

# Field Transformation HMC

Xiao-Yong Jin for Lattice QCD ECP

April 20, 2021

## Contents

<b>1</b>	<b>Introduction</b>	<b>1</b>
<b>2</b>	<b>Generalized field transformation</b>	<b>3</b>
<b>3</b>	<b>Test with 2D U(1) model without explicit neural networks</b>	<b>4</b>
<b>4</b>	<b>Discussions and plans</b>	<b>5</b>

## Versions

- April 20, 2021, version 0  
Some text borrowed from the previous 1D XY model report.

## 1 Introduction

Following the arguments of “Trivializing maps” [1], to evaluate,

$$\langle O \rangle = 1/Z \int dx O(x) e^{-S(x)}, \quad (1)$$

we perform a change of variable,

$$x = F(y) \quad (2)$$

with a vector functions  $F$ , we have

$$\langle O \rangle = \frac{1}{Z} \int dy |\det[J(y)]| O(F(y)) e^{-S(F(y))}, \quad (3)$$

where the Jacobian matrix,

$$J(y) = \frac{\partial F(y)}{\partial y}. \quad (4)$$

$F$  has to satisfy,

- Injective (1 to 1), from the new integration domain to the old.
- Continuously differentiable (or differentiable and have continuous inverse).

Rewrite the integral as,

$$\langle O \rangle = \frac{1}{Z} \int d\mathcal{Y} O(F(\mathcal{Y})) e^{-S(F(\mathcal{Y})) + \ln|\det[J(\mathcal{Y})]|}. \quad (5)$$

$F$  is a trivializing map, when

$$S(F(\mathcal{Y})) - \ln|\det[J(\mathcal{Y})]| = \text{constant} \quad (6)$$

and our expectation value simplifies to

$$\langle O \rangle = \frac{1}{Z} \int d\mathcal{Y} O(F(\mathcal{Y})). \quad (7)$$

In terms of HMC, we add the conjugate momenta,  $\pi$ , and use the equations of motion derived from the Hamiltonian,

$$\mathcal{H}(\mathcal{Y}, \pi) = \frac{1}{2} \pi^2 + S(F(\mathcal{Y})) - \ln|\det[J(\mathcal{Y})]|, \quad (8)$$

as

$$\frac{d}{dt} \pi = -\frac{\partial}{\partial \mathcal{Y}} \mathcal{H} = -J(\mathcal{Y}) S'(F(\mathcal{Y})) + \text{tr} \left[ J^{-1} \frac{d}{d\mathcal{Y}} J \right], \quad (9)$$

$$\frac{d}{dt} \mathcal{Y} = \frac{\partial}{\partial \pi} \mathcal{H} = \pi. \quad (10)$$

This is separable and can use the usual explicit, symplectic and symmetric discrete integrators.

Consider a change of variable for  $\pi$  and  $\mathcal{Y}$ ,

$$\pi = J(\mathcal{Y}) p = J(F^{-1}(x)) p, \quad (11)$$

$$\mathcal{Y} = F^{-1}(x), \quad (12)$$

with the Jacobian matrix of determinant 1,

$$\mathcal{J}(p, x) = \begin{bmatrix} J(F^{-1}(x)) & \frac{\partial}{\partial x} J(F^{-1}(x)) p \\ 0 & \frac{\partial}{\partial x} F^{-1}(x) \end{bmatrix}, \quad (13)$$

$$\det[\mathcal{J}] = 1. \quad (14)$$

We get a new Hamiltonian from equation (8),

$$\tilde{\mathcal{H}}(x, p) = \frac{1}{2} p^\dagger M p + S(x) - \ln|\det[J]| \quad (15)$$

where the positive definite  $M$ ,

$$M(x) = J^\dagger(F^{-1}(x)) J(F^{-1}(x)) \quad (16)$$

is the kernel of the kinetic term considered Duane et al [2, 3].

## 2 Generalized field transformation

For a generic field transformation, we can take inspiration from an MD update, where

$$U' \leftarrow U \exp[dt * \text{Proj}_{\text{TAH}}(M)]. \quad (17)$$

$U$  is covariant and  $M$  (being sum of loops) is invariant under gauge transformation.

