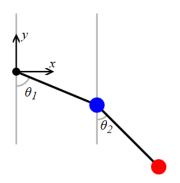
n-kratno nihalo

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1 Dvojno nihalo

Oglejmo si primer na sliki:



Slika 1: Dvojno nihalo [1]

Imamo dve žogici z masama m_1 in m_2 na palčkah dolžine l_1 in l_2 z zanemarljivo maso. Recimo, da je prva žogica na (x_1, y_1) in druga žogica na (x_2, y_2) . Te koordinate dobimo kot:

$$x_1 = l_1 \sin(\theta_1),$$

$$x_2 = l_1 \sin(\theta_1) + l_2 \sin(\theta_2),$$

$$y_1 = -l_1 \cos(\theta_1),$$

$$y_2 = -l_1 \cos(\theta_1) - l_2 \cos(\theta_2).$$

Potencialna energija

Potencialna energija je definirana kot V=mgh, kjer je višina žogice dana zh=y. Za posamezni masi velja:

$$h_1 = y_1, \quad h_2 = y_2.$$

Torej celotna potencialna energija sistema je:

$$V = V_1 + V_2$$

= $m_1 g(-l_1 \cos \theta_1) + m_2 g(-l_1 \cos \theta_1 - l_2 \cos \theta_2)$.

Kinetična energija

Kinetična energija vsake žogice je podana z izrazom

$$T = \frac{1}{2}mv^2.$$

Hitrost žogic dobimo s pomočjo izraza $v=\sqrt{\dot{x}^2+\dot{y}^2}.$

Odvodi koordinat so:

$$\dot{x}_1 = l_1 \dot{\theta}_1 \cos(\theta_1)
\dot{x}_2 = l_1 \dot{\theta}_1 \cos(\theta_1) + l_2 \dot{\theta}_2 \cos(\theta_2)
\dot{y}_1 = l_1 \dot{\theta}_1 \sin(\theta_1)
\dot{y}_2 = l_1 \dot{\theta}_1 \sin(\theta_1) + l_2 \dot{\theta}_2 \sin(\theta_2)$$

Skupna kinetična energija sistema je torej:

$$T = T_1 + T_2 = \frac{m_1}{2}(\dot{x}_1^2 + \dot{y}_1^2) + \frac{m_2}{2}(\dot{x}_2^2 + \dot{y}_2^2)$$

Lotimo se sedaj zapisa sistema Euler-Lagrangeevih enačb [2] za $\mathcal{L} = T - V$:

$$\mathcal{L} = \frac{m_1}{2} (\dot{x}_1^2 + \dot{y}_1^2) + \frac{m_2}{2} (\dot{x}_2^2 + \dot{y}_2^2) + m_1 g (l_1 \cos \theta_1) + m_2 g (l_1 \cos \theta_1 + l_2 \cos \theta_2)$$

$$= \frac{m_1}{2} l_1^2 \dot{\theta}_1^2 + \frac{m_2}{2} (l_1^2 \dot{\theta}_1^2 + l_2^2 \dot{\theta}_2^2 + 2 l_1 \dot{\theta}_1 l_2 \dot{\theta}_2 \cos \theta_1 \cos \theta_2 + 2 l_1 \dot{\theta}_1 l_2 \dot{\theta}_2 \sin \theta_1 \sin \theta_2)$$

$$+ m_1 g l_1 \cos \theta_1 + m_2 g (l_1 \cos \theta_1 + l_2 \cos \theta_2)$$

$$= \frac{1}{2} m_1 l_1^2 \dot{\theta}_1^2 + \frac{1}{2} m_2 (l_1^2 \dot{\theta}_1^2 + l_2^2 \dot{\theta}_2^2 + 2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \cos(\theta_1 - \theta_2)) + (m_1 + m_2) l_1 g \cos \theta_1$$

$$+ m_2 l_2 g \cos \theta_2.$$

Želimo dobiti enačbe oblike:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}_i} \right) = \frac{\partial \mathcal{L}}{\partial \theta_i}, \quad i = 1, 2.$$

Najprej izpeljimo za θ_1 :

