

Fluid Dynamics Equation Sheet

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- Chapter 1: Introduction

- Gage Pressure

- $P_{gage} = P_{abs} - P_{ambient}$

- p_{abs} = absolute pressure

- $p_{ambient}$ = local pressure

- p_{gage} = what engineers measure

- Specific weight

- $\gamma = \rho g$

- Specific Gravity

- $SG = \frac{\rho}{\rho_{H_2O@4^\circ C}}$

- Ideal Gas Law

- $\rho = \frac{p}{RT}$

- p = absolute pressure

- R = a substance specific gas constant

- T = absolute temperature

- Shear Stress

- $\tau = \frac{F}{A}$

- Newtonian Fluid Shear Stress

- $\tau = \mu \frac{\partial u}{\partial y}$

- μ = dynamic viscosity

- u = fluid velocity

- y = the y direction

- $\frac{\partial}{\partial t} (Momentum) = \nabla \cdot \tau + \text{"other stuff"}$

- Divide by density on both sides

- $\frac{\partial}{\partial t} (Velocity) = \nabla \cdot \frac{\tau}{\rho}$

- Where τ is stress

- Chapter 2: Fluid Statics

- Equation of Pressure Gradient

- $\vec{\nabla} p = -\rho \vec{a}_t$

- $\vec{\nabla} p = \frac{\partial p}{\partial x} \hat{i} + \frac{\partial p}{\partial y} \hat{j} + \frac{\partial p}{\partial z} \hat{k}$

- \vec{a}_t = acceleration

■ Special Pressure Equation Cases 1a:

- No motion and incompressible

$$\circ \frac{\partial p}{\partial x} = 0$$

$$\circ \frac{\partial p}{\partial y} = 0$$

$$\circ \frac{\partial p}{\partial z} = -\rho g$$

$$\blacksquare \vec{a}_{t,ext} = 0$$

- External forces are 0 and fluid is incompressible

$$\blacksquare -\rho g \text{ is from gravity}$$

■ Special Pressure Equation Cases 1b:

- No motion and compressible

$$\circ \frac{\partial p}{\partial z} = -\rho(z)g$$

- Most gasses are this way

- Density is a function of position

■ Special Pressure Equation Cases 2a:

- Constant acceleration and incompressible

$$\circ \frac{\partial p}{\partial x} = -\rho a_x$$

$$\circ \frac{\partial p}{\partial y} = -\rho a_y$$

$$\circ \frac{\partial p}{\partial z} = -\rho g - \rho a_z$$

$$\blacksquare \vec{a}_t = \langle a_x, a_y, a_z \rangle$$

- Buoyancy Force

$$\blacksquare F_b = \rho g V = \gamma V$$

- ρ = the density of the fluid
- V = the volume that is submerged
- $\gamma = \rho g$ = specific weight

- Line of constant pressure

■ Multivariable Calculus review:

$$\bullet dp = \frac{\partial p}{\partial x} dx + \frac{\partial p}{\partial y} dy + \frac{\partial p}{\partial z} dz = 0$$

- For constant pressure, derivative $dp = 0$

- This come from the 2d equation:

$$\blacksquare y - y_0 = \frac{dy}{dx} (x - x_0)$$

$$\blacksquare \quad dy = \frac{dy}{dx} dx$$

- Chapter 3: The Bernoulli Equation

- Bernoulli's Equation

$$\blacksquare \quad p + \frac{1}{2}\rho V^2 + \rho g z = \text{Constant}$$

- p = pressure
- ρ = fluids density
- V = Velocity **magnitude**
- $V = \sqrt{x^2 + y^2 + z^2}$
- z = height
- Assumptions
 - No Viscosity
 - Steady flow
 - Incompressible
 - Only along streamlines

