



Min. 22.00

27/03

Robotics 2

Dynamic model of robots: Analysis, properties, extensions, uses

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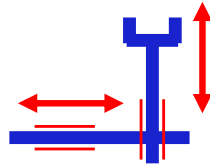
DIPARTIMENTO DI INGEGNERIA INFORMATICA
AUTOMATICA E GESTIONALE ANTONIO RUBERTI



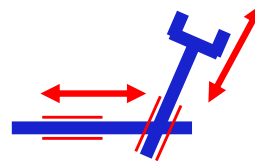
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Analysis of inertial couplings

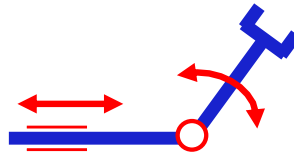
- Cartesian robot



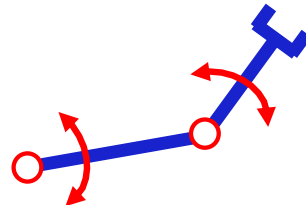
- Cartesian "skew" robot



- PR robot

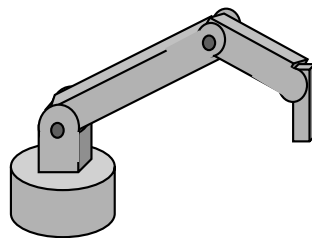


- 2R robot



- 3R articulated robot

(under simplifying assumptions on the CoMs)



Two second order decoupled differ

$$M = \begin{pmatrix} m_{11} & 0 \\ 0 & m_{22} \end{pmatrix}$$

→ NO CORIOLIS, CENTR terms

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gravity terms
↳ g_2 has linear relation to gravity

$$M = \begin{pmatrix} m_{11} & m_{12} \\ m_{12} & m_{22} \end{pmatrix}$$

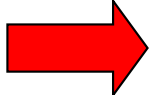
$$M = \begin{pmatrix} m_{11} & m_{12}(q_2) \\ m_{12}(q_2) & m_{22} \end{pmatrix}$$

$$M = \begin{pmatrix} m_{11}(q_2) & m_{12}(q_2) \\ m_{12}(q_2) & m_{22} \end{pmatrix}$$

$$M = \begin{pmatrix} m_{11}(q_2, q_3) & 0 & 0 \\ 0 & m_{22}(q_3) & m_{23}(q_3) \\ 0 & m_{23}(q_3) & m_{33} \end{pmatrix}$$

Analysis of gravity term

- absence of gravity
 - constant U_g (motion on horizontal plane)
 - applications in remote space
- static balancing
 - distribution of masses (including motors)
- mechanical compensation
 - articulated system of springs
 - closed kinematic chains

 $g(q) \approx 0$





Bounds on dynamic terms

- for an open-chain (serial) manipulator, there always exist positive real constants k_0 to k_7 such that, for **any** value of q and \dot{q}

$$k_0 \leq \|M(q)\| \leq k_1 + k_2\|q\| + k_3\|q\|^2 \quad \text{inertia matrix}$$

$$\|S(q, \dot{q})\| \leq (k_4 + k_5\|q\|) \|\dot{q}\| \quad \text{factorization matrix of Coriolis/centrifugal terms}$$

$$\|g(q)\| \leq k_6 + k_7\|q\| \quad \text{gravity vector}$$

- if the robot has only **revolute** joints, these simplify to

$$k_0 \leq \|M(q)\| \leq k_1 \quad \|S(q, \dot{q})\| \leq k_4\|\dot{q}\| \quad \|g(q)\| \leq k_6$$

(the same holds true with bounds $q_{i,min} \leq q_i \leq q_{i,max}$ on prismatic joints)

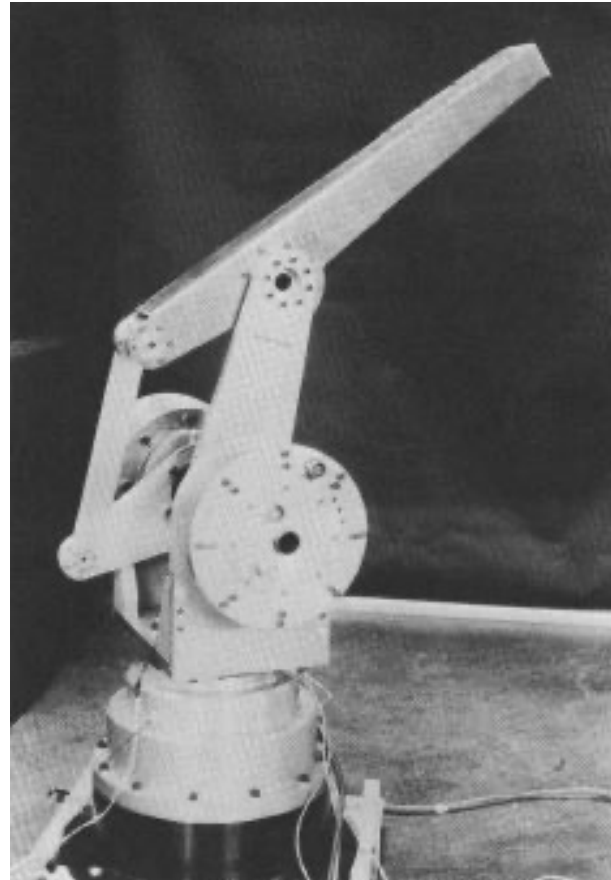
NOTE: norms are either for vectors or for matrices (induced norms)

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Robots with closed kinematic chains - 1

