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UNIVERSITAT POLITÈCNICA DE CATALUNYA

# Cubesat Constellation Astrea

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## Report

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# Contents

<b>List of Tables</b>	<b>ii</b>
<b>List of Figures</b>	<b>iii</b>
<b>I Orbit Design</b>	<b>1</b>
<b>1 Orbit Geometry</b>	<b>2</b>
. . . . .	3
1.1 Keplerian Geometry . . . . .	4
. . . . .	4
1.2 Dynamic equations . . . . .	6
. . . . .	6
<b>2 Orbital Coverage</b>	<b>7</b>
2.1 Satellite Footprint . . . . .	8
2.1.1 Introduction . . . . .	8
2.1.2 Footprint Computation . . . . .	8
2.2 Elevation Angle . . . . .	10
. . . . .	10
2.2.1 Elevation angle cone . . . . .	10
. . . . .	10
2.2.2 Atmospheric restrictive conditions . . . . .	11
. . . . .	11
. . . . .	11
. . . . .	11
2.2.3 Elevation angle of other current constellations . . . . .	13
. . . . .	13
. . . . .	14
2.3 Minimum Plane Inclination . . . . .	15
In conclusion . . . . .	15
2.4 Satellite to Satellite Visibility . . . . .	16

. . . . .	16
2.5 Market Study: Current Nanosatellites in Orbit . . . . .	18
2.5.1 Criteria for the orbital height of the satellites . . . . .	18
Satellites currently in Orbit	
. . . . .	18
The most interesting potential clients	
. . . . .	19
2.5.2 New Space: Adapting to new society needs . . . . .	19
<b>3 Constellation Configuration</b>	<b>20</b>
3.1 Introduction: The Global Positioning System Example . . . . .	21
3.2 Polar Orbit Constellation . . . . .	22
3.2.1 Introduction . . . . .	22
3.2.2 General Configuration . . . . .	23
3.2.2.1 The Streets of Coverage Method . . . . .	23
3.2.3 Results of Streets of Coverage . . . . .	26
3.3 Walker-Delta Constellation . . . . .	28
3.3.1 Full Walker-Delta Constellation . . . . .	28
3.3.1.1 Characteristics . . . . .	28
. . . . .	29
3.3.1.2 Notation . . . . .	29
3.3.1.3 Coverage . . . . .	29
3.3.2 Semi Walker Delta Configuration . . . . .	30
. . . . .	30
3.3.2.1 Advantages . . . . .	30
. . . . .	30
3.3.2.2 Disadvantages . . . . .	31
. . . . .	31
3.3.3 Other Walker Delta Configurations . . . . .	31
. . . . .	31
3.3.3.1 SWDC including an additional polar orbit. . . . .	32
. . . . .	32
3.3.3.2 Mixed Walker Delta. . . . .	32
. . . . .	32
. . . . .	33
3.4 Testing Method . . . . .	35
3.4.1 Introduction . . . . .	35
3.4.2 Second Subsection . . . . .	35
<b>4 Orbit Perturbations</b>	<b>36</b>
4.1 Sources of Perturbation . . . . .	37

4.1.1	Introduction to Orbit Perturbations . . . . .	37
	The Approach of the Perturbations Study . . . . .	38
4.1.2	Gravity Potential of Earth . . . . .	38
4.1.3	Atmospheric Drag . . . . .	39
4.1.4	3rd Body Perturbations . . . . .	40
4.1.5	Other Perturbations . . . . .	40
4.2	Significant Perturbations . . . . .	41
	Propagation Algorithm . . . . .	41
	In conclusion . . . . .	41
4.3	Orbit Decay . . . . .	42
4.3.1	Effects on the Ascention Node . . . . .	42
	4.3.1.1 Introduction . . . . .	42
	4.3.1.2 Second Subsection . . . . .	42
4.3.2	Effects of the Solar Cicle . . . . .	42
	. . . . .	42
4.3.3	Orbital Decay Propagation Results . . . . .	45
	4.3.3.1 Introduction . . . . .	45
	4.3.3.2 Second Subsection . . . . .	45
4.3.4	Dynamic Orbit Decay Computation . . . . .	45
	4.3.4.1 Introduction . . . . .	45
	4.3.4.2 Second Subsection . . . . .	45
4.4	Orbital Station-Keeping . . . . .	46
4.4.1	Raising the orbit height to increase Lifetime . . . . .	46
4.4.2	Using Thrusters to increase Lifetime . . . . .	47
	4.4.2.1 Energy equation . . . . .	47
	. . . . .	47
	. . . . .	48
	4.4.2.2 Delta-V . . . . .	48
	4.4.2.3 Time . . . . .	49
	4.4.2.4 Propellant . . . . .	49
	4.4.2.5 Orbit maintenance . . . . .	49
	4.4.2.6 Results . . . . .	50
	. . . . .	51

# List of Tables

1.1.1	Eccentricity values depending on the shape of the orbit . . . . .	4
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# List of Figures

1.0.1	Geocentric-equatorial frame. Extracted from [?]. . . . .	3
1.1.1	Geocentric-equatorial frame and the Classical Orbital Elements. Extracted from [?]. . . . .	5
2.1.1	Single satellite coverage geometry . . . . .	8
2.2.1	Elevation angle cone. Source: NOAA . . . . .	10
2.2.2	Minimum elevation angle as function of latitude. Source: [a general evaluation criterion] . . . . .	14
2.3.1	Minimum Inclination to provide coverage at different latitude for different orbit apogees. . . . .	15
2.5.1	Distribution of the currently in orbit nanosatellites. . . . .	18
3.1.1	Distribution of the expanded 24-slot GPS constellation. . . . .	21
3.2.1	Distribution of the 66 Iridium constellation satellites. . . . .	22
3.2.2	Distribution of the planes for Polar Orbits design. . . . .	23
3.2.3	Single plain street of coverage. The footprints of the satellites superpose leading to a street. . . . .	24
3.2.4	Two plains streets of coverage. An optimum phasing needs to be obtained.	24
3.2.5	Variation of number of satellites for different heights and elevation angles	26
3.2.6	Variation of number of satellites for different heights between 500 and 600km. . . . .	27
3.3.1	Definition of the inclination $\delta$ . Extracted from [?] . . . . .	28
3.3.2	Delta pattern as seen from the North Pole. Extracted from [?] . . . . .	29
3.3.3	Delta pattern 65°: 30/6/1 . . . . .	30
3.3.4	Minimum altitude for continuous global coverage. Comparison between polar patterns and Walker delta patterns. Extracted from [?] . . . . .	31
3.3.5	12 plane SWDC. Note the gap and the equidistant planes . . . . .	32
3.3.6	This geometry distribution induces a large anti-symmetric gap . . . . .	33
3.3.7	Added polar orbit to the 11 plane based SWDC . . . . .	34
3.3.8	8 plane based MWDC generated for 210 degrees . . . . .	34
4.2.1	Logaritmnic plot of the modulus of the increases in Angular Arguments of the orbit . . . . .	41

## LIST OF FIGURES

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4.3.1	Deviation of densities in the upper atmosphere due to the 19th solar cycle	43
4.3.2	Measured intensities of the 23rd and 24th solar cycles. Source: NOAA	44
4.4.1	Increase in the Lifetime obtained by setting the constellation in a higher orbit . . . . .	46
4.4.2	Hohmann transfer. Extracted from [?] . . . . .	48
4.4.3	Height variation of the satellite . . . . .	51
4.4.4	Height variation of the satellite with a more restrictive minimum height	52



**Part I**

**Orbit Design**

## Chapter 1

# Orbit Geometry

Throughout this chapter, the bases of orbital geometry will be explained in order to correctly understand the parameters that will later be exposed when dealing with the constellation orbits (or the position of the satellites in them). However, long theoretical explanations will be avoided so as not to distract the reader from the main objective of the project.

To understand the movement in space is enough to apply the Newton's laws. These, however, need an inertial non-rotating frame to be correctly described. When dealing with Earth-orbiting, one usually chooses a reference system called *geocentric-equatorial system* which is shown in the figure 1.0.1 As can be seen, the XY plane coincides with the plane Equatorial with the X axis pointing in the direction of the vernal equinox <sup>1</sup>. The Z axis correspond the axis of rotation of the earth and points to the north (following the right-hand rule).

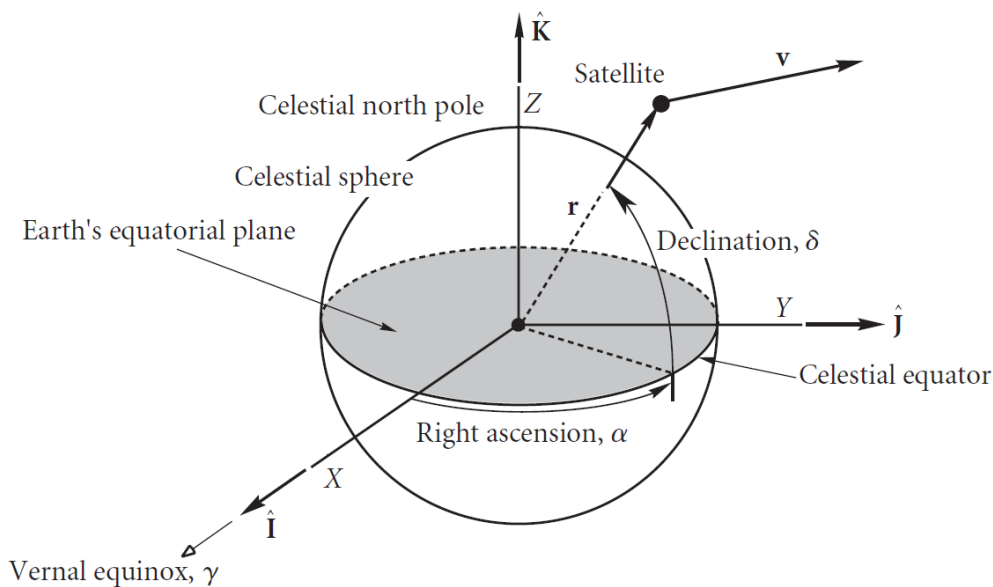


Figure 1.0.1: Geocentric-equatorial frame. Extracted from [?].

By defining this system, any point in the space can be depicted by its position vector  $r$  and we can study its movement by the velocity vector  $\dot{r}$ . These elements are useful especially for computational work but they nearly do not provide information about the orbit. For these reason, the orbital elements were developed.

<sup>1</sup>an imaginary line found by drawing a line from the Earth to the Sun on the first day of spring

## 1.1 Keplerian Geometry

The *Classical Orbital elements*, also known as the *Keplerian elements* as an attribution to Johannes Kepler, are six independent quantities which are sufficient to describe the size, shape and orientation of an orbit. This set of elements are shown in the figure 1.1.1 and are defined as follows:

- **Semi-major axis** ( $a$ ): It is related to the size of the orbit and is defined by the sum of the apogee (furthest point) and the perigee (closest point) divided by two.
- **Eccentricity** ( $e$ ): It defines the shape of the orbit with respect to that of a circle. Thus, the eccentricity of a circular orbit is null while hyperbolic orbits have an eccentricity greater than one.

Circular	$e = 0$
Elliptical	$0 < e < 1$
Parabolic	$e = 1$
Hyperbolic	$e > 1$

Table 1.1.1: Eccentricity values depending on the shape of the orbit

- **Inclination** ( $i$ ): the inclination is the angle between the positive Z axis and the angular momentum vector ( $\mathbf{h}$ ) which is perpendicular to the orbital plane. The inclination of the orbit can take a value from 0 deg to 180 deg. For  $0 \text{ deg} \leq i \leq 90 \text{ deg}$  the motion is *prograde* and for  $90 \text{ deg} \leq i \leq 180 \text{ deg}$  the motion is *retrograde*.
- **Right ascension of the ascending node - RAAN** ( $\Omega$ ): This parameter, along with the inclination define the orientation of the orbital plane. It is the angle between the positive X axis and the intersection of the orbital plane with the equatorial plane XY in counterclockwise direction. The intersection mentioned is called the node line and the point where the orbit passes through the node line (from south to north) is the ascending node ( $0 \text{ deg} \leq \Omega \leq 360 \text{ deg}$ ).
- **Argument of perigee** ( $\omega$ ): Is defined as the angle between the ascending node and the perigee. It describes the orientation of the ellipse with respect to the frame ( $0 \text{ deg} \leq \omega \leq 360 \text{ deg}$ ).
- **True Anomaly** ( $\phi$ ): This last quantity is used to describe the satellite's instantaneous position with respect to the perigee. It is the angle, measured clockwise, between the perigee and the satellite position. From all the orbital elements, the true anomaly is the only that changes continuously. Sometimes, true anomaly is substituted by the mean anomaly, which can be calculated using another auxiliary

angle called the eccentric anomaly.

$$\begin{aligned}\cos E &= \frac{e + \cos \theta}{1 + e \cos \theta} \\ M &= E - e \sin E\end{aligned}\tag{1.1.1}$$

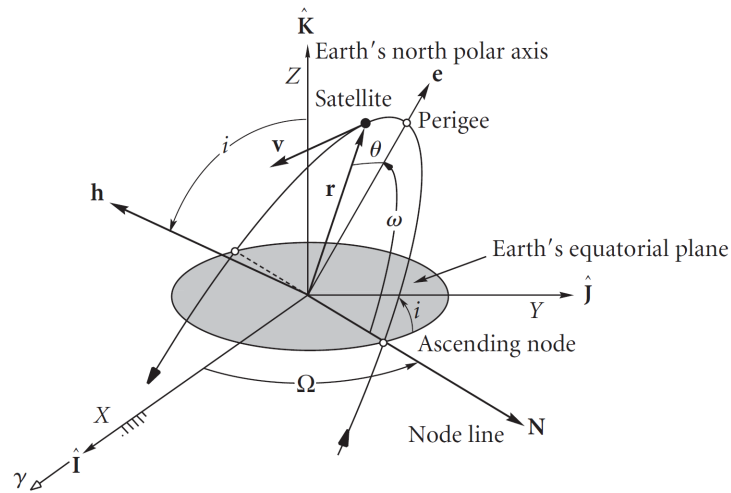


Figure 1.1.1: Geocentric-equatorial frame and the Classical Orbital Elements. Extracted from [?].

