AE837

Advanced Mechanics of Damage Tolerance

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upcoming schedule

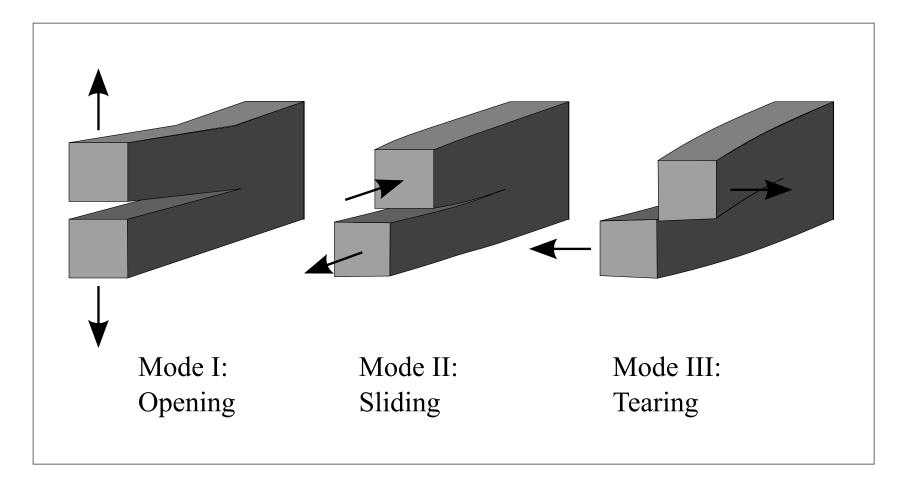
- Sep 3 Elastic Stress Field, Homework 1 Due
- Sep 5 Elastic Stress Field
- Sep 10 Elastic Stress Field
- Sep 12 Elastic Stress Field

outline

- review
- complex airy stress
- westergaard function method
- solutions using westergaard

review

fracture modes



stress field

• The approximate stress field for the various modes can be expressed in terms of the stress intensity factor

$$egin{align} \sigma_{yy} &= rac{K_I}{\sqrt{2\pi x}} + O(\sqrt{x}) & \sigma_{xy} &= \sigma_{yz} &= 0 \ \sigma_{xy} &= rac{K_II}{\sqrt{2\pi x}} + O(\sqrt{x}) & \sigma_{yy} &= \sigma_{yz} &= 0 \ \sigma_{yz} &= rac{K_III}{\sqrt{2\pi x}} + O(\sqrt{x}) & \sigma_{yy} &= \sigma_{xy} &= 0 \ \end{pmatrix}$$

- A stress function technique that can be used to solve many planar problems is known as the *Airy stress function*
- This method reduces the governing equations for a planar problem to a single unknown function

ullet We assume first that body forces are derivable from a *potential* function, V

$$F_x = -rac{\partial V}{\partial x} \ F_y = -rac{\partial V}{\partial y}$$

- How restrictive is this assumption?
- Most body forces are linear (gravity) and can easily be represented this way

Consider the following

$$egin{align} \sigma_{xx} &= rac{\partial^2 \phi}{\partial y^2} + V \ \sigma_{yy} &= rac{\partial^2 \phi}{\partial x^2} + V \ au_{xy} &= -rac{\partial^2 \phi}{\partial x \partial y} \ \end{pmatrix}$$

- The function $\phi = \phi(x, y)$ is known as the Airy stress function
- Equilibrium is automatically satisfied

compatibility

• Substituting the Airy Stress function and potential function into the relationships, we find

$$\frac{\partial^4 \phi}{\partial x^4} + 2 \frac{\partial^4 \phi}{\partial x^2 \partial y^2} + \frac{\partial^4 \phi}{\partial y^4} = -\frac{1 - 2\nu}{1 - \nu} \left(\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} \right) \qquad \text{plane strain}$$

$$\frac{\partial^4 \phi}{\partial x^4} + 2 \frac{\partial^4 \phi}{\partial x^2 \partial y^2} + \frac{\partial^4 \phi}{\partial y^4} = -(1 - \nu) \left(\frac{\partial^2 V}{\partial x^2} + \frac{\partial^2 V}{\partial y^2} \right) \qquad \text{plane stress}$$

compatibility

• If there are no body forces, or the potential function satisfies Laplace's Equation

$$\nabla^2 V = 0$$

• Then both plane stress and plane strain reduce to

$$rac{\partial^4 \phi}{\partial x^4} + 2rac{\partial^4 \phi}{\partial x^2 \partial y^2} + rac{\partial^4 \phi}{\partial y^4} = 0$$