The most generic form of a field transformation could be

$$U(x, \mu)' \leftarrow \text{Proj}_{\text{SU}} \sum_L a_L \prod_{\text{ord}} [L(x, \mu)], \quad (18)$$

where  $L(x, \mu)$  is any line connecting gauge links from  $x$  to  $x + \hat{\mu}$ ,  $a_L$  is a scalar coefficient.  $L(x, \mu)$  may or may not need to go through  $U(x, \mu)$ . we may also have loops in it, and results in a polynomial of such loop after sum.

The  $\text{Proj}_{\text{SU}}$  has to satisfy the constraint that

$$X \text{Proj}_{\text{SU}}(M) Y = \text{Proj}_{\text{SU}}(XMY) \quad (19)$$

for any  $X$  and  $Y$  in the group.

We can, however, rewrite the equation to be similar to the MD update,

$$U(x, \mu)' \leftarrow \text{Proj}_{\text{SU}} \sum_L a_L U(x, \mu) U(x, \mu)^\dagger \prod_{\text{ord}} [L(x, \mu)], \quad (20)$$

such that

$$U(x, \mu)^\dagger \prod_{\text{ord}} [L(x, \mu)] \quad (21)$$

becomes a complete loop, which is gauge invariant. We then have it in a simplified form,

$$U' \leftarrow \text{Proj}_{\text{SU}} \sum_R a_R UR, \quad (22)$$

where  $R$  is any loop that start and stop at  $x + \hat{\mu}$ .

We may write it in terms of

$$U' \leftarrow \text{Proj}_{\text{SU}} \sum_{LR} a_{LR} LUR, \quad (23)$$

such that  $L$  is any loop that start and stop at  $x$ , but since we can always pull the  $U$  to the left by multiplying  $UU^\dagger$ , the most generic form is simply,

$$U' \leftarrow \text{Proj}_{\text{SU}} [U(\sum_R a_R R)], \quad (24)$$

where  $R$  is any loop that start and stop at  $x + \hat{\mu}$ , and may or may not go through  $U$ .

We can multiply  $UU^\dagger$  again, so it becomes

$$U' \leftarrow UU^\dagger \text{Proj}_{\text{SU}} [U(\sum_R a_R R)]. \quad (25)$$

If we had a  $\text{Proj}_{\text{SU}}$  that satisfies

$$X \text{Proj}_{\text{SU}}(M)Y = \text{Proj}_{\text{SU}}(XMY) \quad (26)$$

the above update would become

$$U' \leftarrow U \text{Proj}_{\text{SU}}[\sum_R a_R R] \quad (27)$$

though we need  $\text{Proj}_{\text{SU}}$  satisfy a different constraint,

$$X \text{Proj}_{\text{SU}}(M)X^\dagger = \text{Proj}_{\text{SU}}(XMX^\dagger) \quad (28)$$

If we were allowed to do the above change of constraint to  $\text{Proj}_{\text{SU}}$ , it seems we could just use the projection in MD update,

$$U' \leftarrow U \exp[\text{Proj}_{\text{TAH}}(\sum_R a_R R)] \quad (29)$$

It looks like we are only one step away from

$$U' \leftarrow U \exp \left[ \frac{d}{dU} F(\text{ReTr } A, \text{ReTr } B, \text{ReTr } C, \dots) \right] \quad (30)$$

where  $F$  is any analytical function or neural network, and  $A, B, C, \dots$  are any closed loops may or may not passing  $U$ .

We can simplify it by moving the  $U$  independent loops out and keeping the  $U$  dependent loops simple,

$$U' \leftarrow U \exp \left[ c \text{atan}[F(\text{ReTr } X, \text{ReTr } Y, \dots)] \left[ \frac{d}{dU} \text{ReTr}[W] \right] \right] \quad (31)$$

so that  $F(\text{ReTr } X, \text{ReTr } Y, \dots)$  only depends on  $U$  independent loops, while  $W$  contains  $U$  dependent loops,  $c * \text{atan}(F)$  is for restricting the Jacobian to be positive definite, and  $W$  is a sum of any loop and its symmetrized versions including  $U$ . The arbitrary function  $F$  can take the form of a neural network.