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}_{1}} = m_{1} l_{1}^{2} \dot{\theta}_{1} + m_{2} l_{1}^{2} \dot{\theta}_{1} + m_{2} l_{1} l_{2} \dot{\theta}_{2} \cos(\theta_{1} - \theta_{2}),$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}_{1}} \right) = m_{1} l_{1}^{2} \ddot{\theta}_{1} + m_{2} l_{1} l_{2} [\ddot{\theta}_{2} \cos(\theta_{1} - \theta_{2}) - \dot{\theta}_{2} \sin(\theta_{1} - \theta_{2}) (\dot{\theta}_{1} - \dot{\theta}_{2})]$$

$$\frac{\partial \mathcal{L}}{\partial \theta_{1}} = - m_{2} l_{1} l_{2} \dot{\theta}_{1} \dot{\theta}_{2} \sin(\theta_{1} - \theta_{2}) - (m_{1} + m_{2}) l_{1} g \sin \theta_{1}.$$

Torej za i = 1 dobimo:

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta_1}} \right) - \frac{\partial \mathcal{L}}{\partial \theta_1} = (m_1 + m_2)(l_1^2 \ddot{\theta_1}) + m_2 l_1 l_2 [\ddot{\theta_2} \cos(\theta_1 - \theta_2) + \dot{\theta_2}^2 \sin(\theta_1 - \theta_2)] + (m_1 + m_2) l_1 g \sin \theta_1 = 0 \quad / : l_1.$$

Podobno naredimo za θ_2 in dobimo:

$$\frac{\partial \mathcal{L}}{\partial \dot{\theta}_2} = m_2 l_2 \dot{\theta}_2 + m_2 l_1 l_2 \dot{\theta}_1 \cos(\theta_1 - \theta_2)$$

$$\frac{d}{dt} \left(\frac{\partial \mathcal{L}}{\partial \dot{\theta}_2} \right) = m_2 l_2^2 \ddot{\theta}_2 + m_2 l_1 l_2 [\ddot{\theta}_1 \cos(\theta_1 - \theta_2) - \dot{\theta}_1 \sin(\theta_1 - \theta_2) (\dot{\theta}_1 - \dot{\theta}_2)]$$

$$\frac{\partial \mathcal{L}}{\partial \theta_2} = -m_2 l_1 l_2 \dot{\theta}_1 \dot{\theta}_2 \sin(\theta_1 - \theta_2) - m_2 l_2 g \sin \theta_2$$

Euler-Lagrangeeva enačba za i = 2 se torej glasi:

$$\frac{d}{dt}\left(\frac{\partial \mathcal{L}}{\partial \dot{\theta_2}}\right) - \frac{\partial \mathcal{L}}{\partial \theta_2} = m_2 l_2^2 \ddot{\theta_2} + m_2 l_1 l_2 (\ddot{\theta_1} \cos(\theta_1 - \theta_2) - \dot{\theta_1}^2 \sin(\theta_1 - \theta_2)) + m_2 l_2 g \sin\theta_2 = 0 \quad /: l_2$$

Dobimo sistem diferencialnih enačb:

$$\theta_1: \quad (m_1 + m_2)[l_1\ddot{\theta}_1 + g\sin\theta_1] + m_2l_2[\ddot{\theta}_2\cos(\theta_1 - \theta_2) + \dot{\theta}_2^2\sin(\theta_1 - \theta_2)] = 0$$

$$\theta_2: \quad m_2[l_2\ddot{\theta}_2 + g\sin\theta_2] + m_2l_1[\ddot{\theta}_1\cos(\theta_1 - \theta_2) - \dot{\theta}_1^2\sin(\theta_1 - \theta_2)] = 0$$

2 Trojno nihalo

Kaj pa bi se zgodilo, če vzamemo trojno nihalo? Predpostavimo, da na drugo žogico pripnemo preko vrvice l_3 še eno žogico z maso m_3 .

Ta žogica je na položaju

$$\begin{aligned} x_3 &= l_1 \sin \theta_1 + l_2 \sin \theta_2 + l_3 \sin \theta_3 \\ y_3 &= -l_1 \cos \theta_1 - l_2 \cos \theta_2 - l_3 \cos \theta_3 \\ \dot{x_3} &= l_1 \dot{\theta_1} \cos \theta_2 + l_2 \dot{\theta_2} \cos \theta_2 + l_3 \dot{\theta_3} \cos \theta_3 \\ \dot{y_3} &= l_1 \dot{\theta_1} \sin \theta_2 + l_2 \dot{\theta_2} \sin \theta_2 + l_3 \dot{\theta_3} \sin \theta_3 \end{aligned}$$

