



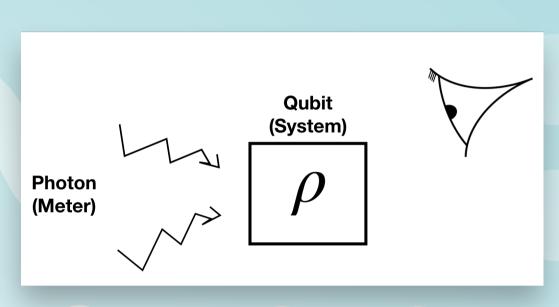
# Completely positive maps for higher-order unraveling of Lindblad master equations

N. Wonglakhon<sup>1</sup>, H. M. Wiseman<sup>1</sup>, A. Chantasri<sup>1,2</sup>

<sup>1</sup>Centre for Quantum Dynamics, Griffith University, Nathan, Queensland 4111, Australia <sup>2</sup>Optical and Quantum Physics Laboratory, Department of Physics, Faculty of Science, Mahidol University, Bangkok 10400, Thailand

## Motivation

In quantum measurement, the system of interest is measured by interacting with a meter (or a bath). For example, one can measure the evolution of a qubit using photons via homodyne detection.

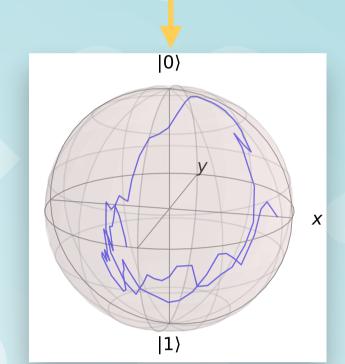


For continuous measurement, system's dynamics conditioned on a diffusive measurement records can be described by a quantum trajectory. The quantum trajectory can be alternatively simulated by using a map or measurement operator which has been proposed in several approaches [1-3].



Measurement records  $y_t$ 

However, the existing measurement operators suffer from finite time step resolution  $\Delta t$ , inducing errors in quantum trajectory simulation. We therefore derive a new approach which yields a measurement operator mitigating the time resolution problem.



Quantum trajectory

## Measurement operator properties

The measurement operator should satisfy the following properties:

Agree with Lindblad master equation, i.e.,

$$\int dy \wp_{\text{ost}}(y) \hat{M}_y \bullet \hat{M}_y^{\dagger} = e^{\Delta t \mathcal{L} \bullet} \approx (\bullet + \sum_{k=1}^n \frac{\Delta t^k}{k!} \mathcal{L}^k \bullet), \quad (1)$$

where  $\wp_{\text{ost}}(y)$  is the "ostensible" probability and  $\mathcal{L}$  is Lindblad superoperator. We will use the order of expansion n to measure this property. Noting that, in Eq. (1), one can change  $\approx$  to = if  $n=\infty$ .

Complete positivity (CP): the measurement operator should guarantee the trace preserving, i.e.,

$$\int dy \wp_{\text{ost}}(y) \hat{M}_y^{\dagger} \hat{M}_y = \hat{1} + o(\Delta t^k). \quad (2)$$

Noting that, for  $o(\Delta t^k) = 0$ , the map is complete positivity.

Closeness to the "exact" quantum trajectory: the exact individual trajectory ( $\rho_{\text{exact}}$ ) can be simulated in literatures for its particular measurement setting. We here calculate the "averaged" trace distance to measure this property, i.e.,

$$\langle D_j \rangle = \frac{1}{2} \langle \text{Tr} | \rho_j - \rho_{\text{exact}} | \rangle_{y_{\text{exact}}}, \quad (3)$$

where we average over all the exact records and  $\rho_j$  is generated by the maps.

## Higher-order completely positive map

Our new measurement operator is derived from a qubit system coupled to a bosonic (harmonic oscillator) bath via a unitary evolution  $\hat{U}(t+\Delta t,t)$ , projected to a homodyne eigenstate  $|\mathbf{e}_y\rangle$ , i.e.,

$$\hat{M}_y = \langle \mathbf{e}_y | \hat{U}(t + \Delta t, t) | \mathbf{e}_0 \rangle.$$

To construct  $\hat{M}_{\rm W}$ , we intuitively select terms which make the map satisfy the first two conditions to  $o(\Delta t^2)$ .















