

# Dynamics of Compact Relativistic Bodies

Todd Oliynyk

School of Mathematical Sciences

Monash University

in collaboration with Lars Andersson and Bernd Schmidt

# Dynamics of fully relativistic, compact elastic bodies



## Dynamics of fully relativistic, compact elastic bodies



- ◊ Bodies  $B_1$  and  $B_2$  are composed of **elastic matter**
  - asteroids, moons, comets, planets, stars, neutron stars

## Dynamics of fully relativistic, compact elastic bodies



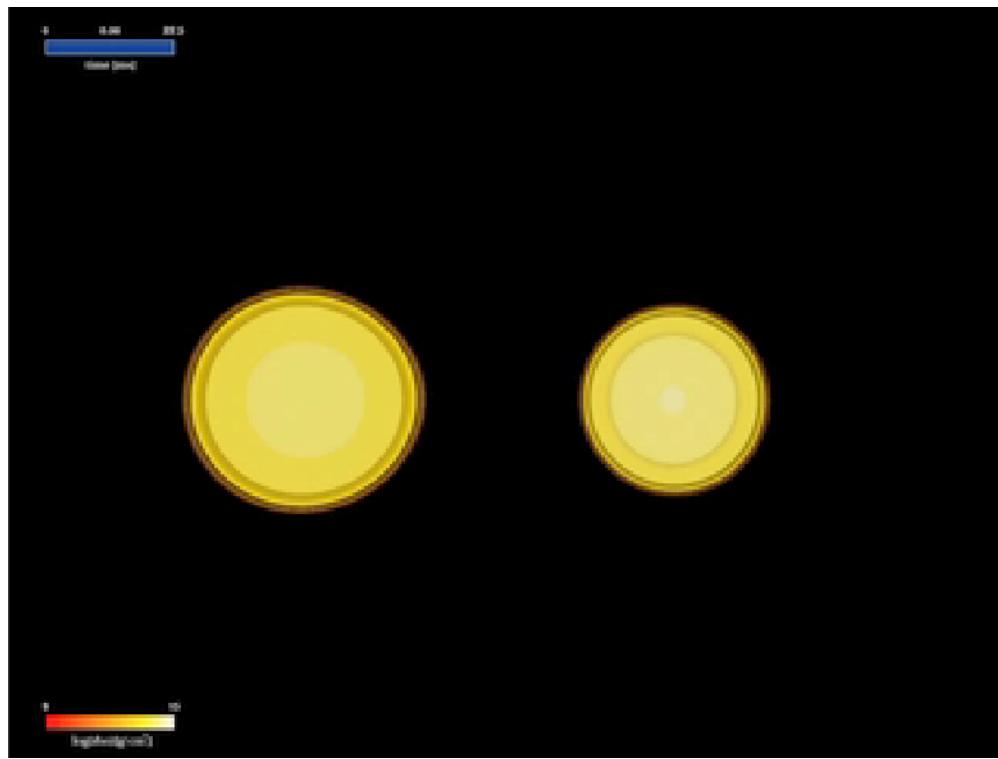
- ◊ Bodies  $B_1$  and  $B_2$  are composed of **elastic matter**
  - asteroids, moons, comets, planets, stars, neutron stars
- ◊ Bodies are **fully relativistic**.

## Dynamics of fully relativistic, compact elastic bodies



- ◊ Bodies  $B_1$  and  $B_2$  are composed of **elastic matter**
  - asteroids, moons, comets, planets, stars, neutron stars
- ◊ Bodies are **fully relativistic**.

Basic Problem : Prove the existence of appropriate solutions that represent the evolution of fully relativistic, compact elastic bodies.

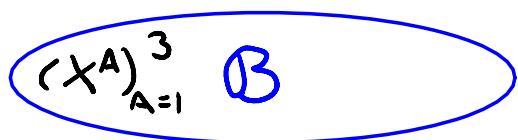


Ref: Rezzolla et. al., Class. Quantum Grav. 27 (2010) 114105

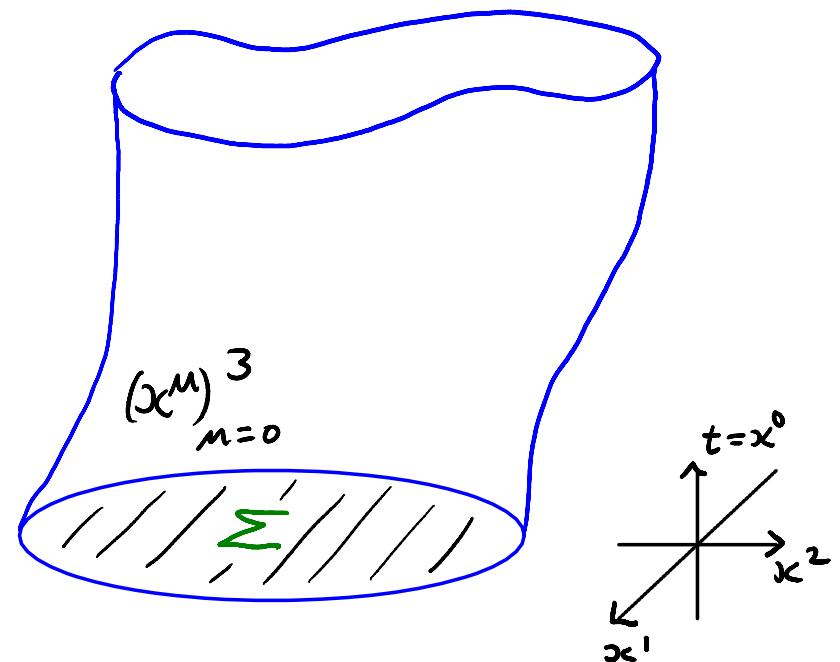
# Relativistic Elasticity

# Relativistic Elasticity

## 3-d Material Manifold



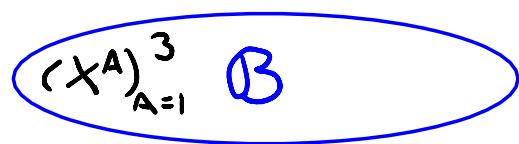
## Spacetime evolution



# Relativistic Elasticity

## 3-d Material Manifold

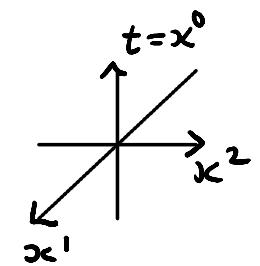
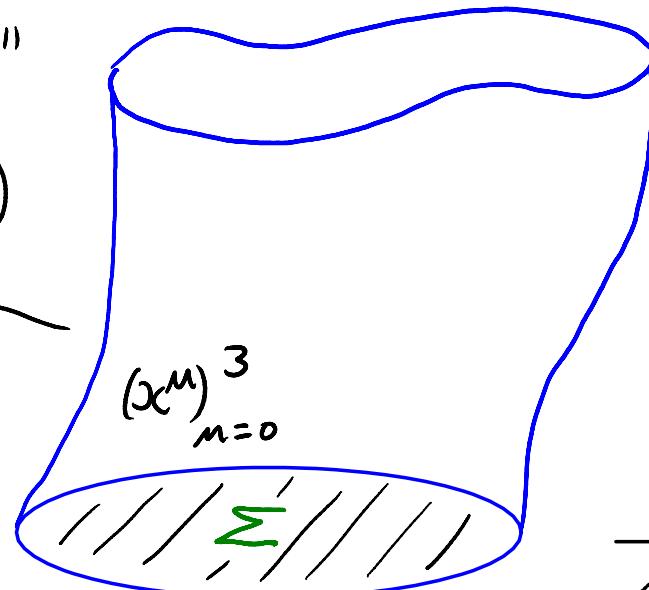
## Spacetime evolution



"Back to Lakes map"

$$f = (f^A(x^m))$$

A curved arrow points from the 3D manifold towards the spacetime diagram.



# Relativistic Elasticity

## 3-d Material Manifold

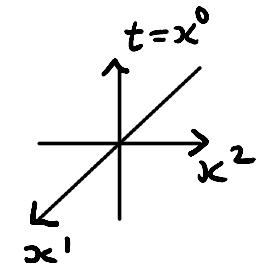
$$(x^A)_{A=1}^3 \quad B$$

Volume form

$$\Omega_{ABC}$$

"Back to Lakes map"  
 $f = (f^A(x^m))$

$$(x^m)_{m=0}^3$$



## Spacetime evolution

Spacetime metric

$$g = g_{\mu\nu} dx^\mu dx^\nu$$

# Relativistic Elasticity

## 3-d Material Manifold

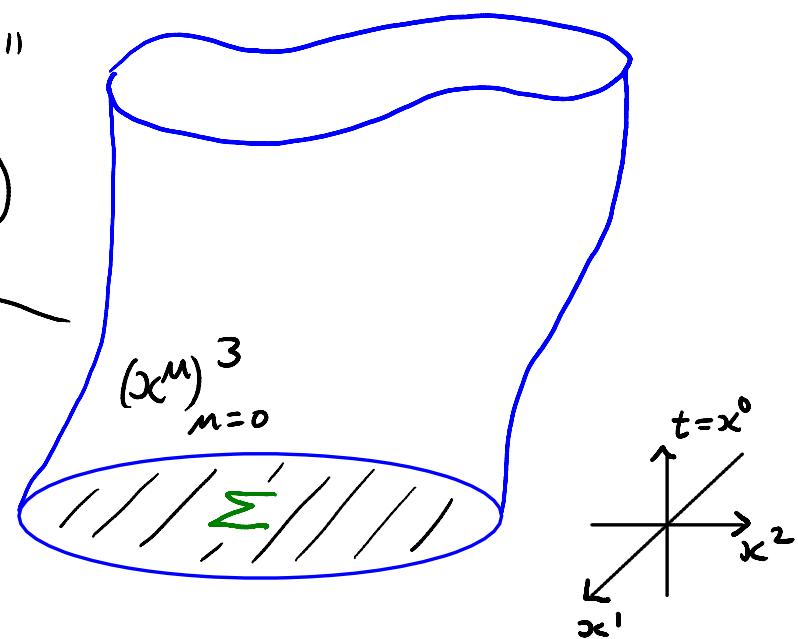
$$(x^A)_{A=1}^3 \quad B \cdot P$$

Volume form

$$\Omega_{ABC}$$

"Back to Lakes map"

$$f = (f^A(x^m))$$



Spacetime metric

$$g = g_{\mu\nu} dx^\mu dx^\nu$$

# Relativistic Elasticity

## 3-d Material Manifold

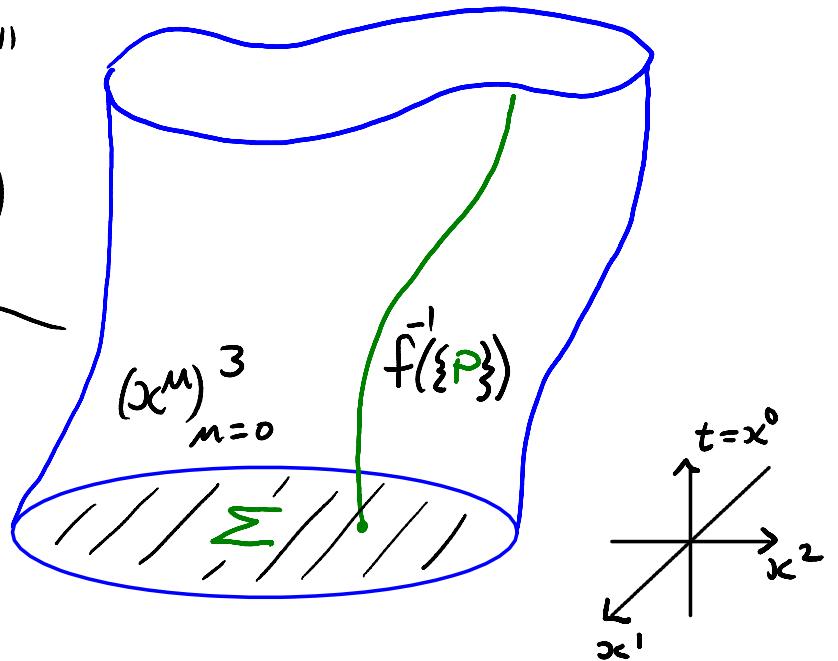
$$(x^A)_{A=1}^3 \quad \mathcal{B} \cdot \rho$$