Za ta sistem imamo:

$$V_{3} = m_{1}gy_{1} + m_{2}gy_{2} + m_{3}gy_{3}$$

$$= -(m_{1} + m_{2} + m_{3})l_{1}g\cos\theta_{1} - \cos\theta_{2}gl_{2}(m_{2} + m_{3}) - \cos\theta_{3}l_{3}m_{3}g$$

$$T_{3} = \frac{1}{2}m_{1}(\dot{x_{1}}^{2} + \dot{y_{1}}^{2}) + \frac{1}{2}m_{2}(\dot{x_{2}}^{2} + \dot{y_{2}}^{2}) + \frac{1}{2}m_{3}(\dot{x_{3}}^{2} + \dot{y_{3}}^{2})$$

$$= \frac{1}{2}m_{1}l_{1}^{2}\dot{\theta_{1}}^{2} + \frac{1}{2}m_{2}[l_{1}^{2}\dot{\theta_{1}}^{2} + l_{2}\dot{\theta_{2}}^{2} + 2l_{1}l_{2}\dot{\theta_{1}}\dot{\theta_{2}}\cos(\theta_{1} - \theta_{2})] + \frac{1}{2}m_{3}[l_{1}^{2}\dot{\theta_{1}}^{2} + l_{2}\dot{\theta_{2}}^{2} + l_{1}^{2}\dot{\theta_{2}}\dot{\theta_{3}}\cos(\theta_{1} - \theta_{2})] + l_{1}^{2}\dot{\theta_{3}}\dot{\theta_{2}}\dot{\theta_{3}}\cos(\theta_{2} - \theta_{3})]$$

$$\mathcal{L} = T_{3} - V_{3}$$

Podobno kot za dvojno nihalo tudi tu dobimo sistem enačb za θ_i , i=1,2,3. Če jih malo preuredimo, dobimo:

$$\theta_{1}: \quad (m_{1}+m_{2}+m_{3})[l_{1}\ddot{\theta}_{1}+g\sin\theta_{1}]+(m_{2}+m_{3})l_{2}[\ddot{\theta}_{2}\cos(\theta_{1}-\theta_{2})+\dot{\theta}_{2}^{2}\sin(\theta_{1}-\theta_{2})] \\ +m_{3}l_{3}[\ddot{\theta}_{3}\cos(\theta_{1}-\theta_{3})+\dot{\theta}_{3}^{2}\sin(\theta_{1}-\theta_{3})]=0 \\ \theta_{2}: \quad (m_{2}+m_{3})[l_{2}\ddot{\theta}_{2}+g\sin\theta_{2}]+(m_{2}+m_{3})l_{1}[\ddot{\theta}_{1}\cos(\theta_{1}-\theta_{2})-\dot{\theta}_{1}^{2}\sin(\theta_{1}-\theta_{2})] \\ +m_{3}l_{3}[\ddot{\theta}_{3}\cos(\theta_{2}-\theta_{3})+\dot{\theta}_{3}^{2}\sin(\theta_{2}-\theta_{3})]=0 \\ \theta_{3}: \quad m_{3}[l_{3}\ddot{\theta}_{3}-g\sin\theta_{3}]+m_{3}l_{1}[\ddot{\theta}_{1}\cos(\theta_{1}-\theta_{3})-\dot{\theta}_{1}^{2}\sin(\theta_{1}-\theta_{3})] \\ +m_{3}l_{2}[\ddot{\theta}_{2}\cos(\theta_{2}-\theta_{3})-\dot{\theta}_{2}^{2}\sin(\theta_{2}-\theta_{3})]=0$$

Zaenkrat sva dobili idejo za reukrzijo, ki izgleda nekako tako?:

Za
$$\theta_1$$
: $(\sum_{i=1}^n m_i)[l_1\ddot{\theta}_1 + g\sin\theta_1] + \sum_{i=2}^n (\sum_{j=i}^n m_j)l_i(\ddot{\theta}_i\cos(\theta_1 - \theta_i) + \dot{\theta}_i^2\sin(\theta_1 - \theta_i)) = 0$
Za θ_n : $m_n[l_n\ddot{\theta}_n + g\sin\theta_n] + m_n\sum_{i=1}^{n-1} l_i[\ddot{\theta}_i\cos(\theta_i - \theta_n) - \dot{\theta}_i^2\sin(\theta_i - \theta_n)] = 0$

Literatura

- [1] Scipython. The double pendulum. Ogled: 24. 10. 2025. URL: https://scipython.com/blog/the-double-pendulum/.
- [2] Wikipedia. Lagrangian mechanics. Ogled: 24. 10. 2025. URL: https://en.wikipedia.org/wiki/Lagrangian_mechanics.