

Created by Prabal Negi

Navier-Stokes in Cylindrical formulation

Weak form and Nek implementation

Formulation of Navier-Stokes 1

Mapping from Cartesian to Cylindrical 1.1

$$x_c = x (1a)$$

$$y_c = r\cos(\theta) \tag{1b}$$

$$z_c = r\sin(\theta) \tag{1c}$$

$$\begin{cases} dx_c \\ dy_c \\ dz_c \end{cases} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -r\sin(\theta) \\ 0 & \sin(\theta) & r\cos(\theta) \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$
 (2)

$$\begin{cases}
 dx_c \\
 dy_c \\
 dz_c
\end{cases} = \begin{pmatrix}
 1 & 0 & 0 \\
 0 & \cos(\theta) & -r\sin(\theta) \\
 0 & \sin(\theta) & r\cos(\theta)
\end{pmatrix} \begin{cases}
 dx \\
 dr \\
 d\theta
\end{cases}$$

$$\Rightarrow \begin{cases}
 dx \\
 dr \\
 d\theta
\end{cases} = \begin{pmatrix}
 1 & 0 & 0 \\
 0 & \cos(\theta) & \sin(\theta) \\
 0 & -\sin(\theta)/r & \cos(\theta)/r
\end{pmatrix} \begin{cases}
 dx_c \\
 dy_c \\
 dz_c
\end{cases}$$
(2)

$$\hat{x} = \hat{x_c} \tag{4a}$$

$$\hat{r} = \cos(\theta)\hat{y_c} + \sin(\theta)\hat{z_c} \tag{4b}$$

$$\hat{\theta} = -\sin(\theta)\hat{y_c} + \cos(\theta)\hat{z_c} \tag{4c}$$

$$\Rightarrow \begin{cases} d\hat{x} \\ d\hat{r} \\ d\hat{\theta} \end{cases} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -\sin(\theta)\hat{y}_c + \cos(\theta)\hat{z}_c \\ 0 & 0 & -\cos(\theta)\hat{y}_c - \sin(\theta)\hat{z}_c \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$

$$\Rightarrow \begin{cases} d\hat{x} \\ d\hat{r} \\ d\hat{\theta} \end{cases} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \hat{\theta} \\ 0 & 0 & -\hat{r} \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$
(5)

1.2 Gradient/Divergence/Laplace Operators

For scalar fields

$$\nabla(\psi) = \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial x_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial x_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial x_c}\right) \hat{x_c} + \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial y_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial y_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial y_c}\right) \hat{y_c} + \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial z_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial z_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial z_c}\right) \hat{z_c}$$

$$\therefore \nabla(\psi) = \frac{\partial \psi}{\partial x}\hat{x} + \frac{\partial \psi}{\partial r}\hat{r} + \frac{1}{r}\frac{\partial \psi}{\partial \theta}\hat{\theta}$$
 (6)

Laplacian of a scalar field

$$\nabla \cdot \nabla(\psi) = \left[\frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial}{\partial \theta} \hat{\theta} \right] \cdot \left[\hat{x} \frac{\partial \psi}{\partial x} + \hat{r} \frac{\partial \psi}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial \psi}{\partial \theta} \right]$$

$$= \begin{cases} \left[\hat{x} \cdot \hat{x} \frac{\partial}{\partial x} \frac{\partial \psi}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial x} \cdot \hat{x} \right] + \left[\hat{x} \cdot \hat{r} \frac{\partial}{\partial x} \frac{\partial \psi}{\partial r} + \frac{\partial \psi}{\partial r} \frac{\partial \hat{r}}{\partial x} \cdot \hat{x} \right] + \left[\hat{x} \cdot \hat{\theta} \frac{\partial}{\partial x} \left(\frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial x} \cdot \hat{x} \right] \right]$$

$$+ \left[\hat{r} \cdot \hat{x} \frac{\partial}{\partial r} \frac{\partial \psi}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial r} \cdot \hat{r} \right] + \left[\hat{r} \cdot \hat{r} \frac{\partial}{\partial r} \frac{\partial \psi}{\partial r} + \frac{\partial \psi}{\partial r} \frac{\partial \hat{r}}{\partial r} \cdot \hat{r} \right] + \left[\hat{r} \cdot \hat{\theta} \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial r} \cdot \hat{r} \right]$$

$$+ \left[\hat{\theta} \cdot \hat{x} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial \psi}{\partial x} + \frac{1}{r} \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial \theta} \cdot \hat{\theta} \right] + \left[\hat{\theta} \cdot \hat{r} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial \psi}{\partial r} + \frac{1}{r} \frac{\partial \psi}{\partial r} \frac{\partial \hat{\theta}}{\partial \theta} \cdot \hat{\theta} \right]$$

$$+ \left[\hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial \theta} \cdot \hat{\theta} \right]$$

$$\therefore \nabla^2(\psi) = \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \psi}{\partial \theta}$$
 (7)