complex airy stress

complex conjugates

• In the cartesian system we can express complex conjugates as

$$z = x + iy$$

$$ar{z} = x - iy$$

• And in polar coordinates we have

$$z=r(\cos heta + i\sin heta) = re^{i heta}$$

$$ar{z} = r(\cos heta - i \sin heta) = re^{-i heta}$$

analytic functions

• A complex function can be written as

$$f(z)=u(x,y)+iv(x,y)$$

• A complex function is said to be analytic if

$$rac{\partial}{\partial x}f(z)=f'(z)rac{\partial z}{\partial x}=f'(z)$$

• and

$$rac{\partial}{\partial y}f(z)=f'(z)rac{\partial z}{\partial y}=if'(z)$$

Cauchy-Riemann

• This means that

$$rac{\partial}{\partial x}f(z) = -irac{\partial}{\partial y}f(z) \ rac{\partial u}{\partial x} + irac{\partial v}{\partial x} = rac{\partial v}{\partial y} - irac{\partial u}{\partial y}$$

Cauchy-Riemann

• From this we obtain the Cauchy-Riemann equations

$$rac{\partial u}{\partial x} = rac{\partial v}{\partial y} \qquad rac{\partial v}{\partial x} = -irac{\partial u}{\partial y}$$

• From which we can easily derive the following

$$\nabla^2 u = \nabla^2 v = 0$$

• The Airy stress function is biharmonic, we can write the following

$$\nabla^2 \phi = P$$

• From which we have

$$\nabla^2 P = \nabla^2 \nabla^2 \phi = 0$$

• And thus *P* will satisfy the compatibility equations

- We can say that *P* is the real part of a complex function $P = Re\{f(z)\}$ where f(z)=P+iQ
- Now we let

$$\psi(z)=rac{1}{4}\int f(z)dz=p+iq$$

• ψ will also be analytic, thus

$$\psi'(z)=rac{1}{4}f(z)$$

• According to the Cauchy-Riemann equations we now have

$$\psi'(z) = rac{\partial p}{\partial x} + irac{\partial q}{\partial x} = rac{\partial q}{\partial y} - irac{\partial p}{\partial y}$$

And we find that

$$P=4rac{\partial p}{\partial x}=4rac{\partial q}{\partial y}$$

• If we now consider $\phi - (xp + yq)$, we can show that

$$\nabla^2[\phi - (xp + yq)] = 0$$

• This means that $\phi - (xp + yq)$ is harmonic, and can be taken as either the real or imaginary portion of some analytic function, $\chi(z)$

$$\phi - (xp + yq) = \text{Re}\{\chi(z)\}$$

• We can now say that

$$xp + yq = \operatorname{Re}\{ar{z}\,\psi(z)\}$$

complex airy stress

• The complex representation of the Airy stress function can now be written as

$$egin{aligned} \phi &= \mathrm{Re}\{ar{z}\psi(z) + \chi(z)\} \ 2\phi(x,y) &= ar{z}\psi(z) + z\psiar{(}z) + \chi(z) + \chiar{(}z) \end{aligned}$$

And from the definition of the Airy stress function we obtain

$$egin{aligned} \sigma_{xx} + i\sigma_{xy} &= rac{\partial^2 \phi}{\partial y^2} - irac{\partial^2 \phi}{\partial x \partial y} &= -irac{\partial}{\partial y} igg(rac{\partial \phi}{\partial x} + irac{\partial \phi}{\partial y}igg) \ \sigma_{yy} - i\sigma_{xy} &= rac{\partial^2 \phi}{\partial x^2} + irac{\partial^2 \phi}{\partial x \partial y} &= rac{\partial}{\partial x} igg(rac{\partial \phi}{\partial x} + irac{\partial \phi}{\partial y}igg) \end{aligned}$$

properties of analytic functions and conjugates

$$egin{aligned} rac{\partial f(z)}{\partial x} &= f' rac{\partial z}{\partial x} &= f'(z) \ rac{\partial f(z)}{\partial x} &= \left(rac{\partial ar{f}(z)}{\partial x}
ight) &= f^{'}(z) \ rac{\partial f(z)}{\partial y} &= f'(z) rac{\partial z}{\partial y} &= if'(z) \ rac{\partial f(z)}{\partial u} &= \left(rac{\partial ar{f}(z)}{\partial u}
ight) &= -if^{'}(z) \end{aligned}$$

complex airy stress functions

The above properties allow us to write

$$rac{\partial \phi}{\partial x} + i rac{\partial \phi}{\partial y} = \psi(z) + z \psi^{'} \overline{(z)} + \chi^{'} \overline{(z)}$$