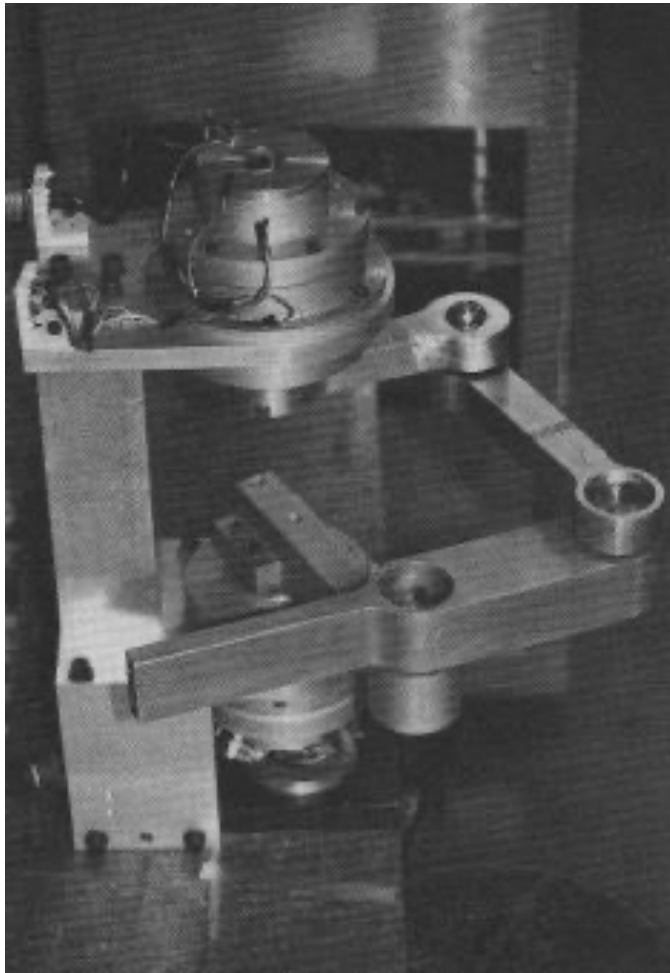


Comau Smart NJ130

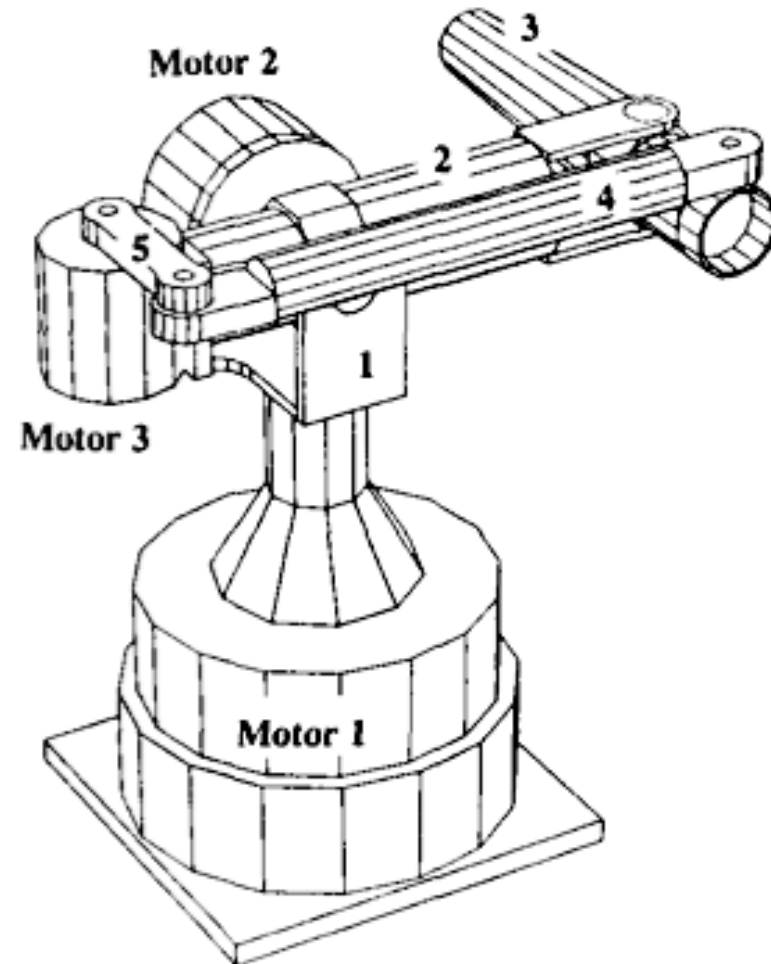


MIT Direct Drive Mark II and Mark III

Robots with closed kinematic chains - 2



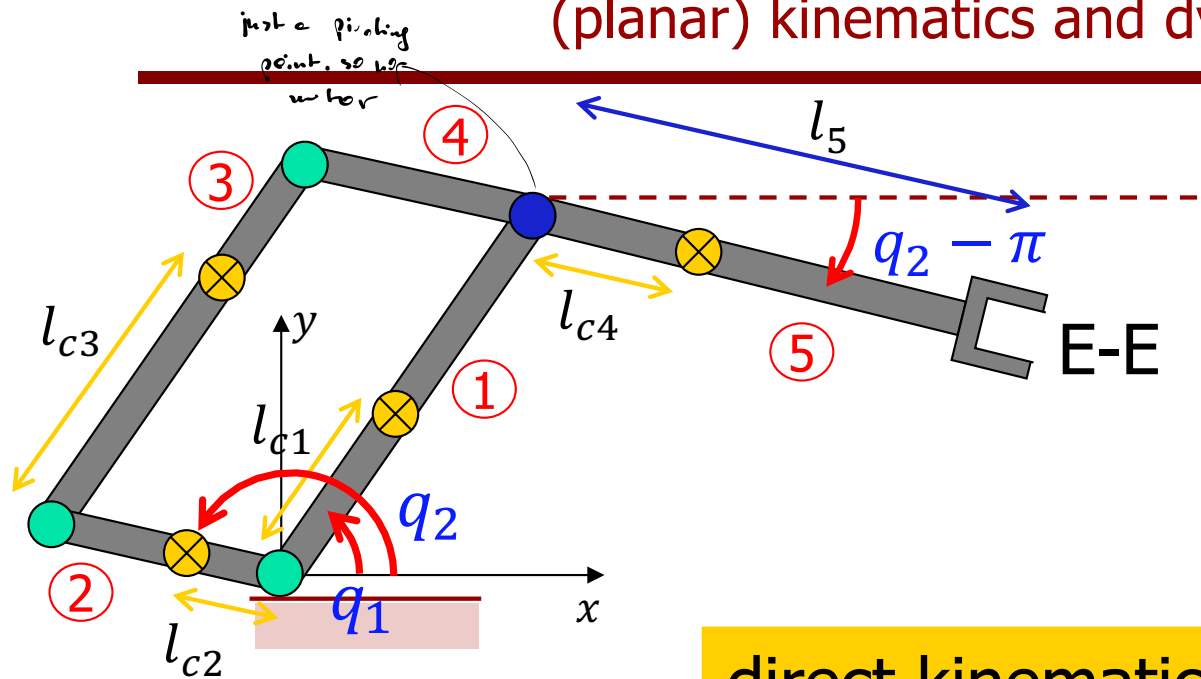
MIT Direct Drive Mark IV
(**planar** five-bar linkage)



UMinnesota Direct Drive Arm
(**spatial** five-bar linkage)

Robot with parallelogram structure

(planar) kinematics and dynamics



⊗ center of mass:
arbitrary l_{ci}

parallelogram:

$$l_1 = l_3$$

$$l_2 = l_4$$

direct kinematics

$$p_{EE} = \begin{pmatrix} l_1 c_1 \\ l_1 s_1 \end{pmatrix} + \begin{pmatrix} l_5 \cos(q_2 - \pi) \\ l_5 \sin(q_2 - \pi) \end{pmatrix} = \begin{pmatrix} l_1 c_1 \\ l_1 s_1 \end{pmatrix} - \begin{pmatrix} l_5 c_2 \\ l_5 s_2 \end{pmatrix}$$

position of center of masses

$$p_{c1} = \begin{pmatrix} l_{c1} c_1 \\ l_{c1} s_1 \end{pmatrix} \quad p_{c2} = \begin{pmatrix} l_{c2} c_2 \\ l_{c2} s_2 \end{pmatrix} \quad p_{c3} = \begin{pmatrix} l_2 c_2 \\ l_2 s_2 \end{pmatrix} + \begin{pmatrix} l_{c3} c_1 \\ l_{c3} s_1 \end{pmatrix} \quad p_{c4} = \begin{pmatrix} l_1 c_1 \\ l_1 s_1 \end{pmatrix} - \begin{pmatrix} l_{c4} c_2 \\ l_{c4} s_2 \end{pmatrix}$$



Kinetic energy

linear/angular velocities

$$\begin{aligned} v_{c1} &= \begin{pmatrix} -l_{c1}s_1 \\ l_{c1}c_1 \end{pmatrix} \dot{q}_1 & v_{c3} &= \begin{pmatrix} -l_{c3}s_1 \\ l_{c3}c_1 \end{pmatrix} \dot{q}_1 + \begin{pmatrix} -l_2s_2 \\ l_2c_2 \end{pmatrix} \dot{q}_2 & \omega_1 &= \omega_3 = \dot{q}_1 \\ v_{c2} &= \begin{pmatrix} -l_{c2}s_2 \\ l_{c2}c_2 \end{pmatrix} \dot{q}_2 & v_{c4} &= \begin{pmatrix} -l_1s_1 \\ l_1c_1 \end{pmatrix} \dot{q}_1 + \begin{pmatrix} l_{c4}s_2 \\ -l_{c4}c_2 \end{pmatrix} \dot{q}_2 & \omega_2 &= \omega_4 = \dot{q}_2 \end{aligned}$$

Note: a (planar) 2D notation is used here!

$$\begin{aligned} T_i \quad T_1 &= \frac{1}{2} m_1 l_{c1}^2 \dot{q}_1^2 + \frac{1}{2} I_{c1,zz} \dot{q}_1^2 & T_2 &= \frac{1}{2} m_2 l_{c2}^2 \dot{q}_2^2 + \frac{1}{2} I_{c2,zz} \dot{q}_2^2 \\ T_3 &= \frac{1}{2} m_3 (l_2^2 \dot{q}_2^2 + l_{c3}^2 \dot{q}_1^2 + 2l_2 l_{c3} c_{2-1} \dot{q}_1 \dot{q}_2) + \frac{1}{2} I_{c3,zz} \dot{q}_1^2 \\ T_4 &= \frac{1}{2} m_4 (l_1^2 \dot{q}_1^2 + l_{c4}^2 \dot{q}_2^2 - 2l_1 l_{c4} c_{2-1} \dot{q}_1 \dot{q}_2) + \frac{1}{2} I_{c4,zz} \dot{q}_2^2 \end{aligned}$$



Robot inertia matrix

$$T = \sum_{i=1}^4 T_i = \frac{1}{2} \dot{q}^T M(q) \dot{q}$$

diagonal terms are constant

$$M(q) = \begin{pmatrix} I_{c1,zz} + m_1 l_{c1}^2 + I_{c3,zz} + m_3 l_{c3}^2 + m_4 l_1^2 & (m_3 l_2 l_{c3} - m_4 l_1 l_{c4}) c_{2-1} \\ (m_3 l_2 l_{c3} - m_4 l_1 l_{c4}) c_{2-1} & I_{c2,zz} + m_2 l_{c2}^2 + I_{c4,zz} + m_4 l_{c4}^2 + m_3 l_2^2 \end{pmatrix}$$

symm

min 5 l. 30

structural condition
in mechanical design

$$m_3 l_2 l_{c3} = m_4 l_1 l_{c4}$$

(*) \rightarrow it becomes constant diagonal



$M(q)$ diagonal and **constant** \Rightarrow centrifugal and Coriolis terms $\equiv 0$

mechanically **DECOUPLED** and **LINEAR**
dynamic model (up to the gravity term $g(q)$)



$$\begin{pmatrix} M_{11} & 0 \\ 0 & M_{22} \end{pmatrix} \begin{pmatrix} \ddot{q}_1 \\ \ddot{q}_2 \end{pmatrix} = \begin{pmatrix} u_1 \\ u_2 \end{pmatrix}$$

big advantage for the design of a motion control law!

each joint
can be treated
independently

just like
prismatic joints



Potential energy and gravity terms

from the y-components of vectors p_{ci}

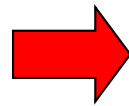
$$\begin{aligned} U_1 &= m_1 g_0 l_{c1} s_1 & U_2 &= m_2 g_0 l_{c2} s_2 \\ U_3 &= m_3 g_0 (l_2 s_2 + l_{c3} s_1) & U_4 &= m_4 g_0 (l_1 s_1 - l_{c4} s_2) \end{aligned}$$

$$U = \sum_{i=1}^4 U_i$$

$$g(q) = \left(\frac{\partial U}{\partial q} \right)^T = \begin{pmatrix} g_0(m_1 l_{c1} + m_3 l_{c3} + m_4 l_1) c_1 \\ g_0(m_2 l_{c2} + m_3 l_2 - m_4 l_{c4}) c_2 \end{pmatrix} = \begin{pmatrix} g_1(q_1) \\ g_2(q_2) \end{pmatrix}$$

gravity components are **always** "decoupled"

in addition,
when (*) holds



$$\begin{aligned} m_{11} \ddot{q}_1 + g_1(q_1) &= u_1 \\ m_{22} \ddot{q}_2 + g_2(q_2) &= u_2 \end{aligned}$$

