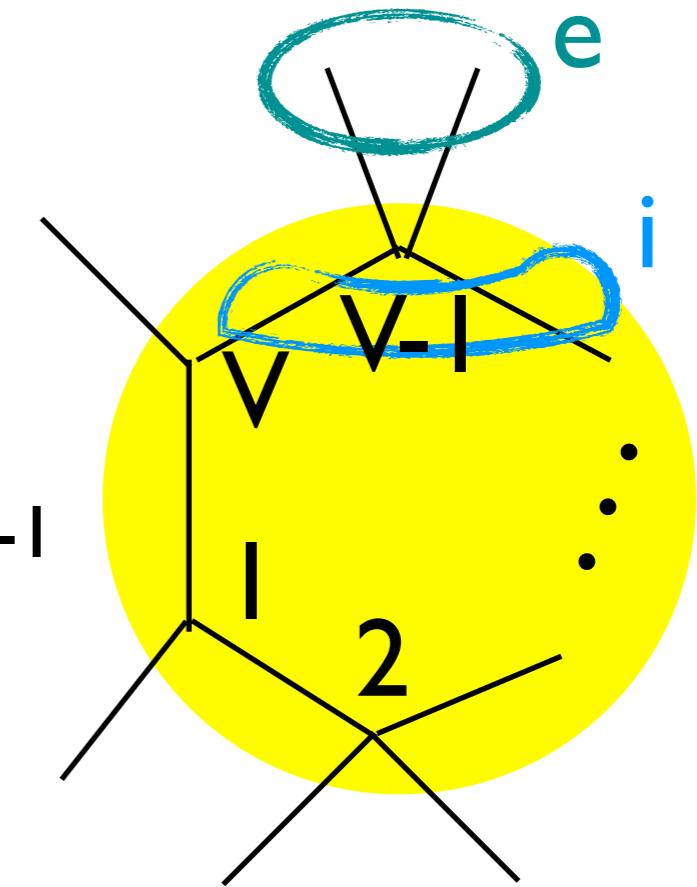
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$$I = \sum i = 3F + 2B \quad (F/B \# \text{ of internal fermions/bosons})$$

$$F+B=V$$

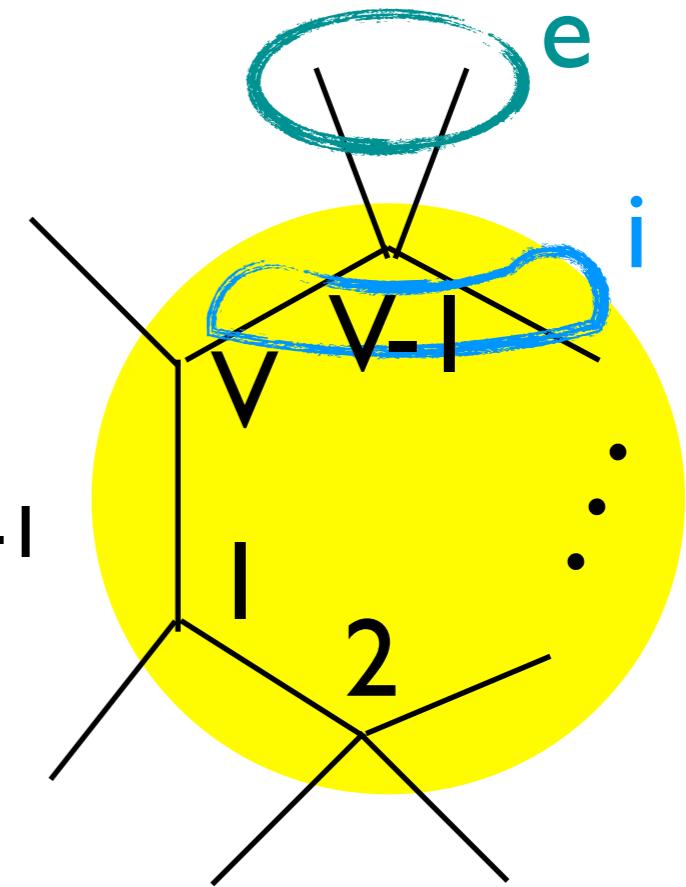
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Each extra loop: one extra propagator

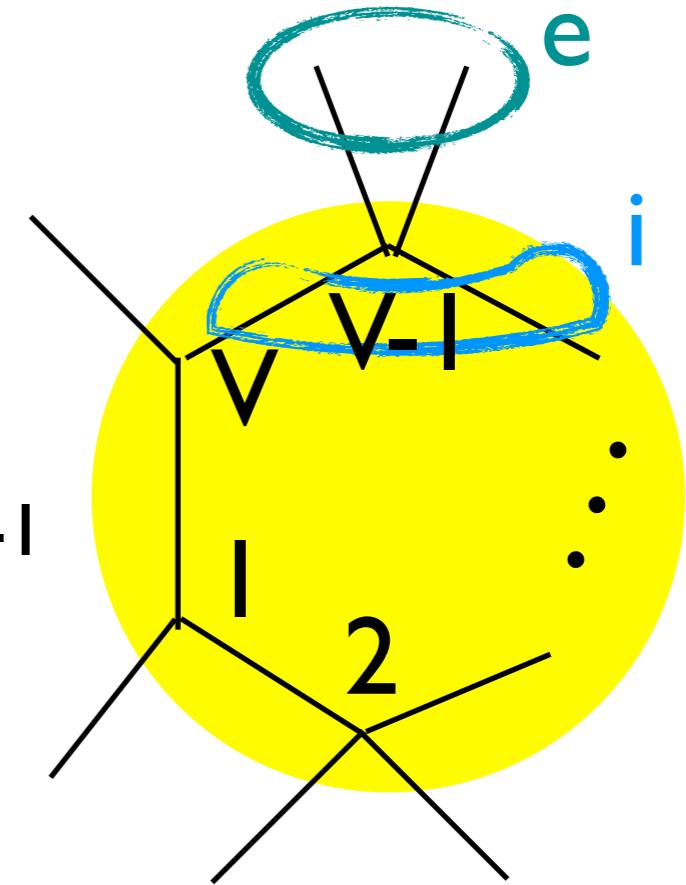
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$$F+B=V+(L-I)$$

Each extra loop: one extra propagator

$$A(D, E, L, V) \sim \left( \int d^4 p \right)^L p^{D-E-I-2B-F} \sim \left( \int d^4 p \right)^L p^{D-E-4V-4(L-1)}$$

# Renormalization

$$A(D, E, L, V) \sim \left( \int d^4 p \right)^L p^{D-E-4V-4(L-1)}$$

Renormalizable model if  $d \leq 4$  for all operators

$D \leq 4V : A \sim \Lambda^{4-E}$  diverges only if  $E \leq 4$

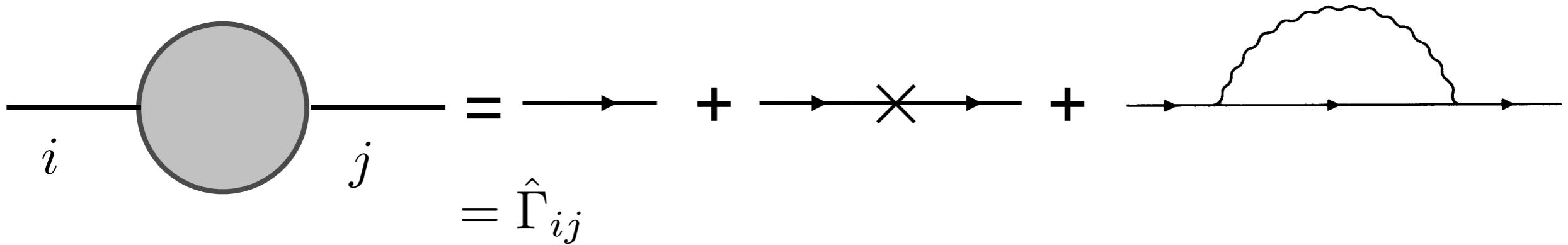
- All the divergence are absorbed by the fields and external parameters redefinitions
- Yang-mills theories
- With spontaneous symmetry breaking

# On-shell scheme

Renormalized mass = Physical mass

Two-point function vanishes on-shell (No loop corrections on the external legs)

Cancel the mixings



$$\tilde{\Re} \hat{\Gamma}_{ij}^f(p) u_j(p) \Big|_{p^2 = m_{f,j}^2} = 0,$$

$$\lim_{p^2 \rightarrow m_{f,i}^2} \frac{p + m_{f,i}}{p^2 - m_{f,i}^2} \tilde{\Re} \hat{\Gamma}_{ii}^f(p) u_i(p) = i u_i(p),$$

