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# Navier-Stokes in Cylindrical formulation

## Weak form and Nek implementation

#### Formulation of Navier-Stokes 1

#### Mapping from Cartesian to Cylindrical 1.1

$$x_c = x (1a)$$

$$y_c = r\cos(\theta) \tag{1b}$$

$$z_c = r\sin(\theta) \tag{1c}$$

$$\begin{cases} dx_c \\ dy_c \\ dz_c \end{cases} = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos(\theta) & -r\sin(\theta) \\ 0 & \sin(\theta) & r\cos(\theta) \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$
 (2)

$$\begin{cases}
 dx_c \\
 dy_c \\
 dz_c
\end{cases} = \begin{pmatrix}
 1 & 0 & 0 \\
 0 & \cos(\theta) & -r\sin(\theta) \\
 0 & \sin(\theta) & r\cos(\theta)
\end{pmatrix} \begin{cases}
 dx \\
 dr \\
 d\theta
\end{cases}$$

$$\Rightarrow \begin{cases}
 dx \\
 dr \\
 d\theta
\end{cases} = \begin{pmatrix}
 1 & 0 & 0 \\
 0 & \cos(\theta) & \sin(\theta) \\
 0 & -\sin(\theta)/r & \cos(\theta)/r
\end{pmatrix} \begin{cases}
 dx_c \\
 dy_c \\
 dz_c
\end{cases}$$
(2)

$$\hat{x} = \hat{x_c} \tag{4a}$$

$$\hat{r} = \cos(\theta)\hat{y_c} + \sin(\theta)\hat{z_c} \tag{4b}$$

$$\hat{\theta} = -\sin(\theta)\hat{y_c} + \cos(\theta)\hat{z_c} \tag{4c}$$

$$\Rightarrow \begin{cases} d\hat{x} \\ d\hat{r} \\ d\hat{\theta} \end{cases} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & -\sin(\theta)\hat{y}_c + \cos(\theta)\hat{z}_c \\ 0 & 0 & -\cos(\theta)\hat{y}_c - \sin(\theta)\hat{z}_c \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$

$$\Rightarrow \begin{cases} d\hat{x} \\ d\hat{r} \\ d\hat{\theta} \end{cases} = \begin{pmatrix} 0 & 0 & 0 \\ 0 & 0 & \hat{\theta} \\ 0 & 0 & -\hat{r} \end{pmatrix} \begin{cases} dx \\ dr \\ d\theta \end{cases}$$
(5)

### 1.2 Gradient/Divergence/Laplace Operators

For scalar fields

$$\nabla(\psi) = \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial x_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial x_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial x_c}\right) \hat{x_c} + \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial y_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial y_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial y_c}\right) \hat{y_c} + \left(\frac{\partial \psi}{\partial x} \frac{\partial x}{\partial z_c} + \frac{\partial \psi}{\partial r} \frac{\partial r}{\partial z_c} + \frac{\partial \psi}{\partial \theta} \frac{\partial \theta}{\partial z_c}\right) \hat{z_c}$$

$$\therefore \nabla(\psi) = \frac{\partial \psi}{\partial x}\hat{x} + \frac{\partial \psi}{\partial r}\hat{r} + \frac{1}{r}\frac{\partial \psi}{\partial \theta}\hat{\theta}$$
 (6)

Laplacian of a scalar field

$$\nabla \cdot \nabla(\psi) = \left[ \frac{\partial}{\partial x} \hat{x} + \frac{\partial}{\partial r} \hat{r} + \frac{1}{r} \frac{\partial}{\partial \theta} \hat{\theta} \right] \cdot \left[ \hat{x} \frac{\partial \psi}{\partial x} + \hat{r} \frac{\partial \psi}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial \psi}{\partial \theta} \right]$$

$$= \begin{cases} \left[ \hat{x} \cdot \hat{x} \frac{\partial}{\partial x} \frac{\partial \psi}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial x} \cdot \hat{x} \right] + \left[ \hat{x} \cdot \hat{r} \frac{\partial}{\partial x} \frac{\partial \psi}{\partial r} + \frac{\partial \psi}{\partial r} \frac{\partial \hat{r}}{\partial x} \cdot \hat{x} \right] + \left[ \hat{x} \cdot \hat{\theta} \frac{\partial}{\partial x} \left( \frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial x} \cdot \hat{x} \right] \\ + \left[ \hat{r} \cdot \hat{x} \frac{\partial}{\partial r} \frac{\partial \psi}{\partial x} + \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial r} \cdot \hat{r} \right] + \left[ \hat{r} \cdot \hat{r} \frac{\partial}{\partial r} \frac{\partial \psi}{\partial r} + \frac{\partial \psi}{\partial r} \frac{\partial \hat{r}}{\partial r} \cdot \hat{r} \right] + \left[ \hat{r} \cdot \hat{\theta} \frac{\partial}{\partial r} \left( \frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial r} \cdot \hat{r} \right] \\ + \left[ \hat{\theta} \cdot \hat{x} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial \psi}{\partial x} + \frac{1}{r} \frac{\partial \psi}{\partial x} \frac{\partial \hat{x}}{\partial \theta} \cdot \hat{\theta} \right] + \left[ \hat{\theta} \cdot \hat{r} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial \psi}{\partial r} + \frac{1}{r} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{r}}{\partial \theta} \cdot \hat{\theta} \right] \\ + \left[ \hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left( \frac{1}{r} \frac{\partial \psi}{\partial \theta} \right) + \frac{1}{r^2} \frac{\partial \psi}{\partial \theta} \frac{\partial \hat{\theta}}{\partial \theta} \cdot \hat{\theta} \right]$$

