

Chapter 1

Introduction, Numbers, and Units

In physics we study the interactions and motion of matter. For our introduction in this class, we will mostly study what I like to call “Physics at the human scale” because it is the physics of daily life that we observe and directly interact with. It describes an apple falling from a tree, a car slamming its brakes, a pendulum swinging, the collision of 2 billiard balls, and more.

1.1 Math and Scientific Notation

This text will assume that you are comfortable with using algebra to solve an equation for a variable and that you have had a class on trigonometry. We will review those topics as they come up but it may be difficult to follow without the math background.

During this text we will use **scientific notation** frequently to write very large or very small numbers in an easier to read format. With it, we write a number in the format

$$a \cdot 10^b \tag{1.1}$$

where $1 \leq a < 10$ and b is an integer. This saves us from having to write and count a lot of zeros that indicate place. The number is a factor multiplied by a power of 10.

Example 1.1

Write the number 38,500,000 in scientific notation.

To write this in the form $a \cdot 10^b$, first let's find a such that $1 \leq a < 10$. Since our number starts with the digits 385, that means $a = 3.85$. From here we have

$$38,500,000 = 3.85 \cdot 10,000,000 \quad (1.2)$$

Now we need to write 10,000,000 as a power of 10. There are 7 zeros, so it will be 10^7 . This means that our number is

$$3.85 \cdot 10^7 \quad (1.3)$$

Example 1.2

Write the number 0.000007018 in scientific notation.

We can use the same reasoning as the above example to find that $a = 7.018$. The number can be written as

$$0.000007018 = 7.018 \cdot 0.000001 \quad (1.4)$$

Now we just need to find the power of 10. Let's start looking at powers of 10.

$$\begin{aligned} 10^0 &= 1 \\ 10^{-1} &= \frac{1}{10} = 0.1 \\ 10^{-2} &= \frac{1}{10^2} = 0.01 \\ 10^{-3} &= \frac{1}{10^3} = 0.001 \end{aligned} \quad (1.5)$$

A negative power means to take 1 divided by the number to that positive power. Looking at 10, we see that as we decrease the power (larger negative value), we move the decimal point 1 to the left. This will give us

$$0.000001 = \frac{1}{10^6} = 10^{-6} \quad (1.6)$$

which makes our number

$$0.000007018 = 7.018 \cdot 10^{-6} \quad (1.7)$$

1.2 Units

Physicists will try to describe the world around us quantitatively, meaning with numbers. In order to ascribe meaning to those numbers we need to use physical **units**. For example, think about if you were baking a loaf of bread and the recipe called for “4 flour”. Would you know how much flour was required? The recipe should say “4 *cups* of flour” so that you know the amount required for the bread.

In physics, like most sciences, we use **SI units**. In the USA, you likely encounter the Imperial system of units in daily life. Common Imperial units you might have seen include the inch for length and gallon for volume. These units are not the standard for scientific disciplines, and as such we will focus on using SI units. There are 3 base units that we will use frequently:

- **kilogram** - the SI base unit of mass. Mass is a measure of how much matter there is. It is related to but not the same as weight. We will discuss weight further in a later chapter, but for now think of mass as how much “stuff” there is.
- **meter** - the SI base unit of length. If you don’t have a good mental picture of a meter, it is a little longer than 3 feet.
- **second** - the SI base unit of time. We will continue to see minutes and hours used throughout, but for doing quantitative work we will want to use seconds for everything.

1.3 SI Prefixes

With these SI base units we use **prefixes** to scale the units relative to our problem. Table 1.1 contains several of the most common prefixes, with the bolded ones being most important for this class.

One thing you should see is that SI prefixes allow us to scale our units by powers of 10. This is much simpler than the Imperial system. For lengths we have 12 inches to a foot, 3 feet to a yard, and 1760 yards to a mile. Compare

this to just multiplying or dividing by powers of 10 and see how much simpler SI units are to work with!

The main purpose to having the prefixes is it allows us to scale the units to our problem. If we were to talk about the distance between San Diego and Los Angeles, it would make sense to use kilometers because that distance is about 194 kilometers. In meters, it would be $194 \cdot 10^3 = 194,000$ meters. Prefixes allow us to use numbers that are closer to 1, which are easier to read and for our brains to process.

There are other prefixes and units that are used within different branches of physics and other fields of science. We will see some of this in CHAPTER HERE when we discuss the solar system. Again, the goal is always to make the numbers we work with closer to 1.

In addition to the base units, we will also have **derived units**, which are a combination of SI base units. On a car speedometer you may have seen the unit km/hr , which means “kilometers per hour”. This is a derived unit because it combines the units of length and time.

1.4 Converting Units

Now that we have SI prefixes, it is important that we learn how to convert between different units. This process is called **dimensional analysis** because the units are the “dimensions” that we are looking at.

To convert units, we will write unit factors that are fractions equal to 1. For example,

$$\frac{1 \text{ mm}}{10^{-3} \text{ m}} = 1 \quad (1.8)$$

Table 1.1: SI Prefixes

| Prefix | Abbreviation | Numerical value |
|---------------|---------------------------------|-----------------|
| nano- | n | 10^{-9} |
| micro- | μ (Greek letter <i>mu</i>) | 10^{-6} |
| milli- | m | 10^{-3} |
| centi- | c | 10^{-2} |
| kilo- | k | 10^3 |
| mega- | M | 10^6 |
| giga- | G | 10^9 |

is a unit factor that equals 1 because the prefix milli- means 10^{-3} , so 1 millimeter is equal to 10^{-3} meter. This fraction can be multiplied by a quantity to change the unit but not the physical value because the fraction is 1. Now let's use that to convert 3.25 meters to millimeters. We can use the unit fraction to cancel the meters from the original value, and have a unit of millimeters left in the numerator.

$$3.25 \cancel{m} \cdot \frac{1 \text{ mm}}{10^{-3} \cancel{m}} = 3,250 \text{ mm} \quad (1.9)$$

The reason this works is that our unit fraction is equal to 1, so multiplying our original quantity by that unit fraction doesn't change the value! It only changes what unit that value is given in.

You could also think of the unit fraction as 1,000 millimeters in a meter.

$$3.25 \cancel{m} \cdot \frac{1000 \text{ mm}}{1 \cancel{m}} = 3,250 \text{ mm} \quad (1.10)$$

Both $\frac{1000 \text{ mm}}{1 \text{ m}}$ and $\frac{1 \text{ mm}}{10^{-3} \text{ m}}$ are valid unit fractions for converting from meters to millimeters. They are equally correct, so use whichever you prefer!

Example 1.3

Convert 500 nanometers to centimeters.

We can convert 500 nanometers to meters first, then convert from meters to centimeters.

$$500 \cancel{nm} = \frac{10^{-9} \cancel{m}}{1 \cancel{nm}} = 5 \cdot 10^{-7} \cancel{m} \quad (1.11)$$

$$5 \cdot 10^{-7} \cancel{m} = \frac{100 \cancel{cm}}{1 \cancel{m}} = 5 \cdot 10^{-5} \cancel{cm} \quad (1.12)$$

Sometimes it is useful to convert to an intermediary unit to help you get to the requested final unit.

Example 1.4

Convert 1 meter to feet. Use the following information:

- 1 foot = 12 inches

- 1 inch = 2.54 cm

The given unit conversions tell us how to solve this problem. First we convert the meter to centimeters, then use the given conversion from centimeters to inches. Finally, we can use the given conversion from inches to feet.

$$1 \cancel{m} \cdot \frac{100 \cancel{cm}}{1 \cancel{m}} \cdot \frac{1 \cancel{in}}{2.54 \cancel{cm}} \cdot \frac{1 \cancel{foot}}{12 \cancel{in}} = 3.28 \text{ feet} \quad (1.13)$$

Example 1.5

A car is driving at 45 kilometers per hour $\frac{km}{hr}$. Convert this to meters per second $\frac{m}{s}$

First let's write out the unit to help us see how to convert this.

$$45 \frac{km}{hr} \quad (1.14)$$

This means that the car moves 45 kilometers in 1 hour. When doing unit conversions it may be helpful to write out the unit like that

$$45 \frac{km}{hr} = \frac{45 km}{1 hr} \quad (1.15)$$

We have kilometers in the numerator and hours in the denominator, so we have to make sure we write our unit fractions correctly to cancel them. First, let's convert the kilometers to meters.

$$\frac{45 \cancel{km}}{1 hr} \cdot \frac{10^3 m}{1 \cancel{km}} = \frac{45,000 m}{1 hr} \quad (1.16)$$

Now we can convert the hour to seconds. It may be helpful to convert to minutes first, then to seconds.

$$\frac{45,000 m}{1 \cancel{hr}} \cdot \frac{1 \cancel{hr}}{60 \cancel{min}} \cdot \frac{1 \cancel{min}}{60 s} = \frac{12.5 m}{1 s} \quad (1.17)$$

The car is driving at $12.5 \frac{m}{s}$.

Chapter 2

Quantities of motion and kinematics

The first thing we must do to describe the motion of objects is define specific terms that we will use. These terms will have clear, specific definitions that may differ from how they are used in everyday language.

2.1 Motion Definitions

The first quantity to talk about is **position**. This is a measurement of where something is along an axis (or multiple axes). This is a length, so the SI unit for position is the meter. Position will be denoted with the variable x in equations.

When we study motion, we often look at an object moving from one position to the other. The first position is the **initial** position, and the second position is the **final** position. In equations we will use subscript i and f for initial final respectively, so that the initial position is x_i and the final position is x_f . To indicate the change in a quantity we will use Δ , which is the Greek letter Delta. The change in position between the initial state and final state can be written as

$$\Delta x = x_f - x_i \tag{2.1}$$

which is a measure of how the position changed moving from the initial state x_i to the final state x_f .

The second quantity we will talk about is **velocity**. This is a measurement of the rate of change in position of a moving object. Velocity will be denoted with the variable v in equations. Velocity has a direction, so it is a vector

quantity. The direction can be denoted with a positive or negative sign to indicate the direction along a position axis.

Average velocity is the average rate of change in position as an object moves from its initial position x_i to its final position x_f over a time interval Δt . For average velocity we will use a subscript “avg”, so it is written as v_{avg} .

$$v_{avg} = \frac{\Delta x}{\Delta t} \quad (2.2)$$

Average velocity does not give us any information about the moment-to-moment velocity throughout the motion, only the average throughout the time interval.

From the equation above, we see that velocity is a length divided by time. This means our SI unit is $\frac{m}{s}$, which is read as *meters per second*.

