Last Modified: 27th Sept, 2021 Acknowledgements: Text

Problem 1.

(i) Find the solution of the following 1st-order PDE:

$$u_t + tu_x = 0, \quad x \in \mathbb{R}, \ t > 0$$

 $u(x, 0) = e^x, \quad x \in \mathbb{R}$

This equation is of the form

$$u_t + p(x, t)u_x = 0$$

whereby the general solution is of the form

$$u(x,t) = f(\phi(x,t))$$

where $\phi(x,t) = C$ is the general solution to the ODE $\frac{\mathrm{d}x}{\mathrm{d}t} = p(x,t)$. Given that p(x,t) = t in this problem, we see that the general solution to $\phi(x,t)$ is

$$x - \frac{1}{2}t^2 = C$$

Whereby u is constant along this characteristic curve $\phi(x,t) = x - \frac{1}{2}t^2$.

The initial condition, given by $u(x,0) = e^x$, therefore implies that the solution to the 1st-order PDE is

$$u(x,0) = f(x) = e^x \Rightarrow u(x,t) = f\left(x - \frac{1}{2}t^2\right) = e^{x - \frac{1}{2}t^2}$$

(ii) Find the solution of the following 1st-order PDE:

$$u_t + 2u_x - u = 0, \quad x \in \mathbb{R}, \ t > 0$$
$$u(x, 0) = x^2, \quad x \in \mathbb{R}$$

We perform a linear change of variables:

$$\alpha = ax + bt$$
$$\beta = cx + dt$$

And we see that

$$\frac{\partial u}{\partial x} = \frac{\partial u}{\partial \alpha} \frac{\partial \alpha}{\partial x} + \frac{\partial u}{\partial \beta} \frac{\partial \beta}{\partial x} = a \frac{\partial u}{\partial \alpha} + c \frac{\partial u}{\partial \beta}$$
$$\frac{\partial u}{\partial t} = \frac{\partial u}{\partial \alpha} \frac{\partial \alpha}{\partial t} + \frac{\partial u}{\partial \beta} \frac{\partial \beta}{\partial t} = b \frac{\partial u}{\partial \alpha} + d \frac{\partial u}{\partial \beta}$$

Therefore

$$u_t + 2u_x = (b+2a)u_\alpha + (d+2c)u_\beta$$

And by choosing $b=1, a=0, d=-2c=-2 \Rightarrow c=1$, we obtain $u_{\alpha}=u$, and therefore, we have that

$$\ln u = \alpha + f(\beta) = t + f(x - 2t)$$

The initial conditions give us

$$ln u(x,0) = f(x) = 2 ln x$$

So that our final form of the general solution is therefore

$$\ln u = \alpha + f(\beta) = t + 2\ln(x - 2t)$$
$$u = (x - 2t)^2 \cdot e^t$$

(iii) We have the following parabolic PDE:

$$\phi_{xx} + 2\phi_{xy} + \phi_{yy} = 1$$

We know that the families of curves $\xi = c$ where c is a constant are the characteristic curves. Therefore, the change of variables is given by

$$\xi = y - \frac{2}{2 \cdot 1}x = y - x$$
$$\eta = x$$

And the transformed equation therefore becomes

$$a\phi_{\xi\xi} + b\phi_{\xi\eta} + c\phi_{\eta\eta} = 1$$

Where (a, b, c) are given by

$$a = A\xi_x^2 + B\xi_x\xi_y + C\xi_y^2 = 2(-1)(1) = -2$$

$$b = 2A\xi_x\eta_x + B(\xi_x\eta_y + \xi_y\eta_x) + 2C(\xi_y\eta_y) = 2(1)(-1) + 2 = 0$$

$$c = 0$$

And therefore the canonical equation is $2\phi_{\xi\xi} + 1 = 0$. Now we solve it as follows:

$$\phi_{\xi\xi} = -\frac{1}{2}$$

$$\phi_{\xi} = -\frac{1}{2}\xi + c_1 + f(\eta)$$

$$\phi = -\frac{1}{4}\xi^2 + (c_1 + f(\eta))\xi + g(\eta) + c_2$$

$$= -\frac{(y-x)^2}{4} + (c_1 + f(x))(y-x) + g(x) + c_2$$

(iv) We are given the elliptic PDE

$$u_{xx} - 6u_{xy} + 12u_{yy} = 0$$

And now let us find a transofmration of independent variables to convert it to the form of the Laplace's equation $\nabla^2 f = 0$.