$$\therefore \nabla^2(\psi) = \frac{\partial^2 \psi}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial \psi}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 \psi}{\partial \theta} \tag{7}$$

Vector gradient. Highlighted in color are the only terms which are non-zero dif-

ferentials of unit vectors.

$$\nabla(\boldsymbol{u}) = \left[ \hat{x} \frac{\partial}{\partial x} + \hat{r} \frac{\partial}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \right] \left[ u_x \hat{x} + u_r \hat{r} + u_\theta \hat{\theta} \right]$$

$$= \begin{cases} \left[ \hat{x} \frac{\partial u_x}{\partial x} \hat{x} + \hat{x} u_x \frac{\partial \hat{x}}{\partial x} \right] + \left[ \hat{x} \frac{\partial u_r}{\partial x} \hat{r} + \hat{x} u_r \frac{\partial \hat{r}}{\partial x} \right] + \left[ \hat{x} \frac{\partial u_\theta}{\partial x} \hat{\theta} + \hat{x} u_\theta \frac{\partial \hat{\theta}}{\partial x} \right] \\ + \left[ \hat{r} \frac{\partial u_x}{\partial r} \hat{x} + \hat{r} u_x \frac{\partial \hat{x}}{\partial r} \right] + \left[ \hat{r} \frac{\partial u_r}{\partial r} \hat{r} + \hat{r} u_r \frac{\partial \hat{r}}{\partial r} \right] + \left[ \hat{r} \frac{\partial u_\theta}{\partial r} \hat{\theta} + \hat{r} u_\theta \frac{\partial \hat{\theta}}{\partial r} \right] \\ + \left[ \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x} + \hat{\theta} \frac{u_x}{r} \frac{\partial \hat{x}}{\partial \theta} \right] + \left[ \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \hat{r} + \hat{\theta} \frac{u_r}{r} \frac{\partial \hat{r}}{\partial \theta} \right] + \left[ \hat{\theta} \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \hat{\theta} + \hat{\theta} \frac{u_\theta}{r} \frac{\partial \hat{\theta}}{\partial \theta} \right] \\ + \left[ \hat{r} \frac{\partial u_x}{\partial x} \hat{x} + \hat{x} \frac{\partial u_r}{\partial r} \hat{r} + \hat{x} \frac{\partial u_\theta}{\partial r} \hat{\theta} \right] \\ + \left[ \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x} + \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \hat{r} + \hat{\theta} \frac{u_r}{r} \hat{\theta} + \hat{\theta} \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \hat{\theta} - \hat{\theta} \frac{u_\theta}{r} \hat{r} \right] \end{cases}$$
(9)

Which, I rearrange as gradients of individual velocity components (and their respective unit vectors)

$$\nabla(u_x \hat{x}) = \hat{x} \frac{\partial u_x}{\partial x} \hat{x} + \hat{r} \frac{\partial u_x}{\partial r} \hat{x} + \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x}$$
(10a)

$$\nabla(u_r \hat{r}) = \hat{x} \frac{\partial u_r}{\partial x} \hat{r} + \hat{r} \frac{\partial u_r}{\partial r} \hat{r} + \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \hat{r} + \hat{\theta} \frac{u_r}{r} \hat{\theta}$$
(10b)

$$\nabla(u_{\theta}\hat{\theta}) = \hat{x}\frac{\partial u_{\theta}}{\partial x}\hat{\theta} + \hat{r}\frac{\partial u_{\theta}}{\partial r}\hat{\theta} + \hat{\theta}\frac{1}{r}\frac{\partial u_{\theta}}{\partial \theta}\hat{\theta} - \hat{\theta}\frac{u_{\theta}}{r}\hat{r}$$
(10c)

$$\therefore \nabla(\boldsymbol{u}) = \begin{bmatrix} \hat{x}\frac{\partial u_x}{\partial x}\hat{x} & \hat{x}\frac{\partial u_r}{\partial x}\hat{r} & \hat{x}\frac{\partial u_\theta}{\partial x}\hat{\theta} \\ \hat{r}\frac{\partial u_x}{\partial r}\hat{x} & \hat{r}\frac{\partial u_r}{\partial r}\hat{r} & \hat{r}\frac{\partial u_\theta}{\partial r}\hat{\theta} \\ \hat{\theta}\frac{1}{r}\frac{\partial u_x}{\partial \theta}\hat{x} & \hat{\theta}\left(-\frac{u_\theta}{r} + \frac{1}{r}\frac{\partial u_r}{\partial \theta}\right)\hat{r} & \hat{\theta}\left(\frac{u_r}{r} + \frac{1}{r}\frac{\partial u_\theta}{\partial \theta}\right)\hat{\theta} \end{bmatrix} \tag{11}$$