The third quantity we will talk about is **acceleration**. This is a measurement of the rate of change of velocity of a moving object. Acceleration will be denoted with the variable a in equations. Acceleration has a direction, so it is a vector quantity.

Average acceleration is the average rate of change in velocity as an object moves from its initial position x_i to its final position x_f over a time interval Δt . Similar to position, the initial velocity is v_i , the final velocity is v_f , and the change in velocity is $\Delta v = v_f - v_i$.

$$a_{avg} = \frac{\Delta v}{\Delta t} \quad (2.3)$$

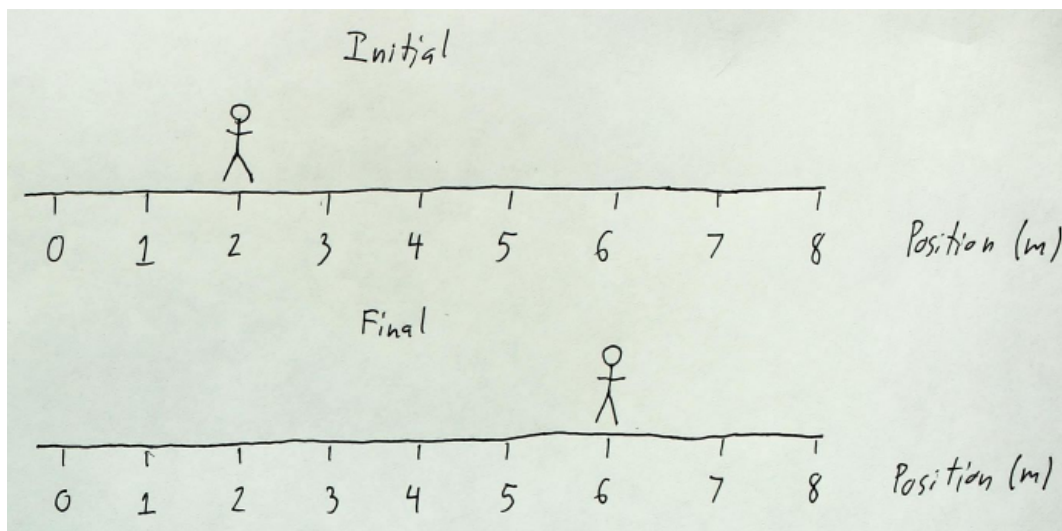
Again, average acceleration gives us the average value over the entire interval of motion but no information about the moment-to-moment value of the acceleration.

From the equation above, we see that acceleration is a velocity divided by time. Our SI unit is $\frac{m/s}{s}$, which is read as *meters per second per second*. The unit can also be written as $\frac{m}{s^2}$, which is read as *meters per second squared*. These two ways of writing the units of acceleration are equivalent and both are correct. For the rest of this text, I will use $\frac{m}{s^2}$.

Example 2.1

Alice starts at a position of 2.0 m and walks to a position of 6.0 m over a duration of 8.0 s . What is her average velocity while she is walking?

We know Alice’s initial position $x_i = 2.0\text{ m}$ and final position $x_f = 6.0\text{ m}$. This means her change in position is



$$\Delta x = x_f - x_i \quad (2.4)$$

$$\Delta x = 6.0 \text{ m} - 2.0 \text{ m} = 4.0 \text{ m} \quad (2.5)$$

In the problem we are told she walks for 8.0 s , so our time duration is $\Delta t = 8.0 \text{ s}$. Now we can use the definition of average velocity

$$v_{avg} = \frac{\Delta x}{\Delta t} \quad (2.6)$$

$$v_{avg} = \frac{4.0 \text{ m}}{8.0 \text{ s}} = 0.50 \frac{\text{m}}{\text{s}} \quad (2.7)$$

Her average velocity is $0.50 \frac{\text{m}}{\text{s}}$. Notice how we get the units of $\frac{\text{m}}{\text{s}}$ by just dividing the units in the equation.

The final quantity we will introduce is **speed**. This is a measurement of how fast an object is moving. You might be familiar with miles per hour (mph) or kilometers per hour ($\frac{\text{km}}{\text{hr}}$) from a car speedometer. The SI unit of speed is $\frac{\text{m}}{\text{s}}$, same as velocity. Speed will be denoted as s in equations.

Average speed can be found by taking the distance traveled d over the time duration Δt .

$$s_{avg} = \frac{d}{\Delta t} \quad (2.8)$$

One important issue we must address is how average speed differs from average velocity. Both are measured in units of $\frac{m}{s}$ and can be used to get information about how fast an object moved.

1. Speed is a scalar quantity, not a vector quantity. Speed does not tell you anything about the direction of motion. There is no such thing as a negative speed because the distance traveled will always be zero or positive.
2. The distance traveled is not always the same as the change in position. Velocity only depends on x_i , x_f , and Δt with no concern for the path taken between x_i and x_f . Speed depends on the distance traveled so it does matter what the path taken is.

The last new definitions we will introduce here are **instantaneous speed** and **instantaneous velocity**. Instantaneous speed is a measure of the speed at an instant in time. Instantaneous velocity is the rate of change in position at an instant in time, which can also be thought of as the instantaneous speed plus the direction of travel.

2.2 Motion Graphs

One helpful tool for studying the motion of objects is to use graphs. This allows us to present data in a visual format rather than text. The x-axis will be the time and the y-axis will be position, velocity, or acceleration.

First, let's start with the following graph of position vs. time (which is also called an x vs t plot).

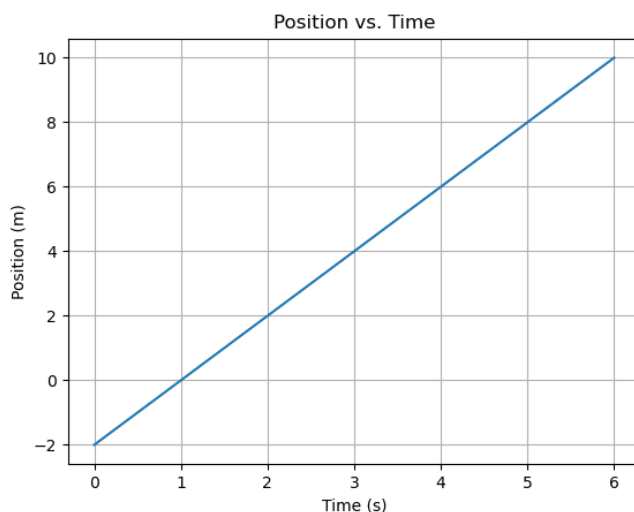


Figure 2.1: Sample position vs. time graph

The graph has a title that tells you what the graph shows and the x- and y-axes are labeled, including units. These features are important to allow somebody reading your graph to easily understand what it is showing.

This graph has time in seconds on the x-axis and position in meters on the y-axis. What this graph is showing is therefore the position of an object as it moves over time. For example, if we wanted to find the position of the object at $t = 4\text{ s}$ we would look at the y-value of the graph at 4 s . The position of the object is therefore $x = 6\text{ m}$.

What if we wanted to find the average velocity of the object throughout this motion? It moves for 6 seconds, so $\Delta t = 6\text{ s}$. The position at $t = 0\text{ s}$ is $x = -2\text{ m}$, so we have in the initial state $x_i = -2\text{ m}$. The position at $t = 6\text{ s}$ is $x = 10\text{ m}$, so we have in the final state $x_f = 10\text{ m}$.

$$v_{avg} = \frac{\Delta x}{\Delta t} = \frac{10\text{ m} - (-2)\text{ m}}{6\text{ s}} = 2 \frac{\text{m}}{\text{s}} \quad (2.9)$$

Since our y-value is position and x-value is time, what we just did above is find the slope of the line! We can extend this in general to recognize that **the slope of a x vs t graph is the velocity**. We can use this to find instantaneous velocities by finding the slope at a point on the graph! Below is a plot of velocity vs. time that describes the same motion.

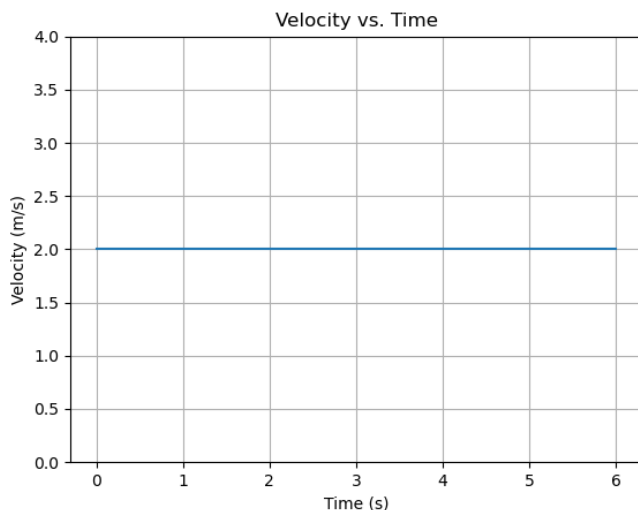


Figure 2.2: Sample velocity vs. time graph for the same motion as 2.1.

The slope of the x vs t plot is constant, so the velocity doesn't change throughout the motion. From 0 s to 6 s we have $v = 2\frac{\text{m}}{\text{s}}$. Next let's look at what we call the “area under the curve”, which is the area of the graph between the x-axis and the graph value. Since the velocity is constant, the area is just a rectangle

$$\text{Area} = 2.0 \frac{\text{m}}{\text{s}} \cdot 6\text{ s} = 12\text{ m} \quad (2.10)$$

Back to the x vs t plot where we had $x_f = 10\text{ m}$ and $x_i = -2\text{ m}$, we have a change in position

$$\Delta x = x_f - x_i = 10\text{ m} - (-2)\text{ m} = 12\text{ m} \quad (2.11)$$

The area under the graph of a v vs t graph is the change in position of the object! Since the y-axis is velocity and the x-axis is time, we are multiplying the velocity by the time to get the change in position. Note that the v vs t graph does not tell us what the value of x_i is, just how much the position changes by.

This is a good time to make a very important point: **motion graphs are not a “picture” of the motion, but instead a visual representation of data.** The x vs t and v vs t look very different from each other despite describing the same motion. That is because they show how the position and velocity change over time, not provide a “picture” of how the object moves. In

almost every case the position, velocity, and acceleration graphs for a moving object will all look very different from each other.

Using the same rational as above, the acceleration is the slope of the v vs t graph and the change in velocity is the area under the a vs t graph. We will try these concepts out in the example below.