- Steady State flow definition

$$\blacksquare \quad \frac{\partial \vec{V}}{\partial t} = \vec{0}$$

- Conservation of mass flow

$$\blacksquare \quad \rho V_1 A_1 = \rho V_2 A_2$$

- V = Velocity
- A = Pipe area

- Volumetric flow rate

$$\blacksquare \quad Q = VA$$

- Q =volumetric flow rate

- Mass flow rate

$$\blacksquare \quad \dot{M} = \rho V A = \rho Q$$

- \dot{M} = mass flow rate

- Chapter 4: Fluid Kinematics

- Equation for Streamlines

$$\blacksquare \quad \frac{dx}{u} = \frac{dy}{v} = \frac{dw}{z}$$

- For two dimensions:
 - $\frac{dy}{dx} = \frac{v}{u}$
 - Because stream line is always **tangent** to velocity field

○ Acceleration

$$\begin{aligned} \blacksquare \quad \vec{a} &= \frac{\partial \vec{V}}{\partial t} + u \frac{\partial \vec{V}}{\partial x} + v \frac{\partial \vec{V}}{\partial y} + w \frac{\partial \vec{V}}{\partial z} \\ &\bullet \quad \vec{a} = a_x \hat{i} + a_y \hat{j} + a_z \hat{k} \\ &\bullet \quad \vec{V} = u \hat{i} + v \hat{j} + w \hat{k} \end{aligned}$$

○ Material Derivative

$$\begin{aligned} \blacksquare \quad \frac{D(\cdot)}{Dt} &= \frac{\partial(\cdot)}{\partial t} + (\vec{V} \cdot \vec{\nabla})(\cdot) \\ &\bullet \quad \text{Acceleration } \vec{a} = \frac{D\vec{V}}{Dt} \end{aligned}$$

○ Extensive/Intensive relationship

$$\begin{aligned} \blacksquare \quad B &= m \cdot b \\ &\bullet \quad B \text{ is an extensive property} \\ &\bullet \quad b \text{ is an intensive property} \\ &\bullet \quad m = \text{mass} \end{aligned}$$

○ Amount of B in V

$$\begin{aligned} \blacksquare \quad B &= \int_V \rho b dV \\ &\bullet \quad V = \text{Volume} \\ &\bullet \quad b = \text{intensive property} \\ &\bullet \quad B = \text{extensive property} \end{aligned}$$

○ Flux of extensive property

$$\begin{aligned} \blacksquare \quad \dot{B} &= \int_S \rho b \vec{V} \cdot \vec{n} dS \\ &\bullet \quad S = \text{a surface} \\ &\bullet \quad b = \text{intensive property} \\ &\bullet \quad B = \text{extensive property} \\ &\bullet \quad \vec{V} = \text{velocity vector} \\ &\bullet \quad \vec{n} = \text{unit vector normal to the surface} \end{aligned}$$

○ Reynolds Transport Theorem

$$\begin{aligned} \blacksquare \quad \frac{\partial B_{cv}}{\partial t} + \dot{B}_{cs} &= 0 \\ &\bullet \quad \frac{\partial B_{cv}}{\partial t} = \text{the unsteady term. It should be zero in steady flow} \\ &\bullet \quad \dot{B} = \text{the fluid flux flowing through the surface} \end{aligned}$$

- Chapter 5: Finite Control Volume Analysis

- Mass Conservation

- $\frac{\partial}{\partial t} \int_{cV} \rho dV + \int_{cs} \rho \vec{V} \cdot \hat{n} ds = 0$

- The first term is the unsteady term
 - cV = is the control Volume
 - cs = control surface
 - You can split the control surface into multiple flat plans and add it together

- Momentum Conservation

- $\frac{\partial}{\partial t} \int_{cV} \vec{V} \rho dV + \int_{cs} \vec{V} \rho \vec{V} \cdot \hat{n} ds = \sum F_{ext}$

- F_{ext} is all external forces
 - Like pressure and stress and such

- External Forces x direction (typically)

- Since $F = pA$

- $\sum F_{ext} = p_{left} A_{left} - p_{right} A_{right} + F_x$

- External Forces y direction (typically)