# Renormalization conditions

$$\tilde{\Re} \bar{u}_i(p') \hat{\Gamma}_{ij}^f(p') \Big|_{p'^2=m_{f,i}^2} = 0,$$

$$\lim_{p'^2 \rightarrow m_{f,i}^2} \bar{u}_i(p') \tilde{\Re} \hat{\Gamma}_{ii}^f(p') \frac{\not{p}' + m_{f,i}}{p'^2 - m_{f,i}^2} = i\bar{u}_i(p').$$

$$\hat{\Gamma}_{ij} = i\delta_{ij}(\not{p} - m_i) + i[f_{ij}^L(p^2)\not{p}\gamma_- + f_{ij}^R(p^2)\not{p}\gamma_+ + f_{ij}^{SL}(p^2)\gamma_- + f_{ij}^{SR}(p^2)\gamma_+]$$

$$\tilde{\Re} [f_{ij}^L(p^2)m_i + f_{ij}^{SR}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\tilde{\Re} [f_{ij}^R(p^2)m_i + f_{ij}^{SL}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\tilde{\Re} \left[ 2m_i \frac{\partial}{\partial p^2} [(f_{ii}^L(p^2) + f_{ii}^R(p^2))m_i + f_{ii}^{SL}(p^2) + f_{ii}^{SR}(p^2)] + f_{ii}^L(p^2) + f_{ii}^R(p^2) \right] \Big|_{p^2=m_i^2} = 0$$

Similar for the vectors and scalars

# Electron renormalization

$$f^L(p^2) = f^R(p^2) = \delta Z_{\psi\psi} - e^2 \int_0^1 dx (2-d)x \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

$$f^{SL}(p^2) = f^{SR}(p^2) = -m\delta Z_{\psi\psi} - \delta m - e^2 \int_0^1 dx \frac{m d}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

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Gamma algebra in d dimension!

# Electron renormalization

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$$\delta m = -\frac{e^2}{(4\pi)^2} \left( \frac{3}{\bar{\epsilon}} + 4 - 6 \log \left( \frac{m}{\mu} \right) \right)$$

Eqs I and 2

# Electron renormalization

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Eqs I and 2

$$\Delta = (1-x)(m^2 - xp^2)$$

$$\frac{\partial \Delta^{-\epsilon}}{\partial p^2} \Big|_{p^2=m^2} = \epsilon x(1-x)\Delta^{-\epsilon-1} \Big|_{p^2=m^2} = \epsilon x(1-x)^{-1-2\epsilon} m^{-2-2\epsilon}$$

# Electron renormalization

$$f^L(p^2) = f^R(p^2) = \delta Z_{\psi\psi} - e^2 \int_0^1 dx (2-d)x \frac{1}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

$$f^{SL}(p^2) = f^{SR}(p^2) = -m\delta Z_{\psi\psi} - \delta m - e^2 \int_0^1 dx \frac{m d}{(4\pi)^{d/2}} \frac{\Gamma(\epsilon)}{\Gamma(2)} \left(\frac{1}{\Delta}\right)^\epsilon$$

$$\delta m = -\frac{e^2}{(4\pi)^2} \left( \frac{3}{\bar{\epsilon}} + 4 - 6 \log \left( \frac{m}{\mu} \right) \right)$$

**Eqs I and 2**

$$\Delta = (1-x)(m^2 - xp^2)$$

$$\frac{\partial \Delta^{-\epsilon}}{\partial p^2} \Big|_{p^2=m^2} = \epsilon x(1-x)\Delta^{-\epsilon-1} \Big|_{p^2=m^2} = \epsilon x(1-x)^{-1-2\epsilon} m^{-2-2\epsilon}$$

After integration over x

$$\frac{\epsilon \Gamma(2) \Gamma(-2\epsilon)}{\Gamma(2-2\epsilon)}$$

$$\delta Z_{\psi\psi} = -\frac{e^2}{(4\pi)^2} \left( \frac{1}{\bar{\epsilon}_{UV}} + \frac{2}{\bar{\epsilon}_{IR}} + 4 - 6 \log \left( \frac{m}{\mu} \right) \right)$$

# Real/Complex masses

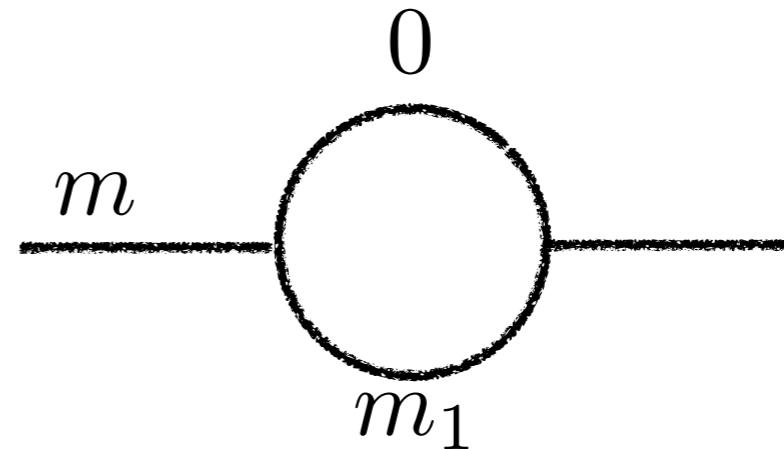
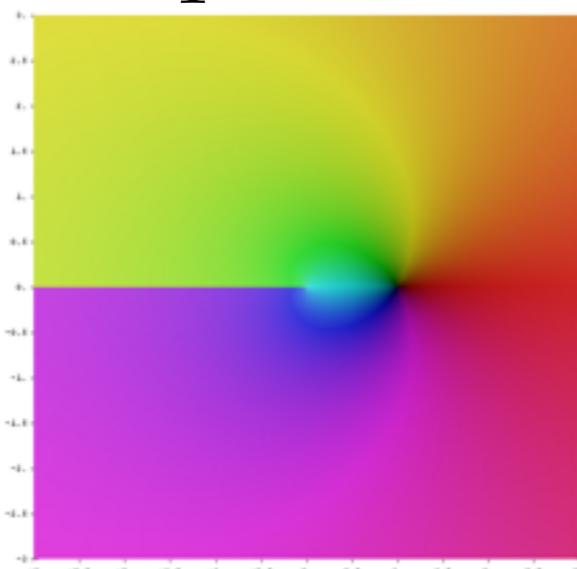
Real masses

$$m_1^2 > m^2$$

$$\Re (\log [m_1^2 - p^2]) \Big|_{p^2=m^2}$$

$$m_1^2 < m^2$$

$$\Re (\log [p^2 - m_1^2] + i\pi) \Big|_{p^2=m^2}$$



Mass corrections are complex if the particle can decay

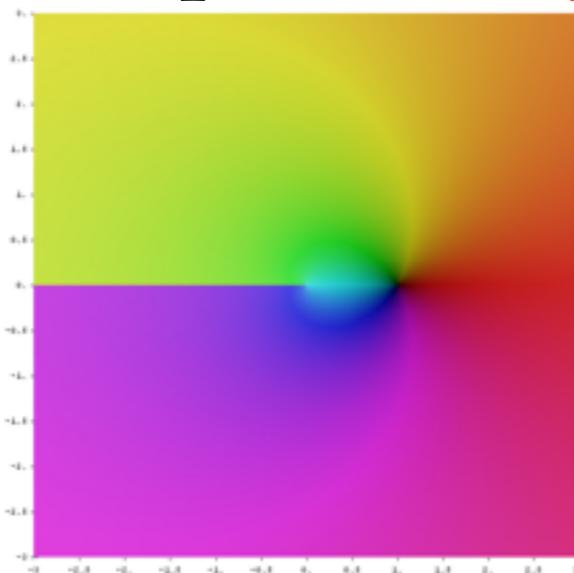
Epsilon prescription of the propagator

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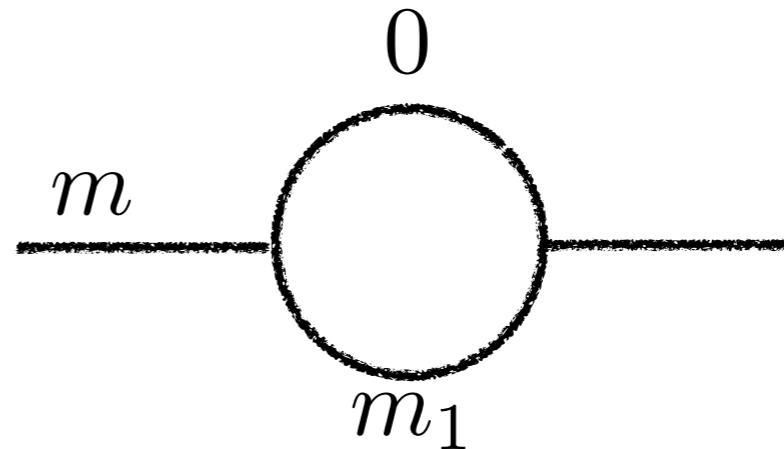
Real masses