## 1.2 Dynamic equations

As aforementioned, the motion of an object in the space can be described using the Newton's laws. The basic idea developed by Newton is to study the Cubesat and the Earth as a spherical bodies in mutual gravitational attraction and neglect the gravitational forces caused by other objects (this is called the *two body* problem). The forces balance is simple since we only have the Earth gravitational attraction, which must compensate the centripetal acceleration of the satellite. Thus, using the law of universal gravitation,

$$-G \frac{M_E m_{sat}}{r^3} \vec{r} = m_{sat} \vec{a}_{sat} \quad (1.2.1)$$

Where  $G$  is the gravitational constant and  $r$  represents the distance between the satellite and the Earth. From the last equation, we only want to obtain the acceleration, therefore:

$$-G \frac{M_E}{r^3} \vec{r} = \vec{a}_{sat} = \frac{d^2 \vec{r}}{dt^2} \quad (1.2.2)$$

For simplicity, it is usual to denote  $\mu = GM_{earth}$  resulting in the following equation:

$$-\frac{\mu}{r^3} \vec{r} = \frac{d^2 \vec{r}}{dt^2} \quad (1.2.3)$$

This expression is a second order equation that models the motion of the Cubesat relative to the Earth and it can be analytically solved. The only problem is that several hypotheses have been applied that make the case different from reality. The formulation should be modified to take into account the effects due to:

- More bodies attracting the satellite (Sun, Moon, Venus, etc.)
- The existence of more forces like the drag, the solar radiation pressure, etc.
- The earth is not an spherical body.

The corrections for considering these things are called perturbations and they are explained in the Chapter ?? of this part of the report.

## Chapter 2

# Orbital Coverage

## 2.1 Satellite Footprint

### 2.1.1 Introduction

The first step to build a satellite network with global coverage is to compute a single satellite footprint.

The footprint of a satellite is defined as the region of Earth where a single satellite can be seen. This Earth coverage surface provided is spherical and depends on some orbital parameters such as:

- Height

When increasing height the footprint of a satellite grows.

- Elevation angle

When increasing the elevation angle, which is the angle between the satellite and the horizontal plane of an arbitrary point of the Earth, the surface seen by the satellites decreases. (This parameter will be later studied in detail)

### 2.1.2 Footprint Computation

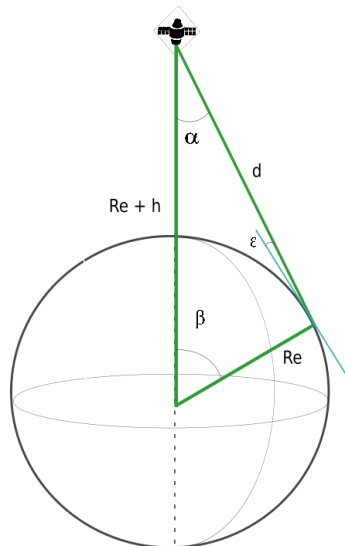


Figure 2.1.1: Single satellite coverage geometry

In order to compute the coverage area we must solve the triangle depicted in figure 2.1.1 where the basic geometry of a satellite footprint is shown.

The most needed parameters are the distance from a random point on Earth (where we can suppose our ground station to be) to the satellite denoted by  $d$  and the central angle, denoted with a  $\beta$ .



## Satellite Footprint

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Applying cosines law to the triangle shown in figure 2.1.1, we obtain the following expression:

$$r^2 = R_{earth}^2 + d^2 - \cos(90 + \epsilon) \quad (2.1.1)$$

Isolating  $d$  from the equation above and changing  $r = R_{earth} + h$ , where  $h$  is the actual height of the satellite regarding the Earth surface, we arrive at:

$$d = R_{earth} \left[ \sqrt{\left( \frac{h + R_{earth}}{R_{earth}} \right)^2 - \cos^2 \epsilon} - \sin \epsilon \right] \quad (2.1.2)$$

From the figure 2.1.1 we can also extract a relation between the central angle, the distance  $d$  and the elevation angle. This relation together with the equation 2.1.2 allow us to find  $\beta$ .

$$\begin{aligned} d \cos \epsilon &= (R_{earth} + h) \sin \beta \\ \beta &= \frac{1}{R_{earth} + h} \arcsin [d(\epsilon) \cos \epsilon] \end{aligned} \quad (2.1.3)$$

Once the central angle  $\beta$  has been computed we are able to obtain the footprint satellite's area using the equation below:

$$S = 2\pi R_{earth}^2 (1 - \cos \beta) \quad (2.1.4)$$

The size of the footprint will determine the level of coverage our constellation provides, therefore when deciding the value of the orbital parameters it has to be a factor to consider.

## 2.2 Elevation Angle

The angle of elevation is essential to calculate the geometry of our constellation. As discussed previously, our aim in this project report is to justify how global coverage will be fulfilled. First, we define for a given groundstation the angle between its beam pointing right to the satellite and the horizontal local plane as the elevation angle. Secondly, a study is conducted in order to relate the height of the satellite, the elevation angle and the coverage of the Earth. Finally, we complete our orbital design by configuring a constellation that will securely define a global coverage fulfillment. Next, we will be defining how these parameters are related.

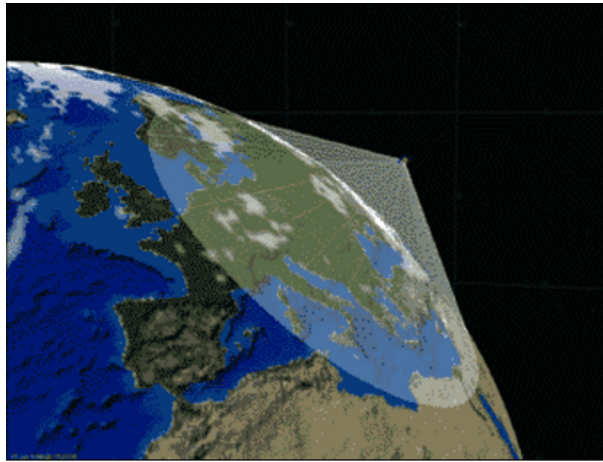


Figure 2.2.1: Elevation angle cone. Source: NOAA

### 2.2.1 Elevation angle cone

Global coverage will be discussed considering the elevation angle and its resulting footprint on Earth. The elevation angle is described by the angular orientation of the antennas in the ground station. However, this angle is also perceived by the satellite in a similar way - it will vary depending on the orientation of the satellite and the angle between horizontal local planes. In order to describe the footprints we must define a cone which vertex is set at the antennas of the satellite, pointing down to Earth, and which generatrix is given by the angle of elevation. This elevation angle based cone is the description of the paths that our communications can take place. In other words, the generatrix of this cone is setting the limits in which the antenna will operate as function of the elevation angle. This implies that our satellite will be able to communicate to all the points contained in the cone. Finally, this cone will be describing a circular surface on top of the Earth which we will call the footprint of the satellite. Additionally, this footprint is the coverage that a single satellite can generate, hence we will be distributing satellites all around the Earth in order to fulfill global coverage.

### 2.2.2 Atmospheric restrictive conditions

In order to obtain the final restrictive angle of elevation needed to contact the ground stations some considerations have to be made. First, a description of the different atmospheric conditions will be defined. Then, we will relate these to our bandwidth in order to analyse if they must be taken into account when communicating with ground stations. [elec2013cantero]

- **Atmospheric gases:** water vapour and oxygen absorptions; important when frequencies are above 3 GHz. More information [64] and [07328546]

- **Precipitations and Clouds:** these conditions are relevant for signals above 10GHz.

- **Cross Polarization Discrimination:** direct consequence of both terrestrial links and rain. Related to non-spherical rain drops which have a polarization rotated towards the component of the major axis, and hence may attenuate a signal wave.

- **Scintillation** is a rapid fluctuation in signal amplitude at low elevations.

- **Radio Refracting Index:** for elevation angles below 3 degrees (especially those below 1 degree) and depending on the latitude of our satellite we may find big signal losses due to the resulting differential ray bending.

- **Ionosphere layers:**

**D layers:** 60-90km. Considerable signal absorptions for 10 MHz and below, with progressively less absorption at higher frequencies and oblique incidencies.

**E layers:** 90-150 km. Absorptions relevant for frequencies lower than 10 MHz, although for sporadic E propagations this value may be increased to 50MHz.

**Sporadic E layers:** Reflections of radio waves in this thin-cloud small layer may reach to frequencies up to 225MHz. These layers are usually formed following the E layers altitudes.

**F layers:** 150-500 km and higher. No absorptions or reflections for these layers. The F2 region allows the longest communication paths, above 210km of altitude.

By means of these physical phenomena we can subtract the elevation angle as function of the latitude. However, we must take into account that these physical conditions give

## **Elevation Angle**

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a value for the elevation angle which may not be the most restrictive. Global coverage conditions, bandwidths, inclination and the final distribution of our constellation will be considering this elevation angle and viceversa, iteratively.

The ASTREA CONSTELLATION was designed and optimized in order to fulfill global coverage for a constant elevation angle - respect to the latitude - of 20 degrees.

Our constellation will be operating at S-band for telemetry and X-band for data relay. Therefore, the satellites need to be operating up to 10 GHz. This directly implies that physical conditions such as atmospheric gases, precipitations and clouds must be studied when determining the elevation angle needed.

The minimum elevation angle is applied in low latitude regions for constellations based on polar orbits whereas this value is also applied out of the low-latitude region for inclined orbit constellations [a general evaluation criterion]. The minimum elevation angle is a specific value which is equivalent to the maximum elevation angle needed to fulfill coverage at a given latitude, considering that the distance between planes is maximum at the equator and that it is reduced for higher latitude positions.

This elevation angle is maximum in a Walker Delta constellation when the latitude is equal to the orbital inclination angle[a general evaluation criterion]. This means that the limiting restrictive elevation angle that we need in order to fulfill global coverage is defined at latitude equal to 72 degrees, which is the inclination of our constellation. Otherwise, we can define a constant elevation angle that will apply to the equator, which will then be, for this model, the restrictive condition.

Accordingly, the approach considered is that of a constant elevation angle to fulfill global coverage at 20 degrees. This implies that our constellation is configured and distributed in order to optimize coverage both at the equator and at the maximum elevation angle latitude. This value has been contrasted and discussed considering the atmospheric conditions and analysing experimental data, which contemplates also the rotation of the Earth among others.

This constant elevation angle model will be very useful in order to analyse and calculate the distribution of the constellation. Nevertheless, we need to describe in an accurate way the minimum elevation angle respect to the latitude. This is why a different model must be approached.

Thus, we need to describe the elevation angle respect to the latitude of our constellation taking into account all considerations above. First, for a latitude of 0 degrees the value of the minimum elevation angle will be of 10 degrees. In our model we have considered that

this value was of constant 20 degrees, so in fact we have redundant global coverage. At latitudes between the equator and 45 degrees our second model increases linearly to 15 degrees. From 45 to 60 degrees the elevation angle also increases linearly to 22 degrees. Then, from 60 to 70 the value increases highly reaching a peak at 70 degrees, where the elevation angle will be of 30 degrees. Finally, from 70 to 80 degrees this model is reduced linearly to 15 degrees, and from 80 to the north and south poles it falls to 0 degrees. This is a simple model that will guarantee global coverage, especially at the latitudes of our ground stations

For the distribution of ground stations we need to guarantee that these will be covered either by one satellite or two at any given moment. As discussed before, the model used was based on a constant 20 degree constant elevation angle. However, for this last model that we have described - which is more realistic - we obtain more coverage than for the constant model except for those regions next to the peak. The most restrictive latitude is now 60 degrees - where all the ground stations are set - and has a 22 degree restriction of the angle of elevation, which is higher than the constant model described previously. These facts imply the following:

- At low latitudes (between 0 and 30 degrees) the constellation fulfills global coverage generously.
- At ground station latitude (60 degrees) the constellation is covering the station successfully. As discussed before, our first model considered a constant 20 degree elevation angle instead of the 22 degrees that now must be corrected. For the previous model coverage was well established with margin. For the latter, the margin has decreased but coverage is still complete. Note: each orbit could be reduced by a number of satellites per plane, but this would endanger the correct and stationary working of the constellation. In this case we would not be able to control possible incidencies such as unoperative satellites with enough margin.
- The ground stations are covered at all time for at least one satellite.

### **2.2.3 Elevation angle of other current constellations**

Analysing the minimum elevation angle needed in order to fulfill global coverage requires, as mentioned before, the understanding first of the restrictive conditions of the atmosphere and how these will alter it. As a consequence of the different physical conditions given before we will be able to determine a relation between latitude and elevation angle. All the same, the elevation angle depends on the bandwidth in which the satellites operate, hence different distributions of this angle respect to the latitude will be

## Elevation Angle

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described depending on the bandwidths used.

- Celestri: 18.8 to 20.2 GHz at 48 degree inclination.
- GlobalStar: 2.4 GHz at 52 degree inclination.
- Iridium: 20 to 30 GHz at 90 degree inclination - polar orbits.

Comparing our configuration to other present constellations some clarifications can be made:

- The minimum elevation angle peak is proportional to the bandwidth at which the satellite is communicating with Earth. For instance, Iridium's peak of elevation angle is the highest relative to the other configurations since it is also working with the highest frequency signals.
- The latitude position of the peaks is related to the inclination of the constellation. Iridium, - a polar orbit based configuration - describes a peak at 90 degrees of latitude whereas Celestri and GlobalStar are near 40 to 50 degrees.

With these tendencies our model can be confirmed as function of the frequencies of the signals and related to the inclination of the orbits.