Example 2.2

The acceleration of an object over time is measured as

Table 2.1: Acceleration vs Time example

| Time (s) | Acceleration (m/s^2) |
|--------------|--------------------------|
| 0 | 9 |
| 1 | 7 |
| 2 | 5 |
| 3 | 3 |
| 4 | 1 |

1. Plot the acceleration vs. time.
2. Use the acceleration vs. time graph to create a velocity vs. time graph if the velocity at $t = 0$ s is $v = -4$ m/s.

First we can create our a vs. t plot with time on the x-axis and acceleration on the y-axis. We can see the plot in 2.3, where the data points have been plotted and fit with a line.

We see that the acceleration changes linearly (as a straight line) with time. The acceleration is positive and decreasing over the entire duration of the graph.

Now we can use the a vs. t graph to create our v vs. t graph. The change in velocity over time is the area under the curve of the a vs. t graph. With the acceleration changing over time, how can we find the area under the curve?

Let's focus first on the interval $t = 0$ s to $t = 1$ s. The acceleration is linear, so the average acceleration over the interval is

$$a_{avg} = \frac{a(t=0) + a(t=1)}{2} \quad (2.12)$$

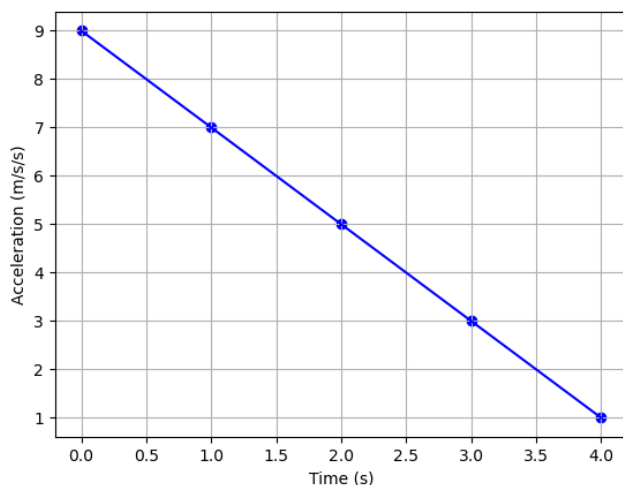


Figure 2.3: Graph created from data in 2.1

The acceleration is a function of time, so $a(t = 0)$ means the acceleration at $t = 0 \text{ s}$ and $a(t = 1)$ means the acceleration at $t = 1 \text{ s}$. The average acceleration between $t = 0 \text{ s}$ and $t = 1 \text{ s}$ is

$$a_{avg} = \frac{9 \text{ m/s}^2 + 7 \text{ m/s}^2}{2} = 8 \frac{\text{m}}{\text{s}^2} \quad (2.13)$$

The area under the curve is the average acceleration multiplied by the time, which is 1 s for this interval

$$\Delta v = a_{avg} \cdot \Delta t \quad (2.14)$$

$$\Delta v = (8 \frac{\text{m}}{\text{s}^2})(1 \text{ s}) = 8 \frac{\text{m}}{\text{s}} \quad (2.15)$$

The velocity at $t = 1 \text{ s}$ is the velocity at $t = 0 \text{ s}$ plus the change in velocity between $t = 0 \text{ s}$ and $t = 1 \text{ s}$.

$$v(t = 1) = v(t = 0) + \Delta v \quad (2.16)$$

$$v(t = 1) = -4 \frac{\text{m}}{\text{s}} + 8 \frac{\text{m}}{\text{s}} = 4 \frac{\text{m}}{\text{s}} \quad (2.17)$$

Try repeating this process for each one second interval. The velocities that you should get are shown in 2.2.

Table 2.2: Velocity vs Time example

| Time (s) | Velocity (m/s) |
|----------|----------------|
| 0 | -4 |
| 1 | 4 |
| 2 | 10 |
| 3 | 14 |
| 4 | 16 |

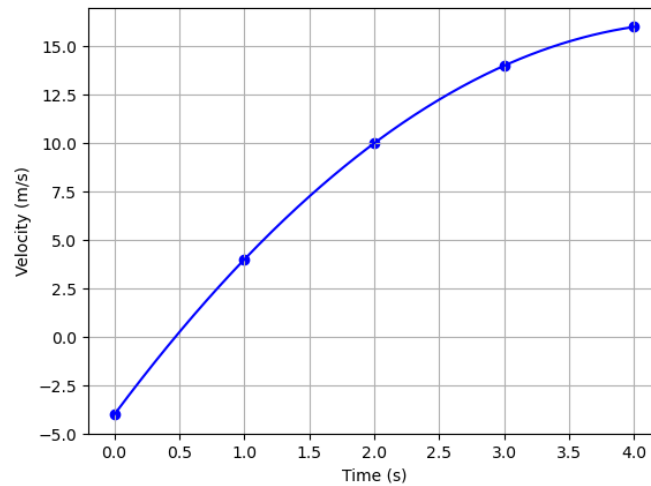


Figure 2.4: Velocity vs time graph created from the data in 2.1 and initial condition of $v = -4 \frac{m}{s}$ at $t = 0 s$.

This graph is a parabola which opens downward. If you are taking or have taken calculus, you would recognize this as being the integral of the a vs. t graph.

2.3 Constant Acceleration Motion

One case of particular interest we will look at is **constant acceleration motion**, where the acceleration does not change over an interval of time. This is useful for 2 reasons:

1. Common. There are several examples of constant acceleration motion or motion that is close enough to having a constant acceleration that we can approximate it as such. This gives us a useful way to describe a variety of physical situations.
2. Easy to work with mathematically.

We will start by saying that acceleration is constant as a function of time

$$\begin{aligned} a(t) &= a \\ a_i &= a_f = a \end{aligned} \tag{2.18}$$

Now we will consider motion from a time of zero ($t = 0$ s) to some later time t . This means that our Δt is

$$\Delta t = t - 0 = t \tag{2.19}$$

Now we can work with our definition of average acceleration. If the acceleration is constant, then the average acceleration is the constant value of acceleration $a_{avg} = a$. Let's see if we can find the velocity v_f at a time t starting with the definition of acceleration.

$$a = \frac{\Delta v}{\Delta t} \tag{2.20}$$

$$a = \frac{v_f - v_0}{t} \tag{2.21}$$

Above we see a subscript zero for the first time. The variable v_0 is pronounced “v naught” and is used to indicate the value of a motion variable at time zero. It is the special case of the initial state being at time zero.

Continuing on with the equation, we can multiply both sides by t .

$$at = v_f - v_0 \tag{2.22}$$

Now we can add v_0 to both sides to get.

$$v_f = v_0 + at \quad (2.23)$$

Which gives us the final velocity v_f at a time t given the velocity at $t = 0$ s is v_0 and the acceleration is a constant value of a . This is the first of the **kinematic equations**, which are used to describe constant acceleration motion.

Constant acceleration motion has a linear velocity, which means we can find the average velocity as

$$v_{avg} = \frac{v_0 + v_f}{2} \quad (2.24)$$

From the definition of average velocity, we also have

$$v_{avg} = \frac{\Delta x}{t} \quad (2.25)$$

Now we can set the right hand sides of 2.24 and 2.25 equal and solve for Δx

$$\frac{\Delta x}{t} = \frac{v_0 + v_f}{2} \quad (2.26)$$

$$\Delta x = \frac{v_0 + v_f}{2} \cdot t \quad (2.27)$$

This is the second kinematic equation, used for finding the change in position if we know the initial and final velocities. We can get a third kinematic equation that does not include v_f if we use the definition of v_f from equation 2.23 and plug that into equation 2.27.

$$\Delta x = \frac{v_0 + v_0 + at}{2} \cdot t \quad (2.28)$$

$$x_f - x_0 = (v_0 + \frac{1}{2}at) \cdot t \quad (2.29)$$

$$x_f = x_0 + v_0t + \frac{1}{2}at^2 \quad (2.30)$$

This is the third kinematic equation and gives us the final position as a function of initial quantities.

So far, these 3 kinematic equations have all included time as a variable. What if we don't know or care about the time but want to study how other quantities are related. Let's return to 2.23 and solve for time t .

$$v_f = v_0 + at \quad (2.31)$$

$$v_f - v_0 = at \quad (2.32)$$

$$t = \frac{v_f - v_0}{a} \quad (2.33)$$

Now we can take this expression for t and plug it into 2.30.

$$x_f = x_0 + v_0 \left(\frac{v_f - v_0}{a} \right) + \frac{1}{2} a \left(\frac{v_f - v_0}{a} \right)^2 \quad (2.34)$$

We can simplify by moving x_0 and working with the terms in parentheses.

$$x_f - x_0 = \frac{1}{a}(v_f v_0 - v_0^2) + \frac{1}{2} a \left(\frac{v_f^2 - 2v_f v_0 + v_0^2}{a^2} \right) \quad (2.35)$$

This may look like a mess but don't worry, we will get a nice equation when we are done. We will replace $x_f - x_0$ with Δx , since that is the definition of the change in position. Also, we will simplify the second term on the right hand side by distributing the $\frac{1}{2}$ and pulling $\frac{1}{a^2}$ out of the parentheses.

$$\Delta x = \frac{1}{a}(v_f v_0 - v_0^2) + \frac{1}{a} \left(\frac{1}{2} v_f^2 - v_f v_0 + \frac{1}{2} v_0^2 \right) \quad (2.36)$$

If we multiply both sides by a , we can remove the parentheses and group the like terms on the right hand side.

$$a\Delta x = v_f v_0 - v_0^2 + \frac{1}{2} v_f^2 - v_f v_0 + \frac{1}{2} v_0^2 \quad (2.37)$$

$$a\Delta x = \frac{1}{2} v_f^2 + v_f v_0 - v_f v_0 + \frac{1}{2} v_0^2 - v_0^2 \quad (2.38)$$

The $v_f v_0$ cross terms actually cancel out!

$$a\Delta x = \frac{1}{2} v_f^2 - \frac{1}{2} v_0^2 \quad (2.39)$$

Now we can solve this equation for v_f^2 .

$$2a\Delta x = v_f^2 - v_0^2 \quad (2.40)$$

$$v_f^2 = v_0^2 + 2a\Delta x \quad (2.41)$$

Here we have the fourth kinematic equation, which does not include time but instead the change in position, initial velocity, final velocity, and acceleration.

All of the equations are collected in 2.4. When presented with a problem, looking at what variable each kinematic equation is missing can help you choose what equation to use. We will go through a few examples and walk through problem solving in more detail. Throughout studying motion with kinematic equations, you need to remember that **we can only use the kinematic equations with constant acceleration motion**. We derived these equations from the starting point of constant acceleration, so if acceleration changes then these equations are no longer valid.