The roots of the characteristic polynomial are given by

$$\lambda_{\pm} = \frac{-6 \pm \sqrt{6^2 - 48}}{2} = -3 \pm \sqrt{3^2 - 12} = -3 \pm i\sqrt{3}$$

And therefore, we see that the characteristics of the equation are:

$$\frac{\mathrm{d}y}{\mathrm{d}x} = -3 \pm i\sqrt{3} \Rightarrow y = (-3 \pm i\sqrt{3})x + c$$

So we see that the two families of the characteristic curves are:

$$c_1 = y + (3 + i\sqrt{3})x$$
, $c_2 = y + (3 - i\sqrt{3})x$

And therefore the required transformation variables ξ and η are

$$\xi = \frac{c_1 + c_2}{2} = y + 3x$$
$$\eta = \frac{c_1 + c_2}{2i} = \sqrt{3}x$$

So the transformed equation therefore becomes

$$a\phi_{\xi\xi} + b\phi_{\xi\eta} + c\phi_{\eta\eta} = 0$$

Where (a, b, c) are given by

$$a = A\xi_x^2 + B\xi_x\xi_y + C\xi_y^2 = (3)^2 - 6(3) + 12 = 3$$

$$b = 2A\xi_x\eta_x + B(\xi_x\eta_y + \xi_y\eta_x) + 2C(\xi_y\eta_y) = 2(1)(3)(\sqrt{3}) - 6\sqrt{3} = 0$$

$$c = 3$$

And therefore the canonical equation is $3\phi_{\xi\xi} + 3\phi_{\eta\eta} = 0 \Rightarrow \phi_{\xi\xi} + \phi_{\eta\eta} = \nabla^2\phi = 0$ which is the Laplacian in (ξ,η) coordinates.

Problem 2.

We consider the partial differential equation:

$$u_t + \alpha u_x - \beta u_{xxx} + \gamma u_{xxxx} = 0$$

along a periodic domain $[0, 2\pi]$ with boundary conditions:

$$u(0) = u(2\pi), \quad u_x(0) = u_x(2\pi), \quad u_{xx}(0) = u_{xx}(2\pi), \quad u_{xxx}(0) = u_{xxx}(2\pi)$$

with the following initial conditions:

$$u(x,0) = \cos^2 x - \sin^2 x$$

Because the periodic is doubly periodic, u(x,t) must take on the form:

$$u(x,t) = \sum_{k=-\infty}^{\infty} u_k(t) \cdot e^{ikx}$$

And now, substituting the above into the initial PDE and eliminating e^{ikx} , we get

$$\frac{\partial u_k}{\partial t} + \alpha i k u_k + \beta i k^3 u_k + \gamma k^4 u_k = 0$$
$$\frac{\partial u_k}{\partial t} + u_k (\alpha i k + \beta i k^3 + \gamma k^4) = 0$$

Doing the separation of variables, we then obtain

$$\frac{\mathrm{d}u_k}{u_k} = -(\alpha ik + \beta ik^3 + \gamma k^4) \,\mathrm{d}t$$

$$\ln|u_k| = -(\alpha ik + \beta ik^3 + \gamma k^4)t + c$$

$$\therefore u_k = Ae^{-(\alpha ik + \beta ik^3 + \gamma k^4)t}$$

So the final form of u(x,t) is

$$u(x,t) = \sum_{k=-\infty}^{\infty} A_k e^{-(\alpha ik + \beta ik^3 + \gamma k^4)t} \cdot e^{ikx}$$