Volume form

$$\Omega_{ABC}$$

"Back to Lakes map"

$$f = (f^A(x^m))$$



Spacetime metric

$$g = g_{\mu\nu} dx^\mu dx^\nu$$

## Relativistic Elasticity

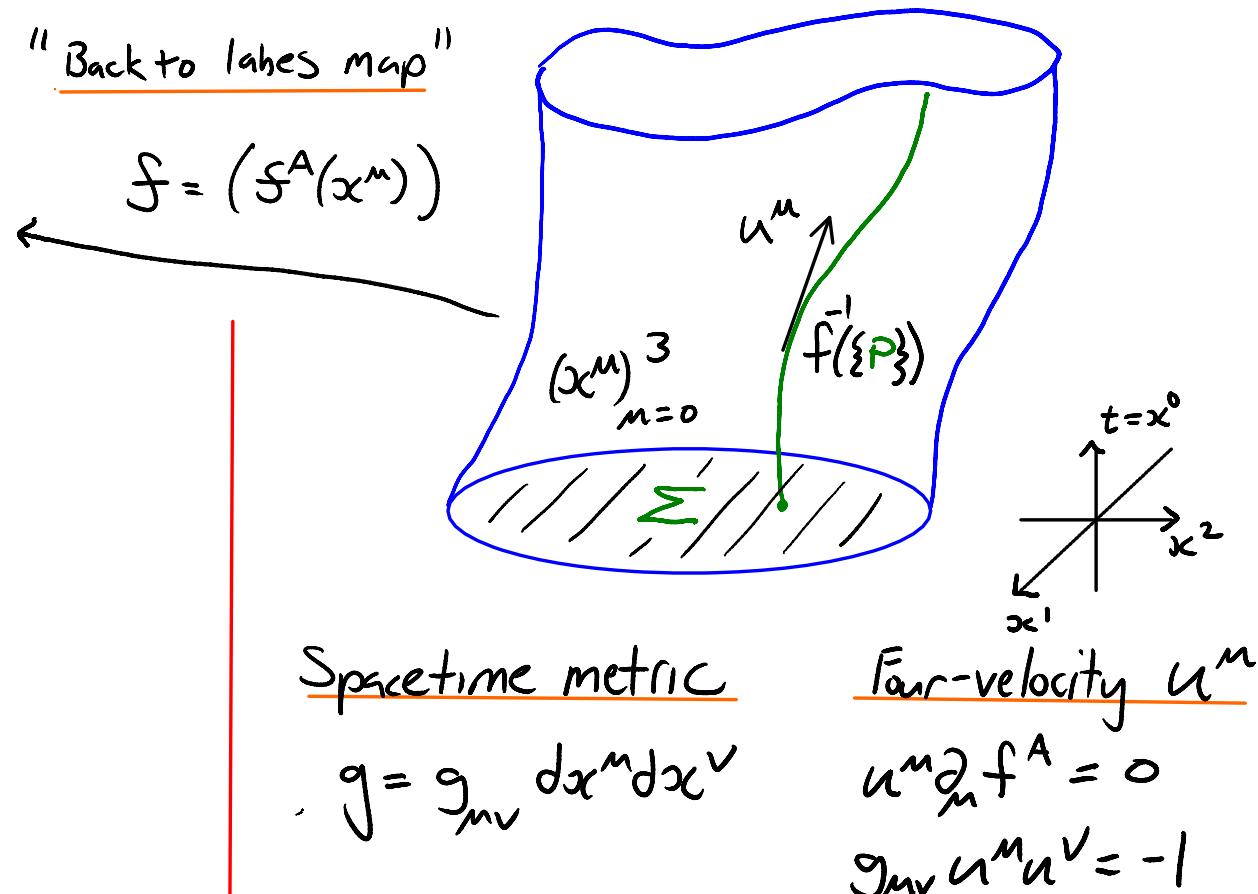
## 3-d Material Manifold

$$(x^4)_{A=1}^3 B \cdot P$$

## Volume form

$$\Omega_{ABC}$$

## Spacetime evolution



# Relativistic Elasticity

## 3-d Material Manifold

$$(x^A)_{A=1}^3 \quad B \cdot P$$

Volume form

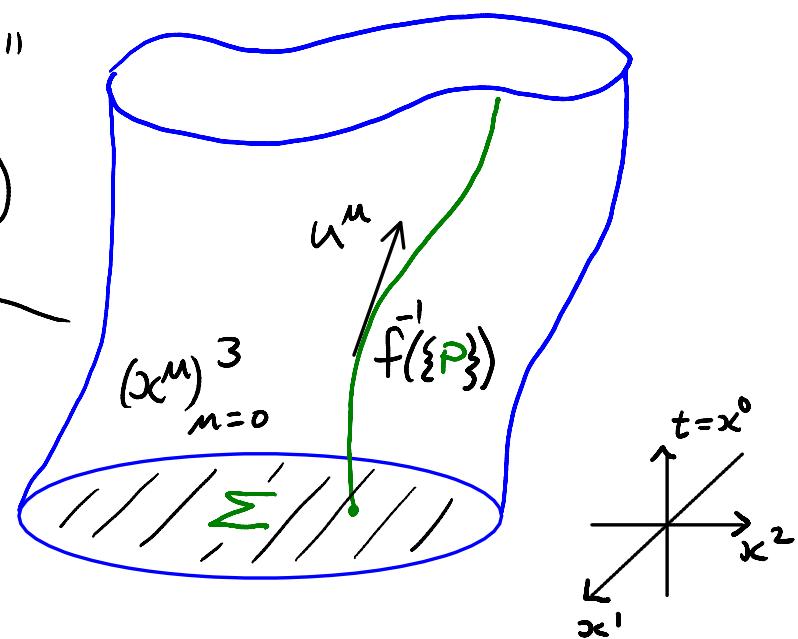
$$\Omega_{ABC}$$

Number density n

$$n^m = \frac{1}{3!n} \epsilon^{m\nu\lambda\rho} \omega_{\nu\lambda\rho}$$

where  $\omega = f^* \omega_0$   $\ddagger$   $\epsilon = \text{vol}(g)$

"Back to Lakes map"  
 $f = (f^A(x^\mu))$



Spacetime metric

$$g = g_{\mu\nu} dx^\mu dx^\nu$$

Four-velocity u^\mu

$$u^\mu \partial_\mu f^A = 0$$

$$g_{\mu\nu} u^\mu u^\nu = -1$$

# Relativistic Elasticity

## 3-d Material Manifold

$$(x^A)_{A=1}^3 \quad B \cdot P$$

Volume form

$$\Omega_{ABC}$$

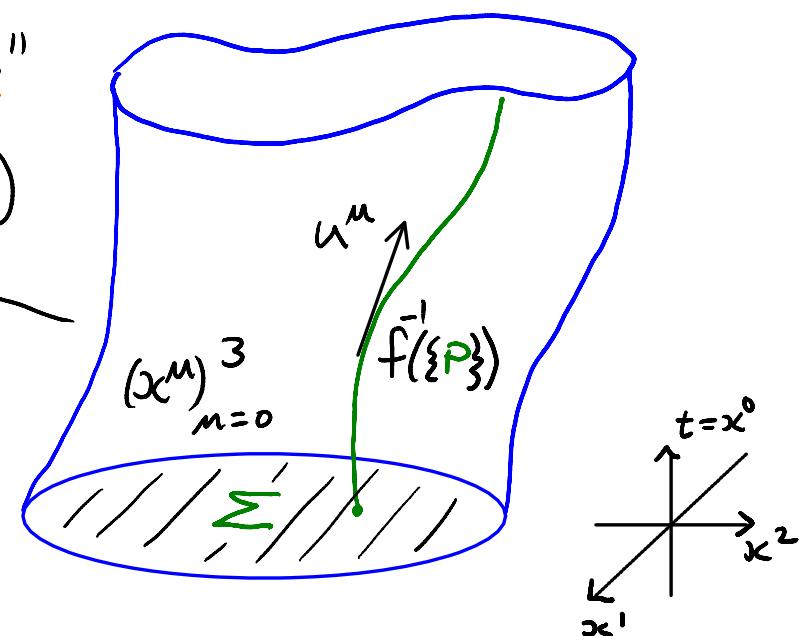
Number density n

$$n^m = \frac{1}{3!n} \epsilon^{m\nu\lambda\rho} \omega_{\nu\lambda\rho}$$

where  $\omega = f^* \omega_0$   $\ddagger$   $\epsilon = \text{vol}(g)$

$$\nabla_m (n u^m) = 0$$

"Back to Lakes map"  
 $f = (f^A(x^m))$



Spacetime metric

$$g = g_{\mu\nu} dx^\mu dx^\nu$$

Four-velocity  $u^m$

$$u^m \partial_m f^A = 0$$

$$g_{\mu\nu} u^\mu u^\nu = -1$$

## Spacetime evolution

# Field Equations

# Field Equations

Lagrangian

$$S[g, \dot{g}] = \int_M P(f, \partial f, g) \sqrt{-|g|} d^4x$$

# Field Equations

Lagrangian

$$S[g, \dot{g}] = \int_M P(f, \partial f, g) \sqrt{-|g|} d^4x$$

Euler-Lagrange equations

$$\frac{\delta S}{\delta f} = 0$$

# Field Equations

## Lagrangian

$$S[g, f] = \int_M P(f, \partial f, g) \sqrt{-g} \, d^4x$$

## Euler-Lagrange equations

$$\frac{\delta S}{\delta f} = 0 \Rightarrow$$

$$-\sum_A := \frac{1}{\sqrt{-g}} \partial_\mu \left( \sqrt{-g} \frac{\partial P}{\partial (\partial_\mu f^A)} \right) - \frac{\partial P}{\partial f^A} = 0$$

## Field Equations

### Lagrangian

$$S[g, f] = \int_M \rho(f, \partial f, g) \sqrt{-|g|} d^4x$$

### Euler-Lagrange equations

$$\frac{\delta S}{\delta f} = 0$$

$$\Rightarrow -\sum_A := \frac{1}{\sqrt{-|g|}} \partial_\mu \left( \sqrt{-|g|} \frac{\partial \rho}{\partial (\partial_\mu f^A)} \right) - \frac{\partial \rho}{\partial f^A} = 0$$

### Diffeomorphism Covariance

$$\rho = \rho(f, H)$$

where

$$H^{AB} = g^{\mu\nu} \partial_\mu f^A \partial_\nu f^B$$

# Stress Energy Tensor

## Stress Energy Tensor

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \left( \rho \sqrt{-g} \right) \delta g^{\mu\nu}$$

## Stress Energy Tensor

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \frac{(\rho \sqrt{-g})}{\delta g^{\mu\nu}} \Rightarrow \left\{ \begin{array}{l} T_{\mu\nu} = 2 \frac{\partial \rho}{\partial g^{\mu\nu}} - \rho g_{\mu\nu} \\ \nabla_\mu T^\mu_\nu = \sum_A \partial_\nu f^A \end{array} \right.$$

## Stress Energy Tensor

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \left( \rho \sqrt{-g} \right) \Rightarrow \left\{ \begin{array}{l} T_{\mu\nu} = 2 \frac{\partial \rho}{\partial g^{\mu\nu}} - \rho g_{\mu\nu} \\ \nabla_\mu T^\mu_\nu = \sum_A \partial_\nu f^A = 0 \Leftrightarrow \sum_A = 0 \end{array} \right.$$

## Stress Energy Tensor

$$T_{\mu\nu} = \frac{2}{\sqrt{-g}} \left( \rho \sqrt{-g} \right) \delta g^{\mu\nu} \Rightarrow \begin{cases} T_{\mu\nu} = 2 \frac{\partial \rho}{\partial g^{\mu\nu}} - \rho g_{\mu\nu} \\ \nabla_m T^m_\nu = \sum_A \partial_\nu f^A = 0 \Leftrightarrow \sum_A = 0 \end{cases}$$