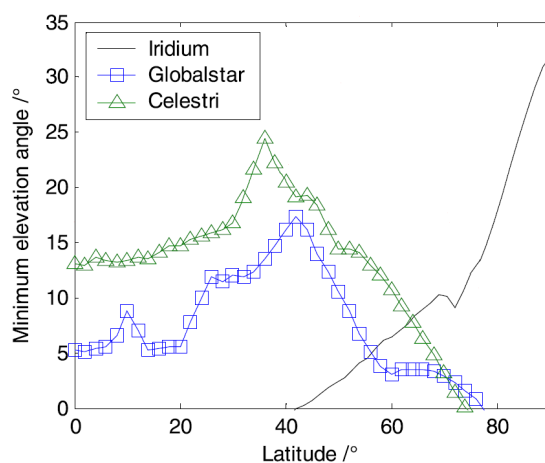


Figure 2.2.2: Minimum elevation angle as function of latitude. Source: [a general evaluation criterion]

## 2.3 Minimum Plane Inclination

As it has been pointed before, there are several factors to take into account in order to design a constellation that provides global coverage on Earth. In this section the minimum inclination to achieve that purpose is assessed. Using the theory previously developed, we can observe the following results:

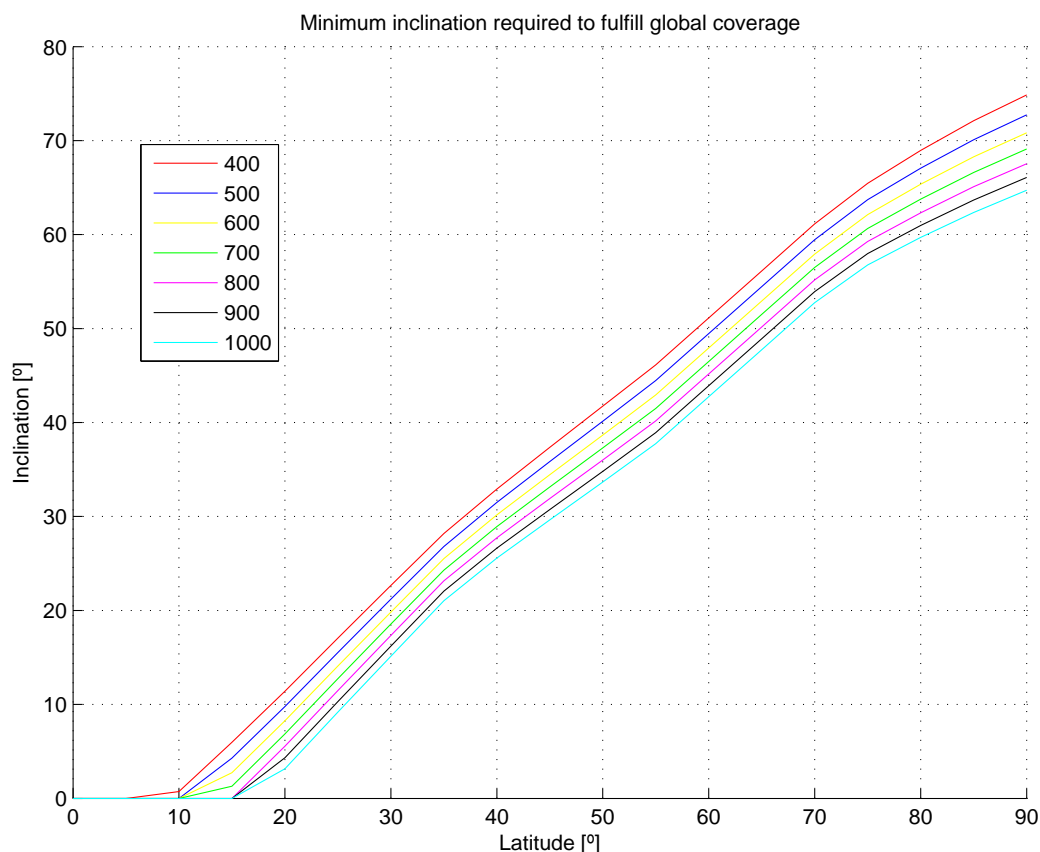


Figure 2.3.1: Minimum Inclination to provide coverage at different latitude for different orbit apogees.

As it can be observed, if the goal of the design is to provide full global coverage, the distribution of elevation angles with latitude is not significant, since the inclination is required to be higher than approximately  $75^\circ$ . In the other cases, the change of minimum elevation angle distribution causes changes of tendency in the distribution of inclination required.

### In conclusion

The main point is that there is a limit inclination for a Walker-Delta constellation configuration in order to provide global coverage at the desired latitude. With this study, this limits in the design algorithms can be set.

## 2.4 Satellite to Satellite Visibility

One of the restrictive conditions that we must take into account is the visibility between satellites. Communications among different satellites is the key point of our constellation. Therefore, this has to be guaranteed considering a model which will represent the conditions of the atmosphere for LEO communications.

In order to fulfill communications among satellites we must consider that a straight beam can be described between two consecutive satellites, which will then communicate with others. These two satellites will need to be at a distance such that the Earth itself doesn't interfere in this straight beam. Depending on the bandwidth of our constellation we will also have to consider that this communication beam will not interfere with a given element of the atmosphere such as the upper layers of the ionosphere. Thus, a model will be developed in order to limit the minimum altitude at which this beam is guaranteed to pass through safely.

This model is a restrictive condition that we need to satisfy when designing our constellation. The highest restrictive conditions are the upper layers of the ionosphere, specifically the E layers at 150 km above the surface of the Earth. Reflections and absorptions can occur for both E layers and sporadic E layers. E layers may reflect signals of frequencies below 10 MHz whereas Sporadic E layers can be a problem up to 225 MHz. Working for S bands and X bands implies that neither of these layers will alter the signals of our constellation.

Operating and computing with these conditions a maximum distance is obtained which defines how far these satellites can be from each other. A simple equation is used to calculate this distance considering the height of the satellites and the height of the E layers in the atmosphere.

$$d = 2\sqrt{(R + h_{sat})^2 - (R + h_{atm})^2}$$

$$h_{sat} = 550 \text{ km}$$

$$h_{atm} = 150 \text{ km}$$

$$R = 6371 \text{ km}$$



## Satellite to Satellite Visibility

---

The final expression for the distance between two satellites indicates that distance between two satellites has to be smaller than 4640 km approximately. For this result we conclude that this restrictive condition is actually less restrictive than the 9 planes needed for our constellation. Thus, satellite to satellite visibility is a parameter which will not affect the design of our constellation after all.

## 2.5 Market Study: Current Nanosatellites in Orbit

### 2.5.1 Criteria for the orbital height of the satellites

#### Satellites currently in Orbit

If only geometric considerations were to be applied in the design of a satellite constellation, it is clear that the higher the orbit the broader is the footprint in the surface leading to a smaller number of satellites. However, if the service of communications is to be offered, the satellites currently in orbit or in design phases need to be at higher orbit than the one of the constellation. The purpose of that requirement is to intersect the field of view of the satellites that nowadays point to Earth.

From source [?] we can study how the currently on orbit satellites are launched and specially, in which orbits. The results of the study of this source is presented below. All of them are in Low Earth Orbits, and half of them above 550km. In total, there are 203 operational satellites.

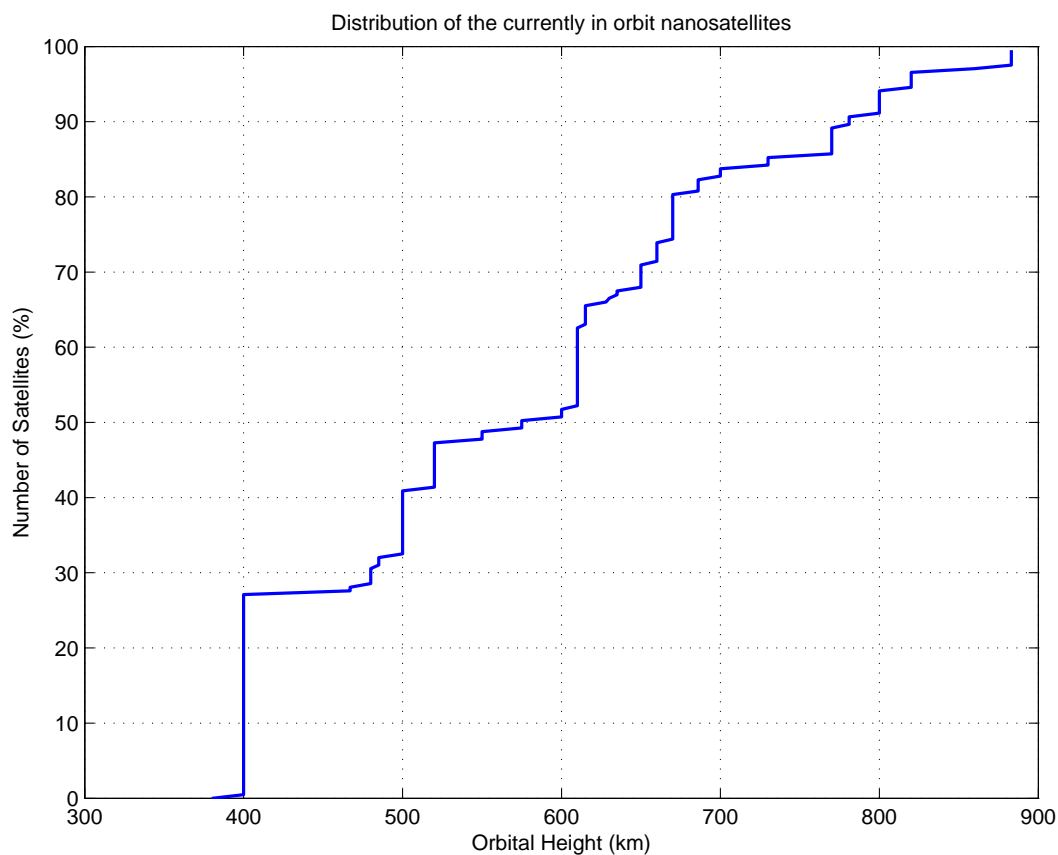


Figure 2.5.1: Distribution of the currently in orbit nanosatellites.

**The most interesting potential clients**

Lots of satellites are orbiting at heights lower than 500km, mainly because one of the most feasible way of launching a small satellite is from the International Space Station. However, this very low LEOs are related to very high speeds and specially to low lifetimes, since drag affects them in a more significant way. To the interest of the constellation, the satellites at higher altitudes are a better commercial target, since they are going to be in orbit for longer missions. In addition, the same orbit decay problems are avoided for the constellation satellites.

**2.5.2 New Space: Adapting to new society needs**

Nowadays new satellites willing to provide services to Earth are being positioned closer than ever. Where closer can be applied in many points of view. Physically, the satellites are placed every time at lower orbits, since the energetic requirement is lower. Technically, the space certified materials and hardware are becoming more feasible, and new launchers are smaller. In the end, everything comes down to an economic approach, launching satellites is becoming cheaper every time and this means closer to the private pocket.

In the future, the possibility of using the Astrea constellation to contact Earth can reduce the requirements for the antennas and AOCSs to communicate with ground, leading to a whole new level of resources for the satellite payload. For instance, by communicating to the constellation pointing to outer space instead of pointing down to Earth. That is just a way in which Astrea is in the New Space Generation.

**In conclusion,** In the decision process one of the statistics considered with certain weight will be the following: the ratio of satellites at which the constellation will be able to provide service considering that nowadays all of them point down to Earth.

## Chapter 3

# Constellation Configuration

*"Our two greatest problems are gravity and paperwork. We can lick gravity, but sometimes the paperwork is overwhelming."*

---

Werner von Braun, 1958

### 3.1 Introduction: The Global Positioning System Example

Depending on the application the Space Segment of a mission can vary in an infinite number of ways. Probably the most famous and widely used satellite constellation is the the Global Positioning System satellite network. In this case, it uses an irregular geometry.

#### **The GPS Constellation: An example of irregular distributed orbits**

The GPS is a constellation property of the U.S. It provides positioning, navigation and timing. The constellation was designed with a 24-slot arrangement to ensure a visibility of at least four satellites from any point on the planet. Nowadays the constellation has expanded to a total operative number of 27-slot since June 2011. Some characteristic parameters of the satellites are the following:

- Orbit: Almost Circular
- Height = 20,200 km (MEO);
- Lifetime = 12.5 years;
- Satellite Cost = 166 million USD;
- Inclination =  $55^\circ$ ;
- Number of planes = 6;
- Phasing:  $30^\circ$ - $105^\circ$ - $120^\circ$ - $105^\circ$ ;

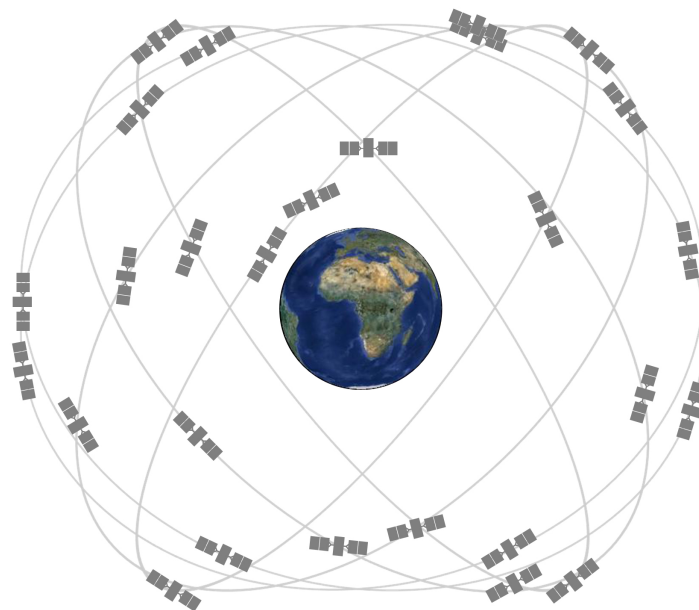


Figure 3.1.1: Distribution of the expanded 24-slot GPS constellation.

## 3.2 Polar Orbit Constellation

### 3.2.1 Introduction

Polar Orbits are probably the simplest way to configure an evenly spaced constellation. As we will see in the section **Orbit Perturbations** when the inclination is the same for all the planes, the deviations tend to be the same for all the satellites. In addition, the computation of the number of satellites required is also easier.