Vector gradient. Highlighted in color are the only terms which are non-zero differentials of unit vectors.

$$\nabla(\boldsymbol{u}) = \left[\hat{x} \frac{\partial}{\partial x} + \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \right] \left[\hat{x} u_x + \hat{r} u_r + \hat{\theta} u_\theta \right]$$

$$= \left\{ \begin{bmatrix} \hat{x} \frac{\partial u_x}{\partial x} \hat{x} + \hat{x} u_x \frac{\partial \hat{x}}{\partial x} \end{bmatrix} + \left[\hat{x} \frac{\partial u_r}{\partial x} \hat{r} + \hat{x} u_r \frac{\partial \hat{r}}{\partial x} \right] + \left[\hat{x} \frac{\partial u_\theta}{\partial x} \hat{\theta} + \hat{x} u_\theta \frac{\partial \hat{\theta}}{\partial x} \right]$$

$$+ \left[\hat{r} \frac{\partial u_x}{\partial r} \hat{x} + \hat{r} u_x \frac{\partial \hat{x}}{\partial r} \right] + \left[\hat{r} \frac{\partial u_r}{\partial r} \hat{r} + \hat{r} u_r \frac{\partial \hat{r}}{\partial r} \right] + \left[\hat{r} \frac{\partial u_\theta}{\partial r} \hat{\theta} + \hat{r} u_\theta \frac{\partial \hat{\theta}}{\partial r} \right]$$

$$+ \left[\hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x} + \hat{\theta} \frac{u_x}{r} \frac{\partial \hat{x}}{\partial \theta} \right] + \left[\hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \hat{r} + \hat{\theta} \frac{u_r}{r} \frac{\partial \hat{r}}{\partial \theta} \right] + \left[\hat{\theta} \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \hat{\theta} + \hat{\theta} \frac{u_\theta}{r} \frac{\partial \hat{\theta}}{\partial \theta} \right]$$

$$(8)$$

$$\therefore \nabla(\boldsymbol{u}) = \begin{bmatrix} \hat{x} \frac{\partial u_x}{\partial x} \hat{x} & \hat{x} \frac{\partial u_r}{\partial x} \hat{r} & \hat{x} \frac{\partial u_\theta}{\partial x} \hat{\theta} \\ \hat{r} \frac{\partial u_x}{\partial r} \hat{x} & \hat{r} \frac{\partial u_r}{\partial r} \hat{r} & \hat{r} \frac{\partial u_\theta}{\partial r} \hat{\theta} \\ \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x} & \hat{\theta} \left(-\frac{u_\theta}{r} + \frac{1}{r} \frac{\partial u_r}{\partial \theta} \right) \hat{r} & \hat{\theta} \left(\frac{u_r}{r} + \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \right) \hat{\theta} \end{bmatrix} \tag{9}$$

Vector Laplacian:

$$\nabla \cdot \nabla (\boldsymbol{u}) = \left(\hat{\boldsymbol{x}} \cdot \frac{\partial}{\partial \boldsymbol{x}} + \hat{\boldsymbol{r}} \cdot \frac{\partial}{\partial r} + \hat{\boldsymbol{\theta}} \cdot \frac{1}{r} \frac{\partial}{\partial \theta} \right) \begin{pmatrix} \hat{\boldsymbol{x}} \frac{\partial u_x}{\partial \boldsymbol{x}} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{x}} \frac{\partial u_r}{\partial \boldsymbol{x}} \hat{\boldsymbol{r}} + & \hat{\boldsymbol{x}} \frac{\partial u_\theta}{\partial \boldsymbol{x}} \hat{\boldsymbol{\theta}} + \\ \hat{\boldsymbol{r}} \frac{\partial u_x}{\partial r} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{r}} \frac{\partial u_r}{\partial r} \hat{\boldsymbol{r}} + & \hat{\boldsymbol{r}} \frac{\partial u_\theta}{\partial r} \hat{\boldsymbol{\theta}} + \\ \hat{\boldsymbol{\theta}} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{\theta}} \left(-\frac{u_\theta}{r} + \frac{1}{r} \frac{\partial u_r}{\partial \theta} \right) \hat{\boldsymbol{r}} + & \hat{\boldsymbol{\theta}} \left(\frac{u_r}{r} + \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \right) \hat{\boldsymbol{\theta}} \end{pmatrix}$$