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$$m \in \mathbb{R}$$



Mass corrections are complex if the particle can decay

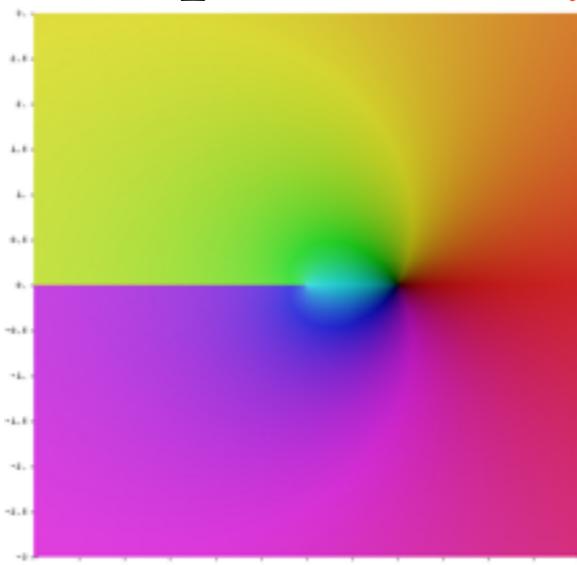
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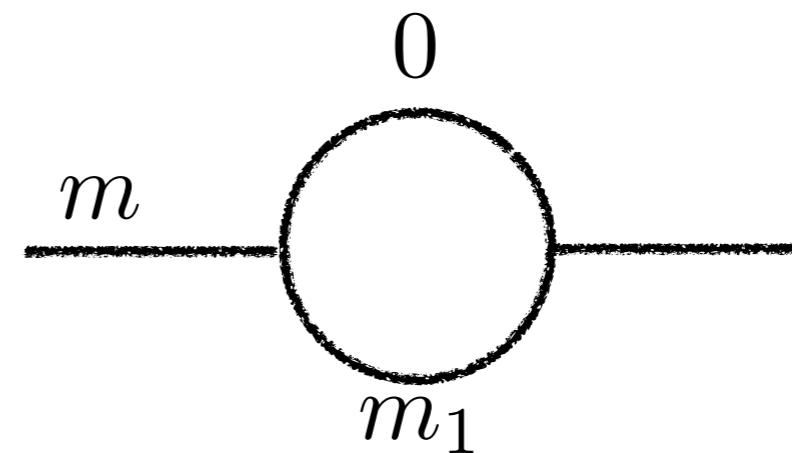
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Mass corrections are complex if the particle can decay

Epsilon prescription of the propagator

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2}$$

$$m^2 \rightarrow m^2 - im\Gamma$$

# Complex mass scheme

$$\left. \begin{array}{l} \phi_0 \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi \\ \phi_0^\dagger \rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi}^*)\phi^\dagger \end{array} \right\} \rightarrow \partial^\mu \phi_0 \partial_\mu \phi_0^\dagger \rightarrow (1 + \Re \delta Z_{\phi\phi}) \partial^\mu \phi \partial_\mu \phi^\dagger$$

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$$m \rightarrow m + \delta m$$

Now with complex masses

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Now with complex masses

Hermitian

$$\mathcal{L}_0 = \mathcal{L} + \delta \mathcal{L}$$

Not hermitian

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$$\cancel{\Re [f_{ij}^L(p^2) m_i + f_{ij}^{SR}(p^2)]} \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\Re [f_{ij}^R(p^2) m_i + f_{ij}^{SL}(p^2)]} \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\Im \left[ 2m_i \frac{\partial}{\partial p^2} [(f_{ii}^L(p^2) + f_{ii}^R(p^2)) m_i + f_{ii}^{SL}(p^2) + f_{ii}^{SR}(p^2)] + f_{ii}^L(p^2) + f_{ii}^R(p^2) \right]} \Big|_{p^2=m_i^2} = 0$$

Similar for the vectors and scalars

# Renormalization conditions

Zero momentum scheme available for the gauge couplings

$$\begin{aligned}\Gamma_{FFV}^\mu(p_1, p_2) = & igT^a \delta_{f_1, f_2} \left[ \gamma^\mu \left( \frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} \right) \right. \\ & + \gamma^\mu \gamma_5 \left( \frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} \right) \\ & \left. + \left( \gamma^\mu h^V(k^2) + \gamma^\mu \gamma_5 h^A(k^2) + \frac{(p_1 - p_2)^\mu}{2m} h^S(k^2) + \frac{k_\mu}{2m} h^P(k^2) \right) \right]\end{aligned}$$



$$\begin{aligned}\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} + h^V(0) + h^S(0) &= 0 \\ \frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} + h^A(0) &= 0.\end{aligned}$$

By gauge invariance



$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{g'_V}{2g} \delta Z_{V'V} + \frac{g'_A}{2g} \delta Z_{V'V} = 0$$

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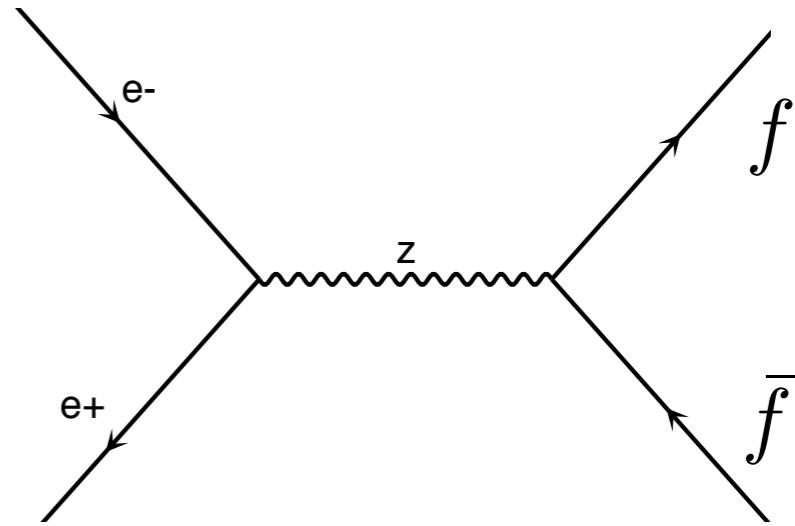
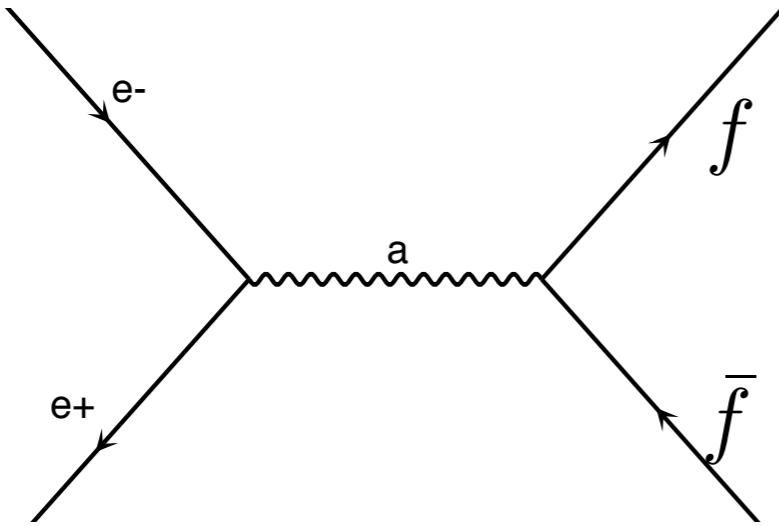
$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{g'_V}{2g} \delta Z_{V'V} + \frac{g'_A}{2g} \delta Z_{V'V} = 0$$

Only from  
two-point  
functions

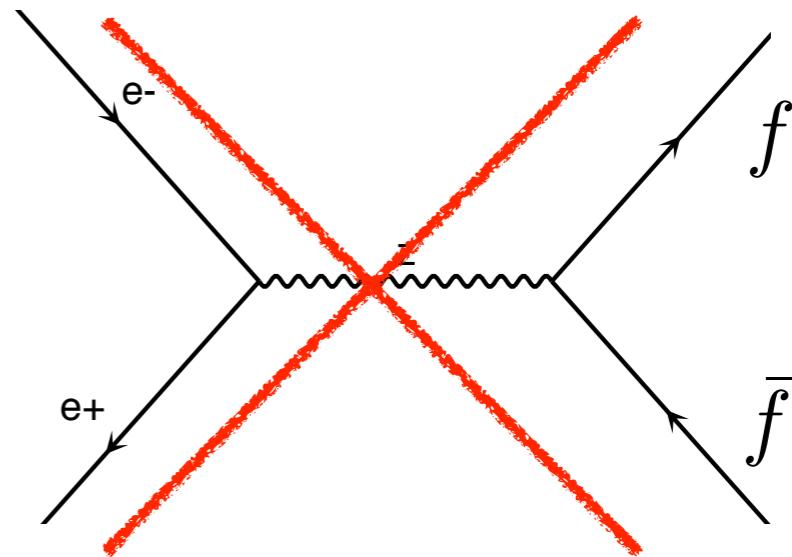
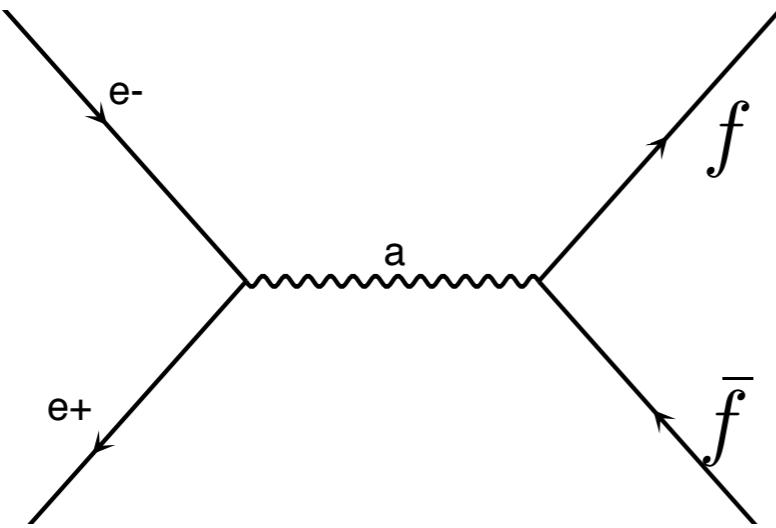
# Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