Vector Laplacian:

$$\nabla \cdot \nabla (\boldsymbol{u}) = \left( \hat{\boldsymbol{x}} \cdot \frac{\partial}{\partial \boldsymbol{x}} + \hat{\boldsymbol{r}} \cdot \frac{\partial}{\partial r} + \hat{\boldsymbol{\theta}} \cdot \frac{1}{r} \frac{\partial}{\partial \theta} \right) \begin{pmatrix} \hat{\boldsymbol{x}} \frac{\partial u_x}{\partial \boldsymbol{x}} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{x}} \frac{\partial u_r}{\partial \boldsymbol{x}} \hat{\boldsymbol{r}} + & \hat{\boldsymbol{x}} \frac{\partial u_\theta}{\partial \boldsymbol{x}} \hat{\boldsymbol{\theta}} + \\ \hat{\boldsymbol{r}} \frac{\partial u_x}{\partial r} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{r}} \frac{\partial u_r}{\partial r} \hat{\boldsymbol{r}} + & \hat{\boldsymbol{r}} \frac{\partial u_\theta}{\partial r} \hat{\boldsymbol{\theta}} + \\ \hat{\boldsymbol{\theta}} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{\boldsymbol{x}} + & \hat{\boldsymbol{\theta}} \left( -\frac{u_\theta}{r} + \frac{1}{r} \frac{\partial u_r}{\partial \theta} \right) \hat{\boldsymbol{r}} + & \hat{\boldsymbol{\theta}} \left( \frac{u_r}{r} + \frac{1}{r} \frac{\partial u_\theta}{\partial \theta} \right) \hat{\boldsymbol{\theta}} \end{pmatrix}$$

First Column

$$\left(\hat{x} \cdot \frac{\partial}{\partial x} + \hat{r} \cdot \frac{\partial}{\partial r} + \hat{\theta} \cdot \frac{1}{r} \frac{\partial}{\partial \theta}\right) \left(\hat{x} \frac{\partial u_x}{\partial x} + \hat{r} \frac{\partial u_x}{\partial r} + \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x}$$

$$= \hat{x} \cdot \left(\hat{x} \frac{\partial^2 u_x}{\partial x^2} + \frac{\partial \hat{x}}{\partial x} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{x} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{\partial \hat{x}}{\partial x}$$

$$+ \hat{x} \cdot \left(\hat{r} \frac{\partial}{\partial x} \frac{\partial u_x}{\partial r} + \frac{\partial \hat{r}}{\partial x} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{x} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{x} \cdot \left(\hat{\theta} \frac{\partial}{\partial x} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) + \frac{\partial \hat{\theta}}{\partial x} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{x} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \frac{\partial \hat{x}}{\partial x}$$

$$+ \hat{r} \cdot \left(\hat{x} \frac{\partial}{\partial r} \frac{\partial u_x}{\partial x} + \frac{\partial \hat{x}}{\partial r} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{r} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial^2 u_x}{\partial r^2} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{r} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{\theta} \frac{\partial}{\partial r} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) + \frac{\partial \hat{\theta}}{\partial r} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \frac{\partial \hat{x}}{\partial r}$$

$$+ \hat{\theta} \cdot \left(\hat{x} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial u_x}{\partial x} + \frac{\partial \hat{x}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial x}\right) \hat{x} + \hat{\theta} \cdot \hat{x} \frac{\partial u_x}{\partial x} \frac{1}{r} \frac{\partial \hat{x}}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{r} \frac{1}{r} \frac{\partial}{\partial \theta} \frac{\partial u_x}{\partial r} + \frac{\partial \hat{r}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial r}\right) \hat{x} + \hat{\theta} \cdot \hat{r} \frac{\partial u_x}{\partial r} \frac{1}{r} \frac{\partial \hat{x}}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \frac{\partial \hat{\theta}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}$$

$$+ \hat{\theta} \cdot \left(\hat{\theta} \frac{1}{r} \frac{\partial}{\partial \theta} \left(\frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \frac{\partial \hat{\theta}}{\partial \theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}\right) \hat{x} + \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta}$$

$$= \left[\frac{\partial^2 u_x}{\partial x^2} + \frac{\partial^2 u_x}{\partial r^2} + \frac{1}{r} \frac{\partial^2 u_x}{\partial r} + \frac{1}{r^2} \frac{\partial^2 u_x}{\partial \theta^2}\right] \hat{x}$$
(12)