Table 2.3: Kinematic Equation

| Equation number | Formula | Missing variable |
|-----------------|--|-------------------|
| 1 | $v_f = v_0 + at$ | No position |
| 2 | $\Delta x = \frac{v_0 + v_f}{2} \cdot t$ | No acceleration |
| 3 | $x_f = x_0 + v_0 t + \frac{1}{2}at^2$ | No final velocity |
| 4 | $v_f^2 = v_0^2 + 2a\Delta x$ | No time |

2.4 Problem Solving with Kinematic Equations

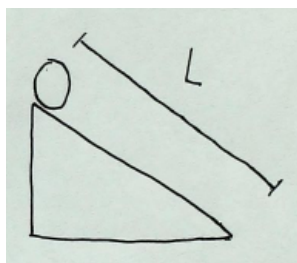
We can follow a basic outline to solving constant acceleration problems using kinematic equations. A typical problem will have a couple sentences describing a physical scenario and asking for a quantity of the motion. When approaching a problem, you should go use the following steps:

1. Identify the known quantities in the problem. This is the information the problem gives you. Drawing a picture is often useful!
2. Identify the quantity that we are trying to solve for.
3. Choose a kinematic equation that contains the known quantities and the one we wish to solve for.
4. Solve the kinematic equation algebraically for the quantity of motion we are trying to find. Do not plug in numbers before doing this!
5. Plug in numbers and solve the problem.

It is important to not plug in numbers early because that makes it difficult to find any mistakes in the solution. Finding algebra mistakes with variables is much easier than finding mistakes with numbers.

Example 2.3

A ball rolling down a ramp accelerates at 3.0 m/s^2 . If the ramp is $L = 2.5 \text{ m}$ long and the ball is released from rest at the top of the ramp, how what is the velocity of the ball when it reaches the bottom of the ramp?



First we can take stock of what we know. We have the acceleration $a = 3.0 \text{ m/s}^2$ given. The ball travels from the top to the bottom of the ramp, so the change in position is the length of the ramp $\Delta x = 2.5 \text{ m}$. That may look

like all of the given information, but there is one more thing we are given. The problem states that the ball “is released from rest”, which means that the ball is stationary when the motion starts (“at rest” means it is still or stationary). Therefore, we also know that $v_0 = 0 \text{ m/s}$.

We want to solve for the velocity at the bottom of the ramp. This is the final state of the problem, so that is our final velocity v_f .

We don’t have time, so it makes sense to look at the 4th kinematic equation.

$$v_f^2 = v_0^2 + 2a\Delta x \quad (2.42)$$

The only value we don’t have here is v_f , so this will work. Let’s take the square root of both sides to solve for velocity.

$$v_f = \sqrt{v_0^2 + 2a\Delta x} \quad (2.43)$$

Plug in numbers.

$$v_f = \sqrt{(0 \text{ m/s})^2 + 2(3.0 \text{ m/s}^2)(2.5 \text{ m})} \quad (2.44)$$

$$v_f = \sqrt{15 \text{ m}^2/\text{s}^2} = 3.9 \text{ m/s} \quad (2.45)$$

The ball has a velocity of 3.9 m/s when it reaches the bottom of the ramp. Notice how tracking the units shows we get units of m/s at the end to help us verify that our work was correct. This is why it is good to track units!

Example 2.4

A car that is driving at 34 m/s begins to brake, slowing down at an acceleration of 1.5 m/s^2 . If the car ends at a velocity of 28 m/s , how much time was it slowing down for?

We know that the car starts at a velocity of $v_0 = 34.0 \text{ m/s}$ and ends at a velocity of $v_f = 28.0 \text{ m/s}$. The acceleration of the car is causing it to slow down, which means the acceleration is in the opposite direction as the velocity. Therefore, the acceleration is negative, with a value of $a = -1.5 \text{ m/s}^2$.

We need to solve for the time t that it takes to change from v_0 to v_f with an acceleration a . This means we want to use the 1st kinematic equation.

$$v_f = v_0 + at \quad (2.46)$$

Now we can solve this algebraically for t .

$$v_f - v_0 = at \quad (2.47)$$

$$t = \frac{v_f - v_0}{a} \quad (2.48)$$

Now we can plug in numbers

$$t = \frac{28.0 \text{ m/s} - 34.0 \text{ m/s}}{-1.5 \text{ m/s}^2} = \frac{-6.0 \text{ m/s}}{-1.5 \text{ m/s}^2} = 4.0 \text{ s} \quad (2.49)$$

The car has to hit the brakes for 4.0 s to slow down as described in the problem.

The general outline to problem solving here will be useful throughout your studies in physics. The variables and equations involved will change, but the method remains the same.

2.5 Gravity and Free-fall

One common scenario where we have constant acceleration is when an object is being acted on by gravity at the Earth's surface. At the Earth's surface, **gravity** causes an object to accelerate straight down toward to the center of the Earth. When an object is accelerating only due to gravity, we call that *free-fall*. Free-fall ignores other effects like air resistance (which we will talk about more later), so it works in cases where those effects are very small. An example is a thrown baseball. While the ball is airborne, the only acceleration it experiences is gravity.

Since free-fall is a special case that we will regularly encounter, we use the variable g to denote free-fall acceleration. The acceleration due to gravity is

$$g = 9.8 \frac{\text{m}}{\text{s}^2} \quad (2.50)$$

directed straight down. If we define up as positive, then $a = -g$. If we define down as positive, then $a = g$.

Example 2.5

Mary drops a rock from a bridge that is 6.0 m above the water. If she releases the rock from rest, how much time does it take for the rock to hit the water?

First we need to define our position coordinates to determine if up or down is positive. Let's say that the water is at a height of 0 m and the bridge is at a height of 6.0 m . This means that $x_0 = 6.0\text{ m}$ and $x_f = 0.0\text{ m}$. Since up is positive, we know that the acceleration is $a = -g = -9.8 \frac{\text{m}}{\text{s}^2}$. Mary releases the rock from rest, so $v_0 = 0\text{ m/s}$.

We are trying to solve for the time t . Based on what we have, it makes the most sense to use the 3rd kinematic equation.

$$x_f = x_0 + v_0 t + \frac{1}{2} a t^2 \quad (2.51)$$

This is a quadratic equation for time, and we are trying to solve for time. In general we would have to use the quadratic formula. However, if we plug in $v_0 = 0\text{ m/s}$, the term that is linear in t (has t to the first power) cancels out and we can solve it more easily.

$$x_f = x_0 + \frac{1}{2} a t^2 \quad (2.52)$$

$$\frac{1}{2} a t^2 = x_f - x_0 \quad (2.53)$$

$$t^2 = \frac{2}{a} (x_f - x_0) \quad (2.54)$$

$$t = \sqrt{\frac{2}{a} (x_f - x_0)} \quad (2.55)$$

$$t = \sqrt{\frac{2}{-9.8\text{ m/s}^2} (0\text{ m} - 6.0\text{ m})} \quad (2.56)$$

$$t = \sqrt{\frac{2}{-9.8\text{ m/s}^2} (-6.0\text{ m})} \quad (2.57)$$

The negative signs cancel, so we have a positive value in the radical.

$$t = \sqrt{\frac{2 \cdot 6.0\text{ m}}{9.8\text{ m/s}^2}} = 1.2\text{ s} \quad (2.58)$$

The rock will hit the water 1.2 s after she drops it.

One thing to remember is that an object that is airborne will always be accelerated at a magnitude g straight down. Magnitude here means the number, so the acceleration is always $9.8 \frac{m}{s^2}$ downward. This is true whether an object is rising or falling. The acceleration is not necessarily in the same direction as the velocity!

One question that we will often look at is the maximum height an object reaches when it is thrown or launched upward. When studying this the final state is when the object reaches its maximum height. **If an object is initially moving straight upward, at the highest point its velocity is zero.** We can think about this by thinking about how the velocity changes. On the way up, the velocity is upward. As it falls back down, the velocity is downward. The highest point is where the velocity crosses zero when switching from upward to downward.

Example 2.6

Rocky tosses a ball straight up in the air at a velocity of $12.0 m/s$. He releases the ball at a height of $h = 1.8 m$ off the ground. What is the maximum height the ball reaches?

We will call up positive so that the initial velocity is $v_0 = 12.0 m/s$. The initial position is the starting height, so $x_0 = 1.8 m$. The final state is the highest point the ball reaches, so its final velocity is zero $v_f = 0 m/s$. The ball is in free-fall and up is positive, so $a = -g = -9.8 \frac{m}{s^2}$.

Finding the highest point means we want the position in the final state, so we are solving for x_f . Looking at the kinematic equations none of them look quite right at first glance, but remember the definition of change in position is $\Delta x = x_f - x_0$. With that, we can use the 4th kinematic equation.

$$v_f^2 = v_0^2 + 2a\Delta x \quad (2.59)$$

$$v_f^2 = v_0^2 + 2a(x_f - x_0) \quad (2.60)$$

Now we can solve for x_f .

$$v_f^2 - v_0^2 = 2a(x_f - x_0) \quad (2.61)$$

$$\frac{v_f^2 - v_0^2}{2a} = x_f - x_0 \quad (2.62)$$

$$x_f = \frac{v_f^2 - v_0^2}{2a} + x_0 \quad (2.63)$$

$$x_f = \frac{(0 \text{ m/s})^2 - (12.0 \text{ m/s})^2}{2 \cdot (-9.8 \text{ m/s}^2)} + 1.8 \text{ m} \quad (2.64)$$

$$x_f = \frac{144 \text{ (m/s)}^2}{19.6 \text{ m/s}^2} + 1.8 \text{ m} = 9.1 \text{ m} \quad (2.65)$$

The ball reaches a maximum height of 9.1 m off the ground.

2.6 Motion in 2 Dimensions

So far we have only talked about motion in one dimension, such as an object falling straight down or sliding along the floor. So how do we handle objects that don't move along a straight line? This is called **2 dimensional (2D) motion**. We can split the motion along 2 perpendicular axes and treat those separately! This will often be a horizontal and vertical split, but can be any perpendicular directions. The standard convention is that horizontal motion is denoted with an x and vertical motion is denoted with a y .

We use the subscript x and y to denote which axis we are looking at. Notice how all positions, velocities (initial and final), and accelerations are split along the axes. The only connection between the 2 axes is the time! For example, we can find the time when an object hits the ground using the y -axis motion and use that time to find how far along the x -axis it moved.