u_i are
(non-conservative) torques
performing work on q_i

further structural conditions in the mechanical design lead to $g(q) \equiv 0!!$

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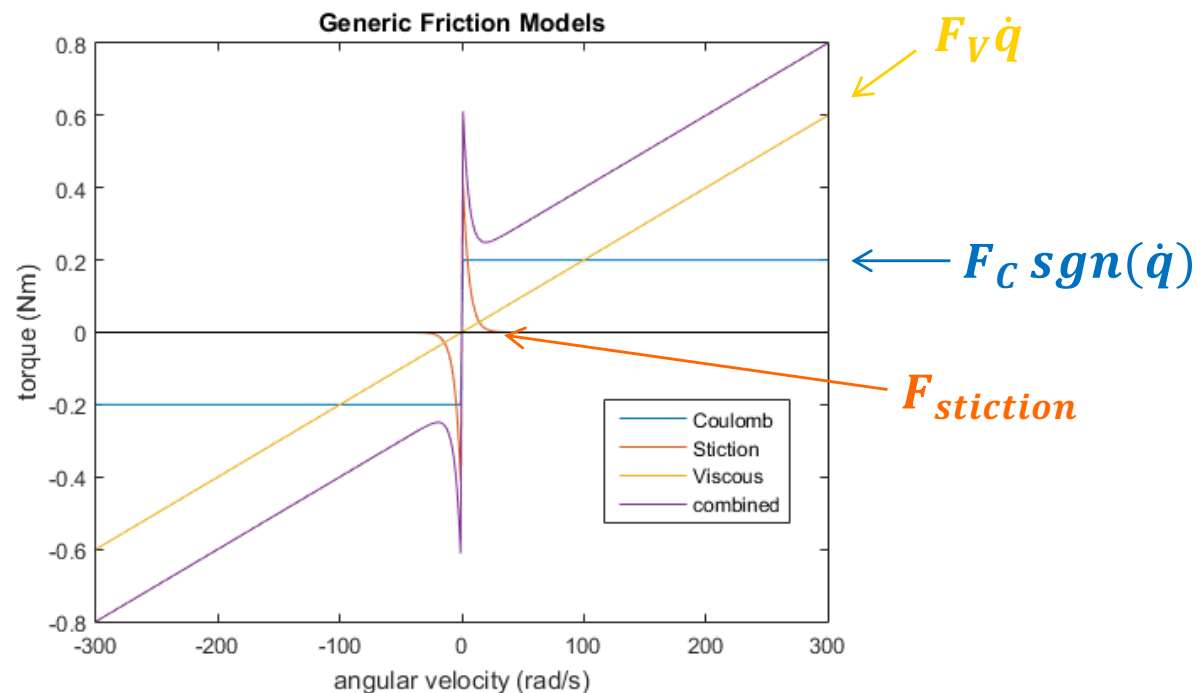
Adding dynamic terms ...

1) dissipative phenomena due to friction at the joints/transmissions

- viscous, Coulomb, stiction, Stribeck, LuGre (dynamic)...
- local effects at the joints
- difficult to model in general, except for:

$$u_{V,i} = -F_{V,i} \dot{q}_i$$

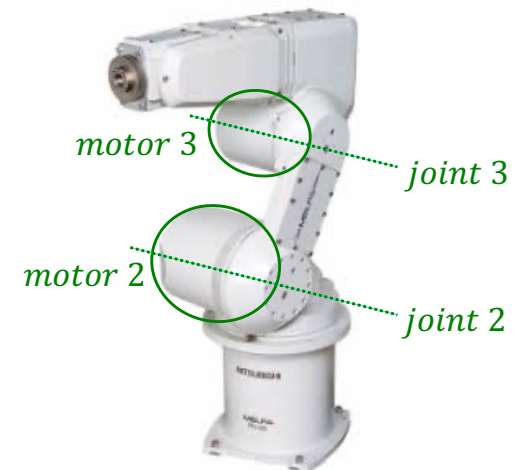
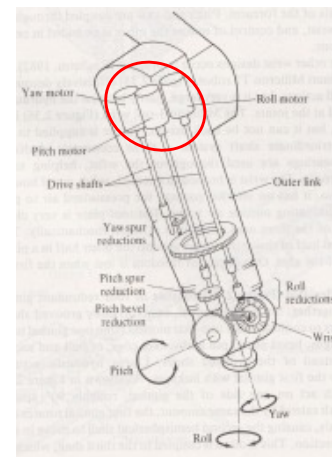
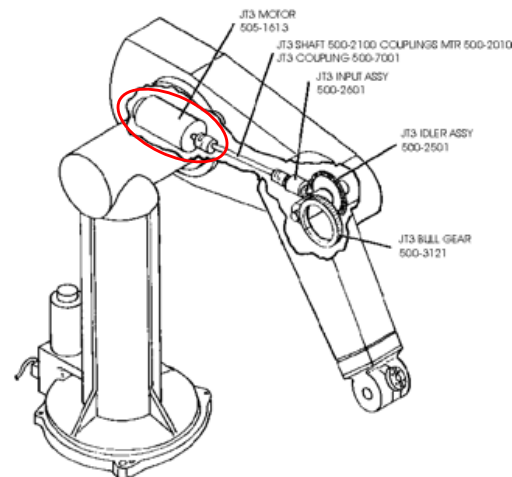
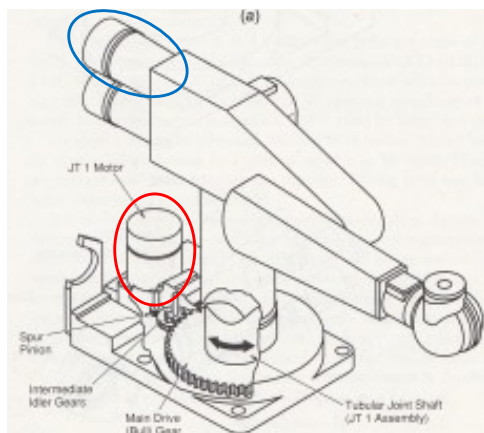
$$u_{C,i} = -F_{C,i} \operatorname{sgn}(\dot{q}_i)$$



Adding dynamic terms ...

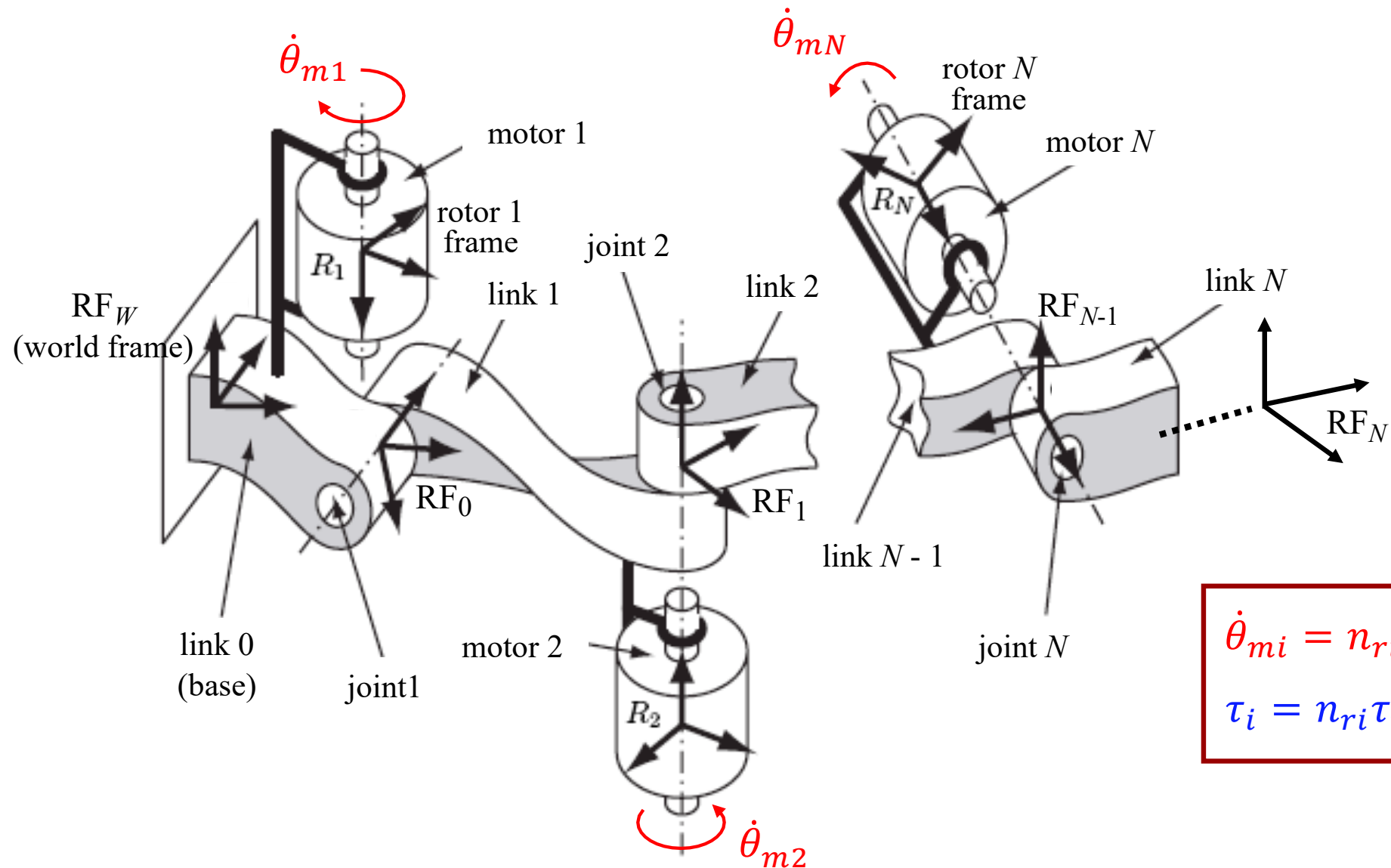
- 2) inclusion of electrical **actuators** (as additional rigid bodies)
 - motor i mounted on link $i - 1$ (or **before**), with very few **exceptions**
 - often with its spinning **axis aligned with joint axis i**
 - (balanced) **mass** of motor included in total mass of carrying link
 - (rotor) **inertia** is to be **added** to robot kinetic energy
 - transmissions with **reduction gears** (often, large reduction ratios)
 - in some cases, multiple motors cooperate in moving multiple links: use a **transmission coupling** matrix Γ (with off-diagonal elements)