- $\sum F_{ext} = p_{bottom} A_{bottom} - p_{top} A_{top} + F_y$

- Chapter 6: Differential Analysis of Fluid Flow

- Volumetric Dilation

- $div(\vec{V})$

- $div(\vec{V}) = \vec{\nabla} \cdot \vec{V}$ watch your coordinate system tho

- Rotation Vector $\vec{\omega}$

- $\frac{1}{2} curl(\vec{V})$

- $curl(\vec{V}) = \vec{\nabla} \times \vec{V}$ watch your coordinate system tho

- Vorticity $\vec{\zeta}$

- $curl(\vec{V})$ or $2\vec{\omega}$

- Gradient Tensor

- $\vec{\nabla} \vec{V} = S_{ij} + \Omega_{ij}$
 - $S = \frac{1}{2} (\vec{\nabla} \vec{V} + (\vec{\nabla} \vec{V})^T)$
 - the rate of stress tensor
 - $\Omega = \frac{1}{2} (\vec{\nabla} \vec{V} - (\vec{\nabla} \vec{V})^T)$
 - Rotation tensor

- $trace(\vec{\nabla} \vec{V}) = div(\vec{V})$

- $trace$ is the sum of the diagonal components of a square matrix

- Mass conservation

- $\vec{\nabla} \cdot \vec{V} = 0$

- “Newton's Second Law”

- $ma = \sum F$
 - $ma = \rho \frac{D\vec{V}}{Dt}$
 - $F_{int} = \vec{\nabla} \cdot T_{ij}$
 - $F_{ext} = \rho \vec{g}$

- Stress tensor

- T_{ij} = isotropic tensor + deviatoric stress tensor
 - isotropic tensor = $-pI_{33}$
 - Where I_{ij} is an $i \times j$ identity matrix
 - deviatoric stress tensor = τ_{ij} the typical stress tensor from mechanics of materials

- Deviatoric Stress Tensor

- $\tau_{ij} = \frac{-2}{3} \mu (\vec{\nabla} \cdot \vec{V}) I_{ij} + 2\mu S_{ij}$
 - $S_{ij} = 0$ in incompressible fluids
 - μ = dynamic viscosity

- Navier Stokes

- $\vec{\nabla} \cdot \vec{V} = 0$

- $\frac{\partial \vec{V}}{\partial t} + (\vec{V} \cdot \vec{\nabla}) \vec{V} = \frac{-1}{\rho} \vec{\nabla} p + \nu \vec{\nabla}^2 \vec{V} + \vec{g}$
 - $\nu = \frac{\mu}{\rho}$ kinematic viscosity
- 2D Navier Stokes
 - Euler Equations
 - Navier stokes without viscous terms
 - $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$
 - $\frac{\partial u}{\partial t} + u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = \frac{-1}{\rho} \frac{\partial p}{\partial x}$
 - $\frac{\partial v}{\partial t} + u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = \frac{-1}{\rho} \frac{\partial p}{\partial y}$
 - Potential functions
 - Linearity applies
 - Assumptions incompressible
 - Inviscid
 - Irrotational
 - Stream functions $\psi(x, y)$
 - $u(x, y) = \frac{\partial \psi(x, y)}{\partial y}$
 - $v(x, y) = \frac{-\partial \psi(x, y)}{\partial x}$
 - In 2D, ψ is a scalar that can generate a vector field
 - ψ is constant along streamlines
 - $\Delta \psi$ is volumetric flow rate
 - Velocity potential ϕ
 - $\vec{V} = \vec{\nabla} \phi$
 - $u = \frac{\partial \phi}{\partial x}$
 - $v = \frac{\partial \phi}{\partial y}$
 - $w = \frac{\partial \phi}{\partial z}$
 - Basic flows
 - Uniform flow magnitude U and angle α
 - $\phi = Ux \cos(\alpha) + Uy \sin(\alpha)$
 - $\psi = Uy \cos(\alpha) - Ux \sin(\alpha)$
 - Source/Sink
 - $\phi = \frac{m}{2\pi} \ln(r)$
 - $\psi = \frac{m}{2\pi} \theta$
 - Free vortex
 - $\phi = \frac{\Gamma}{2\pi} \theta$