Second column:

$$\begin{pmatrix}
\hat{x} \cdot \frac{\partial}{\partial x} + \hat{r} \cdot \frac{\partial}{\partial r} + \hat{\theta} \cdot \frac{1}{r} \frac{\partial}{\partial \theta}
\end{pmatrix} \begin{pmatrix}
\hat{x} \frac{\partial u_r}{\partial x} + \hat{r} \frac{\partial u_r}{\partial r} - \hat{\theta} \frac{u_\theta}{r} + \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta}
\end{pmatrix} \hat{r}$$

$$= \hat{x} \cdot \left(\hat{x} \frac{\partial^2 u_r}{\partial x^2} + \frac{\partial \hat{x}}{\partial x} \frac{\partial u_r}{\partial x}\right) \hat{r} + \hat{x} \cdot \hat{x} \frac{\partial u_r}{\partial x} \frac{\partial \hat{r}}{\partial x}$$

$$+ \hat{x} \cdot \left(\hat{r} \frac{\partial}{\partial x} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial x} \frac{\partial u_r}{\partial r}\right) \hat{r} + \hat{x} \cdot \hat{r} \frac{\partial u_r}{\partial r} \frac{\partial \hat{r}}{\partial x}$$

$$+ \hat{x} \cdot \left(-\hat{\theta} \frac{\partial}{\partial x} \frac{u_\theta}{r} - \frac{\partial \hat{\theta}}{\partial x} \frac{u_\theta}{r}\right) \hat{r} - \hat{x} \cdot \hat{\theta} \frac{u_\theta}{r} \frac{\partial \hat{r}}{\partial x}$$

$$+ \hat{x} \cdot \left(\hat{\theta} \frac{\partial}{\partial x} \left(\frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) + \frac{\partial \hat{\theta}}{\partial x} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{x} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \frac{\partial \hat{r}}{\partial x}$$

$$+ \hat{r} \cdot \left(\hat{x} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial x} + \frac{\partial \hat{x}}{\partial r} \frac{\partial u_r}{\partial x}\right) \hat{r} + \hat{r} \cdot \hat{x} \frac{\partial u_r}{\partial x} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial^2 u_r}{\partial r^2} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_r}{\partial x}\right) \hat{r} + \hat{r} \cdot \hat{r} \frac{\partial u_r}{\partial x} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial x} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_r}{\partial r}\right) \hat{r} + \hat{r} \cdot \hat{r} \frac{\partial u_r}{\partial x} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial x} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_r}{\partial r}\right) \hat{r} + \hat{r} \cdot \hat{r} \frac{\partial u_r}{\partial r} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial x} + \frac{\partial \hat{r}}{\partial r} \frac{\partial u_r}{\partial r}\right) \hat{r} + \hat{r} \cdot \hat{r} \frac{\partial u_r}{\partial r} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial r} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial \theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial \theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \frac{\partial \hat{r}}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial \theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{r} \cdot \hat{\theta} \frac{1}{r} \frac{\partial u_r}{\partial \theta} \frac{\partial r}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}{\partial r} \frac{1}{r} \frac{\partial u_r}{\partial \theta}\right) \hat{r} + \hat{r} \cdot \hat{\theta} \hat{r} \frac{\partial u_r}{\partial r} \frac{\partial r}{\partial r}$$

$$+ \hat{r} \cdot \left(\hat{r} \frac{\partial}{\partial r} \frac{\partial u_r}{\partial r} + \frac{\partial \hat{r}}$$

Third column:

$$\therefore \nabla \cdot \nabla (\boldsymbol{u}) = \begin{cases}
\left[ \frac{\partial^2 u_x}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_x}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_x}{\partial \theta^2} \right] \hat{x} + \\
\left[ \frac{\partial^2 u_r}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_r}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_r}{\partial \theta^2} - \frac{u_r}{r^2} - \frac{2}{r^2} \frac{\partial u_\theta}{\partial \theta} \right] \hat{r} + \\
\left[ \frac{\partial^2 u_\theta}{\partial x^2} + \frac{1}{r} \frac{\partial}{\partial r} \left( r \frac{\partial u_\theta}{\partial r} \right) + \frac{1}{r^2} \frac{\partial^2 u_\theta}{\partial \theta^2} + \frac{2}{r^2} \frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r^2} \right] \hat{\theta}
\end{cases} \tag{15}$$

Which can be written using the Laplacian of individual scalar fields,  $u_x$ ,  $u_r$ ,  $u_\theta$  as,

$$\therefore \nabla \cdot \nabla(\boldsymbol{u}) = \nabla^2(u_x)\hat{x} + \left(\nabla^2(u_r) - \frac{u_r}{r^2} - \frac{2}{r^2}\frac{\partial u_\theta}{\partial \theta}\right)\hat{r} + \left(\nabla^2(u_\theta) + \frac{2}{r^2}\frac{\partial u_r}{\partial \theta} - \frac{u_\theta}{r^2}\right)\hat{\theta}$$
(16)