# Infrared divergences

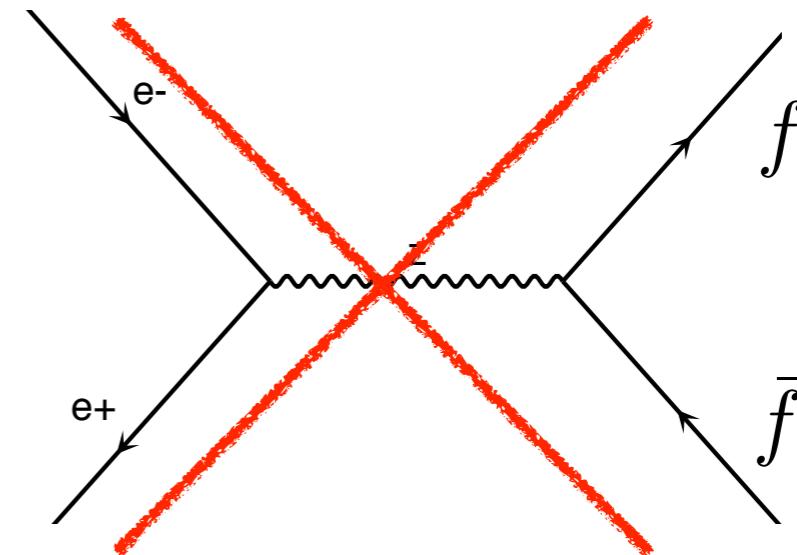
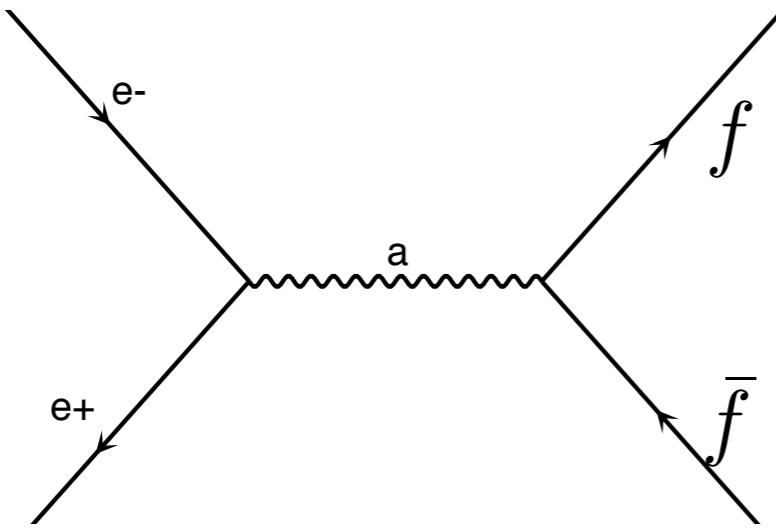


# Infrared divergences



Well below the  $Z$  mass

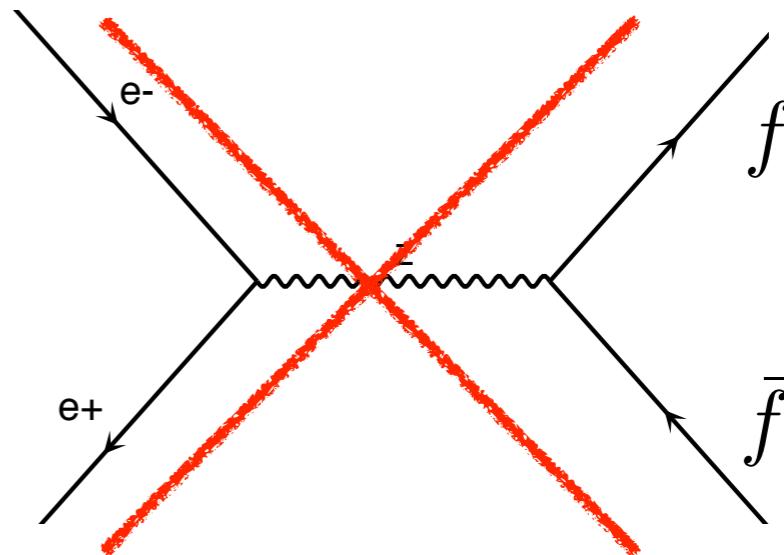
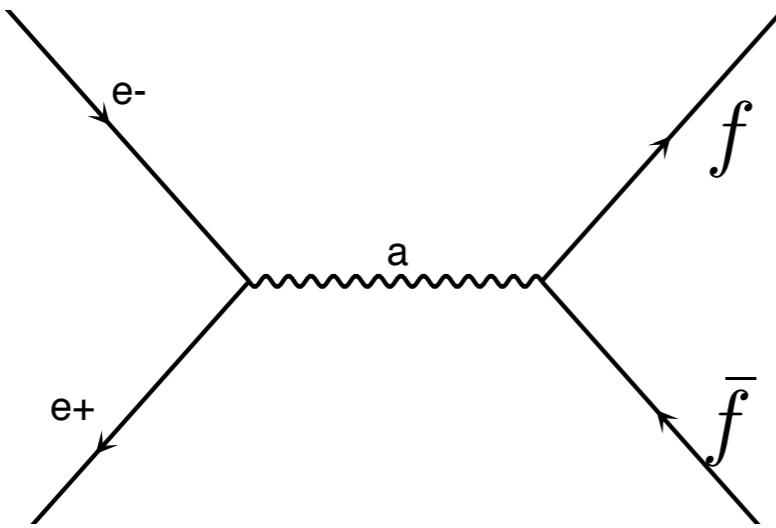
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Well below the  $Z$  mass

$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2$$

# Infrared divergences



Well below the  $Z$  mass

$$\sigma_0 = \frac{4\pi\alpha^2}{3s} Q_f^2$$

$$R \equiv \frac{\sigma(e^+e^- \rightarrow \text{hadrons})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = \frac{\sum_q \sigma(e^+e^- \rightarrow q\bar{q})}{\sigma(e^+e^- \rightarrow \mu^+\mu^-)} = 3 \sum_q Q_q^2$$

# KLN illustration

## Virtual corrections

$$\sigma^V = \sigma_0 \frac{2\alpha_S}{\pi} H(\epsilon) \left[ -\frac{2}{\epsilon^2} + \frac{3}{\epsilon} - 8 + \pi^2 + \mathcal{O}(\epsilon) \right]$$

where  $H(\epsilon) = \frac{3(1+\epsilon)}{(3+2\epsilon)\Gamma(2+2\epsilon)(4\pi)^{2\epsilon}} = 1 + \mathcal{O}(\epsilon)$

## Real corrections

$$\sigma^{q\bar{q}g} = \sigma_0 \frac{2\alpha_S}{\pi} H(\epsilon) \left[ \frac{2}{\epsilon^2} - \frac{3}{\epsilon} + \frac{19}{2} - \pi^2 + \mathcal{O}(\epsilon) \right]$$

$$R = 3 \sum_q Q_q^2 \left\{ 1 + \frac{\alpha_S}{\pi} + \mathcal{O}(\alpha_S^2) \right\}$$

Pole cancellation provides a check

C. Degrande

# Plan

- Renormalization
- KLN Theorem
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# Loop computation

$$\begin{aligned}\mathcal{A}^{1-loop} = & \sum_i \textcolor{red}{d}_i \text{ Box}_i + \sum_i \textcolor{red}{c}_i \text{ Triangle}_i + \sum_i \textcolor{red}{b}_i \text{ Bubble}_i \\ & + \sum_i \textcolor{red}{a}_i \text{ Tadpole}_i + \textcolor{magenta}{R}\end{aligned}$$