#### **The Iridium Constellation: An example of near polar orbits**

The Iridium constellation is a private constellation. It provides voice and data coverage to satellite phones among other services. The constellation was designed with 77 satellites, giving name to the constellation by the chemical element. The constellation was reduced to a number of 66. Sadly, Dysprosium is not such a good commercial name. Some characteristic parameters of the satellites are the following:

- Orbit: Almost Circular
- Height = 781 km (LEO);
- Satellite Cost = 5 million USD;
- Inclination =  $86.4^\circ$ ;
- Number of planes = 11;
- Phasing: Regular;

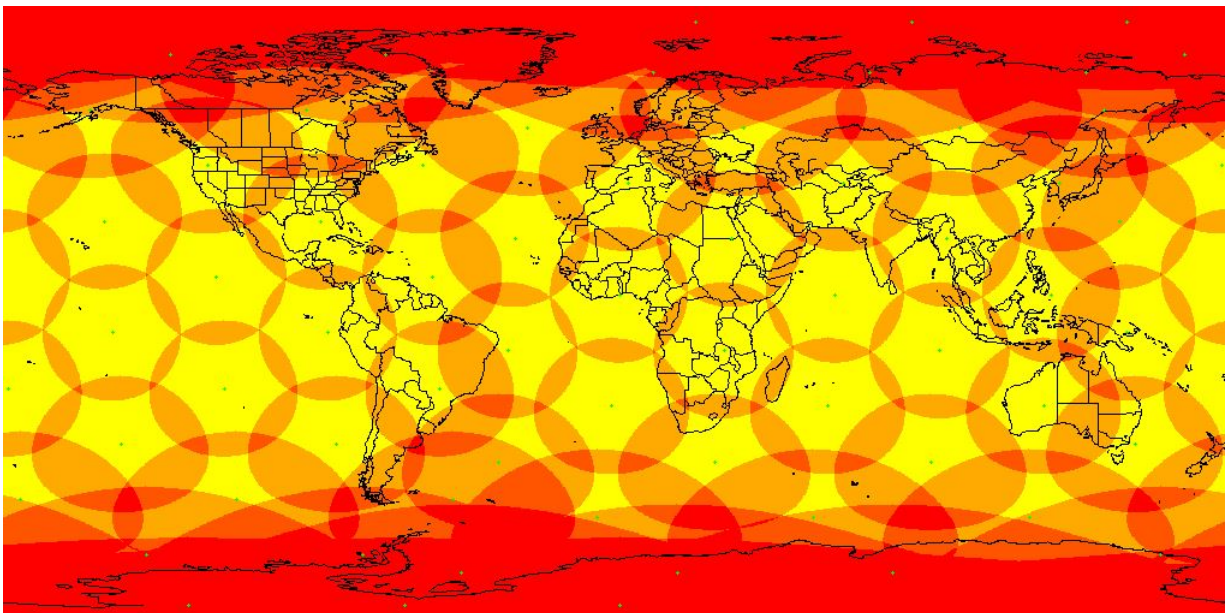


Figure 3.2.1: Distribution of the 66 Iridium constellation satellites.

### 3.2.2 General Configuration

The Polar Orbits configuration consists in the distribution of plains with inclination equal to 90 degrees. Note that the satellites will be travelling parallel to the satellites of the next plain except for the communications between the first and the last plane.

The communications between satellites in antiparallel directions require less space between plains to be fulfilled. In order to solve this inconvenience the separation between the first and the last plain is reduced.

The plains are splitted in the following pattern:

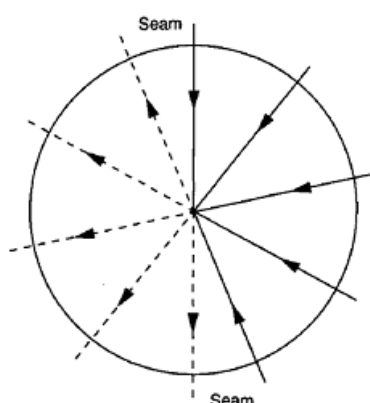


Figure 3.2.2: Distribution of the planes for Polar Orbits design.

#### 3.2.2.1 The Streets of Coverage Method

This Street of Coverage Method is obtained from [?]. As you can see in the figure below, the relations between angles seen from different satellites can be easily computed. The main variables are the following:

Streets of Coverage Method Variables	
N	Number of Satellites
$n_p$	Number of Planes
$N_{pp}$	Number of Satellites per plane
S	Separation between satellites of the same plane
D	General space between planes [°]
$D_0$	Space between antiparallel planes [°]
$\varepsilon$	Elevation angle [°]
$\lambda_{street}$	Street of coverage Width [°]
$\lambda_{max}$	Maximum footprint Radius [°]

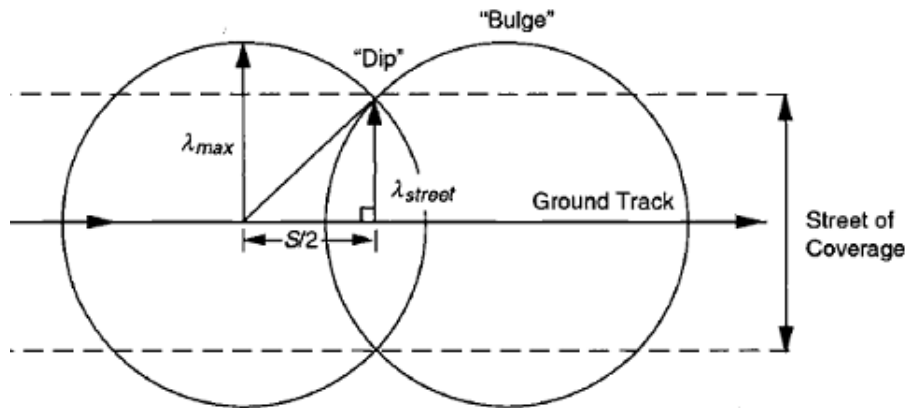


Figure 3.2.3: Single plain street of coverage. The footprints of the satellites superpose leading to a street.

From the figure it can be inferred:

$$S < 2\lambda_{max}$$

$$\cos(\lambda_{street}) = \cos(\lambda_{street}) / \cos(S/2)$$

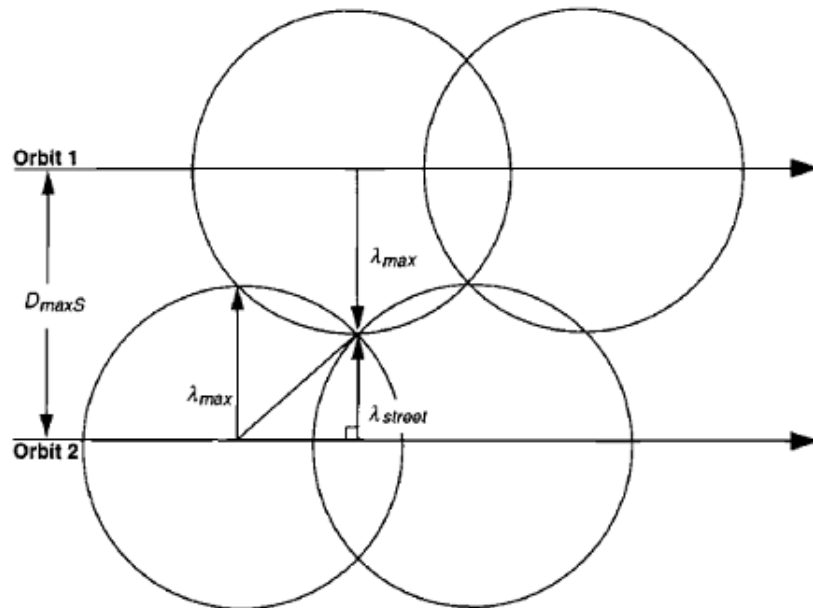


Figure 3.2.4: Two plains streets of coverage. An optimum phasing needs to be obtained.

From this point of view, in general:

$$D = \lambda_{street} + \lambda_{max}$$

n For the antiparallel planes:

$$D_0 = 2\lambda_{street}$$



And the overall relationship between planes sums:

$$180 = (n_p - 1)D + D_0$$

The algorithm for computing the Streets of Coverage Results is defined in the following way:

$$\begin{aligned} \text{Inputs: Height, elevation, inclination...} &\rightarrow \lambda_{max} \rightarrow N_{pp} = \left\lceil \frac{360}{2\lambda_{max}} \right\rceil \rightarrow \\ S = 360/N_{pp} &\rightarrow \lambda_{street} \rightarrow n_p \rightarrow N = N_{pp} * n_p \end{aligned}$$

### 3.2.3 Results of Streets of Coverage

A MATLAB routine has been designed to compute the previously described algorithm. In this conceptual design phase, different heights are computed in order to see the evolution of the number of satellites.

#### General Solution

The program is runned in a broad range of parameters to see the evolution of the number of satellites. As it can be predicted, as the height increases the number of satellites is reduced. The reason is that the footprint of the satellites increases with the height. In addition, as the minimum elevation over the horizon to contact the satellites is reduced, the number of satellites is also reduced for the same reason.

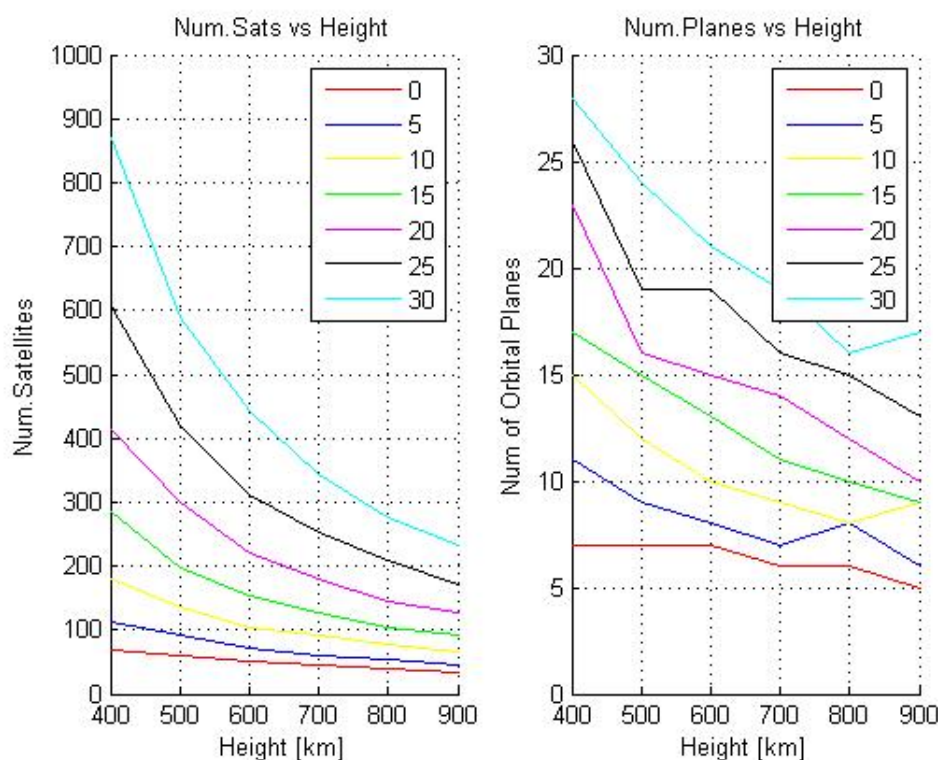


Figure 3.2.5: Variation of number of satellites for different heights and elevation angles

#### Detailed Solution

Given the previously justified assumptions, the same simulation is computed for a more reasonable range of results. In this case, the elevation is set as:

$$\varepsilon = 20^\circ$$

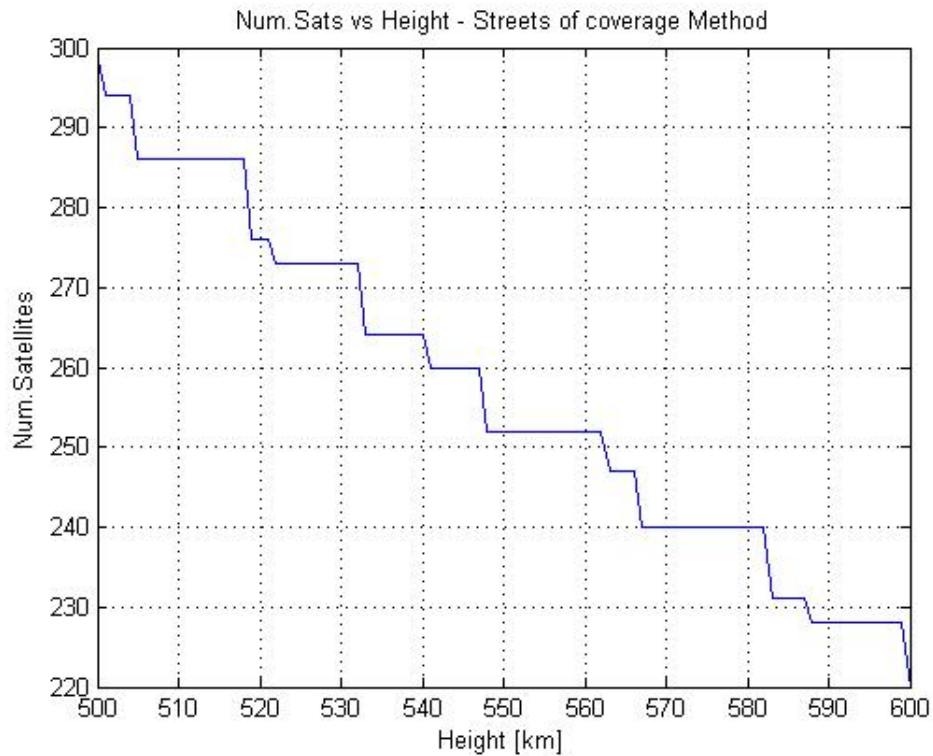


Figure 3.2.6: Variation of number of satellites for different heights between 500 and 600km.