Alternate form

$$\bar{T}_{\mu\nu} = \rho u_\mu u_\nu + t_{\mu\nu}$$

where

$$t_{\mu\nu} = \eta \tilde{\gamma}_{AB} \partial_\mu f^A \partial_\nu f^B$$

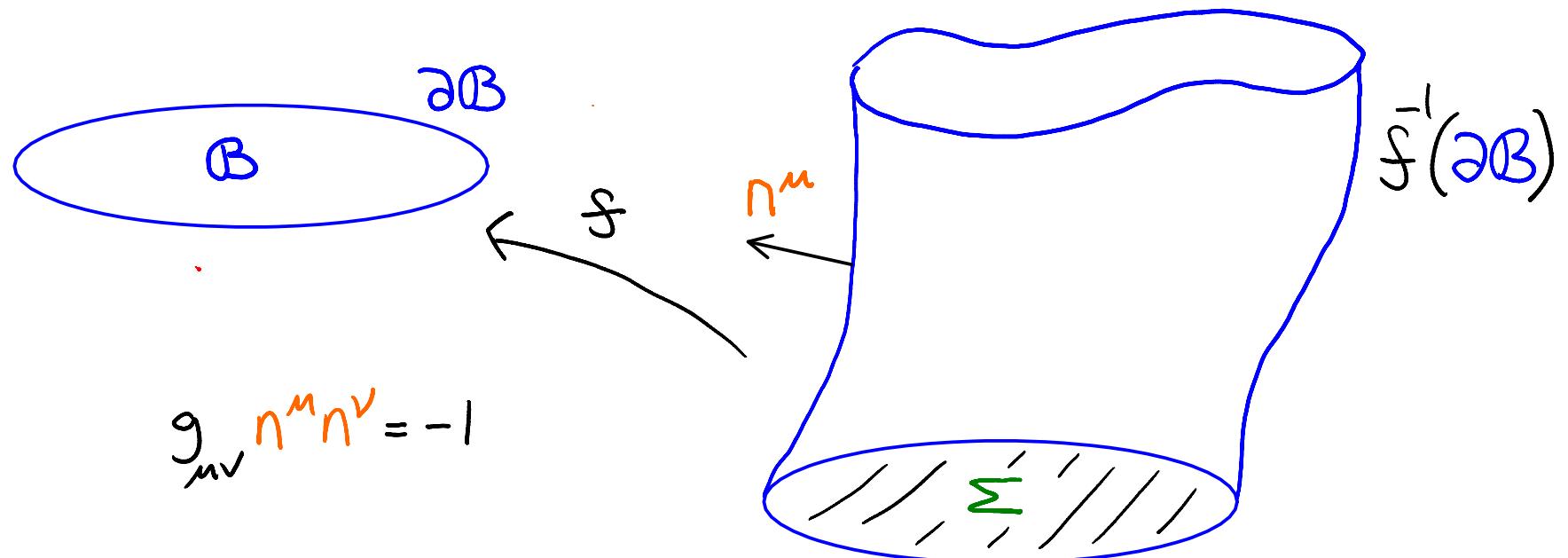
with

$$\epsilon := \frac{P}{n} \quad (\text{stored energy function})$$

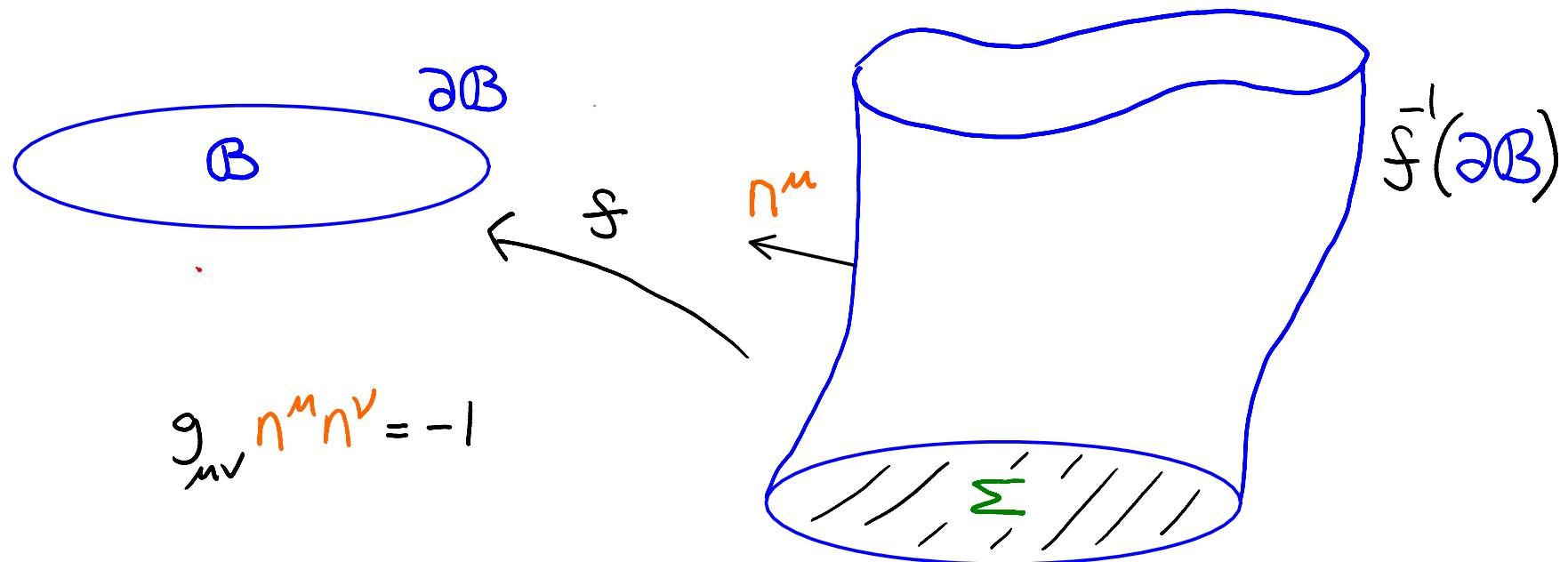
and.

$$\tilde{\gamma}_{AB} := 2 \frac{\partial \epsilon}{\partial H^{AB}} \quad (\text{2}^{\text{nd}} \text{ Piola-Kirchhoff stress tensor})$$