- Box, Triangle, Bubble and Tadpole are known scalar integrals
- Loop computation = find the coefficients
  - Unitarity
  - Multiple cuts
  - Tensor reduction (OPP)

# R<sub>2</sub>

$$\bar{A}(\bar{q}) = \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\bar{N}(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}, \quad \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2$$

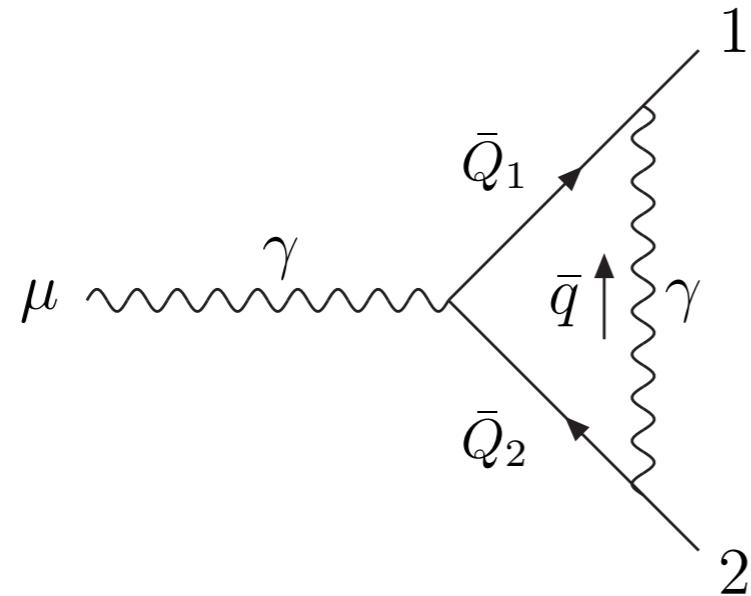
$$\bar{N}(\bar{q}) = N(q) + \tilde{N}(\tilde{q}, q, \epsilon)$$

d      4       $\epsilon$

$$R_2 \equiv \lim_{\epsilon \rightarrow 0} \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\tilde{N}(\tilde{q}, q, \epsilon)}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}$$

Finite set of vertices that can be computed once  
for all

# R<sub>2</sub> example



$$\begin{aligned}\bar{Q}_1 &= \bar{q} + p_1 = Q_1 + \tilde{q} \\ \bar{Q}_2 &= \bar{q} + p_2 = Q_2 + \tilde{q}\end{aligned}$$

$$\bar{D}_0 = \bar{q}^2$$

$$\bar{D}_1 = (\bar{q} + p_1)^2$$

$$\bar{D}_2 = (\bar{q} + p_2)^2$$

**'t Hooft-Veltman scheme**

$$\bar{\eta}^{\bar{\mu} \bar{\nu}} \bar{\eta}_{\bar{\mu} \bar{\nu}} = d,$$

$$\bar{\gamma}^{\bar{\mu}} \bar{\gamma}_{\bar{\mu}} = d \mathbb{1},$$

$$\begin{aligned}\bar{N}(\bar{q}) &\equiv e^3 \left\{ \bar{\gamma}_{\bar{\beta}} (\bar{Q}_1 + m_e) \gamma_{\mu} (\bar{Q}_2 + m_e) \bar{\gamma}^{\bar{\beta}} \right\} \\ &= e^3 \left\{ \gamma_{\beta} (Q_1 + m_e) \gamma_{\mu} (Q_2 + m_e) \gamma^{\beta} \right. \\ &\quad \left. - \epsilon (Q_1 - m_e) \gamma_{\mu} (Q_2 - m_e) + \epsilon \tilde{q}^2 \gamma_{\mu} - \tilde{q}^2 \gamma_{\beta} \gamma_{\mu} \gamma^{\beta} \right\}\end{aligned}$$

$$R_2 = -\frac{ie^3}{8\pi^2} \gamma_{\mu}$$

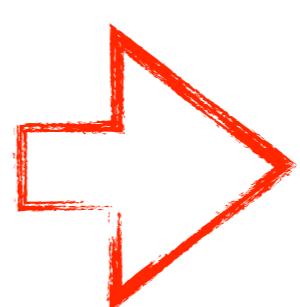
$$\begin{aligned}\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_0 \bar{D}_1 \bar{D}_2} &= -\frac{i\pi^2}{2} + \mathcal{O}(\epsilon), \\ \int d^n \bar{q} \frac{q_{\mu} q_{\nu}}{\bar{D}_0 \bar{D}_1 \bar{D}_2} &= -\frac{i\pi^2}{2\epsilon} g_{\mu\nu} + \mathcal{O}(1)\end{aligned}$$

# R<sub>I</sub>

Due to the  $\epsilon$  dimensional parts of the denominators

Like for the 4 dimensional part but with a different set of integrals

$$\begin{aligned}\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j} &= -\frac{i\pi^2}{2} \left[ m_i^2 + m_j^2 - \frac{(p_i - p_j)^2}{3} \right] + \mathcal{O}(\epsilon), \\ \int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j \bar{D}_k} &= -\frac{i\pi^2}{2} + \mathcal{O}(\epsilon), \\ \int d^n \bar{q} \frac{\tilde{q}^4}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_l} &= -\frac{i\pi^2}{6} + \mathcal{O}(\epsilon).\end{aligned}$$

Only  $R = R_I + R_2$  is gauge invariant  Check

# Plan

- Renormalization
- KLN Theorem
- Rational Terms
- FeynRules at NLO

- Goal : Automate the one-loop computation for BSM models
- Required ingredients :
  - Tree-level vertices
  - R2 vertices (OPP)
  - UV counterterm vertices
- Solution : UFO at NLO

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# UV counterterm vertices

## External parameters

$$\begin{aligned}x_0 &\rightarrow x + \delta x, \\ \phi_0 &\rightarrow (1 + \frac{1}{2}\delta Z_{\phi\phi})\phi + \sum_{\chi} \frac{1}{2}\delta Z_{\phi\chi}\chi.\end{aligned}$$

Same for the conjugate field

One renormalization constant for each fermion chirality

Internal parameters are renormalised by replacing the external parameters in their expressions

$$\mathcal{L}_0 = \mathcal{L} + \delta\mathcal{L} \quad \text{vertices after solving the reno. cond.}$$