Unimation PUMA family



Mitsubishi RV-3S

Placement of motors along the chain



$$\dot{\theta}_{mi} = n_{ri} \dot{\theta}_i$$

$$\tau_i = n_{ri} \tau_{mi}$$



min. 1.15.00 Resulting dynamic model

- **simplifying assumption:** in the **rotational** part of the kinetic energy, only the "spinning" rotor velocity is considered

$$T_{mi} = \frac{1}{2} I_{mi} \dot{\theta}_{mi}^2 = \frac{1}{2} I_{mi} n_{ri}^2 \dot{q}_i^2 = \frac{1}{2} B_{mi} \dot{q}_i^2 \quad T_m = \sum_{i=1}^N T_{mi} = \frac{1}{2} \dot{q}^T B_m \dot{q}$$

diagonal, > 0

- including all added terms, the robot dynamics becomes

$$(M(q) + B_m) \ddot{q} + c(q, \dot{q}) + g(q) + \underbrace{F_V \dot{q} + F_C \operatorname{sgn}(\dot{q})}_{\substack{F_V > 0, F_C > 0 \\ \text{diagonal}}} = \tau$$

constant \rightarrow does NOT contribute to c

moved to the left ...
 motor torques (after reduction gears)

- scaling by the reduction gears, looking **from the motor side**

$$\left(I_m + \operatorname{diag} \left\{ \frac{m_{ii}(q)}{n_{ri}^2} \right\} \right) \ddot{\theta}_m + \operatorname{diag} \left\{ \frac{1}{n_{ri}} \right\} \left(\sum_{j=1}^N \bar{M}_j(q) \ddot{q}_j + f(q, \dot{q}) \right) = \tau_m$$

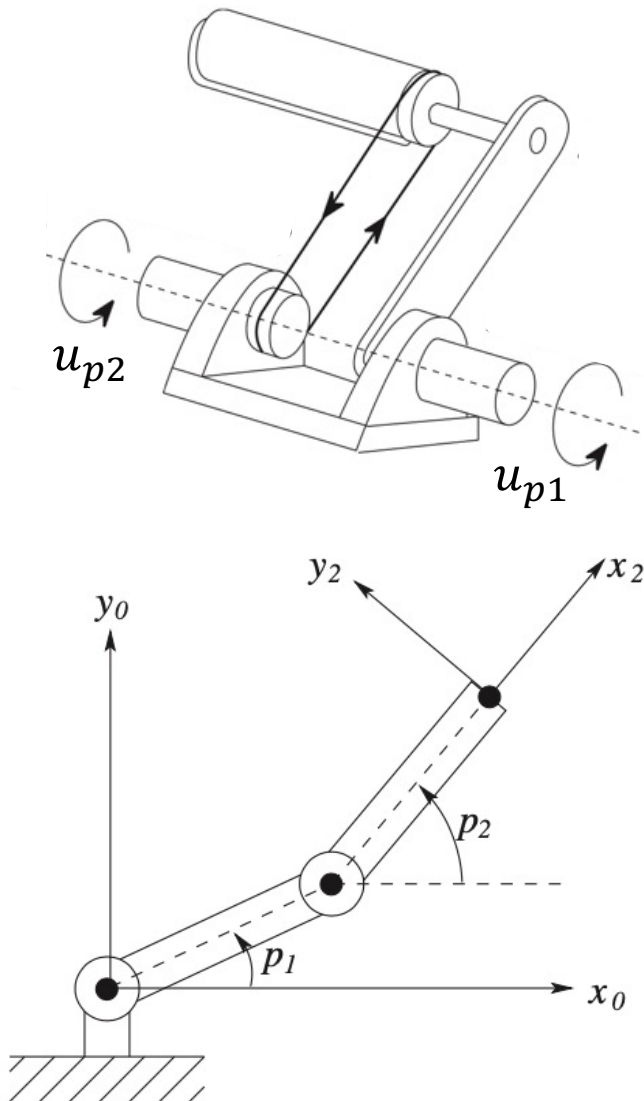
diagonal \rightarrow

except the terms m_{jj}

motor torques (before reduction gears)

Special actuation and associated coordinates

win 1.24.20 planar 2R robot with remotely driven forearm



- motor 1 moves link 1 by p_1
- motor 2 **at the base** moves the **absolute** angle p_2 of link 2
- derive the model in the \mathbf{p} coordinates

$$M(\mathbf{p})\ddot{\mathbf{p}} + \mathbf{c}(\mathbf{p}, \dot{\mathbf{p}}) + \mathbf{g}(\mathbf{p}) = \mathbf{u}_p$$

$$M(\mathbf{p}) = \begin{pmatrix} a_1 - a_3 & a_2 c_{2-1} \\ a_2 c_{2-1} & a_3 \end{pmatrix}$$

$$\mathbf{c}(\mathbf{p}, \dot{\mathbf{p}}) = \begin{pmatrix} -a_2 s_{2-1} \dot{p}_2^2 \\ a_2 s_{2-1} \dot{p}_1^2 \end{pmatrix} \quad \text{no more Coriolis forces!}$$

$$\mathbf{g}(\mathbf{p}) = \begin{pmatrix} a_4 c_1 \\ a_5 c_2 \end{pmatrix}$$

$$c_1 = \cos p_1, c_2 = \cos p_2,$$

$$c_{2-1} = \cos(p_2 - p_1), s_{2-1} = \sin(p_2 - p_1)$$



Including joint elasticity

- in **industrial** robots, use of motion transmissions based on
 - belts
 - harmonic drives
 - long shafts

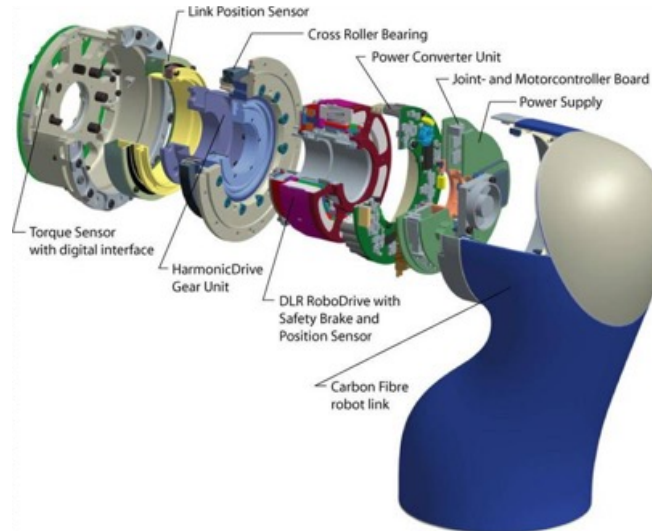
introduces **flexibility** between actuating motors (input) and driven links (output)

- in **research** robots, **compliance** in transmissions is introduced on purpose for **safety** (human collaboration) and/or **energy efficiency**
 - actuator relocation by means of (compliant) cables and pulleys
 - harmonic drives and lightweight (but rigid) link design
 - redundant (macro-mini or parallel) actuation, with elastic couplings
- in both cases, flexibility is modeled as **concentrated at the joints**
- in most cases, assuming small joint deformation (**elastic domain**)

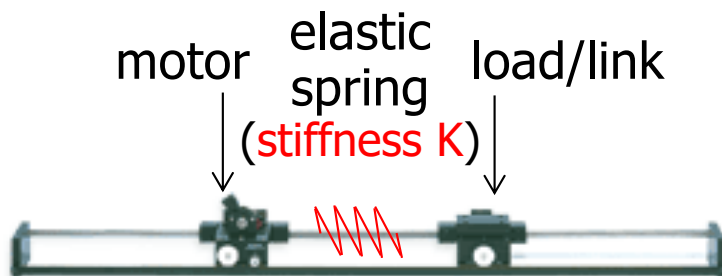
Robots with joint elasticity



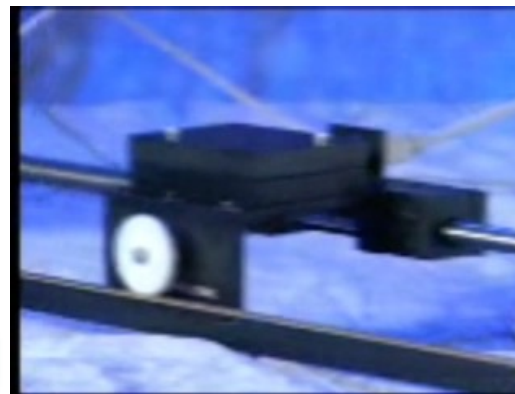
Dexter
with cable transmissions



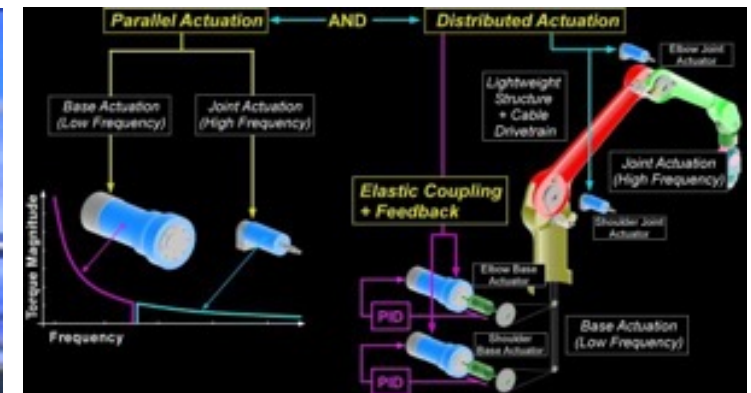
DLR LWR-III
with harmonic drives



Quanser Flexible Joint
(1-dof linear, educational)



video



Stanford DECMMA
with micro-macro actuation



Dynamic model of robots with elastic joints

- introduce $2N$ generalized coordinates

- $q = N$ link positions
- $\theta = N$ motor positions (after reduction, $\theta_i = \theta_{mi}/n_{ri}$)

- add **motor kinetic energy** T_m to that of the links $T_q = \frac{1}{2} \dot{q}^T M(q) \dot{q}$

$$T_{mi} = \frac{1}{2} I_{mi} \dot{\theta}_{mi}^2 = \frac{1}{2} I_{mi} n_{ri}^2 \dot{\theta}_i^2 = \frac{1}{2} B_{mi} \dot{\theta}_i^2 \quad T_m = \sum_{i=1}^N T_{mi} = \frac{1}{2} \dot{\theta}^T B_m \dot{\theta}$$

diagonal, > 0

- add **elastic potential energy** U_e to that due to gravity $U_g(q)$

- K = matrix of **joint stiffness** (diagonal, > 0)

$$U_{ei} = \frac{1}{2} K_i \left(q_i - \left(\frac{\theta_{mi}}{n_{ri}} \right) \right)^2 = \frac{1}{2} K_i (q_i - \theta_i)^2 \quad U_e = \sum_{i=1}^N U_{ei} = \frac{1}{2} (q - \theta)^T K (q - \theta)$$