Example 2.7

Table 2.4: Kinematic Equations in 2 Dimensions

| Equation number | x | y |
|-----------------|--|--|
| 1 | $v_{f,x} = v_{0,x} + a_x t$ | $v_{f,y} = v_{0,y} + a_y t$ |
| 2 | $\Delta x = \frac{v_{0,x} + v_{f,x}}{2} \cdot t$ | $\Delta y = \frac{v_{0,y} + v_{f,y}}{2} \cdot t$ |
| 3 | $x_f = x_0 + v_{0,x} t + \frac{1}{2} a_x t^2$ | $y_f = y_0 + v_{0,y} t + \frac{1}{2} a_y t^2$ |
| 4 | $v_{f,x}^2 = v_{0,x}^2 + 2a_x \Delta x$ | $v_{f,y}^2 = v_{0,y}^2 + 2a_y \Delta y$ |

Chris throws a rock horizontally off a bridge at $4.5 \frac{m}{s}$. If the surface of the water is $12 m$ below where they threw the rock from, how far (horizontally) from the bridge does the rock hit the water?

Let's take stock of what we know along the x and y directions. We know that Chris throws the rock horizontally, so the entire velocity is along the x direction. Define up as positive y , and the direction they throw the rock as positive x .

- $v_{0,x} = 4.5 \frac{m}{s}$
- $v_{0,y} = 0 \frac{m}{s}$
- $a_y = -g = -9.8 \frac{m}{s^2}$
- $a_x = 0 \frac{m}{s^2}$
- $y_0 = 12.0 m$
- $y_f = 0 m$

We can also set $x_0 = 0 m$ and we will want to find x_f . The first thing that we need to do is use the vertical motion to determine how long the rock takes to fall.

$$y_f = y_0 + v_{0,y}t + \frac{1}{2}a_yt^2 \quad (2.66)$$

Plug in $v_{0,y} = 0 \frac{m}{s}$ and $y_f = 0 m$ to simplify the algebra.

$$-y_0 = \frac{1}{2}a_yt^2 \quad (2.67)$$

$$t = \sqrt{\frac{-2y_0}{a_y}} \quad (2.68)$$

$$t = \sqrt{\frac{-2(12.0 m)}{(-9.8 m/s^2)}} = 1.56 s \quad (2.69)$$

Now we can use this time to solve for x_f

$$x_f = x_0 + v_{0,x}t + \frac{1}{2}a_xt^2 \quad (2.70)$$

Since $x_0 = 0 m$ and $a_x = 0 \frac{m}{s^2}$, when we plug in numbers we get:

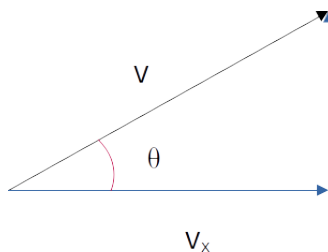


Figure 2.5: Shows the vector representation of v_x and v_y for a velocity vector v at an angle θ above the x -axis. Notice how this is a right triangle with the components as the legs and the velocity as the hypotenuse.

$$x_f = (4.5 \frac{m}{s})(1.56 s) = 7.0 m \quad (2.71)$$

The rock will hit the water at a horizontal distance of $7.0 m$ from where Chris threw it.

To split the motion of an object along 2 axes we will use trigonometry. See how the vector and components form a right triangle in Figure 2.5. The opposite side of the triangle is v_y and the adjacent side of the triangle is v_x . This means we can relate the components to the original velocity as:

$$v_x = v \cos(\theta) \quad (2.72)$$

$$v_y = v \sin(\theta) \quad (2.73)$$

Example 2.8

A potato cannon is a device that uses compressed air to launch a potato out of a tube at a high speed. If such a potato cannon is aimed at an angle of 30.0° above the ground and a potato is fired with an initial velocity of $38.0 \frac{m}{s}$, what are the horizontal and vertical components of the initial velocity?

We are given that $v_0 = 38.0 \frac{m}{s}$ and $\theta = 30.0^\circ$. Since the angle is above horizontal, the horizontal component is the adjacent side of the right triangle and the vertical component is the opposite side of the right triangle. First the horizontal component

$$v_{0,x} = v_0 \cos(\theta) = (38.0 \frac{m}{s}) \cos(30^\circ) = 29.9 \frac{m}{s} \quad (2.74)$$

Now the horizontal component

$$v_{0,y} = v_0 \sin(\theta) = (38.0 \frac{m}{s}) \sin(30^\circ) = 23.4 \frac{m}{s} \quad (2.75)$$

The horizontal component of the velocity is $v_{0,x} = 29.9 \frac{m}{s}$ and the vertical component of the velocity is $v_{0,y} = 23.4 \frac{m}{s}$.

2.7 Projectile Motion

One specific example of 2D motion that we will explore further is **projectile motion**, which is motion where the only acceleration is the free-fall acceleration of gravity. We ignore effects of air resistance so it is an *approximation* of the physical world. This approximation is good for small objects moving at slower speeds. The assumption of free-fall acceleration tells us what the acceleration is (we will treat up as positive for this entire section)

$$a_y = -g = -9.8 \frac{m}{s^2} \quad (2.76)$$

$$a_x = 0 \frac{m}{s^2} \quad (2.77)$$

We will return back to the potato cannon from the end of the previous section for our study of projectile motion. The potato cannon fires a potato at initial velocity v_0 at an angle of θ above horizontal. This gives us

$$v_{0,x} = v_0 \cos(\theta) \quad (2.78)$$

$$v_{0,y} = v_0 \sin(\theta) \quad (2.79)$$

Let's suppose that the potato cannon is fired over level ground and we want to know how far from the cannon the potato will land. Level ground here means that the potato will land at the same height it was launched from. We are trying to solve for Δx under the condition that $\Delta y = 0 m$. First we can use the vertical equation to find the time that the potato is airborne

$$y_f = y_0 + v_{0,y}t + \frac{1}{2}a_yt^2 \quad (2.80)$$

$$y_f - y_0 = (v_{0,y} + \frac{1}{2}a_y t) \cdot t \quad (2.81)$$

Since we are on the condition that $\Delta y = 0 \text{ m}$ we know that that $y_f - y_0 = 0 \text{ m}$.

$$0 = (v_{0,y} + \frac{1}{2}a_y t) \cdot t \quad (2.82)$$

There are two solutions, one where $t = 0 \text{ s}$ and one where $v_{0,y} + \frac{1}{2}a_y t = 0$. The $t = 0 \text{ s}$ solution is when the potato is launched! This is not the solution of interest to us, so we will look at the nonzero solution.

$$0 = v_{0,y} + \frac{1}{2}a_y t \quad (2.83)$$

$$\frac{-1}{2}a_y t = v_{0,y} \quad (2.84)$$

$$t = \frac{2v_{0,y}}{a_y} \quad (2.85)$$

Now let's plug in $v_{0,y} = v_0 \sin(\theta)$ and $a_y = -g$

$$t = \frac{2v_0 \sin(\theta)}{g} \quad (2.86)$$

With the time, we can look at the horizontal component.

$$x_f = x_0 + v_{0,x}t + \frac{1}{2}a_x t^2 \quad (2.87)$$

$$x_f - x_0 = v_{0,x}t + \frac{1}{2}a_x t^2 \quad (2.88)$$

Remember that $\Delta x = x_f - x_0$ so we can switch the left hand side to Δx . This is where we will also replace $v_{0,x} = v_0 \cos(\theta)$. Now we can plug in $a_x = 0 \frac{\text{m}}{\text{s}^2}$ from our assumption about projectile motion and $t = \frac{2v_0 \sin(\theta)}{g}$.

$$\Delta x = v_0 \cos(\theta) \left(\frac{2v_0 \sin(\theta)}{g} \right) \quad (2.89)$$

$$\Delta x = \frac{2v_0^2 \cos(\theta) \sin(\theta)}{g} \quad (2.90)$$

If we use the trigonometry identity

$$\sin(2\theta) = 2 \cos(\theta) \sin(\theta) \quad (2.91)$$

We can write this as

$$\Delta x = \frac{v_0^2 \sin(2\theta)}{g} \quad (2.92)$$

which is known as the **range equation**. This gives the horizontal distance a projectile travels before it lands at the same height it was launched from.

There is one more question of interest we should look at, and that is the maximum height a projectile reaches. The highest point is when the object switches from moving upward to downward, so its vertical velocity is temporarily zero. We want to find Δy for when $v_{f,y} = 0 \text{ m/s}$.

$$v_{f,y}^2 = v_{0,y}^2 + 2a_y \Delta y \quad (2.93)$$

Plug in our highest point condition of $v_{f,y} = 0 \frac{\text{m}}{\text{s}}$

$$0 = v_{0,y}^2 + 2a_y \Delta y \quad (2.94)$$

$$2a_y \Delta y = -v_{0,y}^2 \quad (2.95)$$

$$\Delta y = \frac{-v_{0,y}^2}{2a_y} \quad (2.96)$$

Plug in $v_{0,y} = v_0 \sin(\theta)$ and $a_y = -g$ to get

$$\Delta y = \frac{v_0^2 \sin^2(\theta)}{2g} \quad (2.97)$$

2.8 Uniform Circular Motion

While we did look a lot at constant acceleration motion throughout this chapter, there are some types of motion with varying acceleration that are simple enough for us to study. One example is **uniform circular motion**, which is defined as

- The object moves in a circular path.
- The object moves at a constant speed.

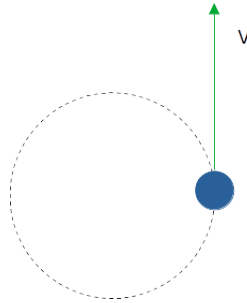


Figure 2.6: In uniform circular motion, the instantaneous velocity is tangential (along the outside) of the circular path the object moves along.

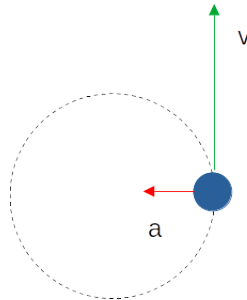


Figure 2.7: In uniform circular motion, the instantaneous acceleration is radial, directed toward the center of the circular path the object moves along.

This might sound like a very specific set of conditions that wouldn't be widely applicable, but it is a valuable *approximation*. A major goal of physics is to find useful simplifications to difficult problems in order to study them easily. There are many physical systems in which an object moves in a way that is very close to being uniform circular motion and we can approximate them as such. This could include cars going around a curve in the road and ferris wheels.

What is the acceleration of an object in uniform circular motion? The first part of answering that is figuring out the direction of the acceleration. In order for an object to move in a circular path, its velocity must be along the circle. That means the instantaneous velocity is always **tangential** to the circle, as seen in Figure 2.6.