First Column

$$\left(\hat{x} \cdot \frac{\partial}{\partial x} + \hat{r} \cdot \frac{\partial}{\partial r} + \hat{\theta} \cdot \frac{1}{r} \frac{\partial}{\partial \theta}\right) \left(\hat{x} \frac{\partial u_x}{\partial x} + \hat{r} \frac{\partial u_x}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x}$$

$$= \hat{x} \cdot \left(\hat{x} \frac{\partial^2 u_x}{\partial x^2} + \frac{\partial \hat{x}}{\partial x} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{x} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{\partial \hat{x}}{\partial x}$$

$$+ \hat{x} \cdot \left(\hat{r} \frac{\partial}{\partial x} \frac{\partial u_x}{\partial r} + \frac{\partial \hat{r}}{\partial x} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{x} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{x} \cdot \left(\hat{\theta} \frac{\partial}{\partial x} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) + \frac{\partial \hat{\theta}}{\partial x} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{x} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \frac{\partial \hat{x}}{\partial x}$$

$$+ \hat{r} \cdot \left(\hat{x} \frac{\partial}{\partial r} \frac{\partial u_x}{\partial x} + \frac{\partial \hat{x}}{\partial r} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{r} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial^2 u_x}{\partial r^2} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{r} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{\theta} \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) + \frac{\partial \hat{\theta}}{\partial r} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{\theta} \cdot \left(\hat{x} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial u_x}{\partial x} + \frac{\partial \hat{x}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{\theta} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{1}{r} \frac{\partial \hat{x}}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{r} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial u_x}{\partial r} + \frac{\partial \hat{r}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{\theta} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{1}{r} \frac{\partial \hat{x}}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \frac{\partial \hat{\theta}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \frac{\partial \hat{\theta}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}$$

$$= \left[\frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial r^2} + \frac{1}{r} \frac{\partial^2 u_x}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u_x}{\partial \theta^2}\right] \hat{x}$$
(10)

Second column:

Third column:

$$\therefore \nabla \cdot \nabla (\boldsymbol{u}) = \begin{cases}
\left[\frac{\partial^2 u_x}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_x}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_x}{\partial \theta^2} \right] \hat{x} + \\
\left[\frac{\partial^2 u_r}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_r}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_r}{\partial \theta^2} - \frac{u_r}{r^2} - \frac{2}{r^2} \frac{\partial u_\theta}{\partial \theta} \right] \hat{r} + \\
\left[\frac{\partial^2 u_\theta}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left(r \frac{\partial u_\theta}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_\theta}{\partial \theta^2} + \frac{2}{r^2} \frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r^2} \right] \hat{\theta}
\end{cases} \tag{13}$$

Which can be written using the Laplacian of individual scalar fields, u_x, u_r, u_θ as,

$$\therefore \nabla \cdot \nabla(\boldsymbol{u}) = \nabla^2(u_x)\hat{x} + \left(\nabla^2(u_r) - \frac{u_r}{r^2} - \frac{2}{r^2}\frac{\partial u_\theta}{\partial \theta}\right)\hat{r} + \left(\nabla^2(u_\theta) + \frac{2}{r^2}\frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r^2}\right)\hat{\theta}$$
(14)

2 Nek Implementations:

2.1 Opgradt

The opgradt subroutine implements the gradient of a field defined on Mesh 2 (pressure mesh), integrated w.r.t the Mesh 1 test functions. This is done with integration by parts which creates a boundary term.

