

52nd—24th INTERNATIONAL-RUDOLF ORTVAY PROBLEM SOLVING CONTEST IN PHYSICS

Problem 19

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In this problem, we set $4\pi\epsilon_0 = c = 1$. The coordinate of electron is performed by (r, θ) in the polar coordinate system.

The electronic potential of electron is: $V(r) = -\frac{Ze^2}{r}$. So we can write the Lagrangian of this electron [1] is:

$$\mathcal{L} = -m\sqrt{1 - \dot{r}^2 - r^2\dot{\theta}^2} + \frac{Ze^2}{r}.$$

This Lagrangian doesn't have any explicit dependence on θ . Therefore, the angular momentum of the electron is a constant L , so we can write:

$$L = \frac{\partial \mathcal{L}}{\partial \dot{\theta}} = m \frac{r^2 \dot{\theta}}{\sqrt{1 - \dot{r}^2 - r^2 \dot{\theta}^2}}. \quad (1)$$

Cause our Lagrangian doesn't have any explicit dependence on time t , conservation the energy of the electron:

$$E = \frac{m}{\sqrt{1 - \dot{r}^2 - r^2 \dot{\theta}^2}} - \frac{Ze^2}{r}. \quad (2)$$

From (1) and (2), we find $\dot{\theta}$ by r :

$$\dot{\theta} = \frac{L}{Er^2 + Ze^2 r}. \quad (3)$$

Let $r'_\theta = \frac{dr}{d\theta}$. So we can perform \dot{r} by r'_θ :

$$\dot{r} = \frac{dr}{d\theta} \dot{\theta}.$$

We write again (1):

$$\frac{\dot{r}}{\dot{\theta}} = \sqrt{\frac{1}{\dot{\theta}^2} - r^2 - \left(\frac{m}{L}\right)^2 r^4}.$$

And put $\dot{\theta}$ and \dot{r} in our new equation, we will have a relationship of r'_θ and r to find the orbital of the

electron:

$$\begin{aligned}
\frac{dr}{d\theta} &= \sqrt{\left(\frac{Er^2 - Ze^2r}{L}\right)^2 - r^2 - \left(\frac{m}{L}\right)^2 r^4} \\
\Rightarrow \theta &= \int \frac{dr}{\sqrt{\frac{E^2 - m^2}{L^2} r^4 - \frac{2EZe^2}{L^2} r^3 + \left[\left(\frac{Ze^2}{L}\right)^2 - 1\right] r^2}} \\
\Rightarrow \theta &= \frac{1}{\sqrt{1 - \left(\frac{Ze^2}{L}\right)^2}} \arccos \frac{\left[1 - \left(\frac{Ze^2}{L}\right)^2\right]^{\frac{1}{r}} - \frac{EZe^2}{L^2}}{\sqrt{\left(\frac{E}{L}\right)^2 - \left(\frac{m}{L}\right)^2 \left[1 - \left(\frac{Ze^2}{L}\right)^2\right]}} + C.
\end{aligned}$$

Choose the coordinate system with $C = 0$, so the Orbital of the electron is:

$$r = \frac{\frac{L}{E} \left(\frac{L}{Ze^2} - \frac{Ze^2}{L} \right)}{1 + \sqrt{\left(\frac{L}{Ze^2}\right)^2 + \left(\frac{m}{E}\right)^2 - \left(\frac{mL}{EZe^2}\right)^2} \cos \left\{ \sqrt{1 - \left(\frac{Ze^2}{L}\right)^2} \theta \right\}} = \frac{p}{1 + \epsilon \cos(\alpha\theta)}.$$

So, the electron will move in a flower orbital.