For a transformation to be usable in HMC, we want a tractable Jacobian determinant. One way to achieve this is updating gauge links in subsets, such that each update to the subset of  $U$  has a Jacobian matrix, where the only nonzero entries are on its diagonal.

### 3 Test with 2D U(1) model without explicit neural networks

This test uses  $F$  in equation (31) as a simple sum of similar loops, specifically plaquettes and rectangles.

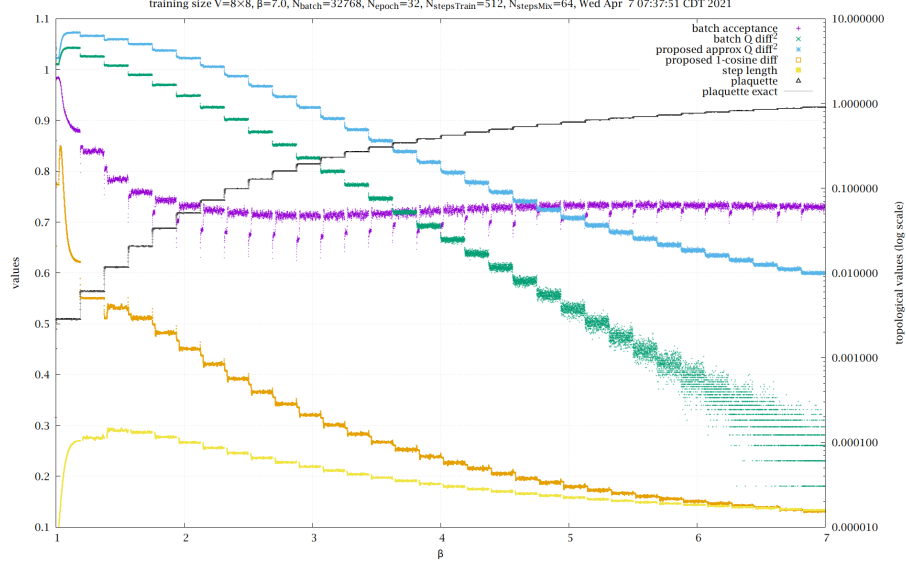


Figure 1: Annealed training steps for  $V = 8 \times 8$  for HMC with 10 leapfrog steps, The only trainable parameter is the step size.

The loss function is

$$l(y', \pi' | y, \pi) = -\frac{1}{N_{\text{batch}}} \sum_{\text{batch}} \max \left\{ 1, e^{H(y, \pi) - H(y', \pi')} \right\} \left( \lambda \frac{1}{N} \sum_x [1 - \cos(P'_x - P_x)] + \rho \left[ \sum_x \sin(P'_x) - \sum_x \sin(P_x) \right]^2 \right), \quad (32)$$

where  $(y', \pi')$  and  $(y, \pi)$  are respectively the proposed and initial configuration in the MD evolution,  $P_x = \theta_{x,1} + \theta_{x+1,2} - \theta_{x+2,1} - \theta_{x,2}$  is the plaquette phase at site  $x$ , and the sum of sin is an approximate of the topological charge. The test here uses  $\lambda = 0.1$  and  $\rho = 1.0$ .

## 4 Discussions and plans

- Test equation (31) with  $F$  parametrized by neural networks.
- Compute and monitor forces during HMC.
- Compute and monitor changes of momenta (integrated forces) during HMC.