- apply **Euler-Lagrange** equations w.r.t. (q, θ)

$2N$ 2nd-order differential equations

$$\begin{cases} M(q) \ddot{q} + c(q, \dot{q}) + g(q) + K(q - \theta) = 0 \\ B_m \ddot{\theta} + K(\theta - q) = \tau \end{cases}$$

no external torques performing work on q

Use of the dynamic model

inverse dynamics



- given a **desired trajectory** $q_d(t)$
 - twice differentiable ($\exists \ddot{q}_d(t)$)
 - possibly obtained from a task/Cartesian trajectory $r_d(t)$, by (differential) kinematic inversion

the **input torque** needed to execute this motion (in **free space**) is

$$\tau_d = (M(q_d) + B_m)\ddot{q}_d + c(q_d, \dot{q}_d) + g(q_d) + F_V \dot{q}_d + F_C \operatorname{sgn}(\dot{q}_d)$$

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(in **contact**, with an external wrench) ... $- J_{ext}^T(q_d) F_{ext,d}$

- useful also for control (e.g., nominal feedforward)
- however, this way of performing the algebraic computation ($\forall t$) is **not efficient** when using the Lagrangian modeling approach
 - symbolic terms grow much longer, quite rapidly for larger N
 - in real time, numerical computation is based on **Newton-Euler** method

State equations

direct dynamics



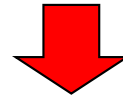
Lagrangian
dynamic model

$$M(q)\ddot{q} + c(q, \dot{q}) + g(q) = u$$

N differential
2nd order
equations

defining the vector of state variables as $x = \begin{pmatrix} x_1 \\ x_2 \end{pmatrix} = \begin{pmatrix} q \\ \dot{q} \end{pmatrix} \in \mathbb{R}^{2N}$

state equations



$$\dot{x} = \begin{pmatrix} \dot{x}_1 \\ \dot{x}_2 \end{pmatrix} = \begin{pmatrix} x_2 \\ -M^{-1}(x_1)[c(x_1, x_2) + g(x_1)] \end{pmatrix} + \begin{pmatrix} 0 \\ M^{-1}(x_1) \end{pmatrix} u$$

$$= f(x) + G(x)u$$

$$\begin{matrix} \uparrow & \uparrow \\ 2N \times 1 & 2N \times N \end{matrix}$$

$2N$ differential
1st order
equations

another choice...

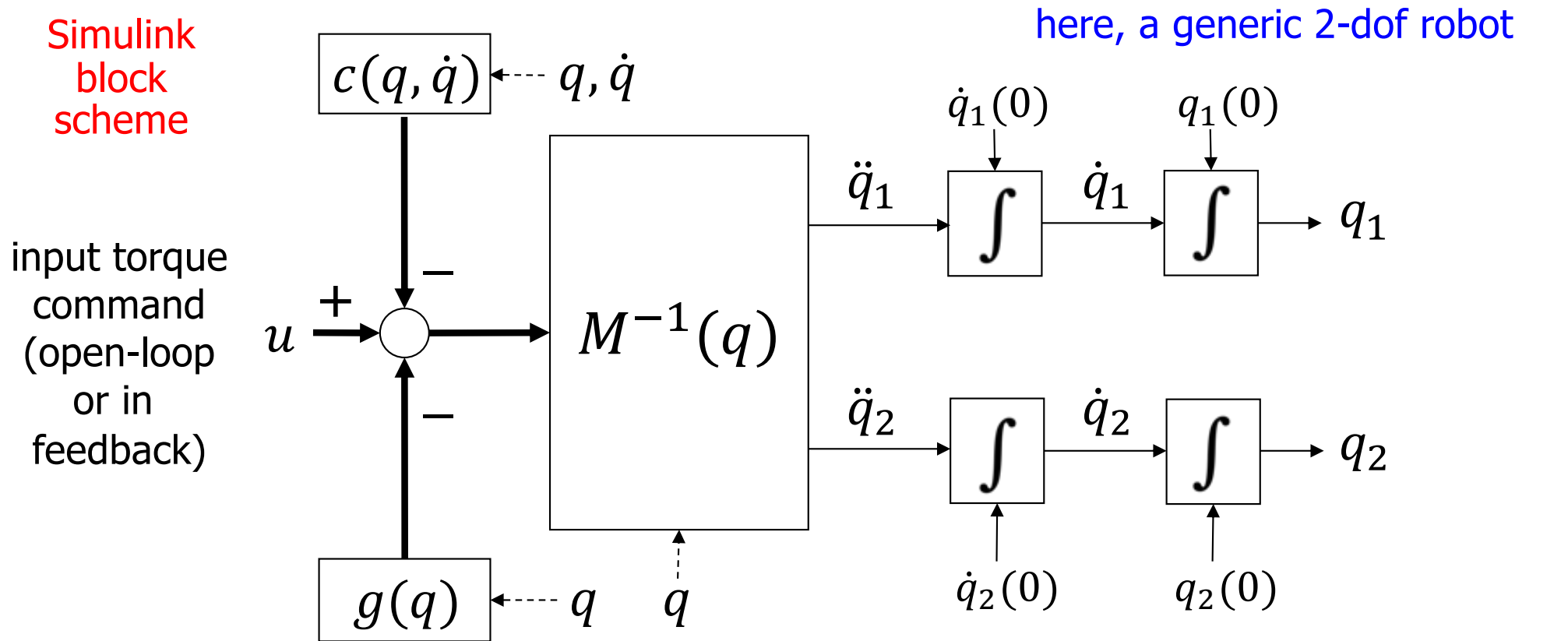
$$\tilde{x} = \begin{pmatrix} q \\ M(q)\dot{q} \end{pmatrix}$$

generalized
momentum

$\dot{\tilde{x}} = \dots$ (do it as exercise)



Dynamic simulation



including "inv(M)"

- initialization (dynamic coefficients and initial state)
- calls to (user-defined) Matlab functions for the evaluation of model terms
- choice of a numerical integration method (and of its parameters)

e.g., 4th-order Runge-Kutta (ode45)



Approximate linearization

- we can derive a **linear** dynamic model of the robot, which is valid **locally** around a given operative condition
 - useful for analysis, design, and **gain tuning** of linear (or, the linear part of) control laws
 - approximation by Taylor series expansion, up to the first order
 - linearization around a (constant) **equilibrium state** or along a (nominal, time-varying) **equilibrium trajectory**
 - usually, we work with (nonlinear) state equations; for mechanical systems, it is more convenient to directly use the **2nd order model**
 - same result, but easier derivation

constant state and input

$$\dot{x} = f(x) \rightarrow \dot{x} = f(x, u)$$

$$(x_e, u_e) : f(x_e, u_e) = 0 \Rightarrow \dot{x} = f(x_e, u_e) +$$

$$\frac{\partial f}{\partial x} \bigg|_{x=x_e, u=u_e} +$$

$$\frac{\partial f}{\partial u} \bigg|_{x=x_e, u=u_e}$$

and impose this

equilibrium **state** $(q, \dot{q}) = (q_e, 0) \quad [\ddot{q} = 0] \quad \longrightarrow \quad g(q_e) = u_e$

equilibrium **trajectory** $(q, \dot{q}) = (q_d(t), \dot{q}_d(t)) \quad [\ddot{q} = \ddot{q}_d(t)]$

$\longrightarrow \quad M(q_d)\ddot{q}_d + c(q_d, \dot{q}_d) + g(q_d) = u_d$



Linearization at an equilibrium state

- variations around an equilibrium state

$$q = q_e + \Delta q \quad \dot{q} = \cancel{\dot{q}_e} + \dot{\Delta q} = \dot{\Delta q} \quad \ddot{q} = \cancel{\ddot{q}_e} + \ddot{\Delta q} = \ddot{\Delta q} \quad u = u_e + \Delta u$$

- keeping into account the **quadratic** dependence of c terms on velocity (thus, neglected around the zero velocity)

$$M(q_e)\ddot{\Delta q} + \underbrace{g(\cancel{q_e}) + \frac{\partial g}{\partial q}\bigg|_{q=q_e}}_{G(q_e)} \Delta q + o(\|\Delta q\|, \|\dot{\Delta q}\|) = \cancel{u_e} + \Delta u$$

infinitesimal terms
of second or higher order

- in state-space format, with $\Delta x = \begin{pmatrix} \Delta q \\ \dot{\Delta q} \end{pmatrix}$

$$\dot{\Delta x} = \begin{pmatrix} 0 & I \\ -M^{-1}(q_e)G(q_e) & 0 \end{pmatrix} \Delta x + \begin{pmatrix} 0 \\ M^{-1}(q_e) \end{pmatrix} \Delta u = A \Delta x + B \Delta u$$



Linearization along a trajectory

- variations around an equilibrium trajectory

$$q = q_d + \Delta q \quad \dot{q} = \dot{q}_d + \dot{\Delta}q \quad \ddot{q} = \ddot{q}_d + \ddot{\Delta}q \quad u = u_d + \Delta u$$

- developing to 1st order the terms in the dynamic model ...

$$M(q_d + \Delta q)(\ddot{q}_d + \ddot{\Delta}q) + c(q_d + \Delta q, \dot{q}_d + \dot{\Delta}q) + g(q_d + \Delta q) = u_d + \Delta u$$

$$M(q_d + \Delta q) \cong M(q_d) + \sum_{i=1}^N \left. \frac{\partial M_i}{\partial q} \right|_{q=q_d} e_i^T \Delta q$$

i -th row of the identity matrix

$$g(q_d + \Delta q) \cong g(q_d) + G(q_d)\Delta q$$

$C_1(q_d, \dot{q}_d)$

$$c(q_d + \Delta q, \dot{q}_d + \dot{\Delta}q) \cong c(q_d, \dot{q}_d) + \underbrace{\left. \frac{\partial c}{\partial q} \right|_{\substack{q=q_d \\ \dot{q}=\dot{q}_d}}}_{C_1(q_d, \dot{q}_d)} \Delta q + \underbrace{\left. \frac{\partial c}{\partial \dot{q}} \right|_{\substack{q=q_d \\ \dot{q}=\dot{q}_d}}}_{C_2(q_d, \dot{q}_d)} \dot{\Delta}q$$