The speed of the object is constant so the magnitude of the velocity is constant. The direction of the velocity is always changing in order to keep the velocity tangential to the circle. As we saw in Section REF HERE, we can consider perpendicular motion separately. That means that if the force

is perpendicular to the velocity, it will change the direction of the velocity without changing the magnitude of the velocity. Therefore, the acceleration is directed radially toward the center of the circle, as seen in Figure 2.7.

Now to find the magnitude of the acceleration. Doing this rigorously requires calculus but we can still get some understanding of the result.

What if the velocity increases? Increasing the velocity will increase the change in velocity. If the velocity vectors are longer, then the Δv vector will be longer as well. Since acceleration is $a_{avg} = \frac{\Delta v}{\Delta t}$, increasing Δv will increase the acceleration. Increasing velocity will also decrease the time interval it takes for the object to travel along the path of the circle. Decreasing Δt will increase the acceleration.

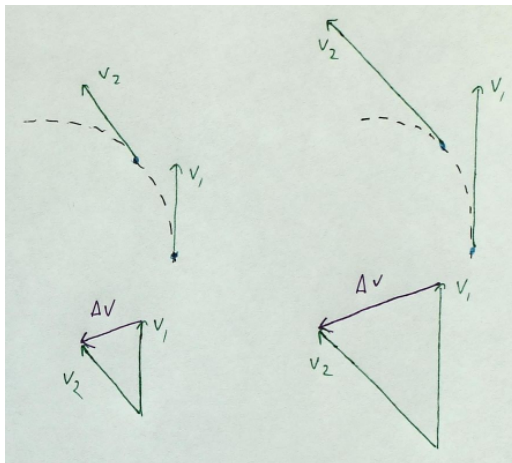


Figure 2.8: A greater velocity will result in a larger change in velocity as the object moves around the circle. That is part of the reason why increasing the velocity results in a larger acceleration.

What if the radius of the circle increases? That increases the distance the object has to move for the same change in velocity Δv . If the velocity is the same, that will increase the time required. This means that increasing the radius decreases the acceleration.

We have 2 ways that increasing velocity increases acceleration and 1 way that increasing radius decreases acceleration. This would lead us to propose

$$a_c = \frac{v^2}{r} \quad (2.98)$$

as the acceleration, which is in fact exactly what we would get if we did the required calculus to solve it! This acceleration is called the **centripetal**

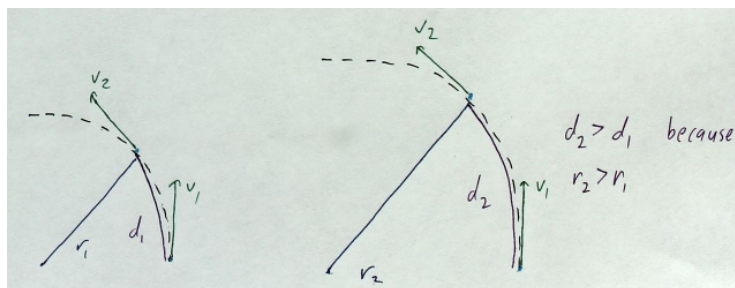


Figure 2.9: If the radius is greater, the object has to move a greater distance around the circle for the same change in velocity. This takes more time, causing the acceleration to decrease.

acceleration, where centripetal means directed toward the center. That's why we denote this acceleration with a subscript c .

The magnitude of the acceleration is constant, so why isn't this constant acceleration motion? The answer lies in the fact that the acceleration is always directed toward the center of the circle. As the object moves around the circle, the direction of the acceleration must be constantly changing so that the acceleration is always pointed toward the center of the circle. That changing direction is why the acceleration isn't constant in uniform circular motion.

Example 2.9

Alex is jogging at $2.5 \frac{m}{s}$ around a circular track of radius $10.0 m$. What is their centripetal acceleration?

We can plug in the velocity and radius into Equation 2.98

$$a_c = \frac{(2.5 m/s)^2}{10.0 m} = 0.625 \frac{m}{s^2} \quad (2.99)$$

Alex has a centripetal acceleration of $a_c = 0.625 \frac{m}{s^2}$.

Example 2.10

An amusement park wants a rollercoaster to have a centripetal acceleration that is "half a g " ($\frac{g}{2} = 4.9 \frac{m}{s^2}$) while going around a circular curve. If rollercoaster is moving at a velocity of $21 \frac{m}{s}$ around the curve, what is the radius of the curve?

First we can start with Equation 2.98 and solve for the radius.

$$a_c = \frac{v^2}{r} \quad (2.100)$$

$$a_c r = v^2 \quad (2.101)$$

$$r = \frac{v^2}{a_c} = \frac{(34 \text{ m/s})^2}{4.9 \text{ m/s}^2} = 90.0 \text{ m} \quad (2.102)$$

The radius of the curve has to be 90.0 m . (You can think of this as if the curve was extended to a full circle, that circle would have a radius of 90.0 m .)

Chapter 3

Forces

We spent the previous chapter talking about motion, but we have yet to address the physical reason *why* objects move. What is the interaction that causes motion? That is the question that will motivate this chapter.

3.1 Newton's Laws of Motion

Isaac Newton (1643 - 1727) formulated 3 laws of motion that are a major foundation of physics. His laws describe how **forces** act on objects to affect their motion. Forces can be thought of as a push or pull on an object. We will go into more detail on specific forces later, but for now we can think of them as a push or pull.

Newton's 1st Law of Motion states that

“An object in motion remains in uniform motion and an object at rest remains at rest, unless acted upon by an external force”

What this means that an object will remain in its current state unless an outside force acts on it. If the object is at rest, it will not start moving unless an outside force acts on it. If it is moving, it will keep moving at a constant velocity (meaning same speed and direction) unless a force acts on it.

Newton's 2nd Law of Motion states that

“The net force acting on an object is equal to the product of its mass and acceleration.”

$$\overrightarrow{F_{net}} = m \overrightarrow{a} \quad (3.1)$$

Newton's 2nd Law of Motion gives the relationship between force as an equation we can use to mathematically describe the motion. Force and acceleration both being vector quantities means that the acceleration is in the same direction as the net force.

Net force means the sum of all forces acting on an object. In most cases objects will have 2 or more forces acting on them, so we will need to add those forces to get the net force acting on the object.

An object's mass determines its resistance to acceleration from a force. This resistance is called **inertia**. The larger the mass, the higher the inertia, and the less an object will accelerate when acted upon by a given force.

Newton's 3rd Law of Motion states that

“When object A exerts a force on object B, object B exerts a force on object A that is equal in magnitude and opposite in direction.

What this means is that when one object pushes on another, the second object pushes back on the first. When the law says “equal in magnitude”, it means that the forces are the same numerical value. If object A exerts a 5 *Newton* force on object B, object B exerts a 5 *Newton* force on object A. This is also called the **action-reaction principle**. For every force, there is an equal and opposite reaction force. The reaction force has the same strength, but is in the opposite direction.

One way you can see this for yourself is to push on a wall. What happens when you do that? You will feel like you are being pushed away from the wall! That is because when you exert a force on the wall, the wall exerts an equal and opposite force on you.

3.2 Force exerted by gravity

In Section 2.5 we looked at the acceleration due to gravity and how it is a constant acceleration $g = 9.8 \frac{m}{s^2}$ for any object in free-fall. Let's apply Newton's 2nd Law of Motion free-fall. Start with 3.1 and plug in g for the acceleration.

$$F_{net} = mg \quad (3.2)$$

For an object in free fall, the net force has magnitude mg direction straight down. Since the only acceleration in free-fall is due to gravity, we know that the only force acting on an object in free-fall is gravity. Therefore, we can say that the force of gravity acting on an object of mass m is

$$F_g = mg \quad (3.3)$$

which will always be directed straight down. This force acts on any object at the Earth's surface, even if the object is not in free-fall. Early on, we said that weight is not the same as mass. **Weight** is the force of gravity acting on an object. Weight is proportional to mass, but depends on the surface gravity. This is why astronauts weighed less on the moon, despite their mass remaining constant!

Example 3.1

Lucy has to move a box that weighs 235 N . What is the mass of the box in kilograms?

Here we will use equation 3.3 and solve for the mass m .

$$m = \frac{F_g}{g} \quad (3.4)$$

The force of gravity is the weight, 235 N .

$$m = \frac{235 \text{ N}}{9.8 \text{ m/s}^2} = 24.0 \text{ kg} \quad (3.5)$$

The box has a mass of $m = 24 \text{ kg}$.

We can say that free-fall is when the only force acting on an object is the gravitational force

$$F_{net} = F_g \quad (3.6)$$

3.3 Normal Force

Set any object on the table in front of you (pencil, book, cell phone, literally anything). Is it accelerating? No! The object set on the table is sitting stationary. That means its acceleration is zero, and from Newton's first law we know that means there is no net force acting on the object.

The force of gravity is acting down on it, so there must be a force acting upward. To determine the source of the force, we can think about why the object isn't falling down. From real life experience, we can say it is obviously the table! The table stops the object from falling by exerting an upward force on the object. The force a surface exerts on an object is called the **normal force**. This is the force that keeps the object from moving "through" the surface. Normal force is always perpendicular to the surface, so on a perfectly horizontal table it is direction straight up. We will denote normal force as F_N .

Example 3.2

A book of mass 0.40 kg is sitting stationary on a table. What is the normal force acting on the book?

Let's call up positive. Since the book is stationary, we know its acceleration is zero and therefore the net force is zero. The gravitational force is F_g directed downward.

$$F_g = -mg \quad (3.7)$$

The normal force is directed upward to make the net force zero.

$$F_{net} = F_N + F_g = 0\text{ N} \quad (3.8)$$

$$F_N = -F_g = -(-mg) \quad (3.9)$$

$$F_N = (0.40\text{ kg})(9.8\frac{m}{s^2}) = 3.9\text{ N} \quad (3.10)$$

The normal force for an object sitting on a horizontal surface with no other forces acting on it is the same as the gravitational force, just directed upward.

3.4 Forces in 2 Dimensions

Many physical situations have forces exerted along 2 dimensions. For example, if we have a box sitting on the floor and somebody pushes it we have 3 forces acting on the box:

1. Gravity acting downward.

2. Normal force acting upward.
3. The person's push acting horizontally.

This scenario has forces that act vertically and horizontally. When studying motion like this, we can take advantage of the fact that **perpendicular forces are independent**. What this means is that a vertical force acting on an object has no impact on the horizontal motion of that object. For the box being pushed along the floor, we can analyze gravity and normal force together since they are vertical forces. Then we can look at the push separately since that is horizontal.