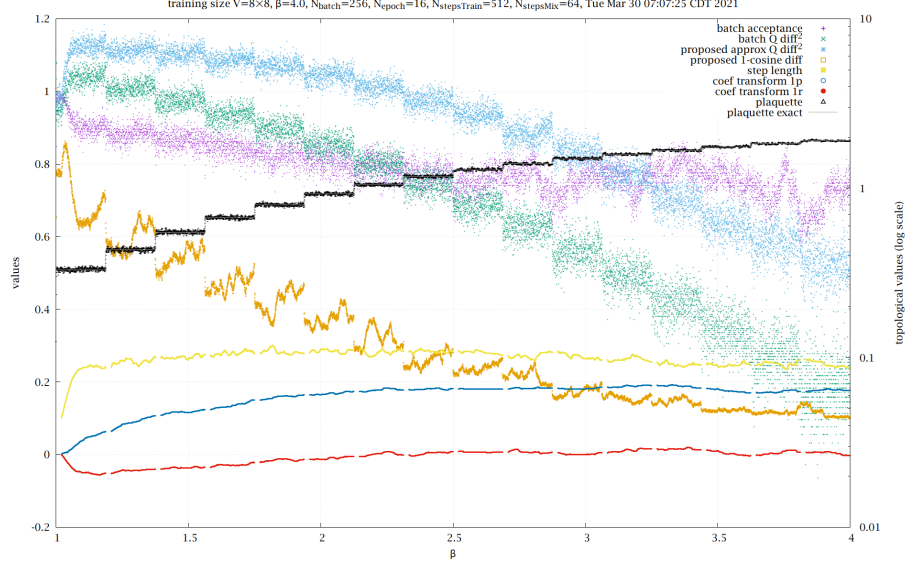


Figure 2: Annealed training steps for  $V = 8 \times 8$  with 10 leapfrog steps.

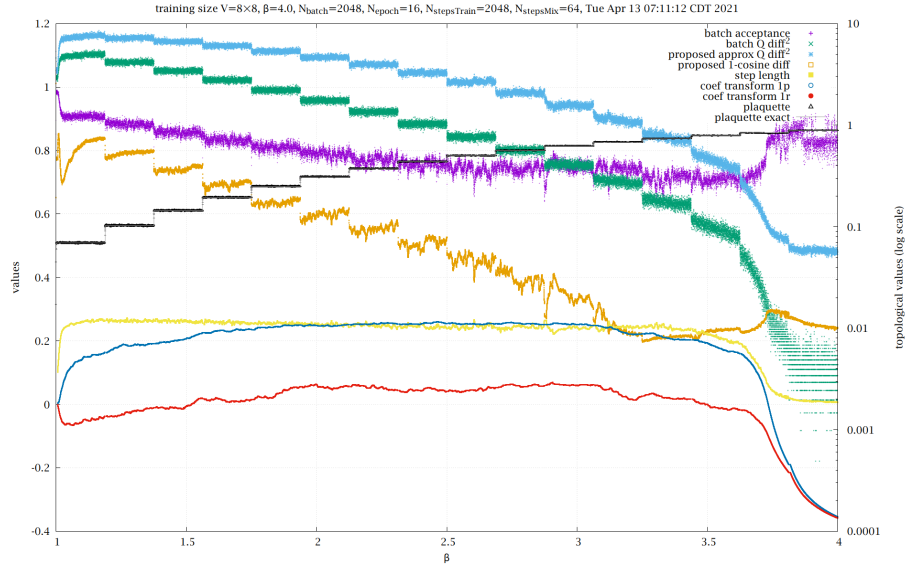


Figure 3: Annealed training steps for  $V = 8 \times 8$  with 10 leapfrog steps.

- Vary the number of leapfrog steps.

## References

- [1] Martin Lüscher. “Trivializing maps, the Wilson flow and the HMC algorithm”. In: *Commun. Math. Phys.* 293 (2010), pp. 899–919. DOI: 10.1007/s00220-009-0953-7. arXiv: 0907.5491 [hep-lat] (cit. on p. 1).
- [2] Simon Duane et al. “Acceleration of Gauge Field Dynamics”. In: *Phys. Lett. B* 176 (1986), p. 143. DOI: 10.1016/0370-2693(86)90940-8 (cit. on p. 2).
- [3] Simon Duane and Brian J. Pendleton. “GAUGE INVARIANT FOURIER ACCELERATION”. In: *Phys. Lett. B* 206 (1988), pp. 101–106. DOI: 10.1016/0370-2693(88)91270-1 (cit. on p. 2).