

OSCILLATORY INTEGRALS IN ONE DIMENSION

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In this note we give an overview of the theory of oscillatory integrals in one dimension, which gives an essentially complete description of the behavior of integrals of the form

$$I(\lambda) = \int_a^b e^{i\lambda\Phi(x)} \Psi(x) dx, \quad \lambda > 0,$$

as $\lambda \rightarrow \infty$, where Φ and Ψ are smooth functions. The behavior of $I(\lambda)$ is governed by three basic principles: *localization*, *scaling*, and *asymptotics*. We shall present these respective principles as three propositions: the first of these can be thought of as a principle of non-stationary phase, the second is one of van der Corput's lemmas, and the third is a formulation of the method of stationary phase; for proofs see [3] and [2].

Proposition 1. *Suppose Ψ has compact support in (a, b) and Φ' is never vanishes, then for all $N \geq 0$ we have*

$$|I(\lambda)| \leq C_{N, \Phi, \Psi} \lambda^{-N}.$$

Remark. If we do not assume that Ψ vanishes near the endpoints of the interval $[a, b]$ then the best estimate we can obtain for $I(\lambda)$ is $O(\lambda^{-1})$. However, in the "periodic" case, i.e, if we have $\Phi^{(k)}(a) = \Phi^{(k)}(b)$ and $\Psi^{(k)}(a) = \Psi^{(k)}(b)$, we again, as in Proposition 1, obtain the rapid decrease of $I(\lambda)$.

Proposition 2. *Suppose Φ is real-valued and $|\Phi^{(k)}(x)| \geq 1$ for all $x \in (a, b)$, then*

$$|I(\lambda)| \leq k C_k \lambda^{-\frac{1}{k}} \left[|\Psi(b)| + \int_a^b |\Psi'(x)| dx \right],$$

whenever (i) $k = 1$ and $\Phi''(x)$ has at most one zero, or (ii) $k \geq 2$.

Remark. We note that the constants in Proposition 2 are $C_1 = 2$, while $C_k \leq 2^{\frac{5}{3}}$ for all $k \geq 2$ and $C_k \rightarrow 4/e$ as $k \rightarrow \infty$; see [1].

Of course if Φ is completely stationary then the best one can do is $|I(\lambda)| \leq (b - a) \|\Psi\|_\infty$.

Proposition 3. *Suppose Φ is real-valued, $\Phi'(x_0) = 0$, while $\Phi''(x_0) \neq 0$. If Ψ is supported in a sufficiently small neighborhood of x_0 , then*

$$I(\lambda) = a_0 \lambda^{-\frac{1}{2}} e^{i\lambda\Phi(x_0)} + O(\lambda^{-\frac{3}{2}})$$

as $\lambda \rightarrow \infty$, where $a_0 = e^{i\frac{\pi}{4}} \left(\frac{2\pi}{\Phi''(x_0)} \right)^{\frac{1}{2}} \Psi(x_0)$, and the bounds occurring in the error term depend on upper bounds for finitely many derivatives of Φ and Ψ on the $\text{supp } \Psi$, the size of this support, and on a lower bound for $|\Phi''(x_0)|$.

Remark 4. If we merely assume that Φ is real-valued, $\Phi'(x_0) = 0$ and $\Phi'(x) \neq 0$ on $\text{supp } \Psi \setminus \{x_0\}$. Then if $\Phi''(x_0) \neq 0$ we may conclude that

$$I(\lambda) = e^{i\lambda\Phi(x_0)}\sigma(\lambda),$$

where σ is a symbol of order $-\frac{1}{2}$, that is $|\sigma^{(\ell)}(\lambda)| \leq c_\ell(1 + \lambda)^{-\frac{1}{2}-\ell}$; see [4].

The constant c_ℓ depends on the $C^{\ell+1}$ norms of Φ and Ψ on the $\text{supp } \Psi$, the size of this support, and on a lower bound for $|\Phi''(x_0)|$.

Example. The *Bessel functions*, defined for $k \in \mathbf{Z}^+$ by

$$J_k(\lambda) = \frac{1}{2\pi} \int_0^{2\pi} e^{i\lambda \sin \theta} e^{ik\theta} d\theta,$$

are a model case for these oscillatory integrals. It follows from Proposition 3 and the remarks following Proposition 1, that for $\lambda \gg 1$

$$(1) \quad J_k(\lambda) = \left(\frac{2}{\pi}\right)^{\frac{1}{2}} \lambda^{-\frac{1}{2}} \cos\left(\lambda - \frac{\pi k}{2} - \frac{\pi}{4}\right) + O(\lambda^{-\frac{3}{2}}).$$

The Bessel function can also be defined for real $k > -\frac{1}{2}$ by the formula

$$J_k(\lambda) = (\pi^{\frac{1}{2}} \Gamma(k + \frac{1}{2}))^{-1} \left(\frac{\lambda}{2}\right)^k \int_{-1}^1 e^{i\lambda t} (1 - t^2)^{k-\frac{1}{2}} dt.$$

These two definitions agree when k is a positive integer and the asymptotic expression (1) is still valid. One can in fact show that

$$(2) \quad J_k(\lambda) = \sigma_1(\lambda)e^{i\lambda} + \sigma_2(\lambda)e^{-i\lambda},$$

where $|\sigma_i^{(\ell)}(\lambda)| \leq c_\ell(1 + \lambda)^{-\frac{1}{2}-\ell}$. See Remark 4.

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