## Boundary Conditions

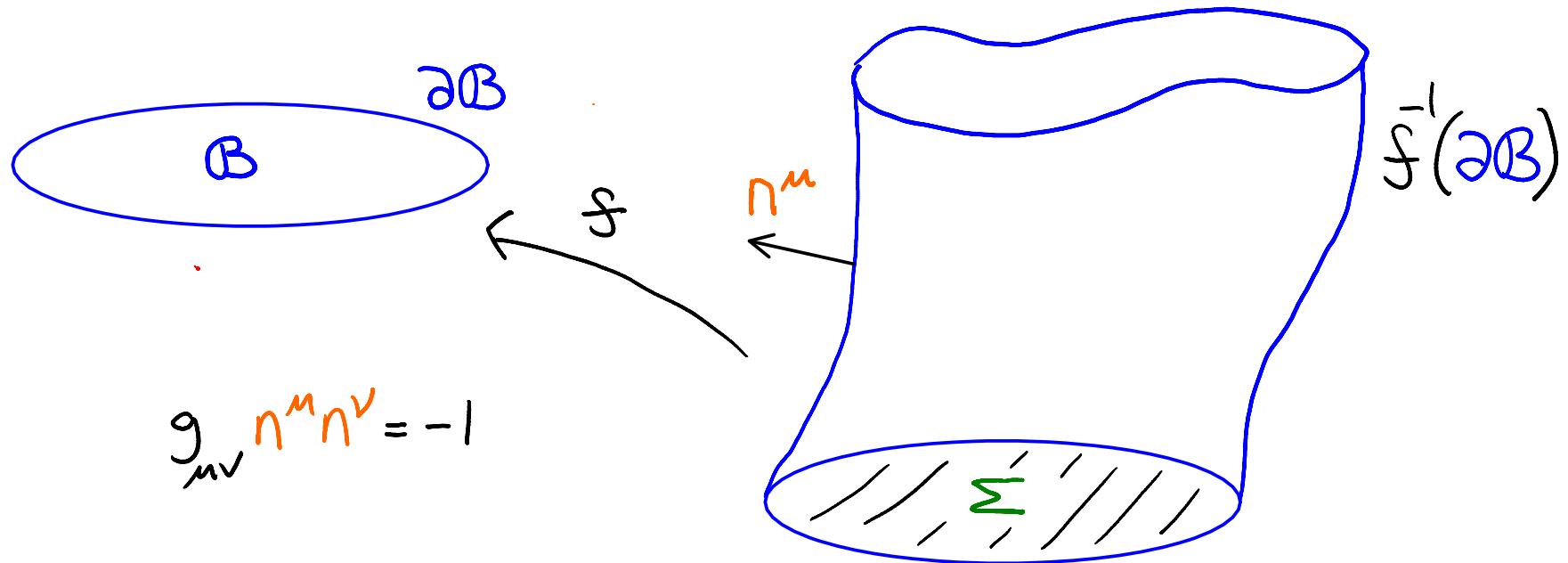


## Boundary Conditions



$$T_\mu^\nu n_\nu \Big|_{f^{-1}(\partial B)} = 0$$

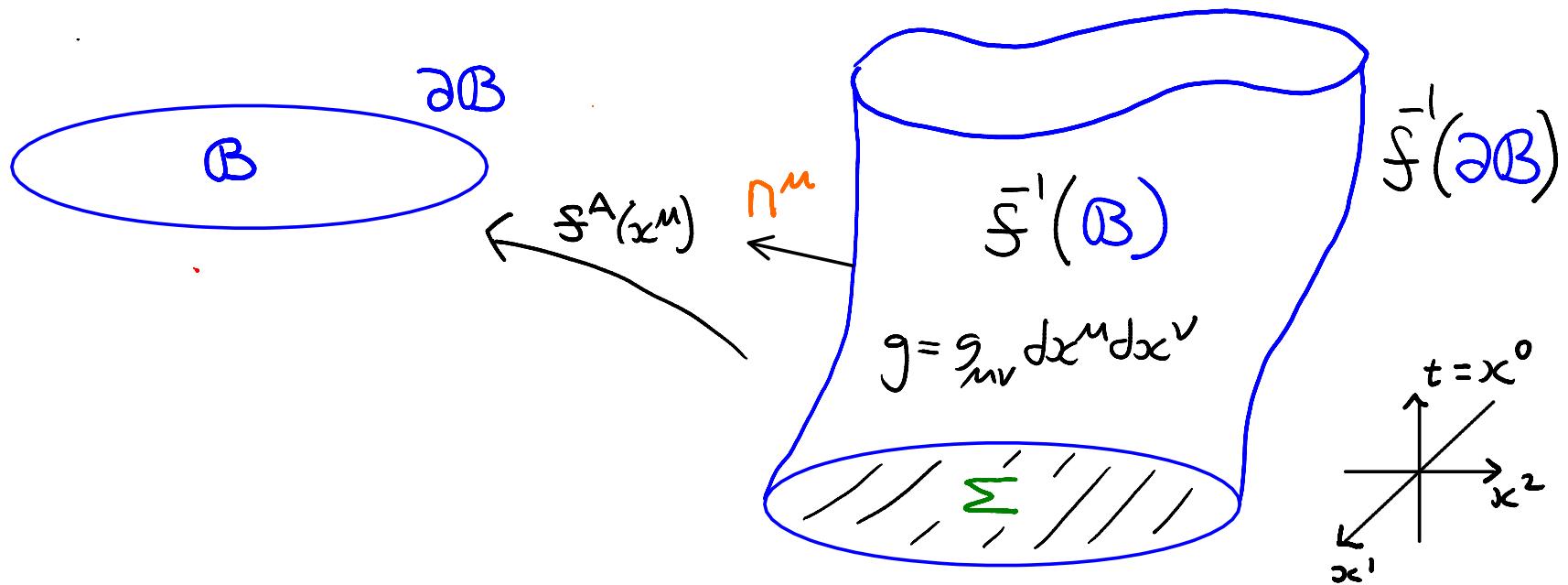
## Boundary Conditions



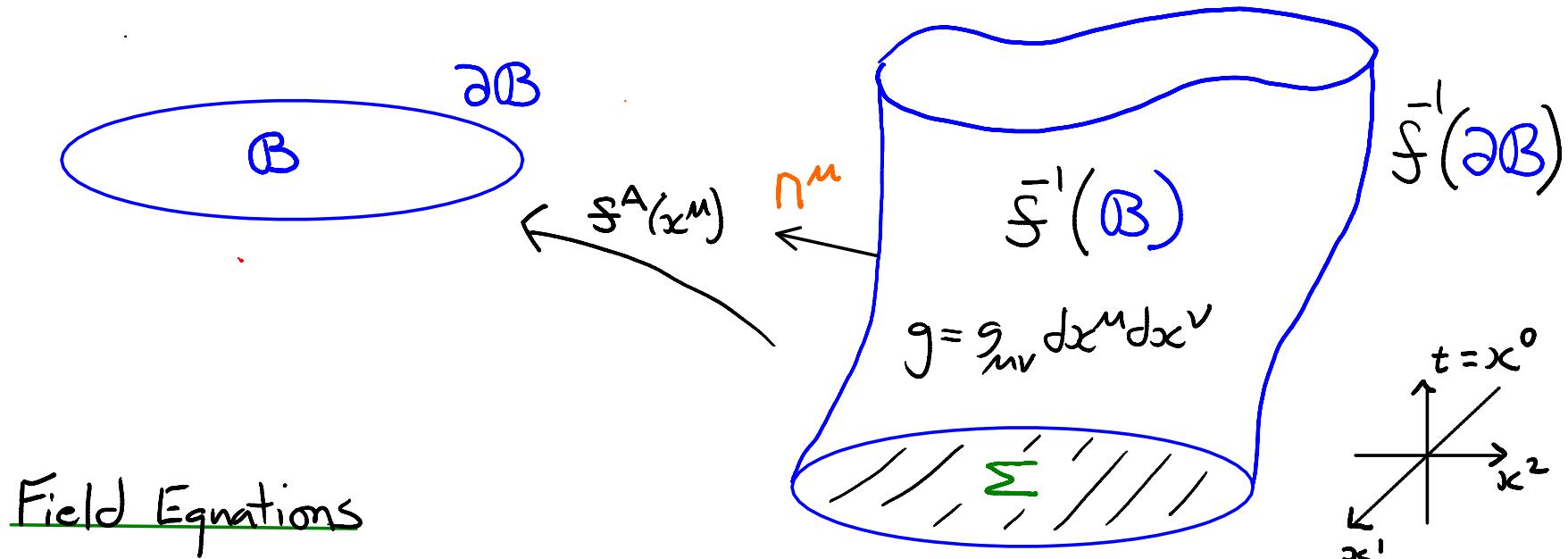
$$g_{\mu\nu} n^\mu n^\nu = -1$$

$$T_\mu^\nu n_\nu \Big|_{f^{-1}(\partial B)} = 0 \iff \sum_{AB} n^\mu \partial_\mu f^B \Big|_{f^{-1}(\partial B)} = 0$$

# Coupling Elasticity to the Einstein Field Equations



# Coupling Elasticity to the Einstein Field Equations



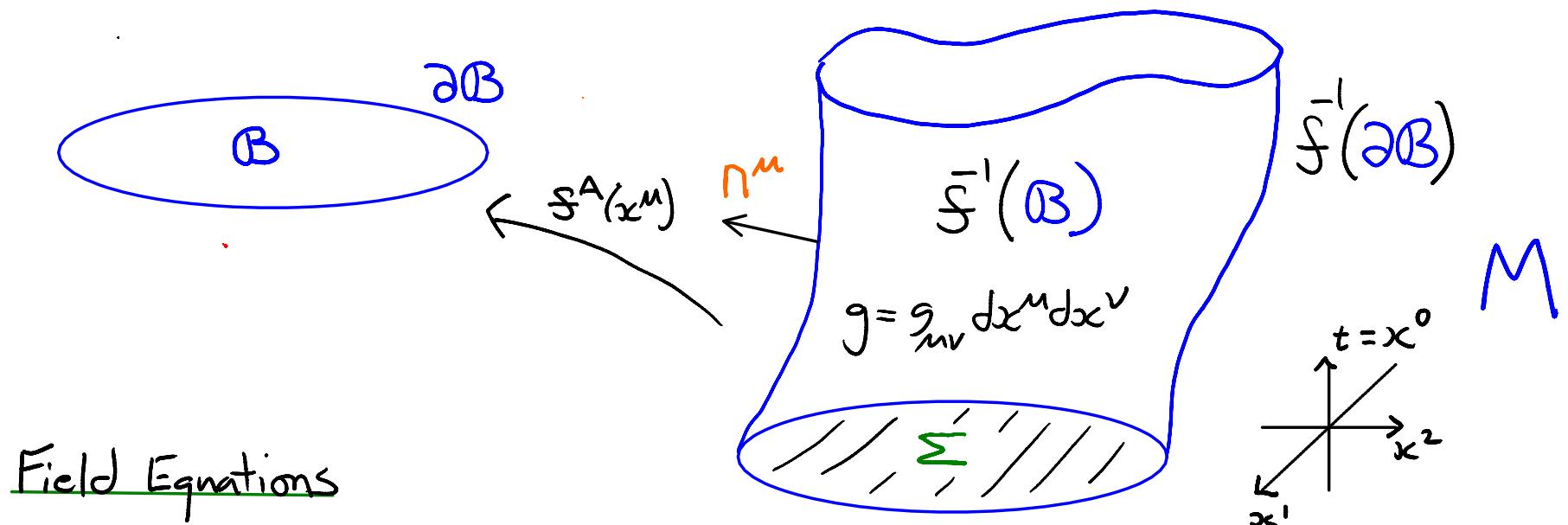
Field Equations

$$R_{\mu\nu} = T_{\mu\nu} - \frac{1}{2} T \delta^{\lambda}_{\mu\nu} g_{\lambda\lambda} \quad \text{in } f^{-1}(B)$$

$$\nabla_\mu T^{\mu\nu} = 0 \quad \text{in } f^{-1}(B)$$

$$\eta^\mu T_{\mu\nu} \Big|_{f^{-1}(\partial B)} = 0$$

# Coupling Elasticity to the Einstein Field Equations



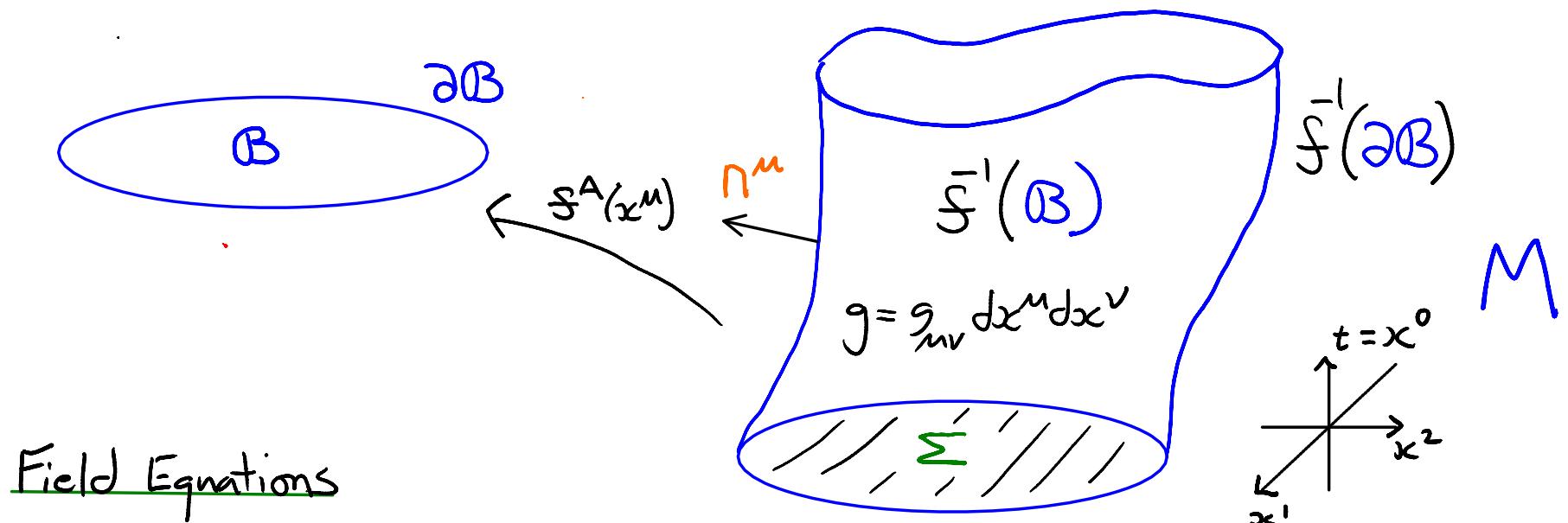
Field Equations

$$R_{\mu\nu} = T_{\mu\nu} - \frac{1}{2} T \delta^{\lambda}_{\mu\nu} g_{\lambda\lambda} \quad \text{in } \bar{f}^{-1}(B) \text{ or } M ?$$

$$\nabla_\mu T^{\mu\nu} = 0 \quad \text{in } \bar{f}^{-1}(B)$$

$$\nabla^\mu T_{\mu\nu} |_{\bar{f}^{-1}(\partial B)} = 0$$

# Coupling Elasticity to the Einstein Field Equations



Field Equations

$$R_{\mu\nu} = T_{\mu\nu} - \frac{1}{2} T \delta^{\lambda}_{\mu\nu} g_{\lambda\lambda} \quad \text{in } M$$

$$\nabla_\mu T^{\mu\nu} = 0 \quad \text{in } f^{-1}(B)$$

$$\nabla^\mu T_{\mu\nu} |_{f^{-1}(\partial B)} = 0$$

## Harmonic Coordinates

## Harmonic Coordinates

$$\diamond \quad -2R_{\mu\nu} + (\nabla_\mu X_\nu + \nabla_\nu X_\mu) = -2T_{\mu\nu} + T^\lambda_\lambda g_{\mu\nu}$$

where

$$X^\mu := g^{\alpha\beta} \nabla^\mu_{\alpha\beta} = \frac{1}{2} g^{\alpha\beta} g^{\mu\lambda} (\partial_\alpha g_{\beta\lambda} + \partial_\beta g_{\alpha\lambda} - \partial_\lambda g_{\alpha\beta})$$