When we have forces that aren't perfectly horizontal or vertical, we can split them into vertical and horizontal components. Figure 3.1 shows how a force angled at an angle θ is split into horizontal and vertical components.

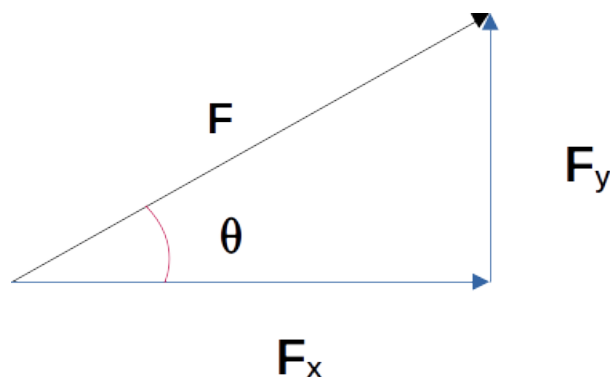


Figure 3.1: A force directed at an angle θ above horizontal split into components. Horizontal is F_x and vertical is F_y .

You can see that the force and components form a right triangle, with the force F as the hypotenuse and the components F_x and F_y are the legs of the right triangle. We can use **trigonometry** to find the sides. From a previous math class you should have seen the trigonometry functions sine, cosine, and tangent. They give the relationships between the opposite leg, adjacent leg, and hypotenuse of a right triangle.

$$\sin(\theta) = \frac{\text{opposite}}{\text{hypotenuse}} \quad (3.11)$$

$$\cos(\theta) = \frac{\text{adjacent}}{\text{hypotenuse}} \quad (3.12)$$

$$\tan(\theta) = \frac{\textit{opposite}}{\textit{adjacent}} \quad (3.13)$$

From the angle θ the horizontal component is the adjacent side, so we can find F_x as

$$F_x = F \cos(\theta) \quad (3.14)$$

The vertical component is the opposite side of the triangle, so we can find F_y as

$$F_y = F \sin(\theta) \quad (3.15)$$

3.5 Free Body Diagrams

If we return to looking at a book on the table, we looked at a situation where there are multiple forces acting in different directions on an object. This can be difficult to keep track of as just text, so in these cases we draw pictures known as **free body diagrams**. These diagrams are simple pictures that show the forces acting on one object. The forces are drawn as arrows in the direction of the force, with the relative length of the arrows in the free body diagram showing the relative magnitudes of the forces. The free body diagram for a book sitting stationary on a table is:

FBD HERE

3.6 Friction

What happens if you slide a book across a table? It will slow down and eventually stop. From Newton's 1st Law of Motion, the change in velocity means there must be a net force. This net force is supplied by **friction**, which is the resistance to 2 surfaces sliding against each other. There are 2 types of friction:

1. **Kinetic friction** - acts when surfaces are sliding against each other. Slows down the moving object(s).
2. **Static friction** - acts when the surfaces are stationary relative to each other. Acts to prevent the surfaces from moving.

Different surfaces will have different friction forces. We know that a hockey puck will slide further on ice than on carpet before stopping, so there is a characteristic of the surfaces that determines the force. The surface interactions are described by the **coefficient of kinetic friction** and **coefficient of static friction**. In equations we will denote these coefficients with the Greek letter μ , pronounced “mu”. Subscripts are used to differentiate kinetic friction μ_k from static friction μ_s . The force of kinetic friction is found by

$$F_k = \mu_k F_N \quad (3.16)$$

which is the normal force multiplied by the coefficient of kinetic friction. This means that if the coefficient of kinetic friction or the normal force increases, the force of kinetic friction increases.

Example 3.3

A box of mass 30.0 kg is sliding along a horizontal floor. The coefficient of kinetic friction between the floor and box is $\mu_k = 0.40$. What is the magnitude of kinetic friction acting on the box?

The normal force here can be found by looking at the vertical axis. The box is not moving up or down, so the gravitational force and normal force are equal in magnitude.

$$F_N = F_g = mg = (30.0 \text{ kg})(9.8 \frac{m}{s^2}) = 294 \text{ N} \quad (3.17)$$

The force of kinetic friction is then

$$F_k = \mu_k F_N = 0.40 \cdot (294 \text{ N}) = 118 \text{ N} \quad (3.18)$$

The force of kinetic friction on the box is $F_k = 118 \text{ N}$

Notice how the coefficient of kinetic friction has no units. That is because it is multiplied by a force to find a force. It is a dimensionless (unitless) factor describing how the 2 surfaces interact when sliding along each other.

The force exerted by static friction is found by

$$F_s \leq \mu_s F_N \quad (3.19)$$

This is very similar to the force of kinetic friction, except that the coefficient is switched out for the coefficient of static friction. However, there is one major difference: *the equal sign is replaced with a less than or equal sign*. That is

because static friction prevents an object from starting to slide, and will not cause one surface to slide past another. In terms of Newton's Laws of Motion, static friction opposes an applied force to make the net force zero. If the applied force is larger than $\mu_s F_N$, then the static friction force is cannot make the net force zero and the object will begin sliding.

Example 3.4

A paperweight of mass 0.75 kg is sitting on a desk. The coefficient of static friction between the desk and paperweight is $\mu_s = 0.54$. There is a force of 3.0 N pushing horizontally on the paperweight, trying to slide it on the desk. Does the paperweight slide on the desk? If not, what is the magnitude of the static friction force acting on the paperweight?

First we need to find out the maximum possible static friction force. First we find the normal force, which is equal to the gravitational force.

$$F_N = F_g = (0.75\text{ kg})(9.8\frac{\text{m}}{\text{s}^2}) = 7.35\text{ N} \quad (3.20)$$

Now we can use Equation 3.19 to find the static friction force.

$$F_s \leq \mu_s F_N = 0.54 \cdot 7.35\text{ N} = 3.97\text{ N} \quad (3.21)$$

The static friction force is less than or equal to 3.97 N . This means that 3.97 N is the maximum possible static friction force. Since the applied force of 3.0 N is less than that, the static friction force can match the applied force and keep the block still.

Now to find the static friction force, we need to find the force required so that the net force is zero.

$$F_{net} = F_{app} + F_s = 0 \quad (3.22)$$

If we say the applied force (F_{app}) is in the positive direction, we can solve the problem as such

$$F_s = -F_{app} = -3.0\text{ N} \quad (3.23)$$

The force of static friction is $F_s = 3.0\text{ N}$ in the opposite direction of the applied force. Note that F_s is not always equal to $\mu_s F_N$!

Chapter 4

Work, Energy, and Power

Newton's Laws of Motion are a powerful tool for describing how forces cause motion, but some problems cannot be easily solved using them. In this chapter we will build off of Newton's Laws and Kinematics to develop new techniques that can be applied more broadly.

4.1 Work and Kinetic Energy

We will start with the 4th kinematic equation

$$v_f^2 = v_0^2 + 2a\Delta x \quad (4.1)$$

The right-hand side has acceleration, so we will multiply both sides by mass so that we can use Newton's 2nd Law to work with force.

$$mv_f^2 = mv_0^2 + 2ma\Delta x \quad (4.2)$$

Now we will subtract mv_0^2 from both sides and divide both sides by 2.

$$ma\Delta x = \frac{1}{2}mv_f^2 - \frac{1}{2}mv_0^2 \quad (4.3)$$

Use Newton's 2nd Law of Motion to replace ma with the force F_{net} to get

$$F_{net}\Delta x = \frac{1}{2}mv_f^2 - \frac{1}{2}mv_0^2 \quad (4.4)$$

The left hand side $F_{net}\Delta x$ is a quantity called **work**. Work, abbreviated W , is defined as a force acting over a displacement. We can find the units of work by multiplying the units of force and displacement. This gives us units of *Newtons · meters*, or $N \cdot m$. This unit is called the **Joule** and is abbreviated

J . When we treat every force acting on the object as a source of work, we call that the external work, which is abbreviated W_{ext} . This is an important distinction that will be important later this chapter.

The right hand side $\frac{1}{2}mv^2$ is a quantity called **kinetic energy**. Kinetic energy, abbreviated KE , is the energy in the motion of an object (energy is a hard quantity to define overall, so we will return to it a little later). The unit of kinetic energy is also the Joule.

This means we can write equation 4.4 as

$$W_{ext} = KE_f - KE_i = \Delta KE \quad (4.5)$$

This is called the **work-kinetic energy theorem**, which states that the sum of all work done on an object is equal to the change in its kinetic energy. We got the initial definitions using kinematic equations, but this result applies to any forces that act over an object while it moves. This means that if we know the change in kinetic energy over some displacement, we can find the **average force** acting on that object while it was moving. The force may vary throughout, but we can find the average using work.

Example 4.1

In the previously mentioned potato cannon, the potato starts at rest and is accelerated through a tube by the compressed.

1. If a potato of mass 0.65 kg leaves the tube at a velocity of $32 \frac{m}{s}$, how much work is done on the potato?
2. If the tube is 1.6 m long, what is the force exerted on the potato?

1. We are given $m = 0.65 \text{ kg}$, $v_0 = 0 \frac{m}{s}$, and $v_f = 32 \frac{m}{s}$. We can use the work-kinetic energy theorem to find the work done on the potato.

$$W_{ext} = \frac{1}{2}(0.65 \text{ kg})(32 \frac{m}{s})^2 - \frac{1}{2}(0.65 \text{ kg})(0 \frac{m}{s})^2 = 312 \text{ J} \quad (4.6)$$

The cannon does $W = 312 \text{ J}$ of work on the potato.

2. Knowing the length of the tube tells us that $\Delta x = 1.6 \text{ m}$. We can use the definition of work $W = F\Delta x$ to find the average force.

$$W = F \Delta x \quad (4.7)$$

$$F = \frac{W}{\Delta x} = \frac{312 \text{ J}}{1.6 \text{ m}} = 195 \text{ N} \quad (4.8)$$

The average force acting on the potato in the tube is $F = 195 \text{ N}$.

4.2 Direction of Force

In the previous section we only talked about accelerations and forces in 1 dimension, where they are aligned with the displacement. But what about forces that aren't along the same axis as the displacement? In Section 3.4 we saw that forces can be split into perpendicular components, where the two components are completely independent. This result tells us that **work is only done by components of the force in the same axis of the motion**. Forces (or the components of forces) perpendicular to the motion do no work!