Weak Laplacian:

$$\nabla(v_x \hat{x}) \cdot \nabla(u_x \hat{x}) = \left(\hat{x} \frac{\partial v_x}{\partial x} \hat{x} + \hat{r} \frac{\partial v_x}{\partial r} \hat{x} + \hat{\theta} \frac{1}{r} \frac{\partial v_x}{\partial \theta} \hat{x}\right) \cdot \left(\hat{x} \frac{\partial u_x}{\partial x} \hat{x} + \hat{r} \frac{\partial u_x}{\partial r} \hat{x} + \hat{\theta} \frac{1}{r} \frac{\partial u_x}{\partial \theta} \hat{x}\right)$$

$$= \hat{x} \frac{\partial v_x}{\partial x} \frac{\partial u_x}{\partial x} \hat{x} + \hat{x} \frac{\partial v_x}{\partial r} \frac{\partial u_x}{\partial r} \hat{x} + \hat{x} \frac{1}{r^2} \frac{\partial v_x}{\partial \theta} \frac{\partial u_x}{\partial \theta} \hat{x}$$

$$(17)$$

$$\nabla(v_{r}\hat{r}) \cdot \nabla(u_{r}\hat{r}) = \left(\hat{x}\frac{\partial v_{r}}{\partial x}\hat{r} + \hat{r}\frac{\partial v_{r}}{\partial r}\hat{r} + \hat{\theta}\frac{1}{r}\frac{\partial v_{r}}{\partial \theta}\hat{r} + \hat{\theta}\frac{v_{r}}{r}\hat{\theta}\right) \cdot \left(\hat{x}\frac{\partial u_{r}}{\partial x}\hat{r} + \hat{r}\frac{\partial u_{r}}{\partial r}\hat{r} + \hat{\theta}\frac{1}{r}\frac{\partial u_{r}}{\partial \theta}\hat{r} + \hat{\theta}\frac{u_{r}}{r}\hat{\theta}\right)$$

$$= \begin{cases} \hat{r}\frac{\partial v_{r}}{\partial x}\frac{\partial u_{r}}{\partial x}\hat{r} + \hat{r}\frac{\partial v_{r}}{\partial r}\frac{\partial u_{r}}{\partial r}\hat{r} + \hat{r}\frac{1}{r^{2}}\frac{\partial v_{r}}{\partial \theta}\frac{\partial u_{r}}{\partial \theta}\hat{r} + \\ \hat{\theta}\frac{v_{r}}{r}\frac{u_{r}}{r}\hat{\theta} + \hat{r}\frac{1}{r}\frac{\partial v_{r}}{\partial \theta}\frac{u_{r}}{r}\hat{\theta} + \hat{\theta}\frac{v_{r}}{r}\frac{1}{\partial \theta}\frac{\partial u_{r}}{\partial \theta}\hat{r} \end{cases}$$

$$(18)$$

$$\nabla(v_{\theta}\hat{\theta}) \cdot \nabla(u_{\theta}\hat{\theta}) = \left(\hat{x}\frac{\partial v_{\theta}}{\partial x}\hat{\theta} + \hat{r}\frac{\partial v_{\theta}}{\partial r}\hat{\theta} + \hat{\theta}\frac{1}{r}\frac{\partial v_{\theta}}{\partial \theta}\hat{\theta} - \hat{\theta}\frac{v_{\theta}}{r}\hat{r}\right) \cdot \left(\hat{x}\frac{\partial u_{\theta}}{\partial x}\hat{\theta} + \hat{r}\frac{\partial u_{\theta}}{\partial r}\hat{\theta} + \hat{\theta}\frac{1}{r}\frac{\partial u_{\theta}}{\partial \theta}\hat{\theta} - \hat{\theta}\frac{u_{\theta}}{r}\hat{r}\right)$$

$$= \begin{cases} \hat{\theta}\frac{\partial v_{\theta}}{\partial x}\frac{\partial u_{\theta}}{\partial x}\hat{\theta} + \hat{\theta}\frac{\partial v_{\theta}}{\partial r}\frac{\partial u_{\theta}}{\partial r}\hat{\theta} + \hat{\theta}\frac{1}{r^2}\frac{\partial v_{\theta}}{\partial \theta}\frac{\partial u_{\theta}}{\partial \theta}\hat{\theta} \\ +\hat{r}\frac{v_{\theta}}{r}\frac{u_{\theta}}{r}\hat{r} - \hat{\theta}\frac{1}{r}\frac{\partial v_{\theta}}{\partial \theta}\frac{u_{\theta}}{r}\hat{r} - \hat{r}\frac{v_{\theta}}{r}\frac{1}{r}\frac{\partial u_{\theta}}{\partial \theta}\hat{\theta} \end{cases}$$
(19)

## 2 Nek Implementations:

## 2.1 Opgradt

The opgradt subroutine implements the gradient of a field defined on Mesh 2 (pressure mesh), integrated w.r.t the Mesh 1 test functions. This is done with integration by parts which creates a boundary term.