- $\psi = \frac{\Gamma}{2\pi} \ln(r)$
 - Doublet
 - $\phi = \frac{K \cos(\theta)}{r}$
 - $\psi = \frac{-K \sin(\theta)}{r}$
 - Converting from cartesian to cylindrical
 - $r = \sqrt{(x - h)^2 + (y - k)^2}$
 - $\theta = \tan^{-1}\left(\frac{y-k}{x-h}\right)$
 - Recall:
 - $u = \frac{\partial \phi}{\partial x} = \frac{\partial \psi}{\partial y}$
 - $v = \frac{\partial \phi}{\partial y} = -\frac{\partial \psi}{\partial x}$
 - $u_r = \frac{\partial \phi}{\partial r} = \frac{1}{r} \frac{\partial \psi}{\partial \theta}$
 - $u_\theta = \frac{1}{r} \frac{\partial \phi}{\partial \theta} = -\frac{\partial \psi}{\partial r}$
 - Where u_r and u_θ are the velocities in those given directions
 - Circulation
 - $\Gamma = \oint_{cs} \vec{V} \cdot d\vec{s}$
- Dimensional Analysis
 - Buckingham-pi theorem
 - A problem can be reduced to $k - r$ nondimensional groups
 - k : Number of variables
 - r : Number of dimensions (units) used
 - Reynolds number
 - $Re = \frac{\rho d V}{\mu} = \frac{d V}{\nu}$
 - Drag coefficient
 - $C_d = \frac{F_d}{\rho V^2 d^2}$ or $\frac{F_d}{\frac{1}{2} \rho V^2 A_c}$
 - Mach number
 - $\frac{V}{\sqrt{\gamma R T}}$
 - γ : Ratio of specific heats = 1.17 for air
- Boundary layer
 - Skin friction
 - $F = \tau A = \mu \frac{\partial u}{\partial y} A$
 - From Newtonian shear stress

- Viscous Flow in a Smooth Laminar Circular Pipe
 - Assumptions
 - Steady flow: $\frac{d}{dt} = 0$
 - Straight lines along pipe: $u_r = 0$
 - No swirls/spirals down the pipe: $u_\theta = 0$
 - Axisymmetric flow: $\frac{\partial}{\partial \theta} = 0$
 - Laminar flow: $Re < 1000$
 - After Navier-Stokes with simplifications
 - Continuity: $\frac{\partial u_z}{\partial z} = 0$
 - r – component: $\frac{\partial p}{\partial r} = 0$
 - z – component: $\frac{\mu}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_z}{\partial r} \right) = \frac{\partial p}{\partial z}$
 - After boundary conditions
 - $u_z(r) = \frac{1}{4\mu} \frac{\partial p}{\partial z} (r^2 - R^2)$
 - Where R is the radius of the pipe
 - $u_z(r) = \frac{1}{4\mu} \frac{\Delta p}{l} (R^2 - r^2)$
 - Pressure drop: $\Delta p = p_1 - p_2$
 - Volumetric flow rate
 - $Q = V_{mean} A = \int u_z dA = \frac{\pi R^4}{8\mu} \frac{\Delta p}{l}$
 - Where Δp is pressure drop ($p_1 - p_2$)
- Viscous Flow in a Rough Circular Pipe
 - Friction factor
 - $f = \frac{\Delta p}{\frac{1}{2} \rho V_{mean}^2} \cdot \frac{D}{l}$
 - Moody diagram

- Shows f as a function of Re and $\frac{\epsilon}{D}$
 - ϵ is the roughness of the pipe
 - If Laminer, $f = \frac{64}{Re}$