Given the initial conditions, we see that

$$u(x,t) = \sum_{k=-\infty}^{\infty} A_k e^{ikx} = \cos^2 x - \sin^2 x$$

$$= \left(\frac{e^{ix} + e^{-ix}}{2}\right)^2 - \left(\frac{e^{ix} - e^{-ix}}{2i}\right)^2 = \frac{e^{2ix} + 2 + e^{-2ix}}{4} + \frac{e^{2ix} - 2 + e^{-2ix}}{4}$$

$$= \frac{1}{2}(e^{2ix} + e^{-2ix})$$

Therefore we see that $A_{-2}=A_2=1/2$ and all other $A_k=0$, such that

$$u(x,t) = \frac{1}{2} \left(e^{-(2\alpha i + 8\beta i + 16\gamma)t} \cdot e^{2ix} + e^{(2\alpha i + 8\beta i - 16\gamma)t} \cdot e^{-2ix} \right)$$

Problem 3.

(i) The solution to a unit impulse at the origin in an infinite domain has the form

$$u(x,t) = \frac{1}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{x^2}{4\kappa t}\right)$$

We wish to verify that for t > 0, the above solution is indeed a solution of the heat conduction equation $u_t = \kappa u_{xx}$.

$$u_t = -\frac{1}{2t} \frac{1}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{x^2}{4\kappa t}\right) + \frac{x^2}{4kt^2} \frac{1}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{x^2}{4\kappa t}\right)$$

$$= \left(-\frac{1}{2t} + \frac{x^2}{4\kappa t}\right) u$$

$$u_x = -\frac{x}{2\kappa t} u$$

$$\therefore u_{xx} = -\frac{1}{2\kappa t} u + \left(\frac{x}{2\kappa t}\right)^2 u = \left(-\frac{1}{2\kappa t} + \frac{x^2}{4\kappa^2 t}\right) u = \frac{u_t}{\kappa}$$

(ii) We next perform a change of independent variables $(x,t) \to (\xi,t')$, where

$$\xi = x - Ut, \quad t' = t$$

Therefore, we see that

$$u_{t} = u_{t'}t'_{t} + u_{\xi}\xi_{t} = u_{t'} - Uu_{\xi}$$

$$u_{x} = u_{t'}t'_{x} + u_{\xi}\xi_{x} = u_{\xi}$$

$$u_{xx} = u_{t't'}t'^{2}_{x} + 2u_{t'\xi}\xi_{x}t'_{x} + u_{\xi\xi}\xi^{2}_{x} + u_{t'}t'_{xx} + u_{\xi}\xi_{xx}$$

$$= 0 + 0 + u_{\xi\xi} + 0 + 0 = u_{\xi\xi}$$

So the linear convection-diffusion equation is transformed into

$$u_t + Uu_x = \kappa u_{xx} \Rightarrow u_{t'} - Uu_{\xi} + Uu_{\xi} = u_{t'} = u_{\xi\xi}$$

The solution of which is

$$u(\xi, t') = \frac{1}{\sqrt{4\pi\kappa t'}} \exp\left(-\frac{\xi^2}{4\kappa t'}\right)$$
$$\therefore u(x - Ut, t) = \frac{1}{\sqrt{4\pi\kappa t}} \exp\left(-\frac{(x - Ut)^2}{4\kappa t}\right)$$

