## Relativity Physics 127 Homework 1

Due Wednesday April 9th 2025, 11:59pm on gradescope.

- 1.) The Einstein summation convention is defined as follows: encountering a repeated spacetime index  $\mu$  (one raised and one lowered) in an expression, we sum over that index from 0 to 3, e.g.  $A^{\mu}B_{\mu} = \sum_{\mu=0}^{3} A^{\mu}B_{\mu}$ .
  - a.) Consider four-vectors  $v^{\mu}$  and  $w^{\mu}$  in Mikowski space, and the Minkowski metric

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & -1 & 0 & 0 \\ 0 & 0 & -1 & 0 \\ 0 & 0 & 0 & -1 \end{pmatrix} .$$

Think of  $v^{\mu}$ ,  $w^{\mu}$  as column vectors v, w, and of  $g_{\mu\nu}$  as a square matrix g.

The Mikowski inner (or dot) product  $v \cdot w$  between v and w is defined by the index-free expression

$$v \cdot w = v^T g w .$$

Here T denotes the transpose. Express this in terms of the components  $v^{\mu}$ ,  $w^{\mu}$ ,  $g_{\mu\nu}$  using Einstein summation convention. Then show that the inner product is symmetric,  $v \cdot w = w \cdot v$ . Recall that you are free to rename pairs of dummy indices that are summed over.

Set v = w = x, where  $x^{\mu}$  is a spacetime point and show that  $x \cdot x$  reduces to the distance-squared in Minkowski space we discussed in class.

b.) Consider a Lorentz transformation  $\Lambda^{\mu}_{\ \nu}$ , which you can think of as a  $4\times 4$  matrix  $\Lambda$  with entries  $\Lambda^{\mu}_{\ \nu}$ . Here the left index  $\mu$  always denotes the row and the right index  $\nu$  the column.

Argue that the transpose matrices  $\Lambda^T$  should have components

$$(\Lambda^T)_{\mu}^{\phantom{\mu}\nu} = \Lambda^{\nu}_{\phantom{\nu}\mu} .$$

The condition that  $\Lambda$  is a Lorentz transformation, expressed in components, takes the form

$$\Lambda^{\mu}{}_{\rho}\Lambda^{\nu}{}_{\sigma}g_{\mu\nu} = g_{\rho\sigma} .$$

Show that this reduces to the index-free matrix equation  $\Lambda^T g \Lambda = g$ .

- c.) Show that the Minkowski inner product  $v \cdot w$  defined in 1a.) above is Lorentz invariant, i.e. if  $v' = \Lambda v$  and  $w' = \Lambda w$ , where  $\Lambda$  is a Lorentz transformation satisfying the condition reviewed in 1b.) above, then show that  $v' \cdot w' = v \cdot w$ .
- 2.) A boost along the x-axis is given and a rotation around the x-axis is given by

$$\Lambda^{\mu}_{\nu}(\alpha) = \begin{pmatrix} \cosh \alpha & \sinh \alpha & 0 & 0 \\ \sinh \alpha & \cosh \alpha & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, \quad R^{\mu}_{\nu}(\theta) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & \cos \theta & \sin \theta \\ 0 & 0 & -\sin \theta & \cos \theta \end{pmatrix},$$

- a) Verify that both matrices are Lorentz transformations, i.e.  $\Lambda^T \cdot g \cdot \Lambda = g$  and  $R^T \cdot g \cdot R = g$
- b) Find the inverse transformations of  $\Lambda$  and R and the parameter transformation of  $\alpha$  and  $\theta$  which gives the inverse.
- c) Calculate  $\Lambda(\alpha_1) \cdot \Lambda(\alpha_2)$  and  $R(\theta_1) \cdot R(\theta_2)$  and show that they can be again written as  $\Lambda(\alpha')$  and  $R(\theta')$  and find  $\alpha'$  and  $\theta'$
- 3.) Consider the action of a boost on a spacetime point  $x^{\mu}$ . For simplicity we consider boosts in the x-direction, and thus also take  $x^{\mu} = (t, x, 0, 0)$  to lie in the t x plane. If  $x^{\mu}$  is timelike, show that you can boost to a frame where x = 0. Similarly, if  $x^{\mu}$  is spacelike, show that you can boost to a frame where t = 0.
- 4.) In class we briefly discussed the fact that the Lorentz group has four disconnected components and is parameterized by six continuous parameters [See also section 1.2.4 in Coleman]. Respeat the analysis for the three dimensional Lorentz group, i.e  $3 \times 3$  matrices which leave the three dimensional Minkoswski metric tensor in variant

$$g_{\mu\nu} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} .$$

i.e. satisfy

$$\Lambda^T \cdot g \cdot \Lambda = g$$

You should find that there are still four disconnected components, but it's not exactly

the same due to the fact that the parity transformation

$$P = \begin{pmatrix} 1 & 0 & 0 \\ 0 & -1 & 0 \\ 0 & 0 & -1 \end{pmatrix} .$$

is actually a rotation (which one?) and hence connected to the identity.

5.) Coleman problem 1.1 (You can either attempt a proof or just do an example path).