# How does it work?

**FeynRules**

Renormalize the Lagrangian

model.mod

model.gen

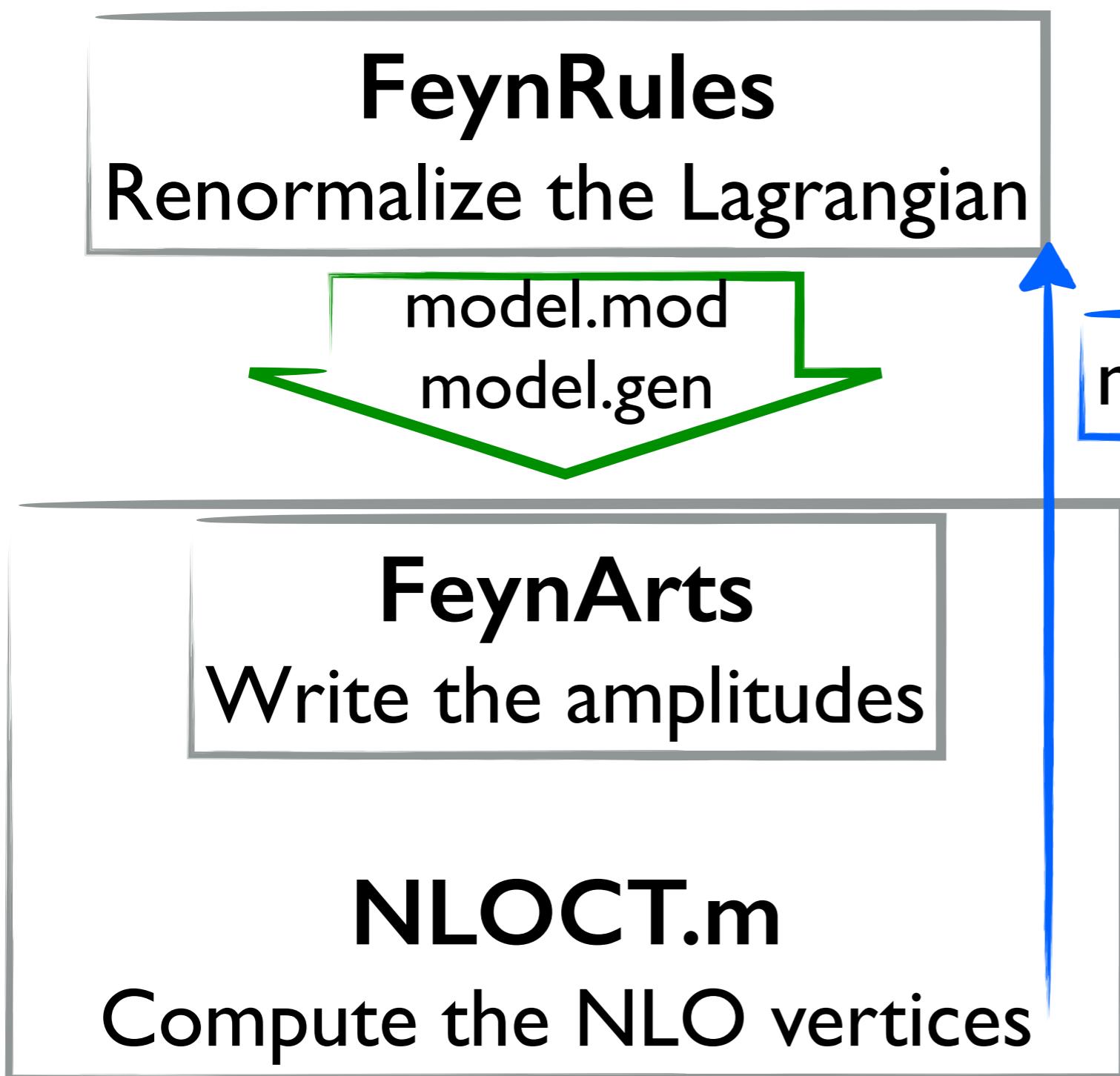
**FeynArts**

Write the amplitudes

**NLOCT.m**

Compute the NLO vertices

# How does it work?



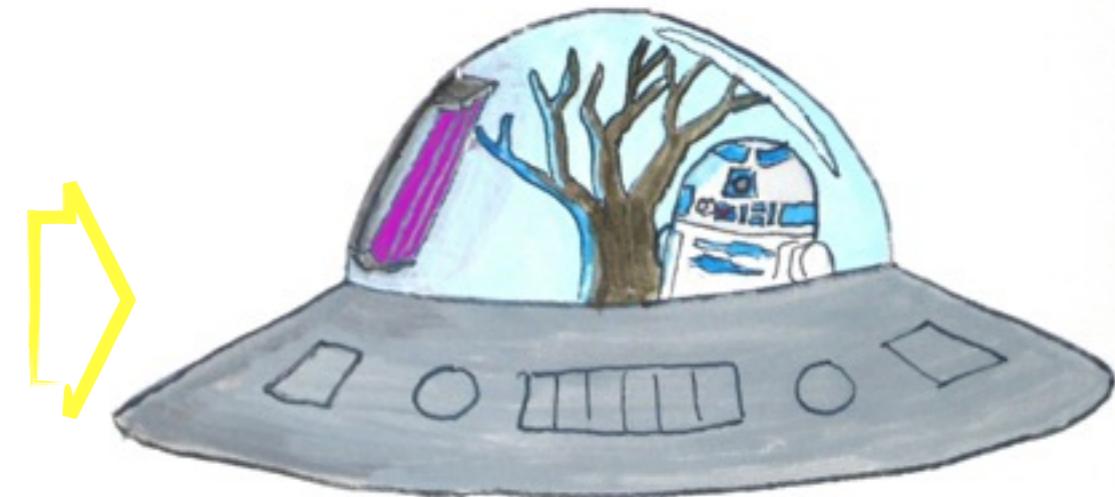
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**FeynRules**  
Renormalize the Lagrangian

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Compute the NLO vertices



model.nlo

# How does it work?

FeynRules :

...

```
Lren = OnShellRenormalization[ LSM , QCDOOnly ->True];
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",
GenericFile -> False]
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FeynRules :

...

```
Get["SMQCDreno.nlo"];
WriteUFO[ LSM , UVCounterterms -> UV$vertlist ,
R2Vertices -> R2$vertlist]
```

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...

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```
Get["SMQCDreno.nlo"];
WriteUFO[ LSM , UVCounterterms -> UV$vertlist ,
R2Vertices -> R2$vertlist]
```

# model.nlo

## Model information (FR+FeynArts model/generic files)

```
R2$vertlist = {  
{{{{anti[u], 1}, {u, 2}}, ((-I/12)*gs^2*  
IndexDelta[Index[Colour, Ext[1]], Index[Colour, Ext[2]]]*IPL[{u, G}]*  
(TensDot[SlashedP[2], ProjM][Index[Spin, Ext[1]], Index[Spin, Ext[2]]] +  
TensDot[SlashedP[2], ProjP][Index[Spin, Ext[1]], Index[Spin, Ext[2]]])/Pi^2},  
...  
}
```

~FeynRules syntax

## UV\$vertlist ( $\epsilon$ is FR\$Eps)

```
FR$InteractionOrderPerturbativeExpansion = {{QCD, 1}, {QED, 0}};
```

NLOCT\$assumptions

QCDOOnly

WriteCT[... , Assumptions -> { ... }]

# UFO@NLO

- **CT\_vertices.py**

```
V_1 = CTVertex(name = 'V_1',
    type = 'R2',
    particles = [ P.g, P.g, P.g ],
    color = [ 'f(1,2,3)' ],
    lorentz = [ L.VVV2 ],
    loop_particles = [ [ [P.b], [P.c], [P.d], [P.s], [P.t], [P.u] ], [ [P.g] ] ],
    couplings = {(0,0,0):C.R2GC_273_53,(0,0,1):C.R2GC_273_54})
```

UV

- **CT\_couplings.py**

```
UVGC_271_34 = Coupling(name = 'UVGC_271_34',
    value = {-1:'( 0 if MB else -(complex(0,1)*G**2)/(24.*cmath.pi**2) ) +',
              '(complex(0,1)*G**2)/(24.*cmath.pi**2)',0:'( -(complex(0,1)*G**2*reglog(MB/MU_R))/',
              '(12.*cmath.pi**2) if MB else 0 )'},
    order = {'QCD':2})
```

Pole

Finite

- In **coupling\_order.py**

```
QCD = CouplingOrder(name = 'QCD',
    expansion_order = 99,
    hierarchy = 1,
    perturbative_expansion = 1)
```

```
QED = CouplingOrder(name = 'QED',
    expansion_order = 99,
    hierarchy = 2)
```

# Restrictions/Assumptions

- Renormalizable Lagrangian, maximum dimension of the operators is 4
- Feynman Gauge
- $\{\gamma_\mu, \gamma_5\} = 0$
- 't Hooft-Veltman scheme
- On-shell/complex mass scheme for the masses and wave functions
- $\overline{\text{MS}}$  by default for everything else (zero-momentum possible for fermion gauge boson interaction)

# NLOCT

- Amplitudes from FeynArts (discard irrelevant diagrams like ghost boxes)
- Compute terms at the generic level

$$\vec{c} \cdot \vec{L} = \sum_i c_i L_i$$

- Feynman parameters
- Remove terms with an odd or too low rank
- Gather loop momentum

$$q^\mu q^\nu q^\rho q^\sigma \rightarrow q^4 \frac{1}{d(d+2)} (\eta^{\mu\nu}\eta^{\rho\sigma} + \eta^{\mu\rho}\eta^{\nu\sigma} + \eta^{\mu\sigma}\eta^{\rho\nu})$$

$$q^\mu q^\nu \rightarrow q^2 \frac{1}{d} \eta^{\mu\nu}.$$

# NLO CT

- Replace momentum integrals

$$\int d^d q \frac{\epsilon}{q^2 - m^2} \Big|_{R_2} = i\pi^2 m^2,$$

$$\int d^d q \frac{\epsilon}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2,$$

$$\int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2 (2a - b)\Delta,$$

$$\int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{R_2} = i\pi^2 \left(a - \frac{1}{2}b\right),$$

$$\int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{R_2} = i\pi^2 \left(a - \frac{5}{6}b\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{q^2 - m^2} \Big|_{UV} = i\pi^2 m^2 \left(\frac{b}{\bar{\epsilon}} + a + b - b \log\left(\frac{m^2}{\mu^2}\right)\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (a\epsilon + b) \left(\frac{1}{\bar{\epsilon}} - \log\left(\frac{\Delta}{\mu^2}\right)\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (2a\epsilon + b\epsilon + 2b) \left(\frac{1}{\bar{\epsilon}} - \log\left(\frac{\Delta}{\mu^2}\right)\right) \Delta,$$

$$\mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},$$

$$\mu^{2\epsilon} \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},$$

- Integrate over the Feynman parameters (but for the two-point UV finite terms)