(iii) We consider discretized forms of the components of the linear convection-diffusion equation as follows:

$$u_{t} = \frac{\partial u}{\partial t} = \frac{P_{i}^{n+1} - P_{i}^{n}}{\tau}$$

$$u_{x} = \frac{\partial u}{\partial x} = \frac{P_{i+1}^{n} - P_{i-1}^{n}}{2\Delta x} = \frac{P_{i+1}^{n} - P_{i-1}^{n}}{2/M} = \frac{M(P_{i+1}^{n} - P_{i-1}^{n})}{2}$$

$$u_{xx} = M^{2}(P_{i+1} + P_{i_{1}} - 2P_{i})$$

Such that we convert the discrete Markov process to:

$$P_i^{n+1} - P_i^n + \left(p - \frac{1}{2}\right) \left(P_{i+1}^n - P_{i-1}^n\right) = \frac{1}{2} \left(P_{i+1} + P_{i_1} - 2P_i\right)$$
$$\tau u_t + \frac{2p - 1}{M} u_x = \frac{2}{M^2} u_{xx} \Rightarrow u_t + \frac{2p - 1}{M\tau} = \frac{1}{2M^2\tau} u_{xx}$$

Comparing against the actual linear convection-diffusion equation, and noting that $t = N\tau$, we have that

$$U = \frac{2p-1}{M\tau} = \frac{2p-1}{Mt/N} \Rightarrow Ut = \frac{(2p-1)N}{M}$$
$$\kappa = \frac{1}{2M^2\tau} \to \kappa t = \frac{N}{2M^2}$$

Let us now transform x and t to x_n and t_n such that

$$x_n = x + c_x$$
, $t_n = t + c_t$

Such that when t = 0)

$$u(x_n, 0) = \frac{1}{\sqrt{4\pi\kappa c_t}} \exp\left(-\frac{(x + c_x - Uc_t)^2}{4\kappa c_t}\right)$$
$$= \frac{1}{\sqrt{2\pi(0.05)^2}} \exp\left(-\frac{(x - 0.2)^2}{2(0.05)^2}\right)$$

So, we see that c_t and c_x are:

$$c_t = \frac{2(0.05)^2}{4\kappa} = \frac{1}{800\kappa}$$
$$c_x = -0.2 + Uc_t$$

And therefore, the final equation therefore becomes

$$u(x_n, t_n) = \frac{1}{\sqrt{4\pi\kappa t_n}} \exp\left(-\frac{(x_n - Ut_n)^2}{4\kappa t_n}\right)$$

$$= \frac{1}{\sqrt{4\pi\kappa (t + c_t)}} \exp\left(-\frac{(x + c_x - U(t + c_t))^2}{4\kappa (t + c_t)}\right)$$

$$= \frac{1}{\sqrt{4\pi\kappa (t + c_t)}} \exp\left(-\frac{(x - 0.2 - Ut)^2}{4\kappa (t + c_t)}\right)$$

$$= \frac{1}{\sqrt{4\pi(\kappa t + 1/800)}} \exp\left(-\frac{(x - 0.2 - Ut)^2}{4(\kappa t + 1/800)}\right)$$

So we have the analytical solution for t=0 and $t=N\tau$

$$u(x,0) = \frac{20}{\sqrt{2\pi}} e^{-200(x-0.2)^2}$$

$$u(x,N\tau) = \frac{1}{\sqrt{4\pi(N/2M^2 + 1/800)}} \exp\left(-\frac{(x-0.2 - (2p-1)N/M)^2}{4(N/2M^2 + 1/800)}\right)$$

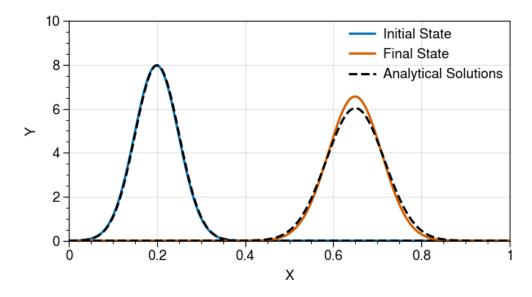


Figure 1: Plots of our analytic solutions (black-dashed), and initial (blue) and final (red) numerical solutions to the linear convection-diffusion equation. Here M = 400, N = 300 and p = 0.8.