Pressure term in the momentum equations:

$$-\int_{\Omega} v \frac{\partial}{\partial x_{j}} (p \delta_{ij}) d\Omega$$
$$-\int_{\Omega} \frac{\partial}{\partial x_{j}} (v p \delta_{ij}) d\Omega + \int_{\Omega} p \delta_{ij} \frac{\partial}{\partial x_{j}} (v) d\Omega$$
$$-\int_{\partial \Omega} (v p \delta_{ij}) n_{j} dA + \int_{\Omega} p \delta_{ij} \frac{\partial}{\partial x_{j}} (v) d\Omega$$

The first term is the boundary condition. The second term is what opgradt evaluates. I represent the determinant of the mapping between reference element and physical coordinates, the Jacobian as  $\mathcal{J}$ . In the continuous formulation, this becomes:

$$w_{i} = \int_{\Omega} p \frac{\partial v}{\partial x_{i}} d\Omega$$

$$= \int_{\Omega} p \frac{\partial v}{\partial x_{i}} \left(\frac{\partial \Omega}{\partial \hat{\Omega}}\right) d\hat{\Omega}$$

$$= \int_{\Omega} p \left(\frac{1}{\mathcal{J}} \frac{\partial r}{\partial x_{i}} \frac{\partial v}{\partial r} + \frac{1}{\mathcal{J}} \frac{\partial s}{\partial x_{i}} \frac{\partial v}{\partial s} + \frac{1}{\mathcal{J}} \frac{\partial t}{\partial x_{i}} \frac{\partial v}{\partial t}\right) \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} p \left(\frac{\partial r}{\partial x_{i}} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial x_{i}} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial x_{i}} \frac{\partial v}{\partial t}\right) d\hat{\Omega}$$

$$\implies w_{i} = \int_{\Omega} p \left(\frac{\partial r_{j}}{\partial x_{i}} \frac{\partial v}{\partial r_{j}}\right) d\hat{\Omega}, \tag{20}$$

where,  $r_j$  represents the reference element coordinate directions r, s, t. After discretization, this becomes:

$$\implies w_i = \sum_k W(x_k) \left( \frac{\partial r_j}{\partial x_i} \frac{\partial v}{\partial r_j} (x_k) \right) p$$

When cross derivative terms like  $\partial s/\partial x$  *etc.* are non-zero, we have the following expression

$$w_{i} = (\mathcal{I}_{t12}^{T} \otimes \mathcal{I}_{s12}^{T} \otimes \mathcal{D}_{r12}^{T})(\frac{\partial r}{\partial x_{i}}. *W. *p) +$$

$$(\mathcal{I}_{t12}^{T} \otimes \mathcal{D}_{s12}^{T} \otimes \mathcal{I}_{r12}^{T})(\frac{\partial s}{\partial x_{i}}. *W. *p) +$$

$$(\mathcal{D}_{t12}^{T} \otimes \mathcal{I}_{s12}^{T} \otimes \mathcal{I}_{r12}^{T})(\frac{\partial t}{\partial x_{i}}. *W. *p)$$

$$(21)$$

In the case of no cross derivative terms, the above expression simplifies to:

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\frac{\partial r}{\partial x} * W * p)$$
 (22a)

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial s}{\partial y} * W * y)$$
(22b)

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial z} * W * p)$$
(22c)

which can be further expressed purely as kronecker products, taking  $\partial r/\partial x$ ,  $\partial s/\partial y$ ,  $\partial t/\partial z$  etc. as diagonal matrices for the respective one dimensional problems.

$$w_1 = (\mathcal{I}_{t12}^T \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 W_{s2} \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial x} W_{r2}) p$$
 (23a)

$$w_2 = (\mathcal{I}_{t12}^T \mathcal{I}_2 W_{t2} \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial u} W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 W_{r2}) p$$
(23b)

$$w_3 = (\mathcal{D}_{t12}^T \frac{\partial t}{\partial z} W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 W_{r2}) p$$
 (23c)

For the cylindrical case we encounter some differences. There is an additional factor of R in front of the integral, i.e.  $\partial\Omega_{x,y,z} \to R\partial\Omega_{x,R,\theta} \to R\mathcal{J}\partial\hat{\Omega}$ . Also, for the  $\theta$  term we have an additional division by R.

$$w_{x} = \int_{\Omega} p \left( \frac{\partial r}{\partial x} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial x} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial x} \frac{\partial v}{\partial t} \right) R d\hat{\Omega}$$

$$w_{R} = \int_{\Omega} p \left( \frac{\partial r}{\partial R} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial R} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial R} \frac{\partial v}{\partial t} \right) R d\hat{\Omega}$$

$$w_{\theta} = \int_{\Omega} p \left( \frac{\partial r}{\partial \theta} \frac{\partial v}{\partial r} + \frac{\partial s}{\partial \theta} \frac{\partial v}{\partial s} + \frac{\partial t}{\partial \theta} \frac{\partial v}{\partial t} \right) d\hat{\Omega}$$

Which will lead to the similar (but not identical) expressions for opgradt. If we assume the grid is Cartesian in the  $x - R - \theta$  space, *i.e.* there are no cross derivatives.