## Harmonic Coordinates

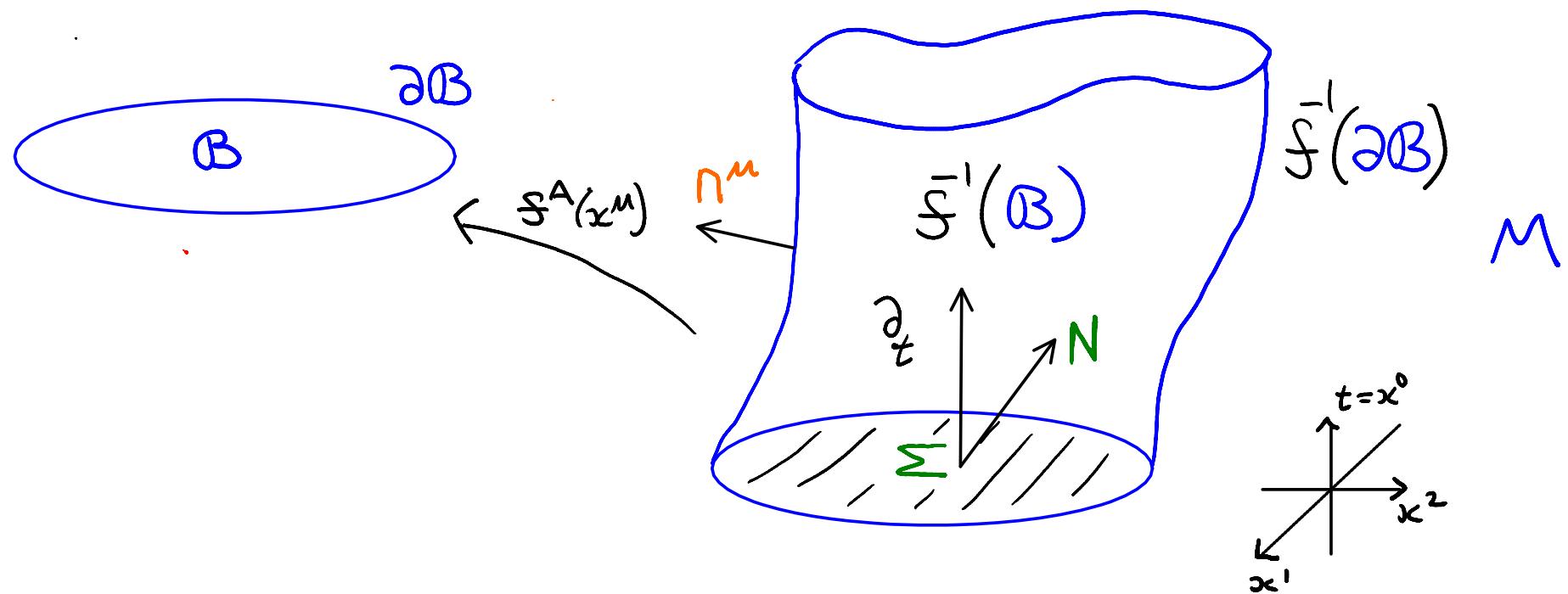
◇  $-2R_{\mu\nu} + (\nabla_\mu X_\nu + \nabla_\nu X_\mu) = -2T_{\mu\nu} + T^\lambda_\lambda g_{\mu\nu}$

where

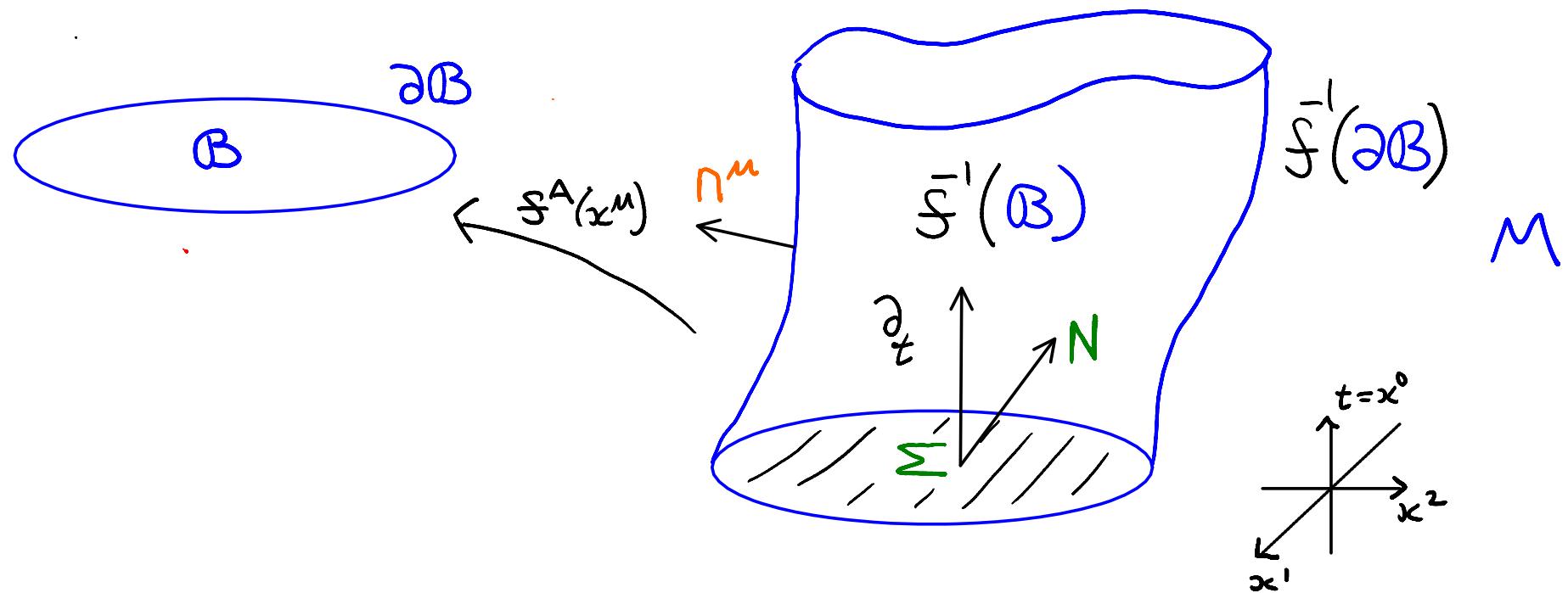
$$X^\mu := g^{\alpha\beta} \nabla^\mu_{\alpha\beta} = \frac{1}{2} g^{\alpha\beta} g^{\mu\lambda} (\partial_\alpha g_{\beta\lambda} + \partial_\beta g_{\alpha\lambda} - \partial_\lambda g_{\alpha\beta})$$

◇  $g^{\alpha\beta} \partial_\alpha \partial_\beta g_{\mu\nu} + Q_{\mu\nu}(g, \partial g) = P_{\mu\nu}(g, f, \partial f)$

## Initial Data Complications I: Gravitational Constraint Equations



# Initial Data Complications I: Gravitational Constraint Equations



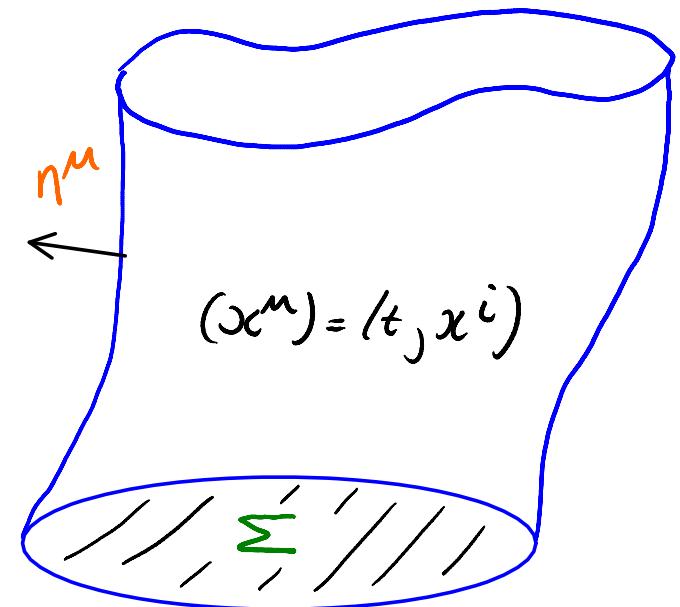
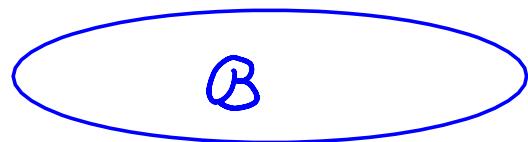
## Constraint Equations

$$(G^{μν} - T^{μν})N_{μ} \Big|_{Σ} = 0$$

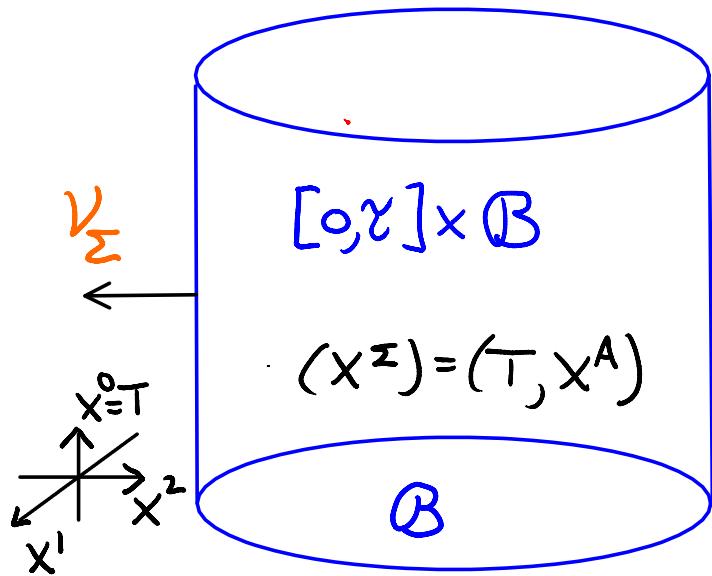
where

$$G_{μν} = R_{μν} - \frac{1}{2} R^λ_λ g_{μν}$$

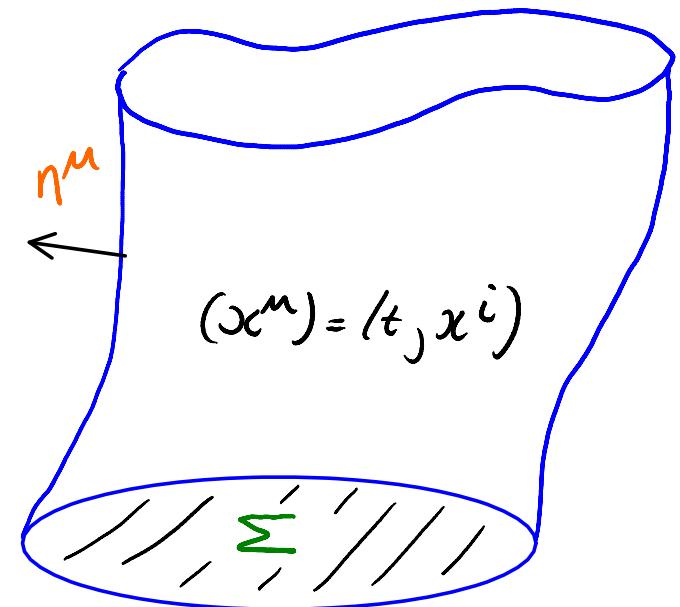
## Material Representation



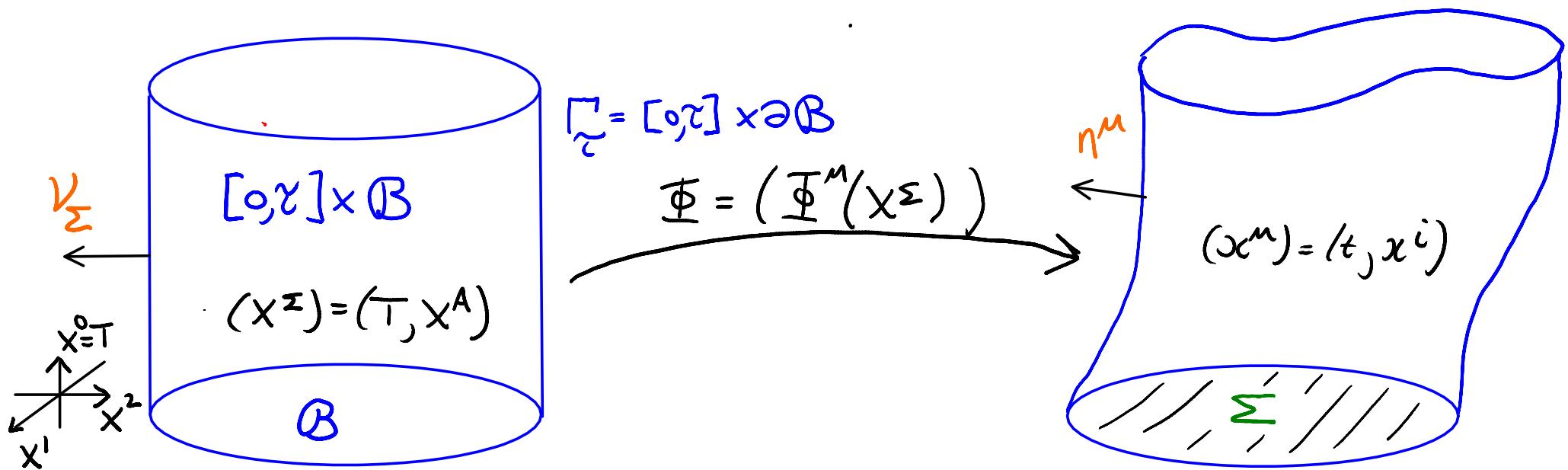
## Material Representation



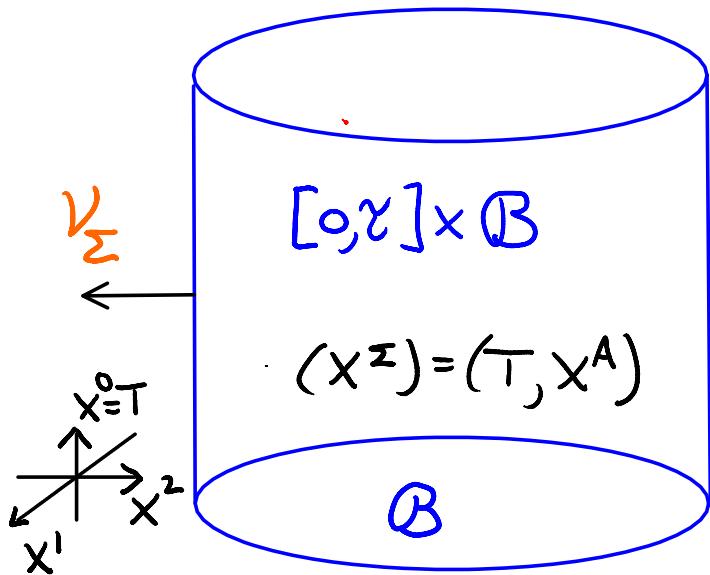
$$\Sigma = [0, \tau] \times \partial \mathbb{B}$$