### Conclusion

The computation and the design of this constellation requires small computational and conceptual effort. However, the number of satellites and planes is greater than expected. Even though the technical complexity can be reduced, the availability of small launchers to reach this particularly inclined orbit is also small. In conclusion, more constellation configurations need to be assessed to compare and select the most feasible one.

### 3.3 Walker-Delta Constellation

Walker Delta Pattern constellations are a type of symmetric, inclined constellation made of equal-radius circular orbits, with an equal number of satellites each one. There are several ways to construct a Walker-Delta Constellation:

- Full Walker-Delta Configuration
- Semi Walker-Delta Configuration
- Custom Walker-Delta Configuration

#### 3.3.1 Full Walker-Delta Constellation

##### 3.3.1.1 Characteristics

A typical delta pattern has the following characteristics:

- The constellation contains a total of  $T$  satellites evenly spaced in each of the  $P$  orbital planes. All planes have the same number of satellites, defined as  $S$ , equally distributed. Thus:

$$T = SP \quad (3.3.1)$$

$$\Delta\varphi = \frac{2\pi}{S} \quad (3.3.2)$$

Where  $\Delta\varphi$  is the angle between satellites in the same plane.

- All orbits have equal inclinations  $\delta$  to a reference plane. If this plane is the Equator (it usually is), then the inclination  $\delta$  equals the orbital parameter inclination  $i$  [?].

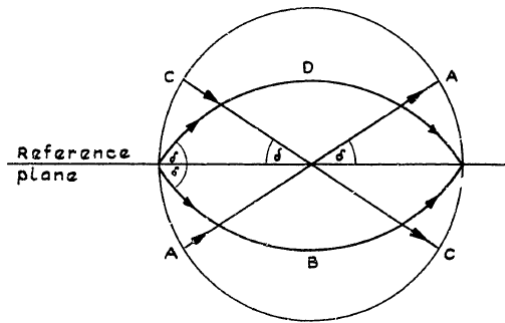


Figure 3.3.1: Definition of the inclination  $\delta$ . Extracted from [?]

- The ascending nodes of the orbits are equally spaced across the full  $2\pi$  ( $360^\circ$  of longitude) at intervals of:

$$\Delta\Omega = \frac{2\pi}{P} \quad (3.3.3)$$

- The position of the satellites in different orbital planes is measured through the factor  $F$ . When a satellite is at its ascending node, a satellite in the most easterly adjacent plane has covered a relative phase difference  $F$ . The real phase difference is defined as:

$$\Delta\Phi = F \frac{2\pi}{P} \quad (3.3.4)$$

In order to have the same phase difference between all orbital planes,  $F$  is defined as an integer, which may have any value from 0 to  $(P-1)$ .

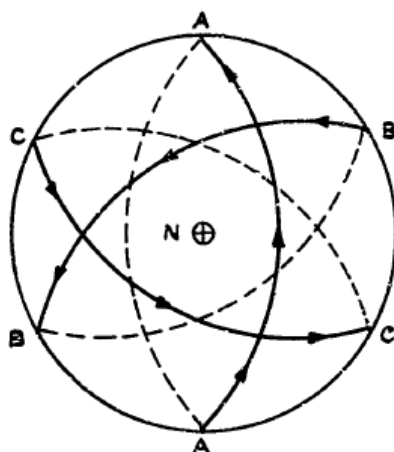


Figure 3.3.2: Delta pattern as seen from the North Pole. Extracted from [?]

With these characteristics, delta constellations are more complex than polar constellations. Because of the inclination of the orbits, the ascending and descending planes and the coverage of the satellites continuously overlap. This characteristic is a constraint on intersatellite networking because the relative velocities between satellites in different orbital planes are larger than in a polar constellation. Consequently, tracking requirements and Doppler shift are increased [?].

### 3.3.1.2 Notation

J.G. Walker developed a notation to define this constellations with only 4 parameters [?]:

$$i : T/P/F$$

Since all satellites are placed at the same altitude, with these notation the shape of the pattern is completely determined. However, to determine all the orbital parameters it is necessary to know the radius of the orbits.

### 3.3.1.3 Coverage

The previous section has shown that in polar orbits the coverage of the constellation could be determined with the streets of coverage method. On the other hand, in delta patterns

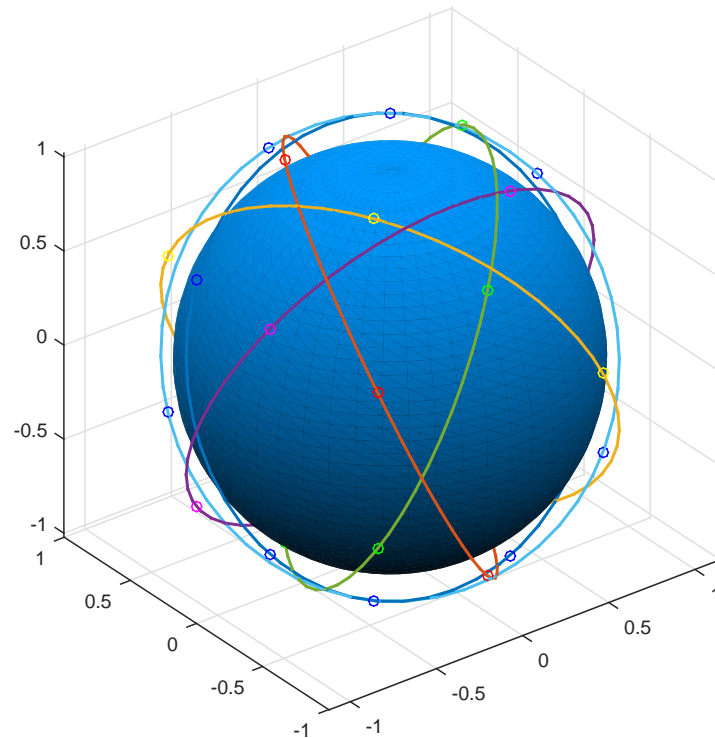


Figure 3.3.3: Delta pattern 65°: 30/6/1

it is necessary to study each configuration to verify its coverage. J.G. Walker determined that delta patterns gave better coverage than polar orbits, but not substantially better in the case of single coverage. This kind of patterns are more useful for double or triple coverage constellations, as it can be seen in Figure 3.3.4. However, his calculations were for a low number of satellites, so it is necessary to compute new results for the number of satellites of the Astrea constellation.

### 3.3.2 Semi Walker Delta Configuration

In order to reduce the necessary costs to design this satellite-based constellation some other configurations will be discussed. The Walker Delta Configuration (WDC) represents the most general constellation for a given inclination different to 90 degrees, i.e. 75 degrees. The WDC is a uniform based 360 degree generated configuration with equidistant orbits, which implies a certain redundant Earth coverage as described in the previous chapter. However, this can and will be solved by generating a 180 degree constellation - Semi Walker Delta Configuration (SWDC) - which will also fulfill global coverage although having some inconvenients.

#### 3.3.2.1 Advantages

- **Distance between planes reduced.** With the SWDC constellation the redundant orbits are directly corrected, thus the distance between planes is reduced to half, as

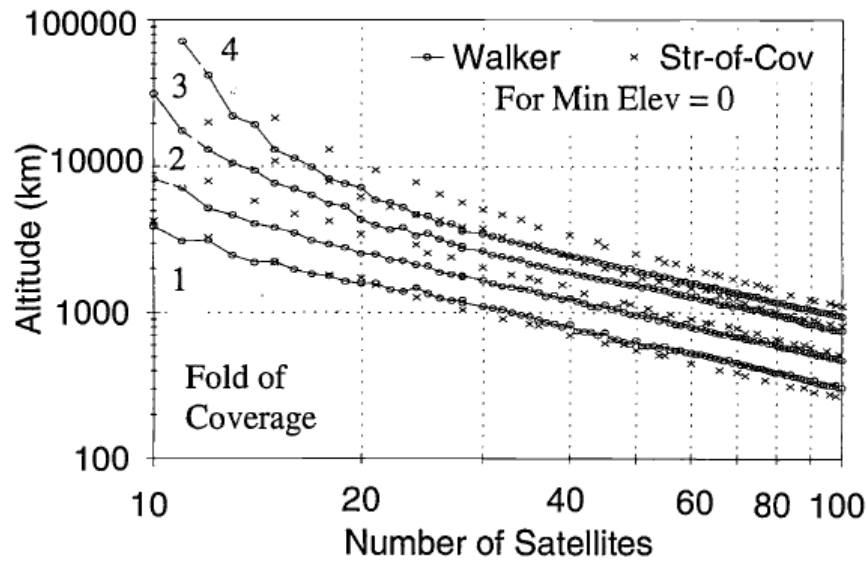


Figure 3.3.4: Minimum altitude for continuous global coverage. Comparison between polar patterns and Walker delta patterns. Extracted from [?]

results from the geometry itself.

- **Less number of planes needed.** This means that in order to approach global coverage fewer planes will be required due to the decrease in distance between planes.
- **Satellites following the same direction - sense** With the SWDC constellation the orbits have no interaction with each other, thus the satellites for each orbit can be set following the same direction. This will significantly improve the communications among satellites from different planes; also, we will be avoiding the Doppler Effect.

### 3.3.2.2 Disadvantages

- **Gap configuration.** With the SWDC constellation the main problem is the gap that results from configuring the constellation at a given inclination and describing equidistant orbits. In order to fulfill global coverage this gap will have to be covered by means of auxiliary orbits.

### 3.3.3 Other Walker Delta Configurations

As we have discussed for the SWDC, the main disadvantage respect to the Walker Delta Configuration is the fact that a gap is obtained, thus a global coverage network cannot be described. In order to cover the entire Earth we have analysed some ways of covering the gap with auxiliary orbits.

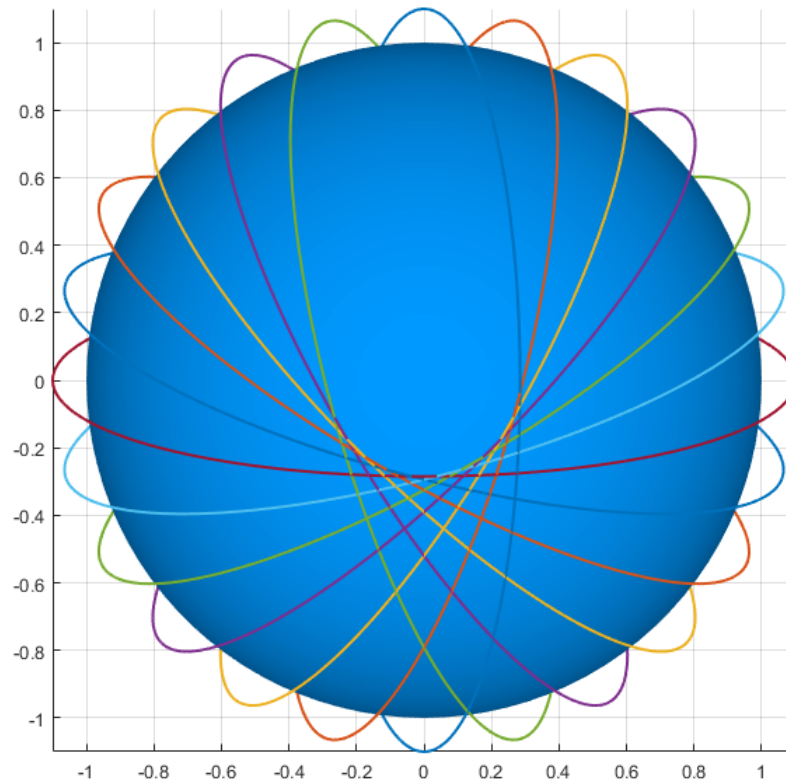


Figure 3.3.5: 12 plane SWDC. Note the gap and the equidistant planes

### 3.3.3.1 SWDC including an additional polar orbit.

This polar orbit would be set directly on top of the gap described by the SWDC. The main issue with polar orbits, as discussed before in this report, is the complex reorientation and decay in inclination that takes place. We must take into account these considerations when covering the entire Earth, especially if we only have one polar orbit in our constellation.

### 3.3.3.2 Mixed Walker Delta.

In order to avoid using polar orbits and their complex reorientations, we can contemplate adding planes to the SWDC. In result, different configurations distributed around the Earth can be described and set in order to fulfill global coverage. As discussed before, the SWDC constellation is generated around 180 degrees whereas the Walker Delta Constellation is a 360 degree generated configuration. This Mixed Walker Delta (MWDC) is the result of adding some planes to the SWDC, thus a constellation can be generated for different degree values, such as 200, 225, 240, etc.



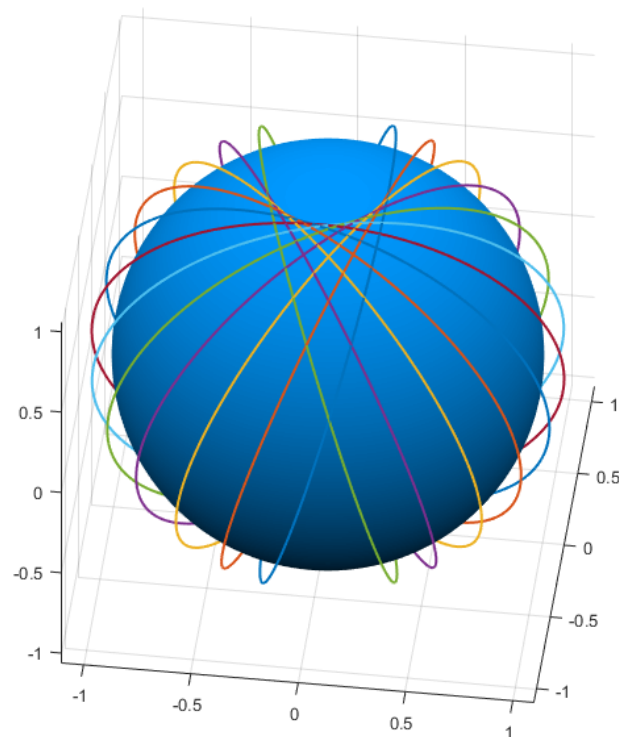


Figure 3.3.6: This geometry distribution induces a large anti-symmetric gap

After different mathematical approaches and optimal solutions, the department of Orbital Design considered that the best option in order to have a global coverage constellation with the least economic and strategic issues - exposed and discussed in previous chapters - would be that of a 210 degree generated MWDC, defined by 8 planes and 21 satellites per plane. This configuration was found optimizing the whole Earth in order to have full coverage without gaps (except for the limitations of this model at high latitudes). An important consideration is that we also analysed other Mixed Walker Delta Configurations for 225 and 240 degrees, but these resulted in a more expensive distribution of satellites.