- Replace masses and couplings by their values for each field insertion

# NLO CT

- Replace momentum integrals

$$\int d^d q \frac{\epsilon}{q^2 - m^2} \Big|_{R_2} = i\pi^2 m^2,$$

$$\int d^d q \frac{\epsilon}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2,$$

$$\int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{R_2} = i\pi^2 (2a - b)\Delta,$$

$$\int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{R_2} = i\pi^2 \left(a - \frac{1}{2}b\right),$$

$$\int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{R_2} = i\pi^2 \left(a - \frac{5}{6}b\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{q^2 - m^2} \Big|_{UV} = i\pi^2 m^2 \left(\frac{b}{\bar{\epsilon}} + a + b - b \log\left(\frac{m^2}{\mu^2}\right)\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{a\epsilon + b}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (a\epsilon + b) \left(\frac{1}{\bar{\epsilon}} - \log\left(\frac{\Delta}{\mu^2}\right)\right),$$

$$\mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^2} \Big|_{UV} = i\pi^2 (2a\epsilon + b\epsilon + 2b) \left(\frac{1}{\bar{\epsilon}} - \log\left(\frac{\Delta}{\mu^2}\right)\right) \Delta,$$

$$\mu^{2\epsilon} \int d^d q \frac{q^2 (a\epsilon + b)}{(q^2 - \Delta)^3} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},$$

$$\mu^{2\epsilon} \int d^d q \frac{q^4 (a\epsilon + b)}{(q^2 - \Delta)^4} \Big|_{UV} = i\pi^2 \frac{b}{\bar{\epsilon}},$$

- Integrate over the Feynman parameters (but for the two-point UV finite terms)

- Replace masses and couplings by their values for each field insertion

# NLOCT

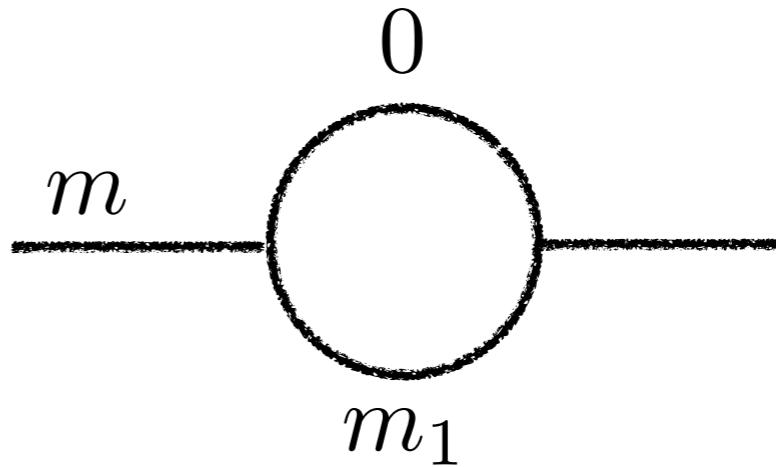
- Perform the color algebra for triplets and octets
- Write the renormalization conditions (fix  $p^2$ ) End R<sub>2</sub>
- Do the integration over the feynman parameters for the UV-finite parts

$$b_0(p^2, m_1, m_2) \equiv \int_0^1 dx \log \left( \frac{p^2(x-1)x + x(m_1^2 - m_2^2) + m_2^2 - i\epsilon_p}{\mu^2} \right)$$

$$b_0(0, 0, 0) = \frac{1}{\bar{\epsilon}}$$

- Solve the renormalization conditions
- Replace the counterterms by their values in the CT vertices

# Real/Complex masses



Real masses

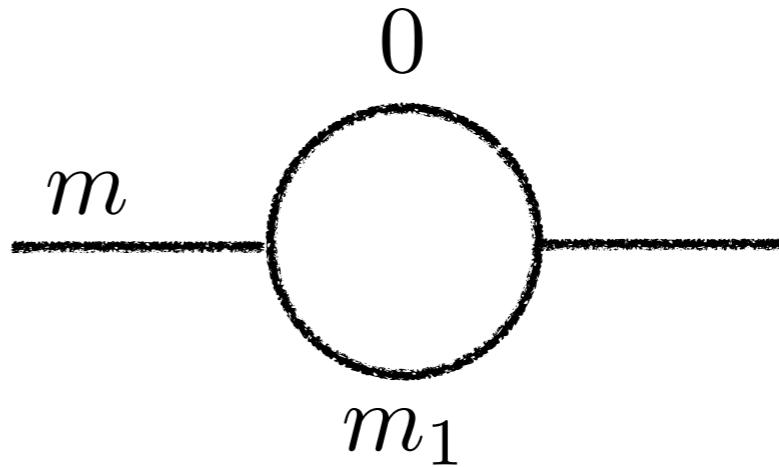
$$m \in \mathbb{R} \quad m_1^2 < m^2 \quad \cancel{\Re(\log [p^2 - m_1^2] + i\pi)} \Big|_{p^2=m^2}$$
$$m_1^2 > m^2 \quad \cancel{\Re(\log [m_1^2 - p^2])} \Big|_{p^2=m^2}$$

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2}$$

All cases are kept unless the users put some assumptions

# Real/Complex masses



Real masses

$$m \in \mathbb{R} \quad m_1^2 < m^2 \quad \cancel{\Re(\log [p^2 - m_1^2] + i\pi)} \Big|_{p^2=m^2}$$
$$m_1^2 > m^2 \quad \cancel{\Re(\log [m_1^2 - p^2])} \Big|_{p^2=m^2}$$

Complex masses

$$\log [m_1^2 - p^2] \Big|_{p^2=m^2} \quad \text{Faster!}$$

All cases are kept unless the users put some assumptions

# Renormalization options

```
FR$LoopSwitches = {{Gf, MW}};
```

Switch internal masses with an external parameter  
appearing in its expression

# Renormalization options

```
FR$LoopSwitches = {{Gf, MW}};
```

Switch int. . . . . before calling `OnShellRenormalization`  
external parameter  
appearing in its `OnShellRenormalization`

# Renormalization options

FR\$LoopSwitches = {{Gf, MW}};

Switch int. . . . . before calling *OnShellRenormalization* external parameter  
appearing in its  $\epsilon_{\mu\nu}$ .

## OnShellRenormalization options

**QCDOnly** : Only the coloured fields and their masses and the couplings with QCD if True

**FlavorMixing** : Forbid all the mixing or allow only some of them

**Exclude4ScalarCT** : No CT for the 4 scalars vertices (but keep the 4 scalars TL)

**Simplify2Point** : Put the quadratic part of the Lagrangian in canonical form (Avoided if False)

# WriteCT options

WriteCT[<model>, <genericfile>, options]

## OnShellRenormalization options

**QCDOOnly** : Only QCD corrections

**Assumptions** : Mass spectrum for the UV counterterms

**Exclude4ScalarCT** : No computation of the CT for the 4 scalars vertices (but keep the 4 scalars TL)

**ZeroMom** : {coupling, vertex} use zero momentum for coupling on vertex

**ComplexMass** : complex mass scheme if True

# R2 :Validation

- tested\* on the SM (QCD:P. Draggiotis et al.  
+QED:M.V. Garzelli et al)
- tested\* on MSSM (QCD:H.-S. Shao,Y.-J. Zhang) : test the Majorana

\*Analytic comparison of the expressions

# UV Validation

- SM QCD : tested\* (W. Beenakker, S. Dittmaier, M. Kramer, B. Plumper)
- SM EW : tested\* (expressions given by H.-S. Shao from A. Denner)

\*Analytic comparison of the expressions

# Tests in event generators

- aMC@NLO
- The SM QCD has been tested by V. Hirschi  
(Comparison with the built-in version)
- The MSSM QCD and SM EW are tested by H.-S. Shao and V. Hirschi
- 2HDM QCD is currently tested ( $p\ p \rightarrow S, H^+ t$ )
  - gauge invariance
  - pole cancelation

# SM tests

==== Finite ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.2565695610e+01	-1.2565705416e+01	-1.2565696276e+01	3.9018817097e-07	Pass
==== Born ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	1.8518318521e-06	1.8518318521e-06	1.8518318521e-06	8.0617231411e-15	Pass
==== Single pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.9397426502e+01	-1.9397426502e+01	-1.9397426504e+01	5.5894073017e-11	Pass
==== Double pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	3.0015206007e-14	Pass
==== Summary ===					
/  passed, 0/  failed==== Finite ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-5.3971186943e+01	-5.3971193753e+01	-5.3971189940e+01	6.3091071914e-08	Pass
==== Born ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	6.4168774056e-05	6.4168764370e-05	6.4168764370e-05	7.5467680882e-08	Pass
==== Single pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-3.7439549398e+01	-3.7439549398e+01	-3.7439549397e+01	6.8122965983e-12	Pass
==== Double pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	2.2443585452e-14	Pass
==== Summary ===					
/  passed, 0/  failed==== Finite ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-5.3769573669e+01	-5.3769573347e+01	-5.3769566412e+01	6.7475496780e-08	Pass