## Material Representation



## Material Representation

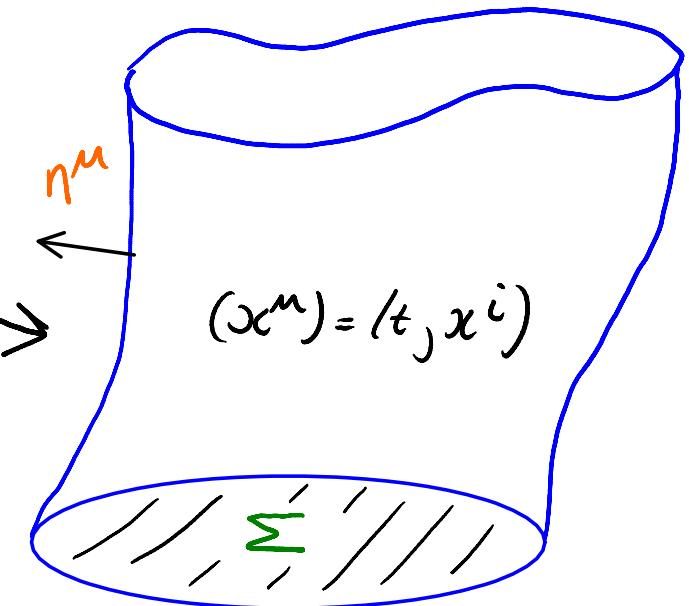


$$\Sigma = [0, \tau] \times \partial \mathcal{B}$$

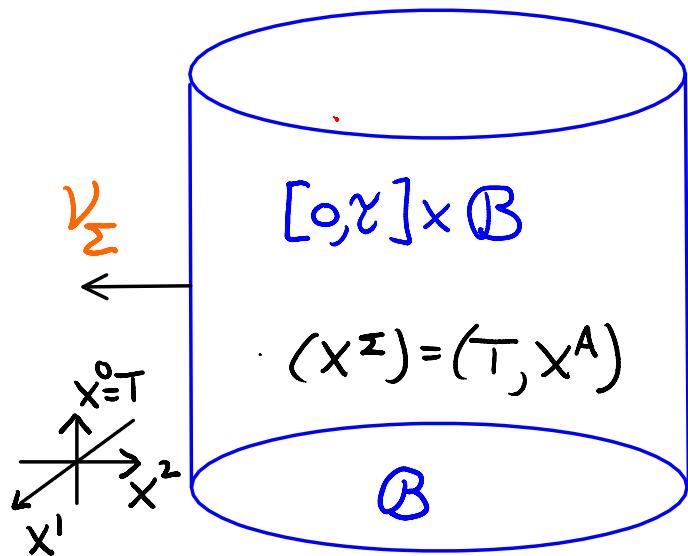
$$\Phi = (\Phi^m(x^\Sigma))$$

$$\Phi^0(T, x) = t \Leftrightarrow T = t$$

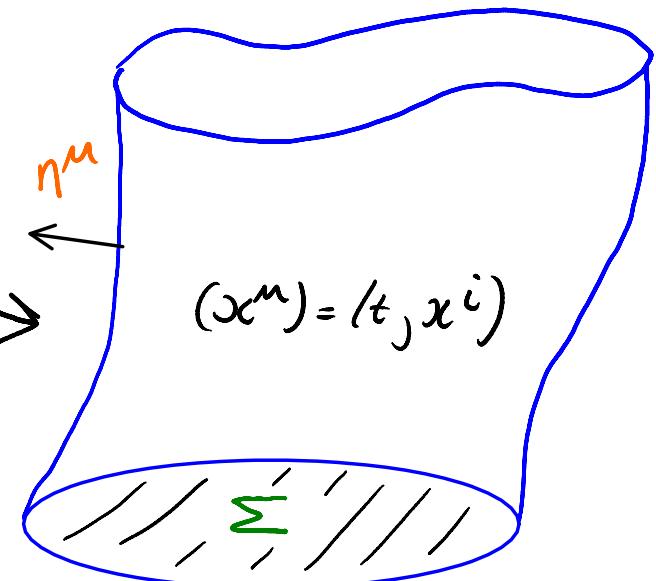
$$f^A(t, \Phi^i(T, x)) = X^A$$



## Material Representation



$$\begin{aligned} \Sigma \hat{\ } &= [0, \tau] \times \partial B \\ \Phi \hat{\ } &= (\Phi^m(x^\Sigma)) \\ \boxed{\begin{aligned} \Phi^0(T, x) &= t \Leftrightarrow T = t \\ f^A(t, \Phi^i(T, x)) &= x^A \end{aligned}} \end{aligned}$$



Unknowns:

$$g_{\mu\nu}(T, X) := g_{\mu\nu}(T, \Phi(T, X)) \quad (\mu, \nu = 0, 1, 2, 3)$$

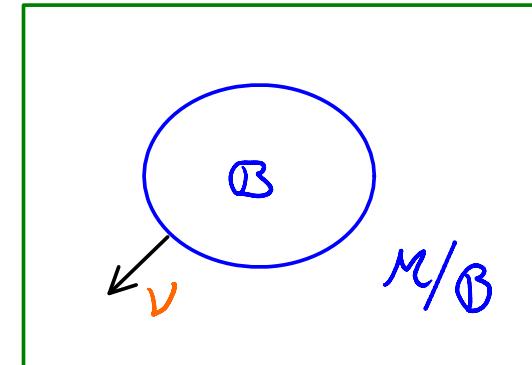
$$\Phi^i(T, X) \quad (i = 1, 2, 3)$$

# Initial Boundary Value Problem

## Initial Boundary Value Problem

Extension of  $\Phi^i$

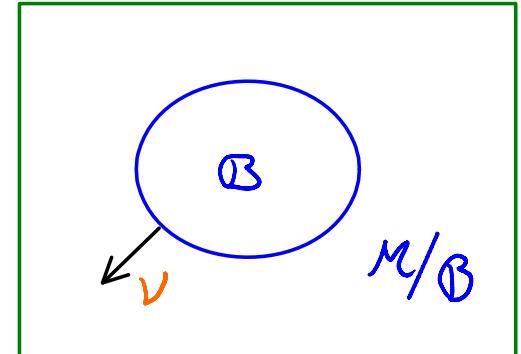
$$\tilde{\Phi}^i \text{ defined on } M \quad \mid \quad \tilde{\Phi}^i|_{\mathcal{B}} = \Phi^i$$



# Initial Boundary Value Problem

Extension of  $\Phi^i$

$$\tilde{\Phi}^i \text{ defined on } M \quad \nmid \quad \tilde{\Phi}^i|_{\mathcal{B}} = \Phi^i$$



IVBP

$$\partial_\Sigma \left( a^\Sigma (\gamma, \partial \tilde{\Phi}) \partial_\gamma \gamma_{\mu\nu} \right) = q_{\mu\nu} (\gamma, \partial \gamma, \partial \tilde{\Phi}) + \chi_B P_{\mu\nu} (\gamma, \partial \gamma, \partial \tilde{\Phi}) \text{ in } [0, \tilde{t}] \times M$$

$$\partial_\Sigma \left( F_i^\Sigma (\gamma, \partial \tilde{\Phi}) \right) = w_i (\gamma, \partial \tilde{\Phi}) \quad \text{in } [0, \tilde{t}] \times B$$

$$\nu \sum F_i^\Sigma (\gamma, \partial \tilde{\Phi}) \Big|_{\Gamma} = 0$$

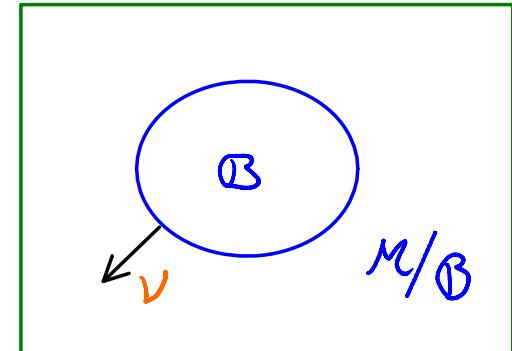
$$(\gamma_{\mu\nu}|_{T=0}, \partial_T \gamma_{\mu\nu}|_{T=0}) = (\gamma_{\mu\nu}^0, \dot{\gamma}_{\mu\nu}^1) \text{ in } B$$

$$(\tilde{\Phi}^i|_{T=0}, \partial_T \tilde{\Phi}^i|_{T=0}) = (\tilde{\Phi}_0^i, \dot{\tilde{\Phi}}_1^i) \text{ in } B$$

# Initial Boundary Value Problem

Extension of  $\Phi^i$

$$\tilde{\Phi}^i \text{ defined on } M \quad \text{if } \tilde{\Phi}^i|_{\mathcal{B}} = \Phi^i$$



IVBP

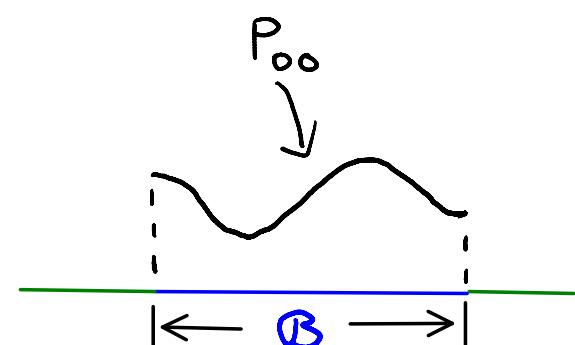
$$\partial_\Sigma (a^\Sigma(\gamma, \partial \tilde{\Phi}) \partial_\gamma \gamma_{\mu\nu}) = q_{\mu\nu}(\gamma, \partial \gamma, \partial \tilde{\Phi}) + \chi_{\mathcal{B}} p_{\mu\nu}(\gamma, \partial \gamma, \partial \tilde{\Phi}) \text{ in } [0, \tilde{t}] \times M$$

$$\partial_\Sigma (F_i^\Sigma(\gamma, \partial \tilde{\Phi})) = w_i(\gamma, \partial \tilde{\Phi}) \text{ in } [0, \tilde{t}] \times \mathcal{B}$$

$$\nu F_i^\Sigma(\gamma, \partial \tilde{\Phi})|_{\Gamma} = 0$$

$$(\gamma_{\mu\nu}|_{T=0}, \partial_T \gamma_{\mu\nu}|_{T=0}) = (\gamma_{\mu\nu}^0, \dot{\gamma}_{\mu\nu}^0) \text{ in } \mathcal{B}$$

$$(\tilde{\Phi}^i|_{T=0}, \partial_T \tilde{\Phi}^i|_{T=0}) = (\Phi_0^i, \dot{\Phi}_1^i) \text{ in } \mathcal{B}$$