Suppose we have a force F and displacement Δx that are separated by an angle θ as shown in Figure 4.1. The component of force along the displacement is the adjacent side of the triangle if we split the force into components (parallel and perpendicular to Δx). This means that the work done by a single force on an object is

$$W = F \Delta x \cos(\theta) \quad (4.9)$$

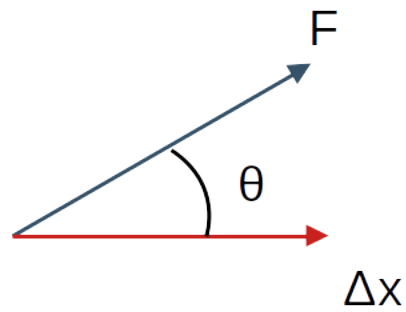


Figure 4.1: Here is a force and displacement separated by an angle of θ . Only the component of force along the displacement does work.

Example 4.2

Yvonne pulls a box along the floor by exerting a force of 70.0 N at an angle of 50.0° above horizontal. If she slides the box 6.50 m , how much work does Yvonne do on the box?

We are given that $F = 70.0\text{ N}$, $\theta = 50.0^\circ$, and $\Delta x = 6.5\text{ m}$. The work done is

$$W = (70.0\text{ N})(6.5\text{ m})\cos(50.0^\circ) = 292\text{ J} \quad (4.10)$$

She does $W = 292\text{ J}$ of work on the box.

4.3 Gravitational Potential Energy

We are going to give special consideration to the work done by gravity on an object. Let's define up as positive, meaning that the force of gravity acting on an object of mass m is

$$F_g = -mg \quad (4.11)$$

This means that the work done by gravity as an object changes height by Δy is

$$W_g = -mg\Delta y \quad (4.12)$$

Example 4.3

Terry throws a rock straight up in the air so it reaches its maximum height at 3.5 m above where they threw it. Use the work-kinetic energy theorem to find the velocity Terry threw the rock at.

We know that $\Delta y = 3.5\text{ m}$ and the final velocity is $v_f = 0\frac{\text{m}}{\text{s}}$. Let's start with the work-kinetic energy theorem.

We can check that this holds true even for objects that don't move just vertically. For example, we will look at a block sliding down a ramp in Figure 4.2. The block slides down the ramp, changing its position by d . The height

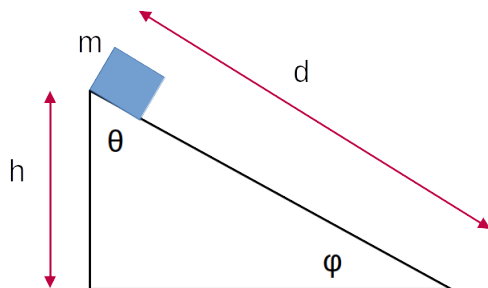


Figure 4.2: A block slides down a ramp of length d , changing its vertical position by $-h$. The angle at the top of the ramp is θ .

of the ramp is h and the angle at the top of the ramp is θ . Defining up as positive, we can write the work done by gravity as

$$W_g = -mg \cdot (-d \cos(\theta)) = mg \cos(\theta) \quad (4.13)$$

The height of the ramp is the adjacent side of the triangle from the angle θ .

$$h = d \cos(\theta) \quad (4.14)$$

This means we can write the work done by gravity as

$$W_g = mgh \quad (4.15)$$

Which means that the work only depends on the vertical movement. Note that the work is positive because the block moves downward, which is in the direction of the force of gravity.

Since gravity from the Earth plays a role in all motion at the Earth's surface, we account for this by describing the work gravity can do with **gravitational potential energy**. An object of mass m at a height of h has a gravitational potential energy of

$$PE_g = mgh \quad (4.16)$$

It is often easier to think about gravitational potential energy rather than the work done by energy, which we will see shortly. When we introduce potential energy, doing work on an object can change both kinetic energy and potential energy.

$$W = \Delta KE + \Delta PE \quad (4.17)$$

Example 4.4

Kevin sees a book on the floor and picks it up to put on a shelf. The book has a mass of 0.30 kg and the shelf is 1.9 m above the floor. How much work does Kevin do on the book?

Let's define the floor as height 0 m , so the shelf is at a height of 1.9 m . The book is at rest in both the final and initial state, so $v_i = v_f = 0\frac{\text{m}}{\text{s}}$. This means that $\Delta KE = 0\text{ J}$, so we have

$$W = \Delta PE = PE_f - PE_i \quad (4.18)$$

$$W = mgh_f - mgh_i \quad (4.19)$$

The book starts on the floor, which we defined as having a height of zero.

$$W = (0.30\text{ kg})(9.8\frac{\text{m}}{\text{s}^2})(1.9\text{ m}) = 5.6\text{ J} \quad (4.20)$$

Kevin does $W = 5.6\text{ J}$ of work on the book.

4.4 Elastic Forces

I am going to guess that everybody reading this has stretched a rubber band before. What did you notice about the band while you were pulling on it? Observations may include:

- The rubber band pulls back against your hand.
- You have to pull harder the further you stretch the rubber band.
- The rubber band will return to its original size and shape.

These are all characteristics of an **elastic force**. Elastic forces are exerted by objects that are stretched or compressed and try to return to their original shape. This includes rubber bands, tennis balls (they compress when they bounce), and the classic physics example of a spring. The force exerted by an ideal spring is stated by Hooke's Law (named after Robert Hooke, who studied springs)

$$F_e = -k\Delta x \quad (4.21)$$

where k is the spring constant and x is how far the spring is stretched or compressed from equilibrium. There is a lot going on here, so let's look at this in detail.

The **spring constant** is a measure of the stiffness of a spring. The units are $\frac{\text{force}}{\text{length}}$, and is typically in $\frac{N}{m}$ or $\frac{N}{cm}$ depending on the spring. A greater spring constant means the spring is stiffer and will have more resistance to being stretched or compressed. The spring constant of each spring is different and it depends on the shape and construction of the spring.

We know that Δx is a displacement, but we need to know what it is a displacement relative to. Hooke's Law needs the displacement away from the **equilibrium length** of the spring. The equilibrium length is the spring's natural resting length, where it does not exert a force. This means that the force is stronger the further from the equilibrium length the spring is stretched or compressed. That is why it felt like you had to pull harder on the rubber band the further you stretched it.

The last thing of note from Hooke's Law is the negative sign, which indicates the force is exerted in the opposite direction that the spring is displaced. The elastic force is a type of **restoring force** because it acts to restore the spring to its original shape.

Example 4.5

A spring with spring constant $60.0 \frac{N}{cm}$ is compressed by a distance of 1.50 cm to the left. What is the magnitude and direction of the force?

Define right as positive. Since our spring is compressed to the left, that makes the displacement negative. We are given values of $k = 60 \frac{N}{cm}$ and $\Delta x = -1.50 \text{ cm}$. We can use Hooke's Law now

$$F_e = -k\Delta x = -(60.0 \frac{N}{cm}) \cdot (-1.50 \text{ cm}) = 90.0 \text{ N} \quad (4.22)$$

The force has a magnitude of 90.0 N . We find that F_e is positive, so the force is directed to the right. That makes sense for the displacement being to the left, since the elastic force is a restoring force.

Example 4.6

A spring is hanging vertically from the ceiling. When a 12.0 kg mass is attached the spring, the spring stretches downward by 5.50 cm . What is the spring constant of the spring in $\frac{N}{m}$?

Define up as positive, so the force of gravity is negative and $\Delta x = -5.50\text{ cm}$. We will convert the displacement to meters because we are asked to give the spring constant in $\frac{N}{m}$.

$$\Delta x = -5.50\text{ cm} \cdot \frac{1\text{ m}}{100\text{ cm}} = -0.0550\text{ m} \quad (4.23)$$

This is a helpful problem to draw a free body diagram. We can see that the block is stationary only if the sum of the force of gravity and the elastic force is zero.

$$F_{net} = F_e + F_g = 0 \quad (4.24)$$

$$F_e = -k\Delta x \quad (4.25)$$

$$F_g = -mg \quad (4.26)$$

$$-k\Delta x - mg = 0 \quad (4.27)$$

$$k = \frac{-mg}{\Delta x} = \frac{-(12.0\text{ kg})(9.8\text{ m/s}^2)}{(-0.0550\text{ m})} = 2140\frac{N}{m} \quad (4.28)$$

The spring constant is $k = 2140\frac{N}{m}$.

While we did talk extensively about about springs, Hooke's Law is a good approximation for many elastic objects at small deformations.

4.5 Elastic Potential Energy

From the example with a block hanging from a vertical spring, let's look at the force that somebody would have to apply to a spring to displace it by a certain amount. Since the force exerted by a spring is $F_e = -k\Delta x$, you have to apply a force of $F_{app} = k\Delta x$ to displace it by Δx . This means the net force is zero at Δx and the spring will be stationary. Let's see this visually in Figure 4.3.

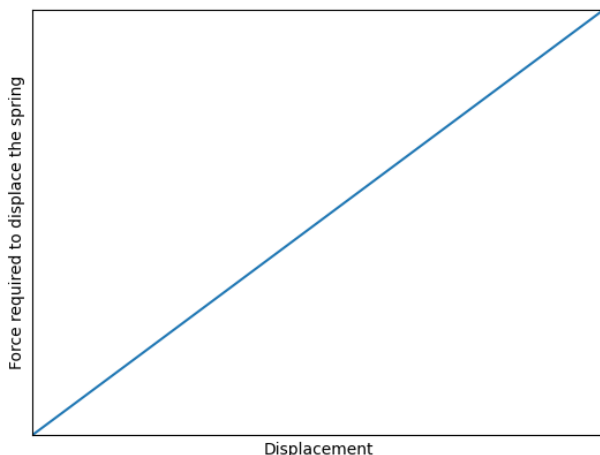


Figure 4.3: The force required to displace a spring increases linearly with displacement.

We know that for a constant force the work is $W = F\Delta x \cos(\theta)$, but here the force isn't constant. With a varying force we can find work by finding the area under the curve of a *Force vs. Displacement* graph. This is very similar to how we found the displacement as the area under the graph of a *v vs. t* plot. Looking at Figure 4.3, the area under the curve is a right triangle with base Δx and height $F = k\Delta x$. The area of a triangle is

$$A = \frac{b \cdot h}{2} \quad (4.29)$$

$$W = \frac{\Delta x \cdot k\Delta x}{2} = \frac{1}{2}k(\Delta x)^2 \quad (4.30)$$

This is the amount of work required to displace the spring by Δx . Similar to how lifting an object converts work into gravitational potential energy, we can think of this work being converted to **elastic potential energy**.

$$PE_e = \frac{1}{2}k(\Delta x)^2 \quad (4.31)$$

This is the potential energy stored in a spring of spring constant k displaced by Δx . We can't use kinematics to study the motion of an object on a spring (since acceleration varies), but work and energy makes it possible!