Linearization along a trajectory (cont)

- after simplifications ...

$$M(q_d)\ddot{\Delta q} + C_2(q_d, \dot{q}_d)\dot{\Delta q} + D(q_d, \dot{q}_d, \ddot{q}_d)\Delta q = \Delta u$$

with

$$D(q_d, \dot{q}_d, \ddot{q}_d) = G(q_d) + C_1(q_d, \dot{q}_d) + \sum_{i=1}^N \left. \frac{\partial M_i}{\partial q} \right|_{q=q_d} \ddot{q}_d e_i^T$$

- in state-space format

$$\begin{aligned} \dot{\Delta x} = & \begin{pmatrix} 0 & I \\ -M^{-1}(q_d)D(q_d, \dot{q}_d, \ddot{q}_d) & -M^{-1}(q_d)C_2(q_d, \dot{q}_d) \end{pmatrix} \Delta x \\ & + \begin{pmatrix} 0 \\ M^{-1}(q_d) \end{pmatrix} \Delta u = A(t) \Delta x + B(t) \Delta u \end{aligned}$$

a linear, but **time-varying** system!!

Coordinate transformation

$$q \in \mathbb{R}^N \quad M(q)\ddot{q} + c(q, \dot{q}) + g(q) = M(q)\ddot{q} + n(q, \dot{q}) = u_q$$

1

if we wish/need to use a **new** set of generalized coordinates p

$$\begin{array}{lll}
 p \in \mathbb{R}^N & p = f(q) & \longrightarrow q = f^{-1}(p) \\
 & \dot{p} = \frac{\partial f}{\partial q} \dot{q} = J(q)\dot{q} & \longrightarrow \dot{q} = J^{-1}(q)\dot{p} \\
 & \ddot{p} = J(q)\ddot{q} + \dot{J}(q)\dot{q} & \longrightarrow \ddot{q} = J^{-1}(q)(\ddot{p} - \dot{J}(q)J^{-1}(q)\dot{p})
 \end{array}$$

by duality
(principle of virtual work)

$$u_q = J^T(q)u_p$$

1

non symmetric

$$M(q)J^{-1}(q)\ddot{p} - M(q)J^{-1}(q)\dot{J}(q)J^{-1}(q)\dot{p} + n(q, \dot{q}) = J^T(q)u_p$$

$$J^{-T}(q) \cdot$$

pre-multiplying the whole equation...



Robot dynamic model after coordinate transformation

$$J^{-T}(q)M(q)J^{-1}(q)\ddot{p} + J^{-T}(q)(n(q, \dot{q}) - M(q)J^{-1}(q)\dot{J}(q)J^{-1}(q)\dot{p}) = u_p$$

$$q \rightarrow p$$

for actual computation,
these inner substitutions
are not necessary

$$(q, \dot{q}) \rightarrow (p, \dot{p})$$

non-conservative
generalized forces
performing work on p



$$M_p(p)\ddot{p} + c_p(p, \dot{p}) + g_p(p) = u_p$$

$$M_p = J^{-T} M J^{-1} \quad \begin{array}{l} \text{symmetric,} \\ \text{positive definite} \\ \text{(out of singularities)} \end{array} \quad g_p = J^{-T} g$$

$$c_p = J^{-T} (c - M J^{-1} \dot{J} J^{-1} \dot{p}) = J^{-T} c - M_p \dot{J} J^{-1} \dot{p} \quad \begin{array}{l} \text{quadratic} \\ \text{dependence on } \dot{p} \end{array}$$

$$c_p(p, \dot{p}) = S_p(p, \dot{p}) \dot{p} \quad \dot{M}_p - 2S_p \quad \text{skew-symmetric}$$

when $p = \text{E-E pose}$, this is the robot **dynamic model** in **Cartesian coordinates**

NOTE: in this case, we have implicitly assumed that $M = N$ (no redundancy!)

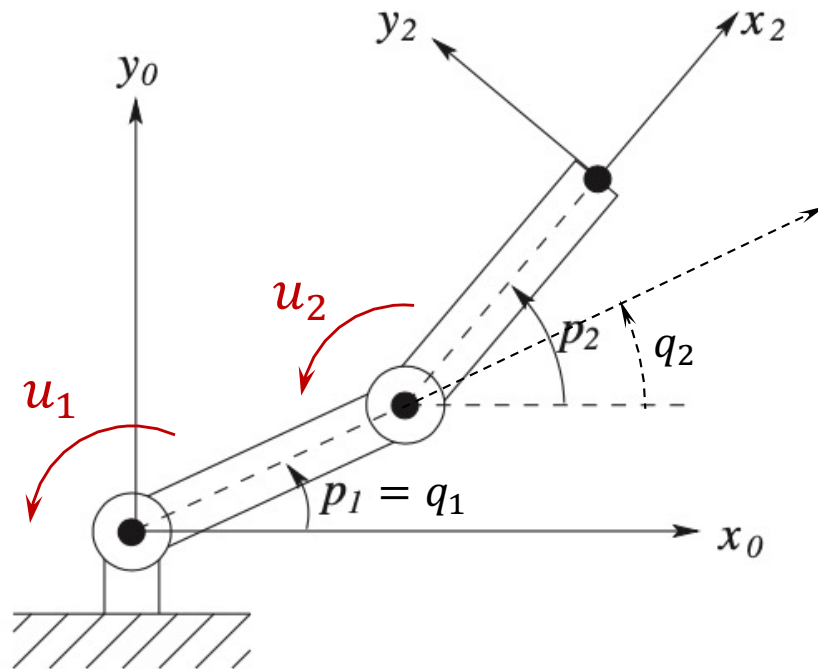
$$\dot{p} = k \dot{q}$$

$$p = k q$$



Example of coordinate transformation

planar 2R robot using absolute coordinates



- motor 1 at joint 1, motor 2 at joint 2
- in place of DH angles q , use the **absolute** angles $p_1 = q_1$ and $p_2 = q_1 + q_2$

$$p = \begin{pmatrix} 1 & 0 \\ 1 & 1 \end{pmatrix} q = J q \quad \text{linear transformation}$$

$$J^{-1} = \begin{pmatrix} 1 & 0 \\ -1 & 1 \end{pmatrix} \quad J^{-T} = \begin{pmatrix} 1 & -1 \\ 0 & 1 \end{pmatrix}$$

- from $M(q)\ddot{q} + c(q, \dot{q}) + g(q) = u_q$ obtained with DH **relative** coordinates

$$M_p = J^{-T} M J^{-1} = \begin{pmatrix} a_1 - a_3 - 2a_2 c_2 & a_2 c_2 \\ a_2 c_2 & a_3 \end{pmatrix} \quad g_p = J^{-T} g = \begin{pmatrix} a_4 c_1 \\ a_5 c_{12} \end{pmatrix}$$

$$c_p = J^{-T} c = \begin{pmatrix} -2a_2 s_2 (\dot{q}_1 + \dot{q}_2)^2 \\ a_2 s_2 \dot{q}_1^2 \end{pmatrix} \quad u_p = J^{-T} u_q = \begin{pmatrix} u_1 - u_2 \\ u_2 \end{pmatrix}$$



Robot dynamic model

in the task/Cartesian space, with redundancy

dynamic model in the joint space

$$M(q)\ddot{q} + n(q, \dot{q}) = \tau$$

$$q \in \mathbb{R}^N$$

$$r = f(q) \in \mathbb{R}^M$$

$$M < N$$

second-order task kinematics

$$\ddot{r} = J(q)\ddot{q} + \dot{J}(q)\dot{q}$$

assumption J is full rank = M

1) isolate the joint acceleration from the dynamics $\Rightarrow \ddot{q} = M^{-1}(q) (\tau - n(q, \dot{q}))$

2) decompose the joint torques in two complementary spaces

$$\tau = J^T(q)F + (I - J^T(q)H(q))\tau_0 \quad H \text{ is any generalized inverse of } J^T$$

$$\in \mathcal{R}(J^T)$$

$$\notin [\mathcal{R}(J^T)]$$

$$J^T H J^T = J^T$$

torques coming from
generalized forces
in the task space ...

... and joint torques τ_0 that do not!

$$\Rightarrow \tau_0 = J^T(q)F, \forall F \in \mathbb{R}^M \Rightarrow (I - J^T(q)H(q))J^T(q)F = 0$$

3) substitute 1) and 2) in the differential task kinematics

$$\Rightarrow \ddot{r} = \underbrace{J(q)M^{-1}(q)J^T(q)}_{\text{task Jacobian}} F + (I - J^T(q)H(q))\tau_0 - n(q, \dot{q}) + \dot{J}(q)\dot{q}$$

4) isolate on the right-hand side the generalized forces F in the task space ...