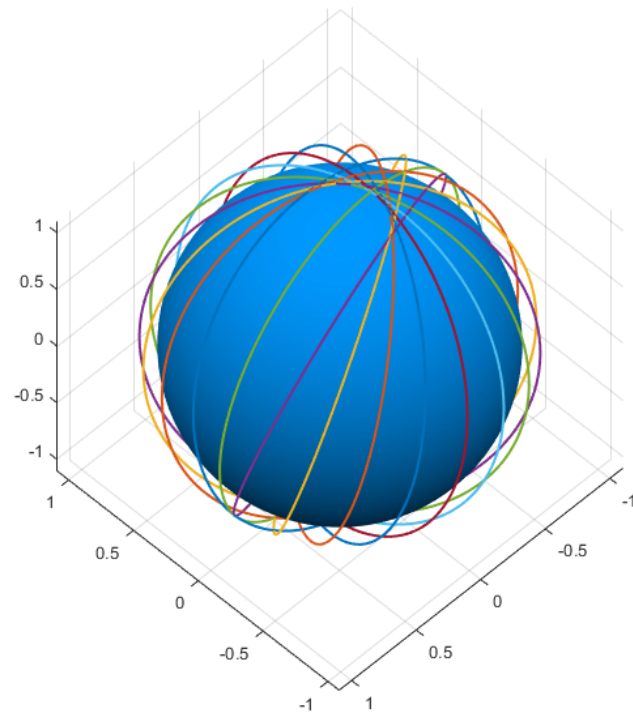


Figure 3.3.7: Added polar orbit to the 11 plane based SWDC

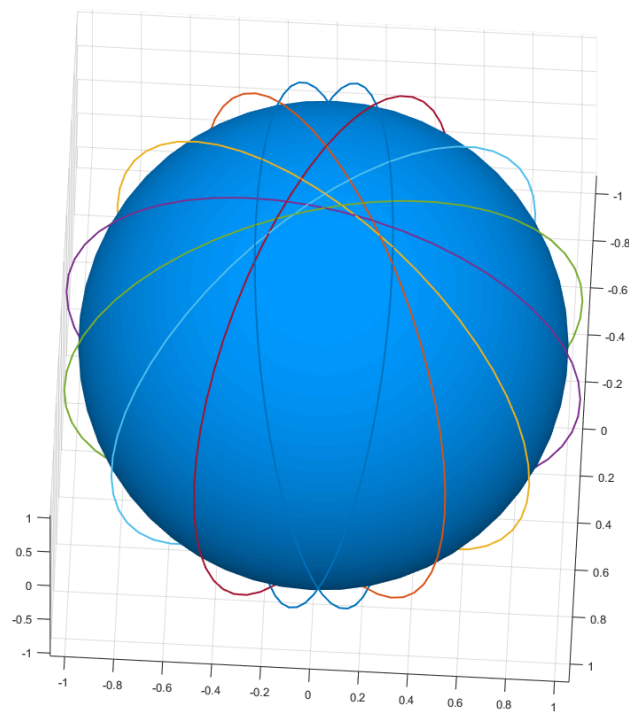


Figure 3.3.8: 8 plane based MWDC generated for 210 degrees

## **3.4 Testing Method**

### **3.4.1 Introduction**

- Primera cosa
- Segunda cosa
- Tercera cosa

### **3.4.2 Second Subsection**

## Chapter 4

# Orbit Perturbations

## 4.1 Sources of Perturbation

### 4.1.1 Introduction to Orbit Perturbations

In this chapter it is seen how the designed orbit configuration varies in time due to external perturbation sources. While some of them can be neglected, there are other of major importance to the future of the constellation. For instance, atmospheric drag determines in plenty of cases the lifetime of the constellation. A first classification of perturbations depending on the time in which their effects are present is the following:

- Secular terms (Sec): They depend on the semimajor axis, the excentricity and the inclination.
- Short Period terms (SP): They depend on the anomalies, this leads to a strong variation in each period.
- Long Period terms (LP): They depend on the argument of the periapsis or the ascendent node.

Even though most of the outter space is vacuum, there ideal models need to consider some factors that escape the typical two body problem. For instance, we can no longer consider Earth as a punctual mass, neither the atmospheric density equal to 0. To enumerate, here is a typical list of the main perturbation sources:

Sources of perturbation:

- Gravity Field of the Central Body
- Atmospheric Drag
- Third Body perturbations
- Solar-Radiation Pressure
- Other Perturbations

All the perturbations can be deeply studied. Consequently, analytical solutions are very hard to find, and even they were found, they do not show clealy a meaning or are not really useful. Instead, there are two mainly used approaches:

- Special Perturbacion: Step-by-step numerical integration of the motion equations with perturbation.
- General Perturbation: Through analytical expansion and integration of the equations of variation of orbit parameters.

**The Approach of the Perturbations Study** For the purposes of these study the different approaches will be assessed. The first analysis will discuss which of the perturbations are the most significant to the study. This analysis will be done considering General Perturbation Techniques. In a deeper second analysis, the two approaches for the perturbations will be assessed and compared considering only the most significant perturbation sources.

#### 4.1.2 Gravity Potential of Earth

Earth's aspherical shape can be modelled as a sum of terms corresponding to the Legendre polynomials. These polynomials can be empirically measured and consider radial symmetry. If one would like to compute also variations in longitude, then should use associated Legendre polynomials.

$$V(r, \delta, \lambda) = -\frac{\mu}{r} \left[ \sum_{n=1}^{\infty} \left( \frac{R_e}{r} \right)^n \sum_{m=0}^n P_{nm} \cos(\delta) (C_{nm} \cos m\lambda + S_{nm} \sin m\lambda) \right] \quad (4.1.1)$$

General Legendre associated polynomials developed Gravitational Potential

$$V(r, \delta) = -\frac{\mu}{r} \left[ 1 - \sum_{n=2}^{\infty} J_n \left( \frac{R_e}{r} \right)^n P_n(\sin \delta) \right] \quad (4.1.2)$$

General Legendre polynomials developed Gravitational Potential

For Earth, the  $J_n$  coefficients are the following:

$$J_2 = 0.00108263 \quad J_3 = -0.00000254 \quad J_4 = -0.00000161$$

Given this distribution, the only significant term  $J_2$ .

$$V(r, \delta) = -\frac{\mu}{r} \left[ 1 - \frac{1}{2} J_2 \left( \frac{R_e}{r} \right)^2 (1 - 3 \sin^2 \delta) \right] \quad (4.1.3)$$

Aproximated Gravitational Potential

If we integrate the force that derives from this potential we can afterwards compute the effect of  $J_2$  On the different orbital elements:

- $\Delta a = 0$

- $\Delta e = 0$

- $\Delta i = 0$

-

$$\Delta \Omega = -3\pi \frac{J_2 R_e^2}{p^2} \cos i \text{ [rad/orbit]} \quad (4.1.4)$$

-

$$\Delta \omega = \frac{3}{2} \pi \frac{J_2 R_e^2}{p^2} (4 - 5 \sin^2 i) \text{ [rad/orbit]} \quad (4.1.5)$$

### 4.1.3 Atmospheric Drag

In order to compute the effect of the remaining atmosphere we use the typical definition of atmospheric drag knowing a drag coefficient:

$$\vec{a}_{drag} = \frac{1}{2} \frac{C_d A}{m} \rho v_{rel}^2 \frac{\vec{v}_{rel}}{|\vec{v}_{rel}|} \quad (4.1.6)$$

The **ballistic coefficient**  $B_c$  is defined as  $\frac{m}{C_d A}$ , characterizing the behaviour of the satellite against atmospheric drag.

### Modelling the Atmosphere

There are several models for the atmosphere. For instance, the most commonly used, the exponential model:

$$\rho = \rho_0 e^{-\frac{h-h_0}{H}} \quad (4.1.7)$$

$$H = \frac{kT}{Mg} \quad (4.1.8)$$

Where:

Exponential Atmosphere Variables	
$\rho$	Density at given height
$\rho_0$	Density at a reference height
$h$	Height over the ellipsoid
$h_0$	Reference height
$H$	Scale Height
$k$	Boltzmann Constant
$T$	Temperature
$M$	Molecular Weight
$g$	Gravity

In addition, other models for the exospheric temperature and the molecular weight need to be used. For this study the ones proposed by The Australian Weather Space Agency are used.

In addition, it is important to note that the following phenomena interfere with the previsions:

- Diurnal Variations
- 27-day solar-rotation cycle
- 11-year cycle of Sun spots
- Semi-annual/Seasonal variations

- Rotating atmosphere
- Winds
- Magnetic Storm Variations
- Others: Tides, Winds,...

Again, if we integrate this force in a period of time, considering the orbit nearly circular, we obtain:

$$\Delta r = -2\pi\rho r^2/B \text{ [/orbit]} \quad (4.1.9)$$

#### 4.1.4 3rd Body Perturbations

The effects of this extra bodies in the system can be computed considering the motion equations. However, some aproximations can be found in the reference as:

$$\dot{\Omega} = \frac{A_m + A_s}{n} \cos i \text{ [}^\circ/\text{day]} \quad (4.1.10)$$

$$\dot{\omega} = \frac{B_m + B_s}{n} (4 - 5 \sin^2 i) \text{ [}^\circ/\text{day]} \quad (4.1.11)$$

Where  $n$  stands for the rate of rotation in orbits/day. In that case, the  $A_m, A_s, B_m$  and  $B_s$  coefficients take as values:

	$A_m + A_s$	$B_m + B_s$
Moon	-0.00338	0.00169
Sun	-0.00154	0.00077

#### 4.1.5 Other Perturbations

In this bag the following low-intensity can be classified:

- Solar Radiation Pressure
- Solid-Earth and Ocean Tides
- Magnetic Field
- South Atlantic Anomaly



## 4.2 Significant Perturbations

### Propagation Algorithm

Given the definitions and approximations to compute perturbations described in the previous section, a propagation in time for the change in orbital parameters is solved. The results are plotted in the graph below:

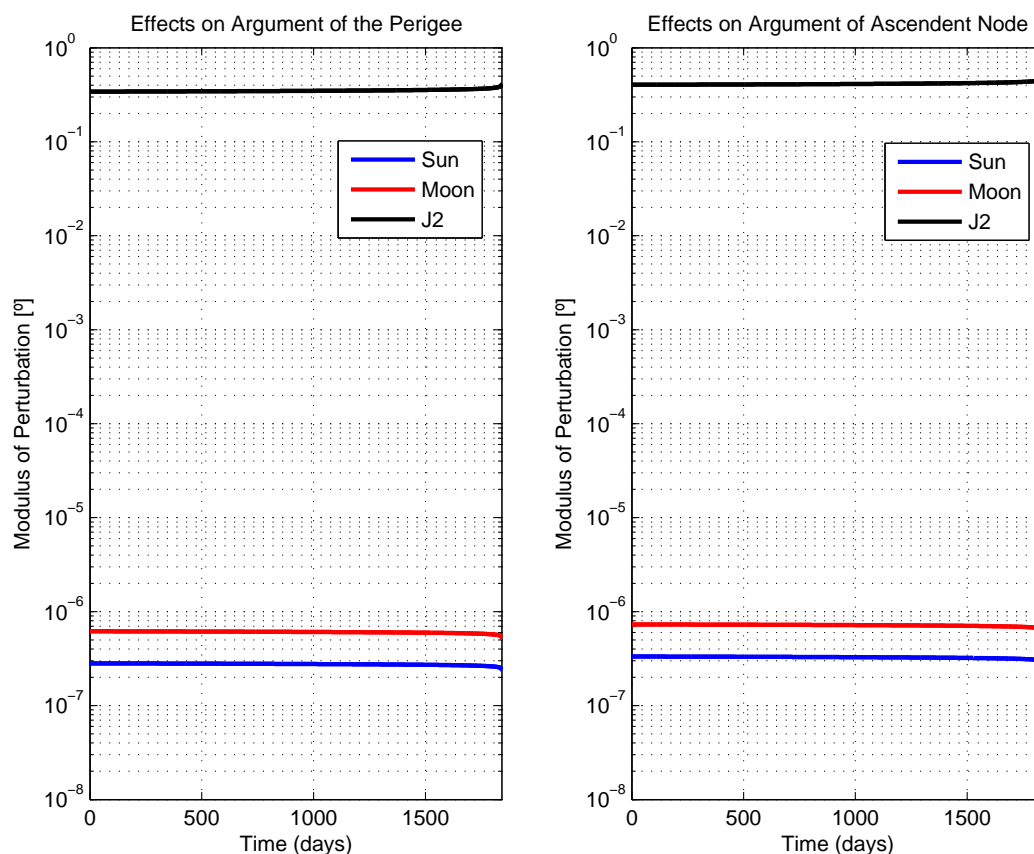


Figure 4.2.1: Logarithmic plot of the modulus of the increases in Angular Arguments of the orbit

As it can be seen, the perturbations caused by 3rd bodies are several orders of magnitude below the order of magnitude of the variation caused by Earth's oblateness. It is also remarkable that the moon has a higher effect than the sun given the relative distance to Earth, even if the sun is way more massive.

Another important observation is that given the very low eccentricity we are considering, the deviation of the argument of the perigee does not affect the performance of the constellation. In other words, since the orbits are considered almost circular there is not a defined Perigee for the orbit.