Which gives

$$egin{aligned} \sigma_{xx}+i\sigma_{xy}&=\psi'(z)+\psi'ar{(}z)-z\psi''ar{(}z)-\chi''ar{(}z)\ \sigma_{yy}-i\sigma_{xy}&=\psi'(z)+\psi'ar{(}z)+z\psi''ar{(}z)+\chi''ar{(}z) \end{aligned}$$

- ullet We can add the two equations to find $\sigma_{xx}+\sigma_{yy}=2(\psi'(z)+\psi'(z))=4\mathrm{Re}\{\psi'(z)\}$
- Similarly we can subtract the equations to find $\sigma_{yy}-\sigma_{xx}-2i\sigma_{xy}=2(z\psi''(z)+\chi''(z))$
- ullet We can re-write using the complex conjugate $\sigma_{yy}-\sigma_{xx}+2i\sigma_{xy}=2(ar{z}\,\psi''(z)+\chi''(z))$

displacement

- We can use Hooke's law to find displacements in terms of the complex functions
- After some algebra, we find

$$2\mu(u_x+iu_y)=\kappa\psi(z)-z\psi^{'}(z)-\chi^{'}(z)$$

• Where

$$\kappa = \frac{\lambda^* + 3\mu}{\lambda^* + \mu}$$

westergaard function method

• If we consider an infinite plate with cracks along the x-axis, and external loads are symmetric with respect to the x-axis, then $\sigma_{xy} = 0$ along y = 0

$$\operatorname{Im}\{ar{z}\psi''(z)+\chi''(z)\}=0 \qquad ext{at } y=0$$

• At y = 0, $z = \bar{z}$, therefore we have (for all y)

$$\chi^{\prime\prime}(z)+z\psi^{\prime\prime}(z)+A=0$$

• Where *A* is some real constant

• After substituting $\chi''(z) = -z\psi''(z) - A$ into the Airy stress relationships we find

$$egin{aligned} \sigma_{xx} &= 2\mathrm{Re}\{\psi'\} - 2y\mathrm{Im}\{\psi''\} + A \ \sigma_{yy} &= 2\mathrm{Re}\{\psi'\} + 2y\mathrm{Im}\{\psi''\} - A \ \sigma_{xy} &= -2y\mathrm{Re}\{\psi''\} \end{aligned}$$

• We now define

$$\psi' = \frac{1}{2}(Z_I + A)$$

ullet where $\hat{Z_I'}\equiv Z_I$ and thus

$$\psi=rac{1}{2}(\hat{Z_I}+Az) \ \psi''=rac{1}{2}Z_I'$$

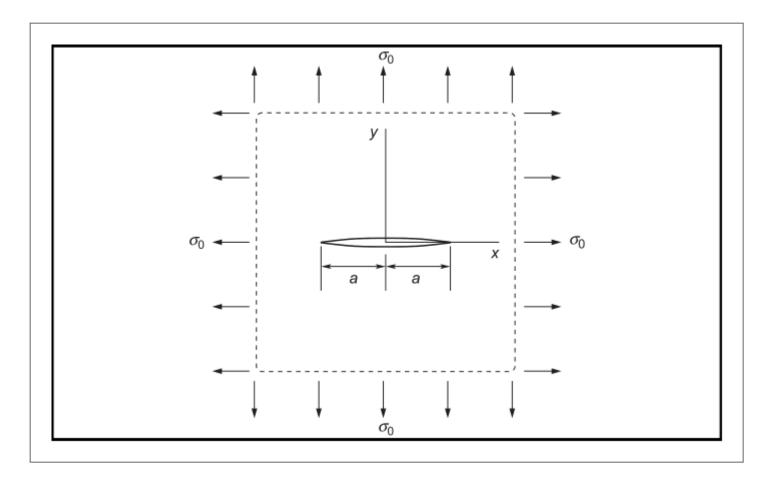
• This gives the following result for the stress field

$$egin{aligned} \sigma_{xx} &= \mathrm{Re}\{Z_I\} - y\mathrm{Im}\{Z_I'\} + 2A \ \sigma_{yy} &= \mathrm{Re}\{Z_I\} + y\mathrm{Im}\{Z_I'\} \ \sigma_{xy} &= -y\mathrm{Re}\{Z_I'\} \end{aligned}$$

• and Z_I is the so-called Westergaard function

solutions using westergaard

mode I crack



8/23/2019 Lecture 5 - elastic stress field

mode I crack