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\frac{\partial r}{\partial x} \cdot *R. *W. *p)$$
(24a)

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial s}{\partial y} \cdot *R. *W. *p)$$
(24b)

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial z} * W * p)$$
(24c)

The factors of R, W and the geometric factors can also be expressed in kronecker product form:

$$w_1 = (\mathcal{I}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{D}_{r12}^T)(\mathcal{I}_2 \otimes \mathcal{I}_2 \otimes \frac{\partial r}{\partial r})(\mathcal{I}_2 \otimes R_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p$$
 (25a)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 R_2 \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial r} \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (25b)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 R_2 W_{s2} \otimes \mathcal{D}_{r12}^T \frac{\partial r}{\partial x} \mathcal{I}_2 W_{r2}) p$$
 (25c)

*R*-direction:

$$w_2 = (\mathcal{I}_{t12}^T \otimes \mathcal{D}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\mathcal{I}_2 \otimes \frac{\partial s}{\partial R} \otimes \mathcal{I}_2)(\mathcal{I}_2 \otimes R_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p$$
 (26a)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial R} R_2 \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (26b)

$$= (\mathcal{I}_{t12}^T \mathcal{I}_2 \mathcal{I}_2 W_{t2} \otimes \mathcal{D}_{s12}^T \frac{\partial s}{\partial R} R_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2 W_{r2}) p$$
 (26c)

 $\theta$ -direction:

$$w_3 = (\mathcal{D}_{t12}^T \otimes \mathcal{I}_{s12}^T \otimes \mathcal{I}_{r12}^T)(\frac{\partial t}{\partial \theta} \otimes \mathcal{I}_2 \otimes \mathcal{I}_2)(\mathcal{I}_2 \otimes \mathcal{I}_2 \otimes \mathcal{I}_2)(W_{t2} \otimes W_{s2} \otimes W_{r2})p$$
 (27a)

$$= (\mathcal{D}_{t12}^T \frac{\partial t}{\partial \theta} \mathcal{I}_2 \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 \mathcal{I}_2 \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2) (W_{t2} \otimes W_{s2} \otimes W_{r2}) p$$
 (27b)

$$= (\mathcal{D}_{t12}^T \frac{\partial t}{\partial \theta} \mathcal{I}_2 W_{t2} \otimes \mathcal{I}_{s12}^T \mathcal{I}_2 \mathcal{I}_2 W_{s2} \otimes \mathcal{I}_{r12}^T \mathcal{I}_2 \mathcal{I}_2 W_{r2}) p$$
 (27c)

### 2.2 Opdiv

The opdiv subroutine implements the divergence operation for a vector field u, defined on Mesh 1 (velocity mesh), integrated w.r.t the Mesh 2 test functions. We can represent  $\partial r/\partial x, \partial s/\partial x, \ldots$  as matrices  $D_x r, D_x s, \ldots$ . For a one dimensional case,  $D_x r$  is a diagonal matrix. Therefore, in one dimension,  $D_x = D_x r D_r$  and  $D_x^T = D_r^T (D_x r)^T$  etc.

$$\int_{\Omega} q \nabla \cdot u = \int_{\Omega} q \frac{\partial}{\partial x_{i}}(u_{i}) d\Omega$$

$$= \int_{\Omega} q \frac{1}{\mathcal{J}} \frac{\partial r_{j}}{\partial x_{i}} \frac{\partial u_{i}}{\partial r_{j}} \mathcal{J} d\hat{\Omega}$$

$$= \sum_{k} q(x_{k}) W(x_{k}) \left( \frac{\partial r_{j}}{\partial x_{i}} \frac{\partial u_{i}}{\partial r_{j}}(x_{k}) \right)$$

$$= \sum_{k} q(x_{k}) W(x_{k}) (D_{x_{i}} r_{j}) (D_{r_{j}} u_{i})$$

$$= q_{k} W_{k} (D_{x_{i}} r_{j}) (D_{r_{j}} u_{i})$$

In the absence of cross geometric factors this becomes

$$\int_{\Omega} q \nabla \cdot u = q_k W_k (D_x r D_r u + D_y s D_s v + D_z t D_t w)$$

$$\int_{\Omega} q \nabla \cdot u = \begin{cases}
q W (\mathcal{I}_{t12} \otimes \mathcal{I}_{s12} \otimes D_x r D_r) u \\
+q W (\mathcal{I}_{t12} \otimes D_y s D_s \otimes \mathcal{I}_{r12}) v \\
+q W (D_z t D_t \otimes \mathcal{I}_{s12} \otimes \mathcal{I}_{r12}) w
\end{cases}$$

$$\int_{\Omega} q \nabla \cdot u = \begin{cases}
q (W_t \mathcal{I}_{t12} \otimes W_s \mathcal{I}_{s12} \otimes W_r D_x r D_r) u \\
+q (W_t \mathcal{I}_{t12} \otimes W_s D_y s D_s \otimes W_r \mathcal{I}_{r12}) v \\
+q (W_t D_z t D_t \otimes W_s \mathcal{I}_{s12} \otimes W_r \mathcal{I}_{r12}) w
\end{cases}$$
(29)