Robot dynamic model

in the task/Cartesian space, with redundancy

$$\rightarrow (J(q)M^{-1}(q)J^T(q))^{-1}\ddot{r} = F + (J(q)M^{-1}(q)J^T(q))^{-1}(J(q)M^{-1}(q)((I - J^T(q)H(q))\tau_0 - n(q, \dot{q})) + \dot{J}(q)\dot{q})$$

5) choose as generalized inverse $H = (JM^{-1}J^T)^{-1}JM^{-1} = (J_M^\#)^T$, i.e., the transpose of the **inertia-weighted pseudoinverse** of the task Jacobian (see block of slides #2)

\rightarrow in this way, the joint torque component τ_0 will **NOT** affect the task acceleration \ddot{r}

$$(J(q)M^{-1}(q)J^T(q))^{-1}\ddot{r} = F + (J(q)M^{-1}(q)J^T(q))^{-1}(\dot{J}(q)\dot{q} - J(q)M^{-1}(q)n(q, \dot{q}))$$

6) the resulting (M –dimensional) task dynamics is then

$$M_r(q)\ddot{r} + n_r(q, \dot{q}) = F \dots + F_{ext}$$

external forces can be added on the rhs of the equations in a **dynamically consistent** way!

with

$$\left. \begin{aligned} M_r(q) &= (J(q)M^{-1}(q)J^T(q))^{-1} \text{ task inertia matrix} \\ n_r(q, \dot{q}) &= M_r(q)(J(q)M^{-1}(q)n(q, \dot{q}) - \dot{J}(q)\dot{q}) \end{aligned} \right\} \text{ for } M = N, \text{ these terms are identical to slide \#27}$$

7) an additional ($N - M$)-dimensional dynamics is needed to describe the full robot!



Dynamic scaling of trajectories

uniform time scaling of motion

- given a smooth **original trajectory** $q_d(t)$ of motion for $t \in [0, T]$
 - suppose to **rescale** time as $t \rightarrow r(t)$ (a strictly **increasing** function of t)
 - in the new time scale, the scaled trajectory $q_s(r)$ satisfies

$$q_d(t) = q_s(r(t)) \quad \xrightarrow{\text{same path executed (at different instants of time)}} \quad \dot{q}_d(t) = \frac{dq_d}{dt} = \frac{dq_s}{dr} \frac{dr}{dt} = q'_s(r) \dot{r}(t)$$
$$\ddot{q}_d(t) = \frac{d\dot{q}_d}{dt} = \left(\frac{dq'_s}{dr} \frac{dr}{dt} \right) \dot{r} + q'_s \frac{d\dot{r}}{dt} = q''_s(r) \dot{r}^2(t) + q'_s(r) \ddot{r}(t)$$

- **uniform scaling** of the trajectory occurs when $r(t) = kt$

$$\dot{q}_d(t) = k q'_s(kt) \quad \ddot{q}_d(t) = k^2 q''_s(kt)$$

Q: what is the new **input torque** needed to execute the **scaled** trajectory?
(suppose **dissipative** terms can be **neglected**)



Dynamic scaling of trajectories

inverse dynamics under uniform time scaling

- the new torque could be recomputed through the inverse dynamics, for every $r = kt \in [0, T'] = [0, kT]$ along the scaled trajectory, as

$$\tau_s(kt) = M(q_s)q_s'' + c(q_s, q_s') + g(q_s)$$

- however, being the dynamic model **linear** in the acceleration and **quadratic** in the velocity, it is

$$\begin{aligned}\tau_d(t) &= M(q_d)\ddot{q}_d + c(q_d, \dot{q}_d) + g(q_d) = M(q_s)k^2q_s'' + c(q_s, kq_s') + g(q_s) \\ &= k^2(M(q_s)q_s'' + c(q_s, q_s')) + g(q_s) = k^2(\tau_s(kt) - g(q_s)) + g(q_s)\end{aligned}$$

- thus, saving separately the total torque $\tau_d(t)$ and gravity torque $g_d(t)$ in the inverse dynamics computation along the **original** trajectory, the **new input torque** is obtained **directly** as

$$\tau_s(kt) = \frac{1}{k^2} (\tau_d(t) - g(q_d(t))) + g(q_d(t))$$

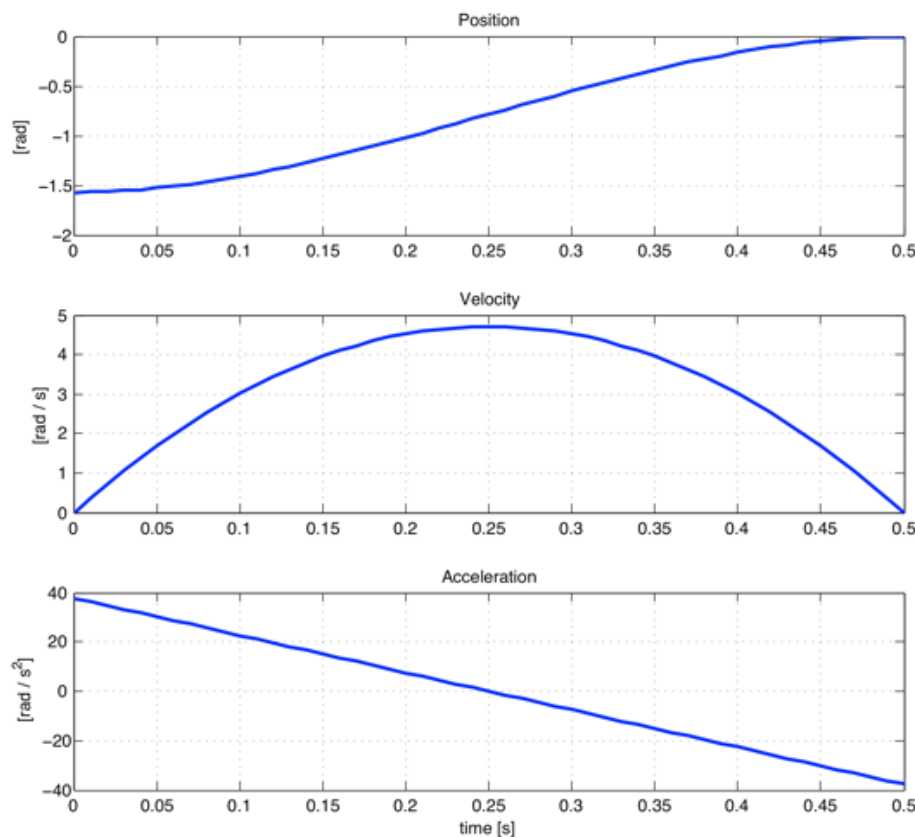
$k > 1$: slow down
⇒ reduce torque
 $k < 1$: speed up
⇒ increase torque

gravity term (only position-dependent): does **NOT** scale!

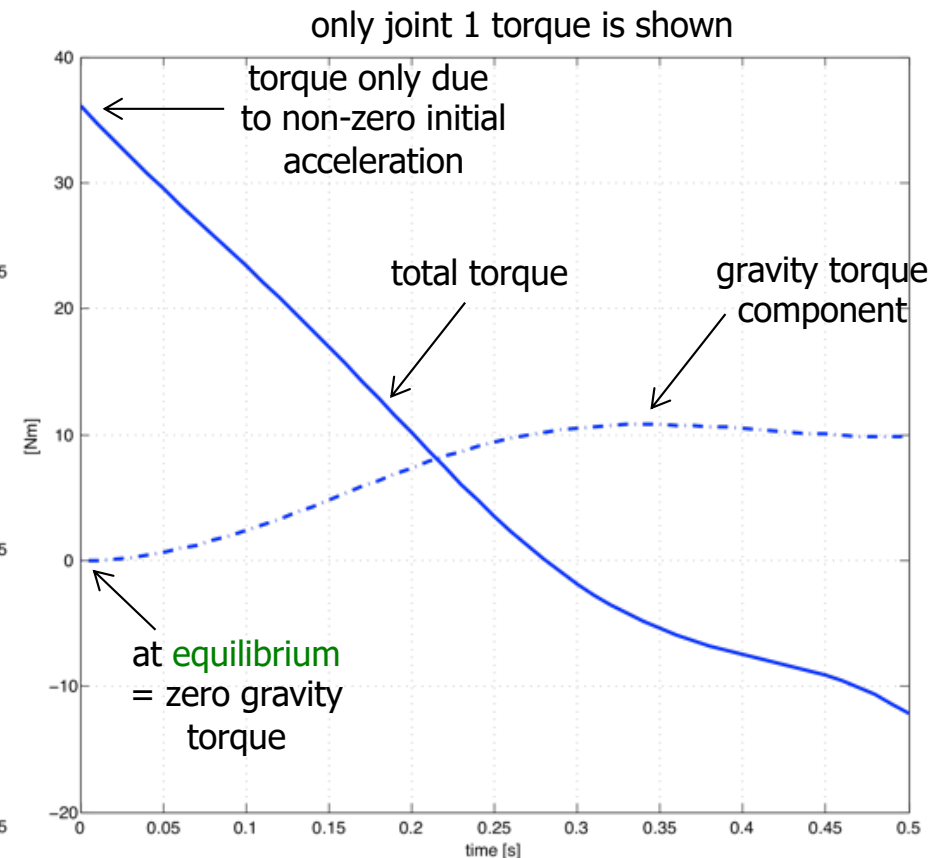
Dynamic scaling of trajectories

numerical example

- rest-to-rest motion with cubic polynomials for planar 2R robot under gravity (from downward **equilibrium** to horizontal link 1 & upward vertical link 2)
- original trajectory lasts $T = 0.5$ s (but maybe violates the torque limit at joint 1)