Pressure term in the momentum equations:

$$-\int_{\Omega} v \frac{\partial}{\partial x_{j}} (p \delta_{ij}) d\Omega$$
$$-\int_{\Omega} \frac{\partial}{\partial x_{j}} (v p \delta_{ij}) d\Omega + \int_{\Omega} p \delta_{ij} \frac{\partial}{\partial x_{j}} (v) d\Omega$$
$$-\int_{\partial \Omega} (v p \delta_{ij}) n_{j} dA + \int_{\Omega} p \delta_{ij} \frac{\partial}{\partial x_{j}} (v) d\Omega$$

The first term is the boundary condition. The second term is what opgradt evaluates. I represent the determinant of the mapping between reference element and physical coordinates, the Jacobian as \mathcal{J} . In the continuous formulation, this becomes:

$$w_{i} = \int_{\Omega} p \frac{\partial v}{\partial x_{i}} d\Omega$$

$$= \int_{\Omega} p \frac{\partial v}{\partial x_{i}} \left(\frac{\partial \Omega}{\partial \hat{\Omega}}\right) d\hat{\Omega}$$

$$= \int_{\Omega} p \left(\frac{1}{\mathcal{J}} \frac{\partial r}{\partial x_{i}} \frac{\partial v}{\partial r} + \frac{1}{\mathcal{J}} \frac{\partial s}{\partial x_{i}} \frac{\partial v}{\partial s} + \frac{1}{\mathcal{J}} \frac{\partial t}{\partial x_{i}} \frac{\partial v}{\partial t}\right) \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} p \left(\frac{\partial r}{\partial x_{i}} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial x_{i}} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial x_{i}} \frac{\partial v}{\partial t}\right) d\hat{\Omega}$$

$$\implies w_{i} = \int_{\Omega} p \left(\frac{\partial r_{j}}{\partial x_{i}} \frac{\partial v}{\partial r_{j}}\right) d\hat{\Omega}, \tag{15}$$

where, r_j represents the reference element coordinate directions r, s, t. After discretization, this becomes:

$$\implies w_i = \sum_k W(x_k) \left(\frac{\partial r_j}{\partial x_i} \frac{\partial v}{\partial r_j} (x_k) \right) p$$

When cross derivative terms like $\partial s/\partial x$ *etc.* are non-zero, we have the following expression

$$w_{i} = (\mathcal{I}_{t12}^{T} \otimes \mathcal{I}_{s12}^{T} \otimes \mathcal{D}_{r12}^{T})(\frac{\partial r}{\partial x_{i}}. *W. *p) +$$

$$(\mathcal{I}_{t12}^{T} \otimes \mathcal{D}_{s12}^{T} \otimes \mathcal{I}_{r12}^{T})(\frac{\partial s}{\partial x_{i}}. *W. *p) +$$

$$(\mathcal{D}_{t12}^{T} \otimes \mathcal{I}_{s12}^{T} \otimes \mathcal{I}_{r12}^{T})(\frac{\partial t}{\partial x_{i}}. *W. *p)$$

$$(16)$$

In the case of no cross derivative terms, the above expression simplifies to:

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\frac{\partial r}{\partial x} * W. * p)$$
(17a)

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial s}{\partial y} * W. * p)$$
(17b)

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial z} * W. * p)$$
(17c)

which can be further expressed purely as kronecker products, taking $\partial r/\partial x$, $\partial s/\partial y$, $\partial t/\partial z$ *etc.* as diagonal matrices for the respective one dimensional problems.

$$w_1 = (\mathcal{I}_{t12}^T \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 W_{s2} \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial x} W_{r2}) p$$
(18a)

$$w_2 = (\mathcal{I}_{t12}^T \mathcal{I}_2 W_{t2} \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial y} W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 W_{r2}) p$$
(18b)

$$w_3 = (\mathcal{D}_{t12}^T \frac{\partial t}{\partial z} W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 W_{r2}) p$$
(18c)

For the cylindrical case we encounter some differences. There is an additional factor of R in front of the integral, i.e. $\partial\Omega_{x,y,z}\to R\partial\Omega_{x,R,\theta}\to R\mathcal{J}\partial\hat{\Omega}$. Also, for the θ term we have an additional division by R.

$$w_{x} = \int_{\Omega} p \left(\frac{\partial r}{\partial x} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial x} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial x} \frac{\partial v}{\partial t} \right) R d\hat{\Omega}$$

$$w_{R} = \int_{\Omega} p \left(\frac{\partial r}{\partial R} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial R} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial R} \frac{\partial v}{\partial t} \right) R d\hat{\Omega}$$

$$w_{\theta} = \int_{\Omega} p \left(\frac{\partial r}{\partial \theta} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial \theta} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial \theta} \frac{\partial v}{\partial t} \right) d\hat{\Omega}$$

Which will lead to the similar (but not identical) expressions for opgradt. If we assume the grid is Cartesian in the $x-R-\theta$ space, *i.e.* there are no cross derivatives.