### In conclusion

The effects of the Moon and the Sun are neglected in comparison with the effects of J2 for the Argument of the ascendent node as well as for the argument of the Perigee.

## 4.3 Orbit Decay

In this chapter the effects of the main perturbations are deeply studied. Firstly, an introduction on the effects of Earth's oblateness on the orbital parameters. Secondly and in more detail, the effects of Atmospheric drag. This is significant because it deviates the power and mass budget to engines and propellant.

### 4.3.1 Effects on the Ascention Node

#### 4.3.1.1 Introduction

- Primera cosa
- Segunda cosa
- Tercera cosa

#### 4.3.1.2 Second Subsection

### 4.3.2 Effects of the Solar Cicle

It is important to consider many parameters when calculating the orbital decay of a satellite. The most important of these parameters for LEO based constellations is drag. As discussed in other chapters, the drag of a satellite depends on the coefficient of drag, its surface, the density of the air and the velocity at which operates. Solar cycles will directly affect the density of the upper atmosphere. This phenomena is relevant when calculating the drag of the satellite and therefore is essential to compute the orbital decay.

Solar cycles are periodic changes in the Sun's activity of approximately 11 years. In each period a solar maximum and minimum can be determined, referring to the amount of periods of sunspot counts. The intensities for these periods vary from cycle to cycle.

Different studies have been made throughout the 20th century cycles. In order to understand the change density of the air changes as consequence of these solar cycles we considered the result data of an old study regarding the 19th solar cycle, which had a duration of 10.5 years between 1958 and 1968. This solar cycle had the highest maximum smoothed sunspot number ever recorded (since 1755), which was of 201.3. This maximum value was recorded in March 1958. This value is high in comparison to other cycles, especially when comparing it to the current 24th solar cycle. In this chapter an analysis will be developed in order to study the influence of the solar cycles on the drag of our satellites.

At 550 km:

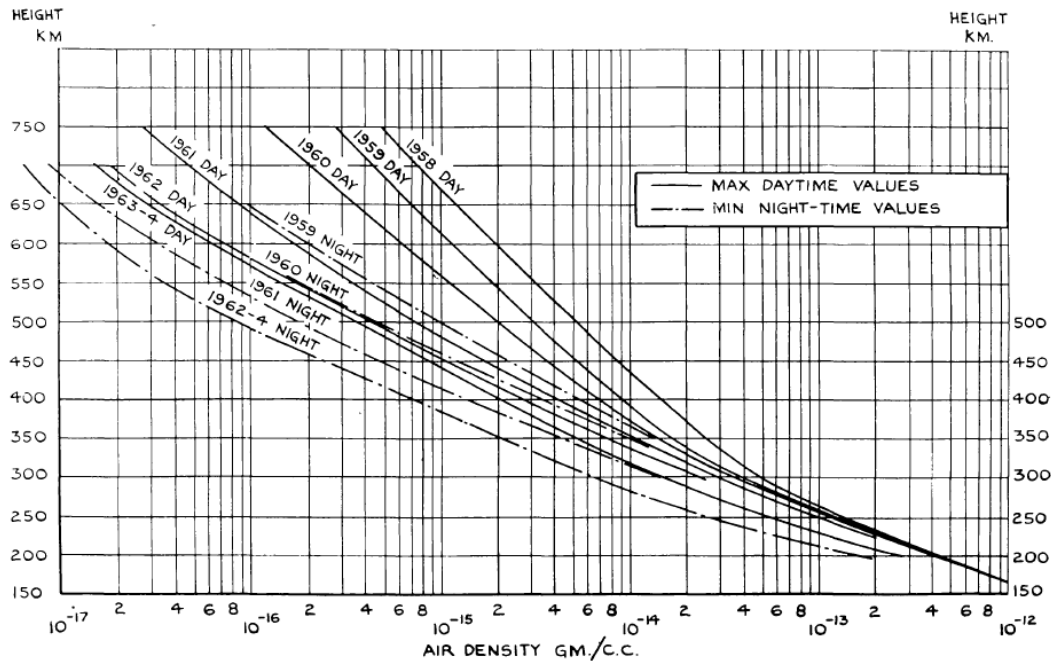


Figure 4.3.1: Deviation of densities in the upper atmosphere due to the 19th solar cycle

Year	D/N	Density at 550km [g/cc]
1958	Day	3.2E-14
1958	Night	5.0E-15
1964	Day	1.35E-15
1964	Night	3.35E-16

These values referring to day and night are the densities of the upper atmosphere at 550 km of altitude respect to the surface of the Earth. The upper atmosphere densities rise during the day following the increase of temperature caused by the radiation of the Sun whereas these values are reduced at night. The orbital decay is on the order of several years whereas these deviations appear every few hours. Thus, in order to compute the orbital decay we will not be taking into account these daily deviations but rather a main value. Therefore the mean density for 1958 will be of  $1.85\text{E-}14$  g/cc and the solar minimum's density of 1964 will be of  $8.4\text{E-}16$  g/cc.

In order to analyse how these values may apply to our constellation we first must adjust these - which belong to the 19th solar cycle - to those of the current 24th cycle, which is noticeable less intense. A way of operating this adjustment is comparing the mean solar maximum achieved by each cycle. The maximum monthly smoothed sunspot number of the 19th cycle had a value of 201.3 and a minimum of 9.6 whereas the current 24th ranges between 11.7 and 81.9 approximately. This means that for the 19th cycle a total

deviation of 191.7 was measured whilst for the 24th cycle this deviation was only of 70.2. This is crucial if we want to analyse the solar maximum densities.

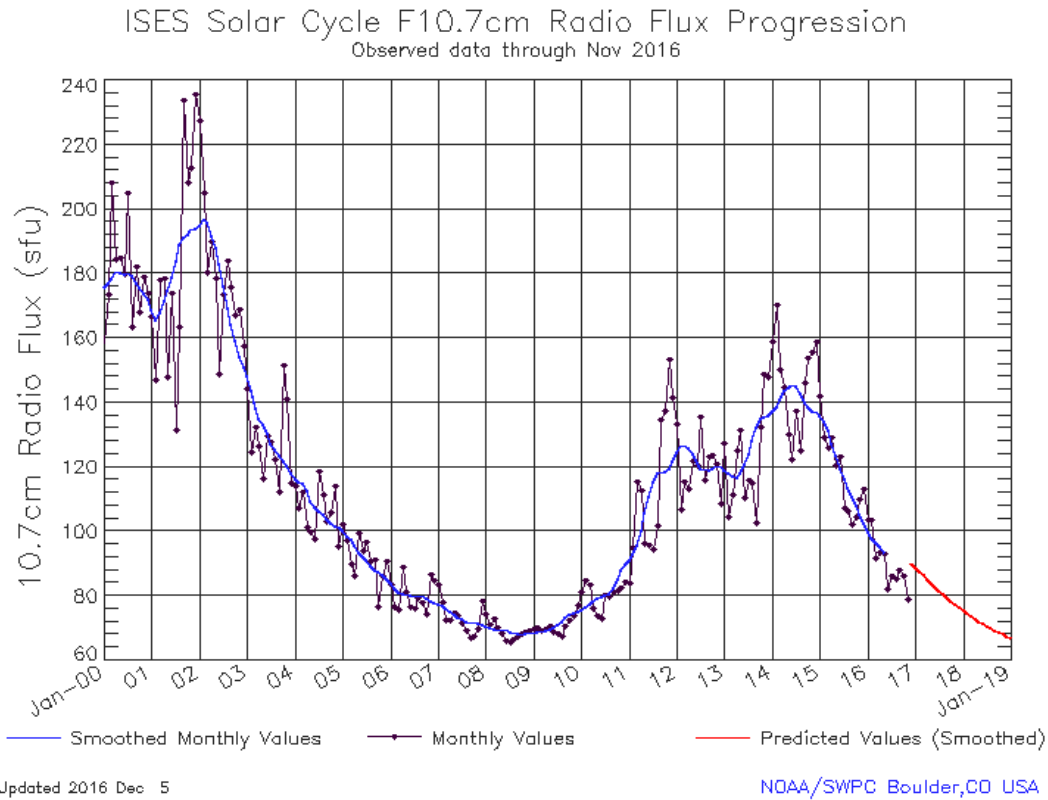


Figure 4.3.2: Measured intensities of the 23rd and 24th solar cycles. Source: NOAA

We must now adjust the mean constant density defined initially to the conditions that this 24th cycle imposes. It is important to note that our satellites will be launched in 2017, and that the 24th cycle is currently decreasing its intensity. Thus, our calculations will be near the conditions of solar minimum, meaning that the drag of our satellite will be smaller than first considered.

Our new approach to the density of the atmosphere at 550 km is near the first approximation, but will consider that we are now entering the solar minimum which will remain more or less constant until 2022. As discussed before, the solar minimum represents a singularity with a minimum density of  $8.4\text{E-}16$  g/cc. The approximation taken will be the resulting constant value which represents the mean smoothed densities between 2017 and 2022.

The final density at 550 km considering the solar minimum during 2017 to 2022 will be of  $2.0\text{E-}15$  g/cc.

### **4.3.3 Orbital Decay Propagation Results**

#### **4.3.3.1 Introduction**

- Primera cosa
- Segunda cosa
- Tercera cosa

#### **4.3.3.2 Second Subsection**

### **4.3.4 Dynamic Orbit Decay Computation**

#### **4.3.4.1 Introduction**

- Primera cosa
- Segunda cosa
- Tercera cosa

#### **4.3.4.2 Second Subsection**

## 4.4 Orbital Station-Keeping

We will study:

- Increased height
- Thrusters

### 4.4.1 Raising the orbit height to increase Lifetime

The key to understand this solution is to see from another point of view the atmospheric drag phenomena. Once we have designed the constellation to provide certain coverage to specific points of the globe, the action of increasing the height of the orbit has the effect of increasing the footprint area on the surface of the earth. As the constellation is set, the time that take the satellites to reach the design height is extra lifetime.

From this point of view, the atmospheric drag phenomena can be recomputed and plotted it in this new way:

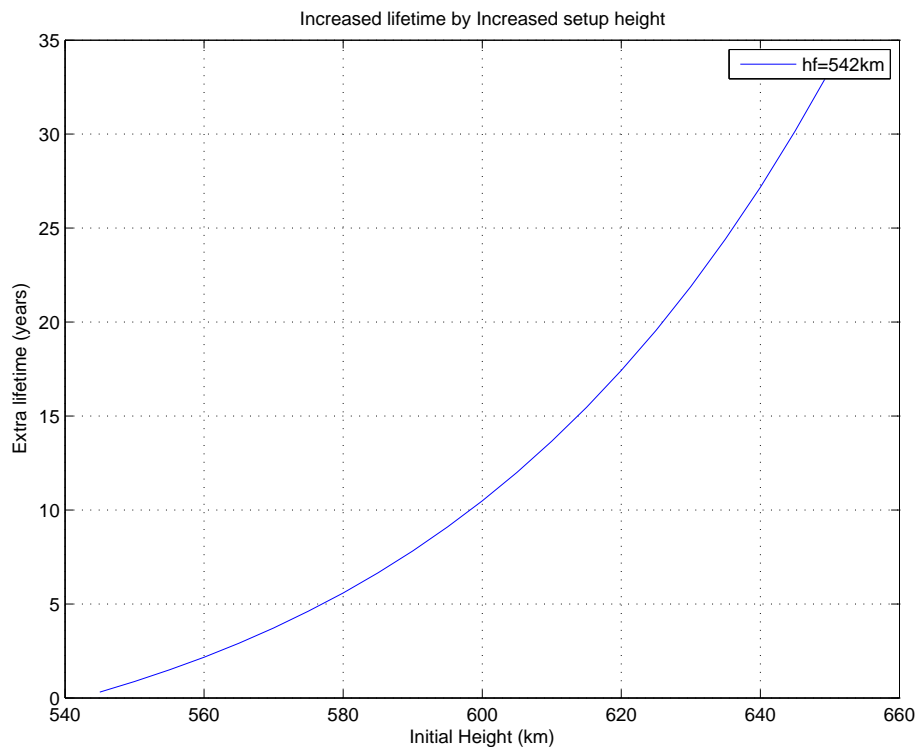


Figure 4.4.1: Increase in the Lifetime obtained by setting the constellation in a higher orbit

As it can be seen, the lifetime increases radically with time. However, this is a dangerous solution, since the coupling with another design parameters is compromised. To list the complications that can lead to:

- **Clients:** With the current technology, the satellites currently in orbit are set to point towards Earth. This means, if the constellation's satellites are at a higher

orbit, the contact is impossible. As the market study reveals, it is important to place the satellites as low as possible.

- **Spacecraft Subsystems:** A higher orbit means a higher gain for the antennas and therefore an increase in the required power.
- **Constellation Reconfiguration:** The overall time to reconfigure the constellation increases with height, since the period of the transition orbits is higher.

### **In conclusion**

This tool is a very powerful option to deal with the orbit decay, even though it is not exactly an operation of Station Keeping itself. Given the high correlation it shows with another subsystems, the possibility of using it needs to be considered while the other design decisions are taken.

## **4.4.2 Using Thrusters to increase Lifetime**

In order to maintain the configuration of the constellation for a longer time, a thruster is installed in each satellite to correct the decrease in altitude due to the orbit decay. The most optimal way to maintain the altitude is through a low-thrust maneuver. However, since this is a preliminar study, the calculations will be computed for a Hohmann transfer maneuver, which is simpler and more effective, but requires more propellant and greater increases of velocity. That is, by computing the velocity and propellant needed for a Hohmann maneuver, the results will be safe for a low-thrust maneuver, because the late one requires less energy.

### **4.4.2.1 Energy equation**

The deduction of the equations needed to solve the Hohmann maneuver begins with the energy equation:

$$\frac{V^2}{2} - \frac{\mu}{r} = -\frac{\mu}{2a} \quad (4.4.1)$$

where  $V$  is the orbital velocity of the satellite,  $r$  is the distance from the focus,  $a$  the semimajor axis of the orbit and  $\mu$  the gravitational constant of the attracting body, in this case, the Earth. This expression shows that the total energy of the satellite equals the sum of its kinetic and potential energy (per mass unit).