# Elasticity Tensor

## Elasticity Tensor

Definition:

$$L_{AB}^{\mu\nu} := \frac{\partial P(f, \partial f)}{\partial(\partial_\nu f^B) \partial(\partial_\mu f^A)}$$

## Elasticity Tensor

Definition:

$$L_{AB}^{\mu\nu} := \frac{\partial P(f, \partial f)}{\partial (\partial_\nu f^B) \partial (\partial_\mu f^A)}$$

Properties:

$$L_{AB}^{\mu\nu} = L_{BA}^{\nu\mu}$$

## Elasticity Tensor

Definition:

$$L_{AB}^{\mu\nu} := \frac{\partial P(f, \partial f)}{\partial (\partial_\nu f^B) \partial (\partial_\mu f^A)}$$

Properties:

$$L_{AB}^{\mu\nu} = L_{BA}^{\nu\mu}$$

Assumptions:  $\exists K_0, K_1, \mu > 0$  such that

$$(i) \left\langle (PL)_{AB}^{ij} \partial_i \psi^A | \partial_j \psi^B \right\rangle_{L^2(\Sigma)} \geq K_0 \|\psi\|_{H^1(\Sigma)} - \mu \|\psi\|_{L^2(\Sigma)}$$

where

$$\Sigma = \left( S|_{t=0} \right)^{-1} (\partial B)$$

$$(ii) L_{AB}^{00} \phi^A \phi^B \leq -K_1 |\phi|$$

# Function Spaces

## Function Spaces

◊ WLOG assume  $\mathcal{M} \cong \mathbb{T}^3$

## Function Spaces

◊ WLOG assume  $M \cong \mathbb{T}^3$

◊ Define

$$\mathcal{H}^{k,s}(\mathbb{T}^3) = H^s(\mathcal{B}) \cap H^s(\mathbb{T}^3 \setminus \mathcal{B}) \cap H^k(\mathbb{T}^3) \quad (s \geq k)$$

with norm

$$\|u\|_{\mathcal{H}^{k,s}(\mathbb{T}^3)}^2 = \|u\|_{\mathcal{B}}^2 + \|u\|_{\mathbb{T}^3 \setminus \mathcal{B}}^2 + \|u\|_{\mathbb{T}^3}^2$$

## Function Spaces

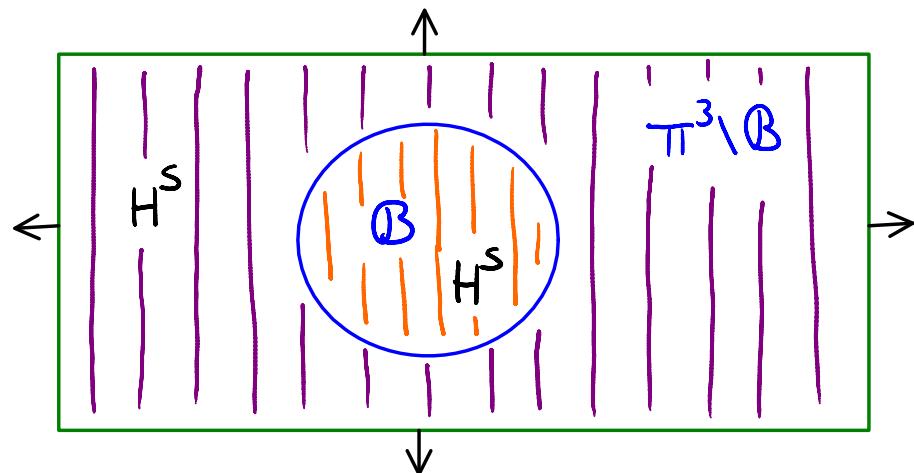
◊ WLOG assume  $\mathcal{M} \cong \mathbb{T}^3$

◊ Define

$$\mathcal{H}^{k,s}(\mathbb{T}^3) = H^s(\mathcal{B}) \cap H^s(\mathbb{T}^3 \setminus \mathcal{B}) \cap H^k(\mathbb{T}^3) \quad (s \geq k)$$

with norm

$$\|u\|_{\mathcal{H}^{k,s}(\mathbb{T}^3)}^2 = \|u\|_{H^s(\mathcal{B})}^2 + \|u\|_{\mathbb{T}^3 \setminus \mathcal{B}}^2 + \|u\|_{H^k(\mathbb{T}^3)}^2$$



## Function Spaces

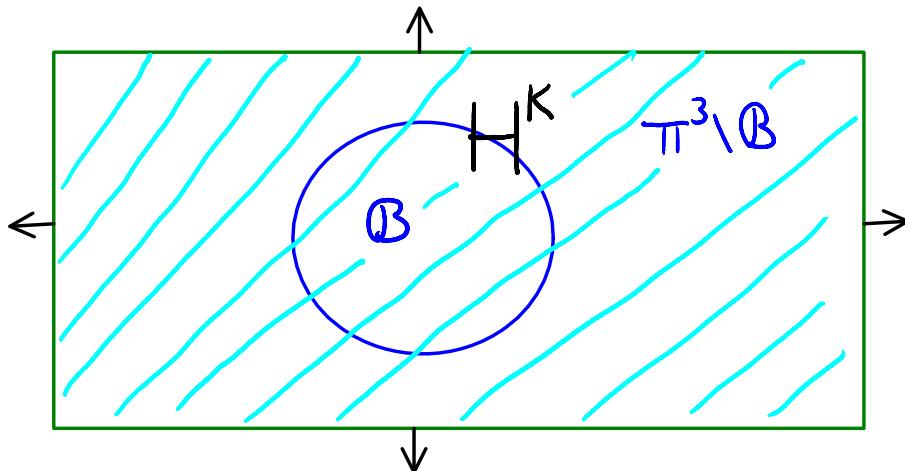
◊ WLOG assume  $\mathcal{M} \cong \mathbb{T}^3$

◊ Define

$$\mathcal{H}^{k,s}(\mathbb{T}^3) = H^s(\mathcal{B}) \cap H^s(\mathbb{T}^3 \setminus \mathcal{B}) \cap H^k(\mathbb{T}^3) \quad (s \geq k)$$

with norm

$$\|u\|_{\mathcal{H}^{k,s}(\mathbb{T}^3)}^2 = \|u\|_{H^s(\mathcal{B})}^2 + \|u\|_{\mathbb{T}^3 \setminus \mathcal{B}}^2 + \|u\|_{H^k(\mathbb{T}^3)}^2$$



## Initial Data Complications II: Compatibility conditions

## Initial Data Complications II: Compatibility conditions

◇  $\frac{\partial^l}{T} \gamma_{\mu\nu} \Big|_{T=0} \in \mathcal{J}^{M_{S+1-l}, S+1-l}(\mathbb{T}^3) \quad l=0, 1, \dots, S+1$

where

$$M_j = \begin{cases} 2 & \text{if } j \geq 2 \\ j & \text{otherwise} \end{cases}$$

and

$$s \in \mathbb{Z}_{>\frac{5}{2}}$$

## Initial Data Complications II: Compatibility conditions

◇  $\left. \frac{\partial^l}{\partial T} \gamma_{\mu\nu} \right|_{T=0} \in \mathcal{H}^{m_{s+1-l}, s+1-l}(\mathbb{T}^3) \quad l=0, 1, \dots, s+1$

where

$$m_j = \begin{cases} 2 & \text{if } j \geq 2 \\ j & \text{otherwise} \end{cases}$$

and

$$s \in \mathbb{Z}_{>\frac{s}{2}}$$

◇  $\left. \frac{\partial^l}{\partial T} \left( \nu_\Sigma F^\Sigma(\gamma, \partial \bar{\Phi}) \right) \right|_{T=0} \in H^{s-l}(\mathcal{B}) \cap H_0^1(\mathcal{B}) \quad l=0, \dots, s-1$

## Local Existence

## Local Existence

Suppose  $s > \frac{n}{2} + 1$  and the initial data

$$(\gamma_{\mu\nu}^0, \gamma_{\mu\nu}^1) \in \mathcal{H}^{2,s+1}(\mathbb{T}^3) \times \mathcal{H}^{2,s}(\mathbb{T}^3)$$

$$(\Phi_0^i, \Phi_1^i) \in \mathcal{H}^{s+1}(\mathbb{B}) \times \mathcal{H}^s(\mathbb{B})$$

satisfies the

## Local Existence

Suppose  $s > \frac{n}{2} + 1$  and the initial data

$$(\gamma_{\mu\nu}^0, \gamma_{\mu\nu}^1) \in \mathcal{H}^{2,s+1}(\mathbb{T}^3) \times \mathcal{H}^{2,s}(\mathbb{T}^3)$$

$$(\Phi_0^i, \Phi_1^i) \in \mathcal{H}^{s+1}(\mathbb{B}) \times \mathcal{H}^s(\mathbb{B})$$

satisfies the

- (i) the gravitational constraint equations, and
- (ii) the compatibility conditions,

## Local Existence

Suppose  $s > \frac{n}{2} + 1$  and the initial data

$$(\gamma_{\mu\nu}^0, \gamma_{\mu\nu}^1) \in \mathcal{H}^{2,s+1}(\mathbb{T}^3) \times \mathcal{H}^{2,s}(\mathbb{T}^3)$$

$$(\Phi_0^i, \Phi_1^i) \in \mathcal{H}^{s+1}(\mathbb{B}) \times \mathcal{H}^s(\mathbb{B})$$

satisfies the

- (i) the gravitational constraint equations, and
- (ii) the compatibility conditions,

then ...

there exists a  $\gamma > 0$  and maps

$$(\chi_{\mu\nu}, \bar{\Phi}^i) \in \bigcap_{l=0}^{s+1} C^l([0, \tilde{\gamma}], \mathcal{H}^{m_{s+1-l}, s+1-l}(\mathbb{T}^3)) \times C^l([0, \tilde{\gamma}], H^{s+1-l}(\mathcal{B}))$$

there exists a  $\gamma > 0$  and maps

$$(\gamma_{\mu\nu}, \bar{\Phi}^i) \in \bigcap_{\ell=0}^{s+1} C^\ell([0, \tilde{\gamma}], \mathcal{H}^{m_{s+1-\ell}, s+1-\ell}(\mathbb{R}^3)) \times C^\ell([0, \tilde{\gamma}], H^{s+1-\ell}(\mathcal{B}))$$

such that  $(\gamma_{\mu\nu}, \bar{\Phi}^i)$  is the unique solution of the IVP:

$$\partial_\Sigma (a^\Sigma(\gamma, \partial \bar{\Phi}) \partial_\gamma \gamma_{\mu\nu}) = q_{\mu\nu}(\gamma, \partial \gamma, \partial \bar{\Phi}) + \chi_{\mathcal{B}} p_{\mu\nu}(\gamma, \partial \gamma, \partial \bar{\Phi}) \text{ in } [0, \tilde{\gamma}] \times \mathcal{M}$$

$$\partial_\Sigma (F_i^\Sigma(\gamma, \partial \bar{\Phi})) = w_i(\gamma, \partial \bar{\Phi}) \text{ in } [0, \tilde{\gamma}] \times \mathcal{B}$$

$$\left. \sum F_i^\Sigma(\gamma, \partial \bar{\Phi}) \right|_{\Gamma} = 0$$

$$(\gamma_{\mu\nu}|_{T=0}, \partial_T \gamma_{\mu\nu}|_{T=0}) = (\gamma_{\mu\nu}^0, \gamma_{\mu\nu}^1) \text{ in } \mathcal{B}$$

$$(\bar{\Phi}^i|_{T=0}, \partial_T \bar{\Phi}^i|_{T=0}) = (\bar{\Phi}_0^i, \bar{\Phi}_1^i) \text{ in } \mathcal{B}$$

# Perfect Fluids

## Perfect Fluids

Assumption:

$$\rho = \rho(n)$$

where

$$n = \left( \frac{1}{6} H^{AB} H^{CD} H^{EF} \Omega_{ACE} \Omega_{BDF} \right)$$

and

$$H^{AB} = g^{\mu\nu} \partial_\mu f^A \partial_\nu f^B$$

# Perfect Fluids

Assumption:

$$\rho = \rho(n)$$

where

$$n = \left( \frac{1}{6} H^{AB} H^{CD} H^{EF} \Omega_{ACE} \Omega_{BDF} \right)$$

and

$$H^{AB} = g^{\mu\nu} \partial_\mu f^A \partial_\nu f^B$$

Consequences:



$$T_{\mu\nu} = (\rho + p) u_\mu u_\nu + p g_{\mu\nu}$$

where

$$p = n \frac{\partial \rho}{\partial n} - \rho$$



$$\nabla_\mu T^{\mu\nu} = 0$$

## A brief (partial) history of the problem

## A brief (partial) history of the problem

---

1757

Euler's  
Eqns.