# SM tests

==== Born ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z g g	3.1531233900e-04	3.1531235770e-04	3.1531235770e-04	2.9654886777e-08	Pass	
==== Single pole ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z g g	-3.7464897007e+01	-3.7464897007e+01	-3.7464897007e+01	4.2333025503e-12	Pass	
==== Double pole ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	2.1316282073e-14	Pass	
==== Summary ===						
	I/I passed, 0/I failed == Finite ==					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z z g	-5.9990384275e+00	-5.9990511729e+00	-5.9990379587e+00	1.1013604745e-06	Pass	
==== Born ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z z g	2.2616997126e-06	2.2617000449e-06	2.2617000449e-06	7.3450366526e-08	Pass	
==== Single pole ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z z g	-1.5469587040e+01	-1.5469587040e+01	-1.5469587040e+01	1.5226666708e-11	Pass	
==== Double pole ===						
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
d~ d > z z g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	2.6645352591e-15	Pass	
==== Summary ===						
	I/I passed, 0/I failed == Finite ==					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result	
g g > h t t~	2.9740187004e+01	2.9740187005e+01	2.9740187036e+01	5.3265970697e-10	Pass	

# SM tests

==== Born ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	1.1079653971e-07	1.1079653974e-07	1.1079653974e-07	1.3190849004e-10	Pass
==== Single pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	-7.0825709000e+00	-7.0825709000e+00	-7.0825709000e+00	5.0901237085e-13	Pass
==== Double pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > h t t~	-6.0000000000e+00	-6.0000000000e+00	-6.0000000000e+00	1.7023419711e-15	Pass
==== Summary ===					
I/I passed, 0/I failed == Finite ==					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	3.6409017466e+01	3.6409021125e+01	3.6409021117e+01	5.0242920154e-08	Pass
==== Born ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	7.0723041711e-07	7.0723046101e-07	7.0723046101e-07	3.1039274206e-08	Pass
==== Single pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	-7.1948086812e+00	-7.1948086773e+00	-7.1948086773e+00	2.7349789963e-10	Pass
==== Double pole ===					
Process	Stored MadLoop v4	ML5 opt	ML5 default	Relative diff.	Result
g g > z t t~	-6.0000000000e+00	-6.0000000000e+00	-6.0000000000e+00	2.5165055225e-15	Pass
==== Summary ===					
I/I passed, 0/I failed == Finite ==					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.2565695610e+01	-1.2565705416e+01	-1.2565696276e+01	3.9018817097e-07	Pass

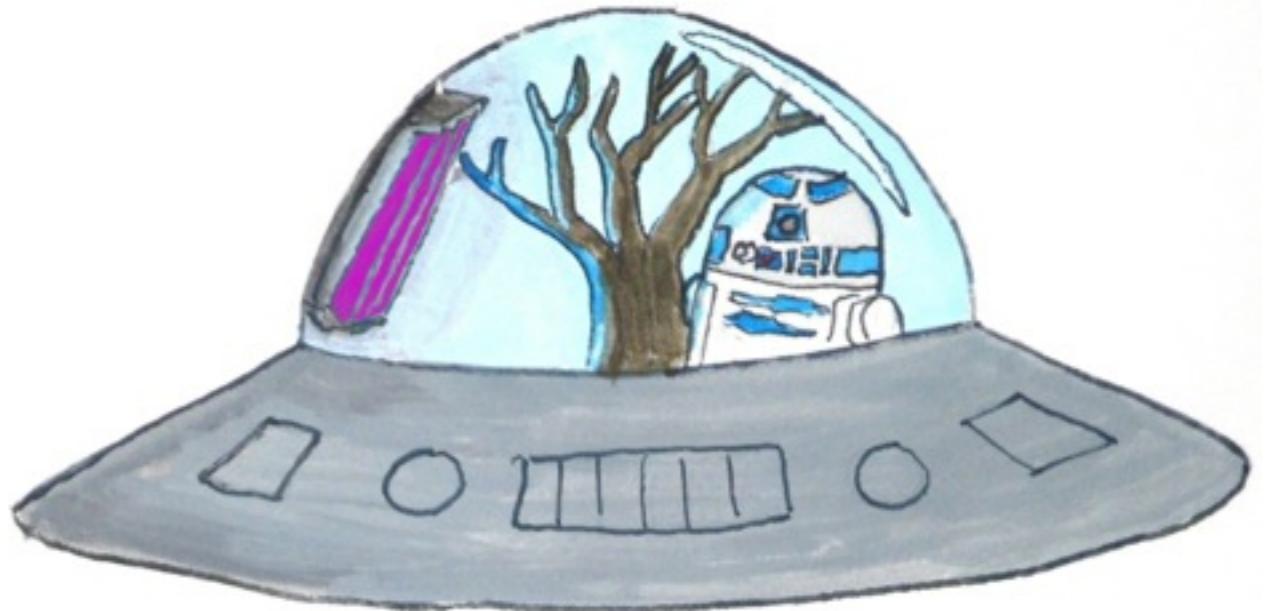
# SM tests

==== Born ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	1.8518318521e-06	1.8518318521e-06	1.8518318521e-06	8.0617231411e-15	Pass
==== Single pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-1.9397426502e+01	-1.9397426502e+01	-1.9397426504e+01	5.5894073017e-11	Pass
==== Double pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d d~ > w+ w- g	-5.6666666667e+00	-5.6666666667e+00	-5.6666666667e+00	3.0015206007e-14	Pass
==== Summary ===					
I/I passed, 0/I failed == Finite ==					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-1.1504816412e+01	-1.1504816557e+01	-1.1504815497e+01	4.6089385415e-08	Pass
==== Born ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	2.3138920858e-06	2.3138920858e-06	2.3138920858e-06	4.3012538015e-15	Pass
==== Single pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-2.8637049838e+01	-2.8637049838e+01	-2.8637049838e+01	1.5718407645e-13	Pass
==== Double pole ===					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > a g g	-8.6666666667e+00	-8.6666666667e+00	-8.6666666667e+00	1.7421961310e-15	Pass
==== Summary ===					
I/I passed, 0/I failed == Finite ==					
Process	Stored ML5 opt	ML5 opt	ML5 default	Relative diff.	Result
d~ d > z g g	-1.0306105482e+01	-1.0306105654e+01	-1.0306102645e+01	1.4600800434e-07	Pass

+2/3

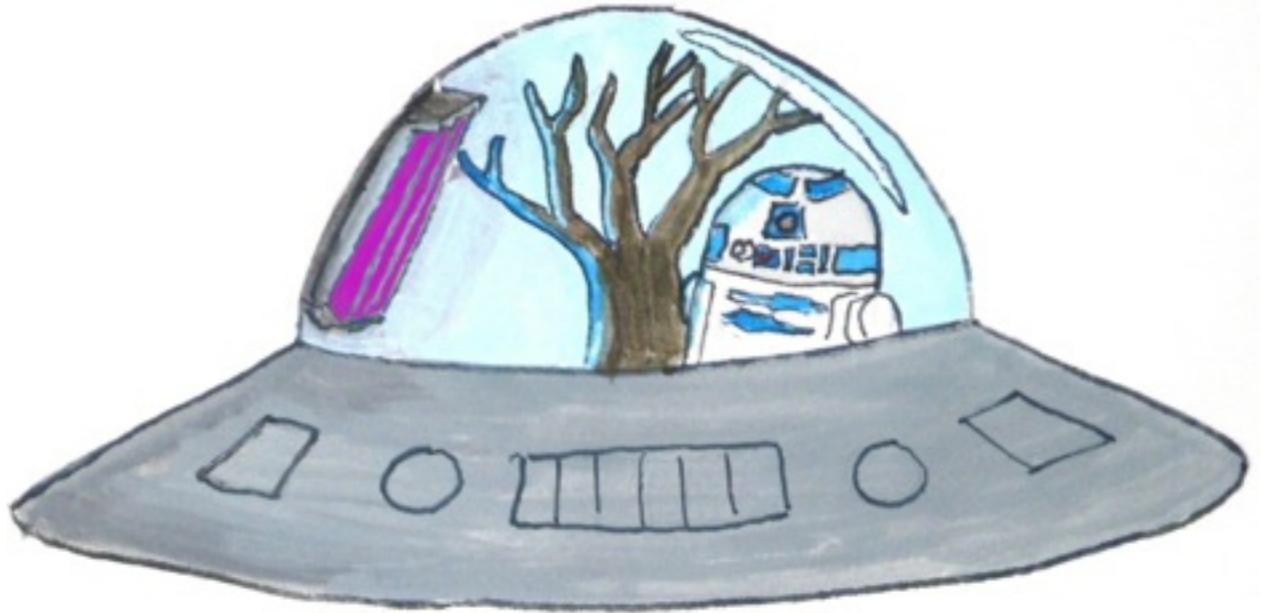
# Summary/Prospect

- Automatic BSM@NLO
  - renormalizable
  - Feynman gauge



# Summary/Prospect

- Automatic BSM@NLO
  - renormalizable
  - Feynman gauge
- Next version
  - EFT
  - Any gauge
  - other renormalization scheme (EW)



# Summary/Prospect

- Automatic BSM@NLO
  - renormalizable
  - Feynman gauge
- Next version
  - EFT
  - Any gauge
  - other renormalization scheme (EW)
- With the help of the FeynRules and Madgraph\_aMC@NLO teams