Here  $D_r$ ,  $D_s$ ,  $D_t$  is essential  $D_{r12}$ ,  $D_{s12}$ ,  $D_{t12}$ , since we are evaluating the derivative of a Mesh 1 field on Mesh 2 points.

For the cylindrical case we have

$$\int_{\Omega} q \overline{\partial}_{x_{i}}(u_{i}) d\Omega$$

$$= \int_{\Omega} q \frac{\partial u_{i}}{\partial x_{i}} R \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} q \left(\frac{\partial u}{\partial x} + \frac{1}{R} \frac{\partial (Rv)}{\partial R} + \frac{1}{R} \frac{\partial w}{\partial \theta}\right) R \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} q \left(\frac{\partial u}{\partial x} + \frac{\partial v}{\partial R} + \frac{v}{R} + \frac{1}{R} \frac{\partial w}{\partial \theta}\right) R \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} q \left(\frac{1}{\mathcal{J}} \frac{\partial r_{j}}{\partial x} \frac{\partial u}{\partial r_{j}} + \frac{1}{\mathcal{J}} \frac{\partial r_{j}}{\partial R} \frac{\partial v}{\partial r_{j}} + \frac{v}{R} + \frac{1}{R} \frac{1}{\mathcal{J}} \frac{\partial r_{j}}{\partial \theta} \frac{\partial w}{\partial r_{j}}\right) R \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} q \left(\frac{\partial r_{j}}{\partial x} \frac{\partial u}{\partial r_{j}} + \frac{\partial r_{j}}{\partial R} \frac{\partial v}{\partial r_{j}} + \frac{\mathcal{J}v}{R} + \frac{1}{R} \frac{\partial r_{j}}{\partial \theta} \frac{\partial w}{\partial r_{j}}\right) R \mathcal{J} d\hat{\Omega}$$

$$= \int_{\Omega} q \left(\frac{\partial r_{j}}{\partial x} \frac{\partial u}{\partial r_{j}} + \frac{\partial r_{j}}{\partial R} \frac{\partial v}{\partial r_{j}} + \frac{\mathcal{J}v}{R} + \frac{1}{R} \frac{\partial r_{j}}{\partial \theta} \frac{\partial w}{\partial r_{j}}\right) R d\hat{\Omega}$$

$$= W_{k} q_{k} \left(R \frac{\partial r_{j}}{\partial x} \frac{\partial u}{\partial r_{j}} + R \frac{\partial r_{j}}{\partial R} \frac{\partial v}{\partial r_{j}} + \frac{R \mathcal{J}v}{R} + \frac{\partial r_{j}}{\partial \theta} \frac{\partial w}{\partial r_{j}}\right) \tag{30}$$

For undeformed elements,  $\mathcal J$  is a constant throughout the element. Therefore in kronecker notation we have

$$\int_{\Omega} q \nabla \cdot u = \begin{cases}
Wq(\mathcal{I}_{12} \otimes R \otimes D_{x12}rD_r)u \\
+Wq(\mathcal{I}_{12} \otimes (RD_{R12}sD_s + \mathcal{J}\mathcal{I}_{12}) \otimes \mathcal{I}_{12})v \\
+Wq(D_{\theta 12}tD_t \otimes \mathcal{I}_{12} \otimes \mathcal{I}_{12})w
\end{cases}$$
(31)

$$\int_{\Omega} q \nabla \cdot u = \begin{cases}
q(W_t \mathcal{I}_{12} \otimes W_s R \otimes W_r D_{x12} r D_r) u \\
+q(W_t \mathcal{I}_{12} \otimes W_s (R D_{R12} s D_s + \mathcal{J} \mathcal{I}_{12}) \otimes W_r \mathcal{I}_{12}) v \\
+q(W_t D_{\theta 12} t D_t \otimes W_s \mathcal{I}_{12} \otimes W_r \mathcal{I}_{12}) w
\end{cases}$$
(32)

### 2.3 Pressure Pseudo-Laplacian

$$S_{\Delta t} = DQD^T \tag{33}$$

where, if  $Q = H^{-1}$  there is no decoupling error and we have the Uzawa algorithm. Alternately,  $Q = B^{-1}$  in which case we incur a decoupling error but avoid the nested iterations since B, being diagonal, can be trivially inverted.