# Comparisons of Maps

Lindblad and CP properties

We make comparisons with 3 different existing maps which has been proposed in literatures in different intuitions, namely, the conventional Ito map  $(\hat{M}_{\rm I})$  [1], Rouchon- Ralph map  $(\hat{M}_{\rm R})$ , having terms quadratic in y [2], and Guevara-Wiseman map  $(\hat{M}_{\rm G})$ , designed for high-order CP [3], to our new map  $(\hat{M}_{\rm W})$ . By considering the properties in Eqs. (1)-(2), we have the analytical results of order of satisfaction as the following:

Map	Lindblad $\mathcal{L}^n$ : $n$ up to 2	CP: k  up to  2
$\hat{M}_{ m I}$	X	×
$\hat{M}_{ m R}$	X	×
$\hat{M}_{ m G}$	X	
$\hat{M}_{ m W}$		

### Quantum trajectory closeness (qubit examples)

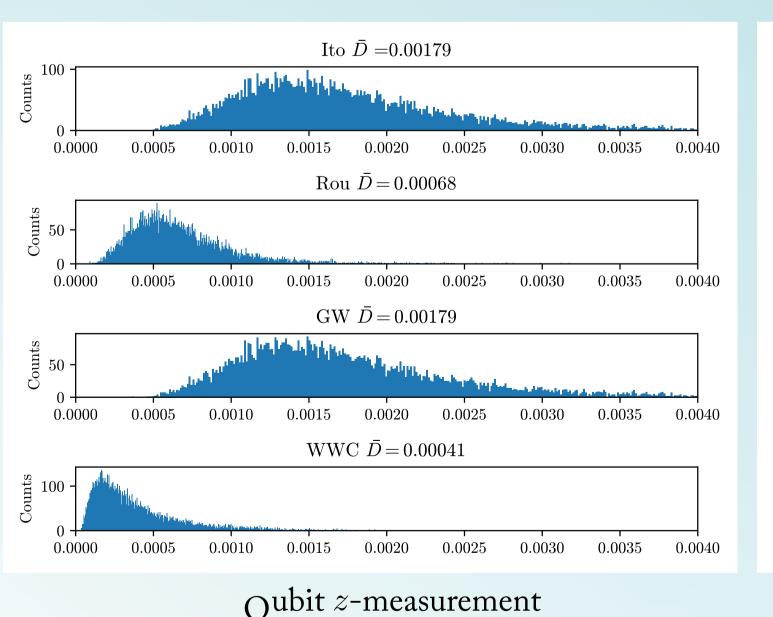
We calculate the averaged trace distance in Eq. (3) of individual trajectories generated by each map  $(\rho_j)$  compared to the trajectories generated by the exact map for two particular examples, namely qubit z-measurement with the Lindblad operator  $\hat{c} \propto \hat{\sigma}_z$  and the measurement of qubit's fluorescence  $(\hat{c} \propto \hat{\sigma}_-)$ .

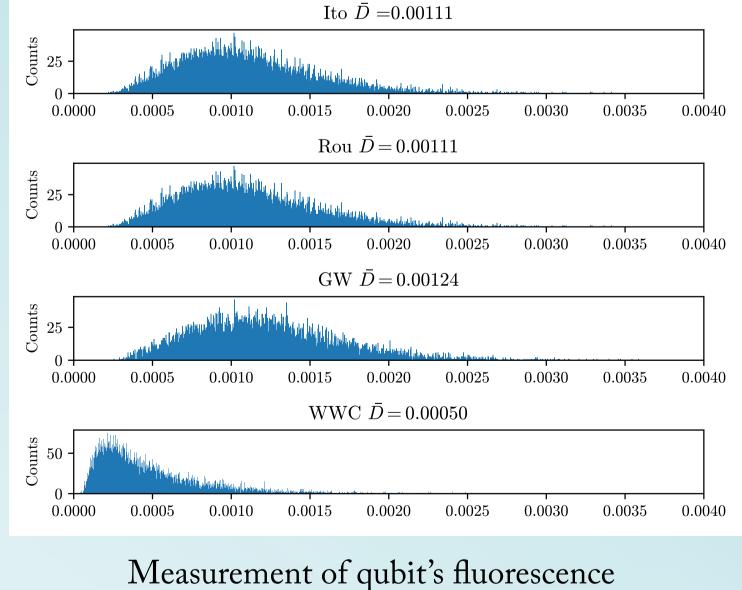
Map	Qubit z-measurement	Qubit's fluorescence
$\langle D_{ m I}  angle$	$0.15\alpha$	0.50 eta
$\langle D_{ m R}  angle$	$0.09\alpha$	$0.50\beta$
$\langle D_{ m G}  angle$	$0.15\alpha$	$0.50\beta$
$\langle D_{ m W}  angle$	$0.04\alpha$	$0.25 \beta$

Here  $\alpha, \beta$  are the first order constant with  $\alpha, \beta \sim \mathcal{O}(\Delta t^{3/2})$ .

#### Simulation results

The histograms show the numerical calculations of the trace distance of the trajectories generated by each map compared to the exact trajectories for two qubit experiments.





#### Discussion

We construct the new map which satisfies analytically both the Lindblad master equation and CP to  $\Delta t^2$ . Interestingly, although the new map derived to satisfy the first two conditions, it gives the closest individual trajectory to the exact one both analytically and numerically despite not high order in  $\Delta t$ .

#### References

- [1] H. M. Wiseman and G. J. Milburn, Quantum measurement and control (Cambridge University Press UK, 2010).
- [2] P. Rouchon and J. F. Ralph, Phys. Rev. A 91, 012118(2015).
- [3] I. Guevara and H. M. Wiseman, Phys. Rev. A 102, 052217 (2020).