This equation can be arranged to obtain the velocity of the satellite. In the case of a circular orbit, the radius is constant, and equal to the semimajor axis. Replacing  $a = r$  in the energy equation and after some operations, the expression of the velocity of a circular orbit is obtained:

$$V_c = \sqrt{\frac{\mu}{r}} \quad (4.4.2)$$

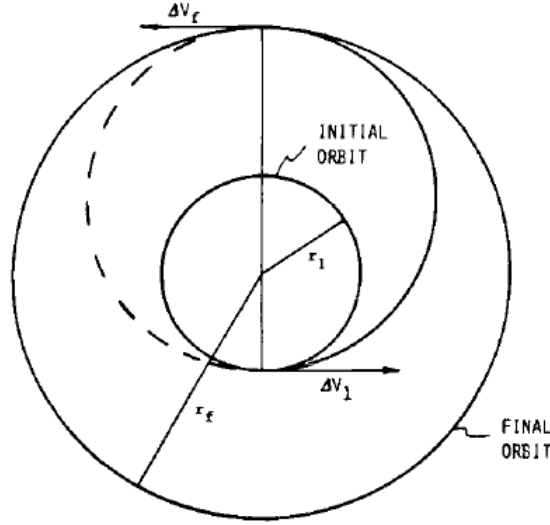


Figure 4.4.2: Hohmann transfer. Extracted from [?]

As it can be deduced from the energy equation, a change in orbital velocity leads to a change in the value of the semimajor axis. This property is used in satellites to change their orbit through a velocity increment  $\Delta V$ . This process is called an orbital maneuver.

#### 4.4.2.2 Delta-V

If the velocity increment  $\Delta V$  is done instantaneously, the maneuver is called an impulsive maneuver. The Hohmann transfer is a two-impulse transfer between coplanar circular orbits. From an initial circular orbit, a tangential velocity increment  $\Delta V_1$  is applied to change the orbit to an ellipse. This ellipse is the transfer orbit, in which the perigee radius is the radius of the initial circular orbit and the apogee radius equals the radius of the final circular orbit. When the satellite reaches the apogee, a second velocity increment  $\Delta V_2$  is applied, so that the satellite reaches the final circular orbit with the apogee radius. If this second velocity is not applied, the satellite will remain in the elliptic orbit.

With the energy equation defined above, it is easy to determine the velocity of the satellite in each orbit. The first orbit and the final ones are circular:

$$V_1 = \sqrt{\frac{\mu}{r_1}} \quad (4.4.3)$$

$$V_f = \sqrt{\frac{\mu}{r_f}} \quad (4.4.4)$$

The velocity in the transfer orbit can be easily calculated with the energy equation applying the definition of the semimajor axis of an ellipse:

$$a = \frac{r_1 + r_f}{2} \quad (4.4.5)$$



The velocities in the perigee and apogee are:

$$V_p = \sqrt{\frac{2\mu r_f}{r_1(r_1 + r_f)}} \quad (4.4.6)$$

$$V_a = \sqrt{\frac{2\mu r_1}{r_f(r_1 + r_f)}} \quad (4.4.7)$$

Therefore the velocity increments are:

$$\Delta V_1 = V_p - V_1 = \sqrt{\frac{2\mu r_f}{r_1(r_1 + r_f)}} - \sqrt{\frac{\mu}{r_1}} \quad (4.4.8)$$

$$\Delta V_2 = V_f - V_a = \sqrt{\frac{\mu}{r_f}} - \sqrt{\frac{2\mu r_1}{r_f(r_1 + r_f)}} \quad (4.4.9)$$

#### 4.4.2.3 Time

It is also necessary to know the time needed to do the maneuver. This time is equal to half of the period of the transfer ellipse:

$$t = \frac{T}{2} = \frac{1}{2} \sqrt{\frac{4\pi^2 a^3}{\mu}} \quad (4.4.10)$$

#### 4.4.2.4 Propellant

In order to know the mass of propellant needed in the maneuver, the Tsiolkovsky rocket equation is applied:

$$\Delta V = g_0 I_{sp} \ln \frac{m_1}{m_f} = g_0 I_{sp} \ln \frac{m_1}{m_1 - m_{prop}} \quad (4.4.11)$$

where  $\Delta V = \Delta V_1 + \Delta V_2$  is the total velocity increment of the maneuver,  $g_0$  is the Earth's gravity,  $I_{sp}$  is the specific impulse of the thruster used,  $m_1$  is the initial mass of the satellite,  $m_f$  is its final mass and  $m_{prop}$  is the mass of propellant used in the maneuver.

$$m_{prop} = m_1 \left( 1 - \exp \left( - \frac{\Delta V}{g_0 I_{sp}} \right) \right) \quad (4.4.12)$$

#### 4.4.2.5 Orbit maintenance

As explained at the beginning of the section, the orbital maneuvers exposed are intended to maintain the altitude of the satellite for a longer time and, consequently, lengthen its life. The method proposed begins when the satellite is deployed at a given height. This height will decrease due to the orbit decay, reaching a critical value, the limit altitude in which the constellation provides global coverage or another given height. Once this critical altitude is achieved, the satellite is put once again at its initial height through a Hohmann maneuver. The process is repeated several times until the satellite runs out of propellant

or until it reaches its desired lifetime.

In reality the satellite will perform a low-thrust maneuver, which is more practical for an electric thruster. In this non-impulsive maneuvers, the thruster is constantly providing a velocity increment to the satellite, but it is so small that the whole transfer maneuver requires a lot of time. This means that it is not necessary to wait until the satellite reaches the critical altitude. The maneuver will start when the satellite is deployed or when it reaches a given altitude (higher than the critical altitude) so that it counteracts the effect of the orbital decay.

#### 4.4.2.6 Results

The results are computed for a 3U CubeSat with an ion thruster. The characteristics of the thruster are the following ones (for more characteristics of the thruster refer to the section ??):

Thrust	100 $\mu\text{N}$
Specific Impulse	2150 s

The first parameters to be defined are the maximum and minimum height of the orbit, mesured from the surface of the Earth. The maximum height is the altitude at which the satellite is deployed, and minimum height is the altitude at which the Hohmann transfer maneuver is applied. The satellite has to be above the minimum height to be functional. Figure 4.4.3 is an example of the height variation of the satellite using the Hohmann maneuver to reach the maximum height once the satellite is in the minimum height. The results of this maneuver are:

Maximum height	545 km
Minimum height	542 km
Number of Hohmann Maneuvers	32
Maximum $\Delta V_1$	0.8237 m/s
Maximum $\Delta V_2$	0.8236 m/s
Total $\Delta V$ Budget	52.7116 m/s
Propellant mass	10 g
Lifetime of the satellite	33.3288 years

Since the thruster used is an ion thruster, the specific impulse is big, and the mass propellant is very low. In this case, the variation of height due to the orbit decay is approximately 3 km per year, so the thruster needs to do a Hohmann maneuver per year. With only 10 g of propellant, the lifetime of the satellite is over 30 years.

Figure 4.4.4 is another example of the Hohmann maneuver with the same amount of propellant but with a more restrictive range of operational heights, only 80 m. It should have the same shape as Figure 4.4.3, but since a lot of maneuvers are applied, the lines have overlapped. The characteristics of this maneuver are:

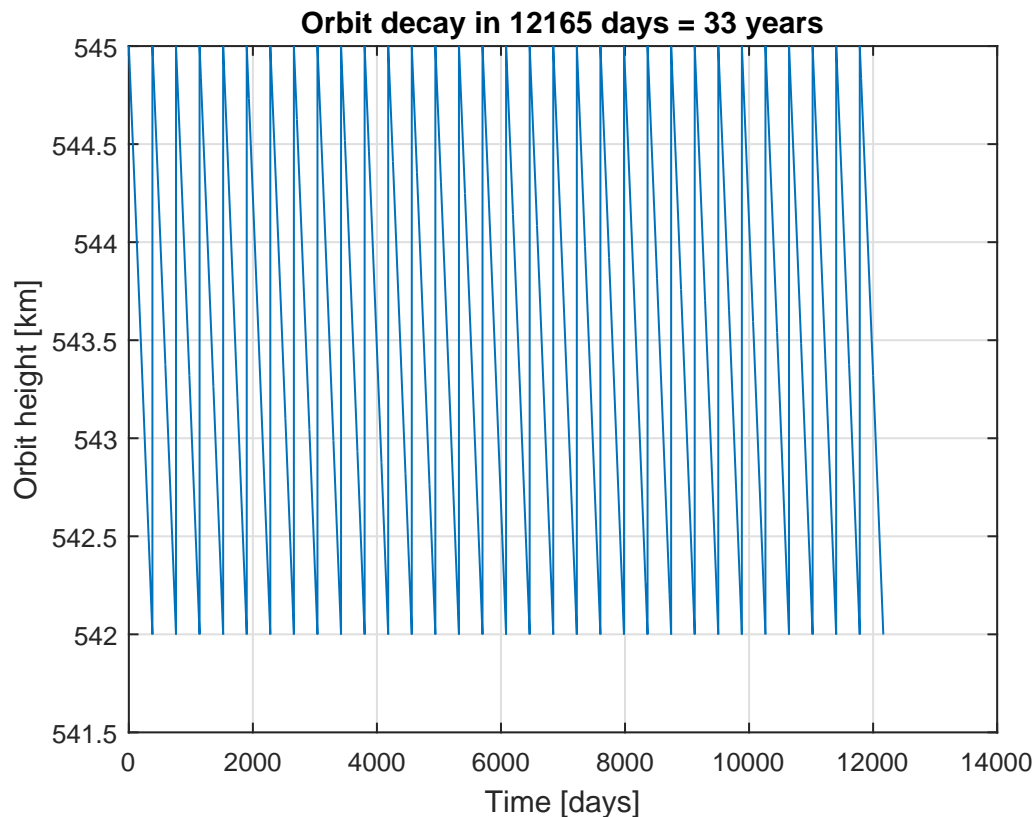


Figure 4.4.3: Height variation of the satellite

Maximum height	545 km
Minimum height	544.92 km
Number of Hohmann Maneuvers	1200
Maximum $\Delta V_1$	0.0221 m/s
Maximum $\Delta V_2$	0.0221 m/s
Total $\Delta V$ Budget	52.7570 m/s
Propellant mass	10 g
Lifetime of the satellite	34.5726 years

Comparing these results with the previous ones, it can be seen that with a more restrictive range of heights, the lifetime of the satellite is practically the same. The velocity increments are lower because the difference in the heights is extremely low, but at the same time, the satellite reaches before the minimum height and the maneuvers needed to maintain the satellite in this range are many more than on the other case. Since the  $\Delta V$  budget is practically the same in both cases, it can be assured that the only difference between them is the number of maneuvers computed.

As mentioned earlier, the results obtained are for a Hohmann maneuver when in reality the satellite will compute a low-thrust maneuver, that requires less velocity increments

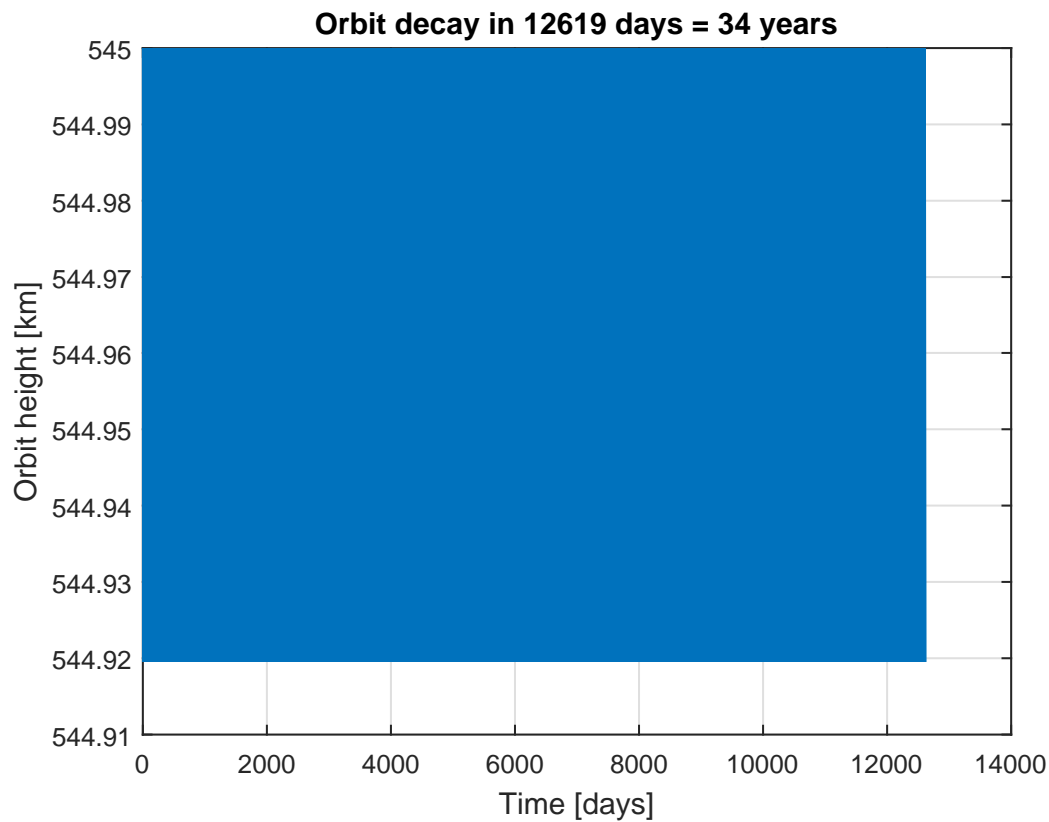


Figure 4.4.4: Height variation of the satellite with a more restrictive minimum height

and less propellant. In conclusion, taking into account these results, it can be stated that the lifetime of the satellite will not be determined by its orbit decay but for the failure of its systems or other external causes. It can also be assured that the satellite is capable of carrying enough propellant to maintain its altitude and to compute other maneuvers if necessary.