## A brief (partial) history of the problem

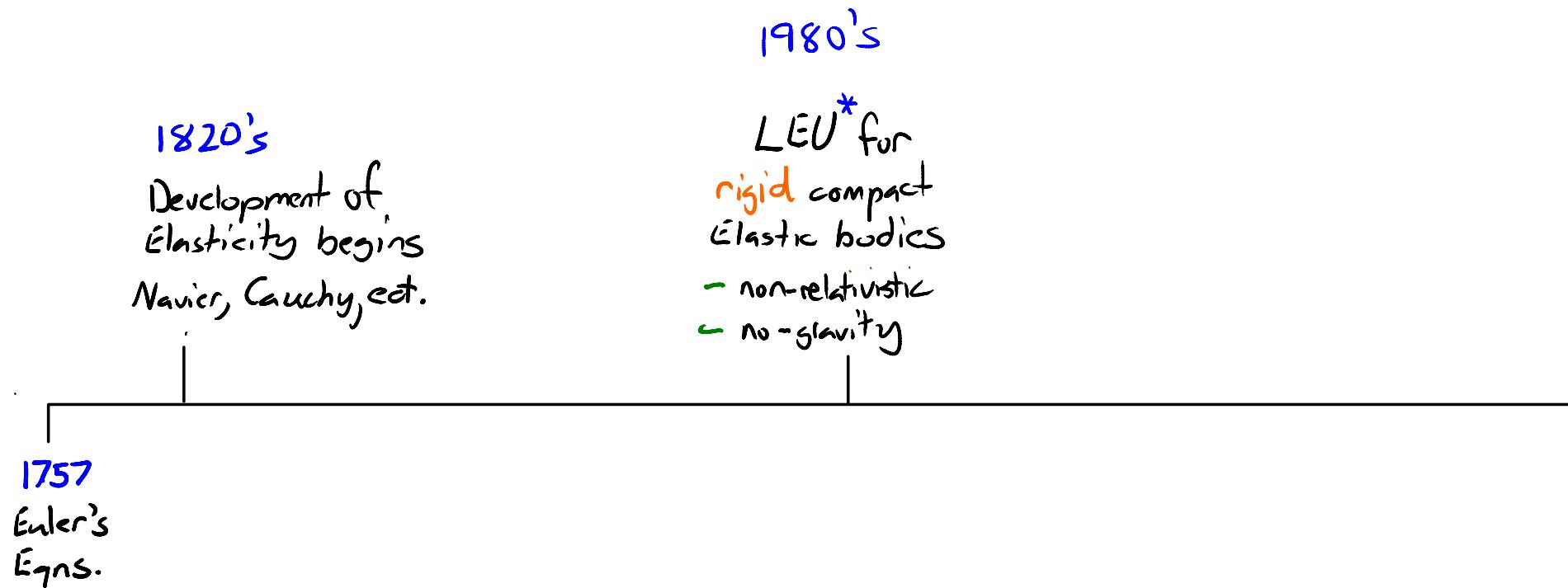
1820's

Development of  
Elasticity begins  
Navier, Cauchy, etc.

1757

Euler's  
Eqns.

## A brief (partial) history of the problem



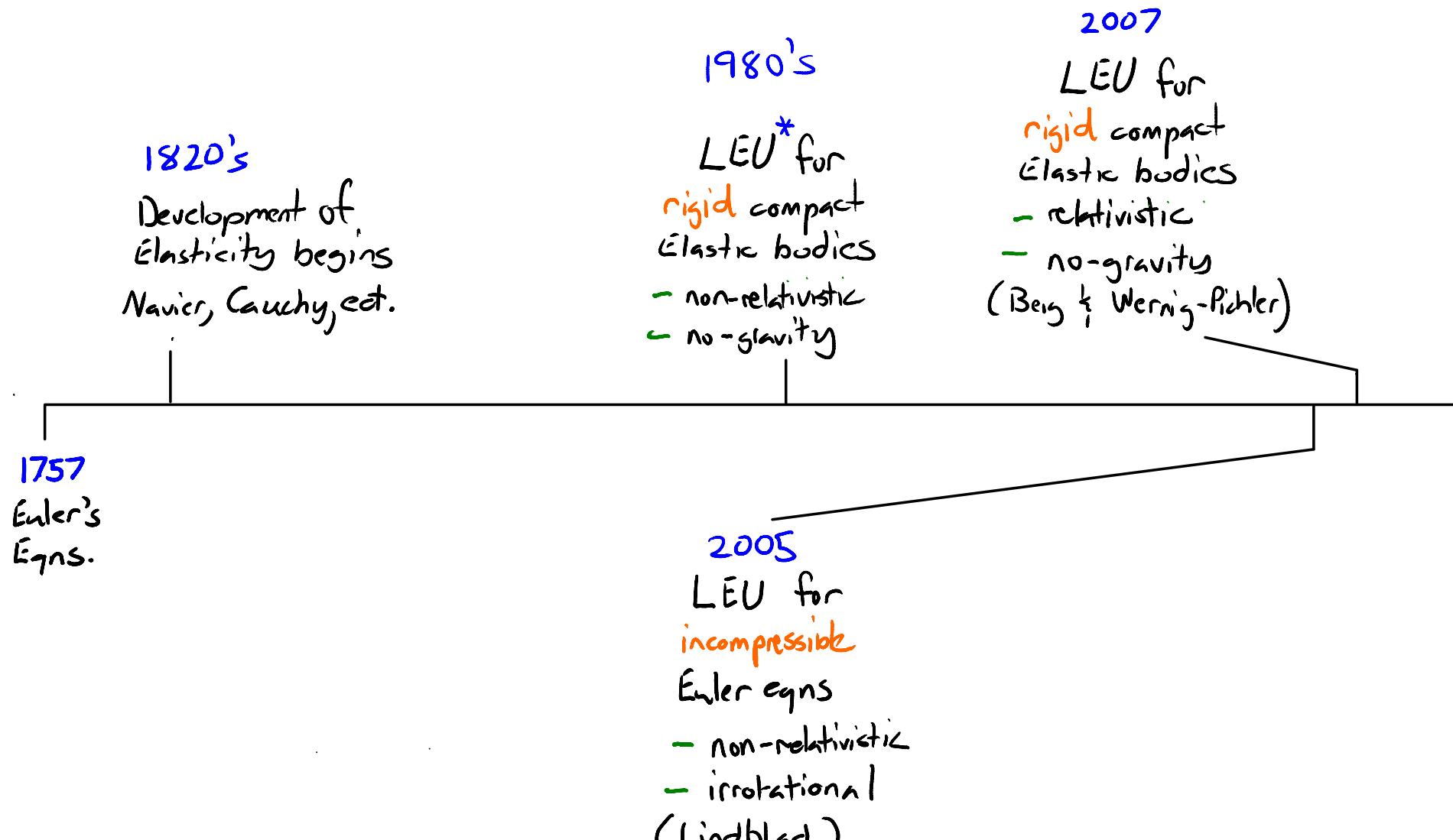
\* Local Existence and Uniqueness (LEU)

# A brief (partial) history of the problem



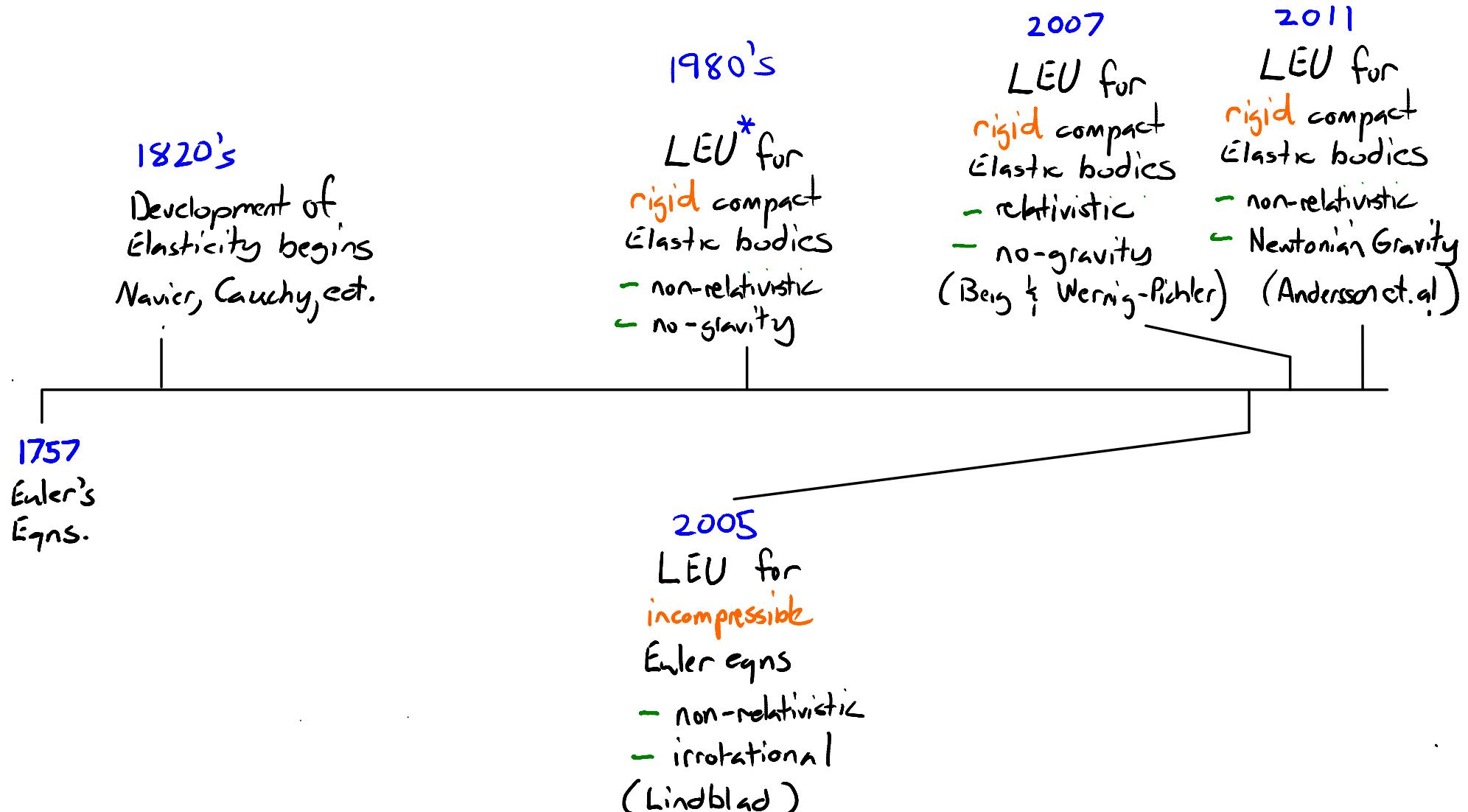
\* Local Existence and Uniqueness (LEU)

# A brief (partial) history of the problem



\* Local Existence and Uniqueness (LEU)

# A brief (partial) history of the problem



\* Local Existence and Uniqueness (LEU)

# A brief (partial) history of the problem