In the SI units ¹:

$$r = \frac{\frac{Lc}{E} \left(\frac{4\pi\epsilon_0 cL}{Ze^2} - \frac{Ze^2}{4\pi\epsilon_0 cL} \right)}{1 + \sqrt{\left(\frac{4\pi\epsilon_0 cL}{Ze^2}\right)^2 + \left(\frac{mc^2}{E}\right)^2 - \left(\frac{4\pi\epsilon_0 c^3 mL}{EZe^2}\right)^2} \cos \left\{ \sqrt{1 - \left(\frac{Ze^2}{4\pi\epsilon_0 cL}\right)^2} \theta \right\}}.$$

In reality, with $L \sim \hbar$, $Z \sim 31$, then $\left(\frac{Ze^2}{4\pi\epsilon_0 cL}\right)^2 \sim 0.05$ and with bigger Z , we can't use the classical approximation or define the semi-axes of ellipse. Therefore, we can't make a new Kepler's 3rd law with the semi-axes. Thus, we will use $p = \frac{L}{E} \left(\frac{L}{Ze^2} - \frac{Ze^2}{L} \right)$, $\epsilon = \sqrt{\left(\frac{L}{Ze^2}\right)^2 + \left(\frac{m}{E}\right)^2 - \left(\frac{mL}{EZe^2}\right)^2}$ and $\alpha = \sqrt{1 - \left(\frac{Ze^2}{4\pi\epsilon_0 cL}\right)^2}$ to make a new Kepler's 3rd law.

In this case, the orbiting time is the time T from $\theta = 0$ to $\theta = \frac{2\pi}{\alpha}$.

a) From (3):

$$\begin{aligned}
dt &= \left[\left(\frac{E}{L}\right) r^2 + \left(\frac{Ze^2}{L}\right) r \right] d\theta \\
&= \frac{p}{\sqrt{1 - \alpha^2}} \frac{1 + (1 - \alpha^2) \cos(\alpha\theta)}{[1 + \epsilon \cos(\alpha\theta)]^2} d\theta \\
\Rightarrow T &= \frac{p}{\sqrt{1 - \alpha^2}} \int_0^{\frac{2\pi}{\alpha}} \frac{1 + (1 - \alpha^2) \epsilon \cos(\alpha\theta)}{[1 + \epsilon \cos(\alpha\theta)]^2} d\theta \\
\Leftrightarrow T &= \frac{p}{\alpha \sqrt{1 - \alpha^2}} \int_0^{2\pi} \frac{1 + (1 - \alpha^2) \epsilon \cos x}{[1 + \epsilon \cos x]^2} dx \\
&= 2\pi p \frac{1 - (1 - \alpha^2) \epsilon^2}{\alpha \sqrt{1 - \alpha^2} (1 - \epsilon^2)^{\frac{3}{2}}}.
\end{aligned}$$

¹We only use the SI units to calculate some values and compare them, after that, we will back to our units

b) In the electron frame:

$$\begin{aligned}
dt' &= dt \sqrt{1 - \dot{r}^2 - r^2 \dot{\theta}^2} \\
&= \frac{m}{L} r^2 d\theta \\
&= p\alpha \sqrt{\frac{(1 - \alpha^2) \epsilon^2 - 1}{1 - \alpha^2}} \frac{d\theta}{[1 + \epsilon \cos(\alpha\theta)]^2} \\
\Rightarrow T' &= p \sqrt{\frac{(1 - \alpha^2) \epsilon^2 - 1}{1 - \alpha^2}} \int_0^{2\pi} \frac{dx}{[1 + \epsilon \cos x]^2} \\
&= 2\pi p \frac{\sqrt{1 - (1 - \alpha^2) \epsilon^2}}{\sqrt{1 - \alpha^2} (1 - \epsilon^2)^{\frac{3}{2}}}.
\end{aligned}$$

References

- [1] A. S. Kompaneyets (1978), *A Course Of Theoretical Physics, Vol. 1 Fundamental Laws*, Mir titles.
- [2] Lim Yung-Kou (1994), *Problems and Solutions on Mechanics*, WORLD SCIENTIFIC, pp. 738-741.