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\frac{\partial r}{\partial x} \cdot *R. *W. *p)$$
(19a)

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial s}{\partial y} \cdot *R. *W. *p)$$
(19b)

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial z} * W. * p)$$
(19c)

The factors of R, W and the geometric factors can also be expressed in kronecker product form:

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\mathcal{I}_2 \otimes \mathcal{I}_2 \otimes \frac{\partial r}{\partial x})(\mathcal{I}_2 \otimes R_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p \quad (20a)$$

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 R_2 \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial x} \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (20b)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 R_2 W_{s2} \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial x} \mathcal{I}_2 W_{r2}) p$$
 (20c)

R-direction:

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\mathcal{I}_2 \otimes \frac{\partial s}{\partial R} \otimes \mathcal{I}_2)(\mathcal{I}_2 \otimes R_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p$$
 (21a)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial R} R_2 \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (21b)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 W_{t2} \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial R} R_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2 W_{r2}) p$$
 (21c)

 θ -direction:

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial \theta} \otimes \mathcal{I}_2 \otimes \mathcal{I}_2)(\mathcal{I}_2 \otimes \mathcal{I}_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p \quad (22a)$$

$$= (\mathcal{D}_{t12}^T \frac{\partial t}{\partial \theta} \mathcal{I}_2 \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (22b)

$$= (\mathcal{D}_{t12}^T \frac{\partial t}{\partial \theta} \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 \mathcal{I}_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2 W_{r2}) p$$
 (22c)

2.2 Opdiv

The opdiv subroutine implements the divergence operation for a vector field u, defined on Mesh 1 (velocity mesh), integrated w.r.t the Mesh 2 test functions. We can represent $\partial r/\partial x, \partial s/\partial x, \ldots$ as matrices $D_x r, D_x s, \ldots$. For a one dimensional case, $D_x r$ is a diagonal matrix. Therefore, in one dimension, $D_x = D_x r D_r$ and $D_x^T = D_r^T (D_x r)^T$ etc.

$$\int_{\Omega} q \nabla \cdot u = \int_{\Omega} q \frac{\partial}{\partial x_i} (u_i) d\Omega$$

$$= \int_{\Omega} q \frac{1}{\mathcal{J}} \frac{\partial r_j}{\partial x_i} \frac{\partial u_i}{\partial r_j} \mathcal{J} d\hat{\Omega}$$

$$= \sum_{k} q(x_k) W(x_k) \left(\frac{\partial r_j}{\partial x_i} \frac{\partial u_i}{\partial r_j} (x_k) \right)$$

$$= \sum_{k} q(x_k) W(x_k) (D_{x_i} r_j) (D_{r_j} u_i)$$

$$= q_k W_k (D_{x_i} r_j) (D_{r_j} u_i)$$

In the absence of cross geometric factors this becomes

$$\int_{\Omega} q \nabla \cdot u = q_k W_k (D_x r D_r u + D_y s D_s v + D_z t D_t w)$$

$$\int_{\Omega} q \nabla \cdot u = \begin{cases}
q W (\mathcal{I}_{t12} \otimes \mathcal{I}_{s12} \otimes D_x r D_r) u \\
+q W (\mathcal{I}_{t12} \otimes D_y s D_s \otimes \mathcal{I}_{r12}) u \\
+q W (D_z t D_t \otimes \mathcal{I}_{s12} \otimes \mathcal{I}_{r12}) u
\end{cases} \tag{23}$$

Here D_r , D_s , D_t is essential D_{r12} , D_{s12} , D_{t12} , since we are evaluating the derivative of a Mesh 1 field on Mesh 2 points.

2.3 Pressure Pseudo-Laplacian

$$S_{\Delta t} = DQD^T \tag{24}$$

where, if $Q=H^{-1}$ there is no decoupling error and we have the Uzawa algorithm. Alternately, $Q=B^{-1}$ in which case we incur a decoupling error but avoid the nested iterations since B being diagonal can be trivially inverted.

References