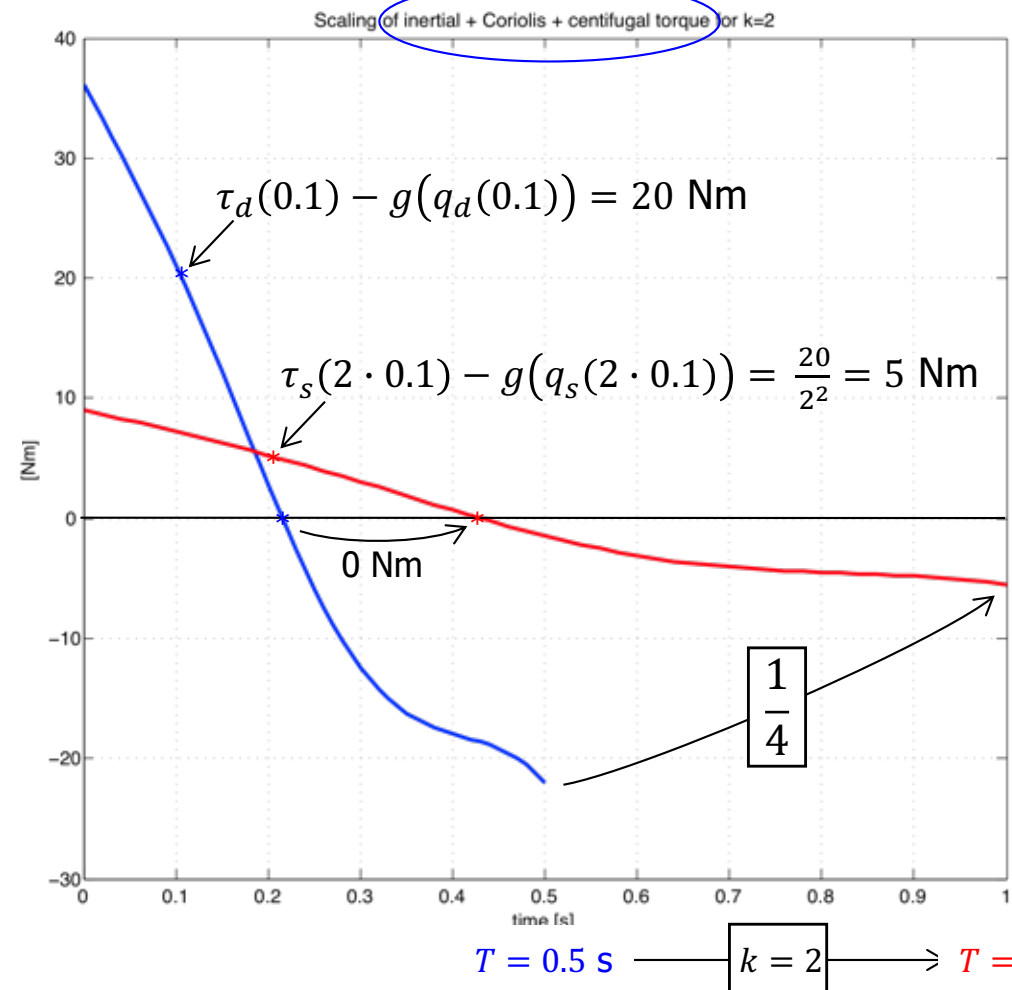
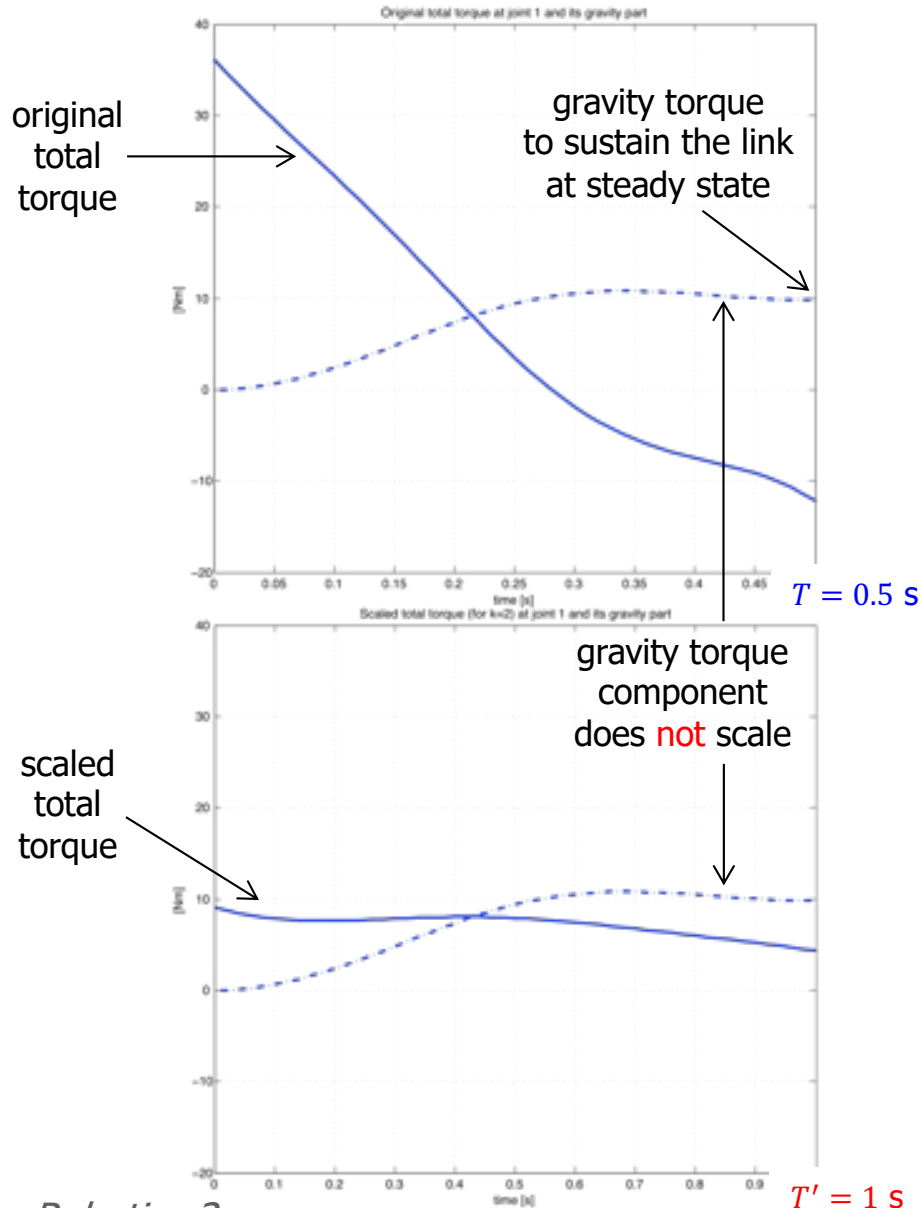
for **both** joints



Dynamic scaling of trajectories

numerical example

- scaling with $k = 2$ (slower) $\rightarrow T' = 1$ s

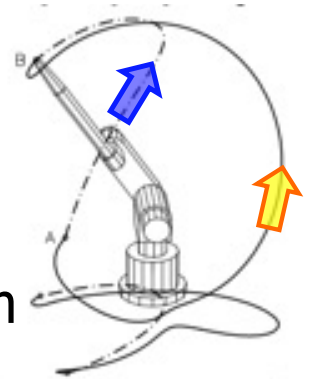


Optimal point-to-point robot motion

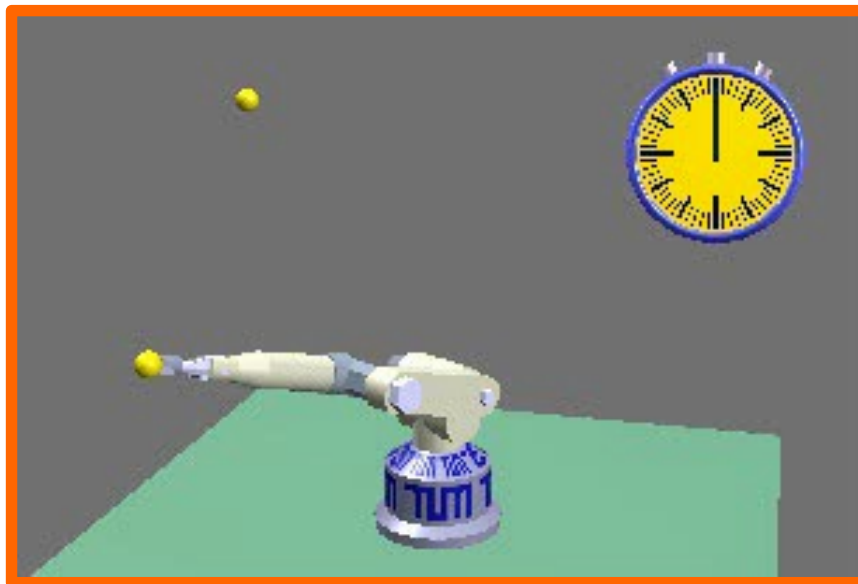
considering the dynamic model



- given the initial and final robot configurations (at rest) and actuator torque bounds, find
 - the **minimum-time** T_{\min} motion
 - the (global/integral) **minimum-energy** E_{\min} motionand the associated **command torques** needed to execute them
- a complex nonlinear optimization problem solved **numerically**

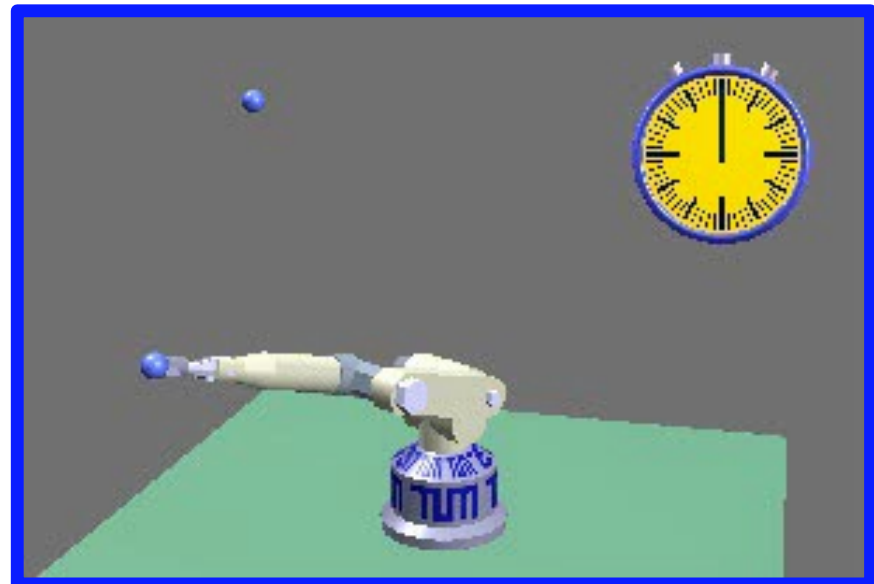


video



$T_{\min} = 1.32 \text{ s}$, $E = 306$

video



$T = 1.60 \text{ s}$, $E_{\min} = 6.14$