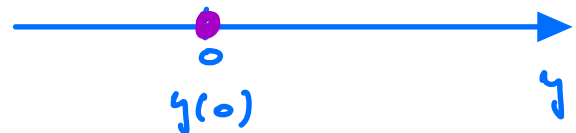


## Lecture 13

Let's assume we have a mechanical/electrical system at rest subject to an external force  $f(t)$

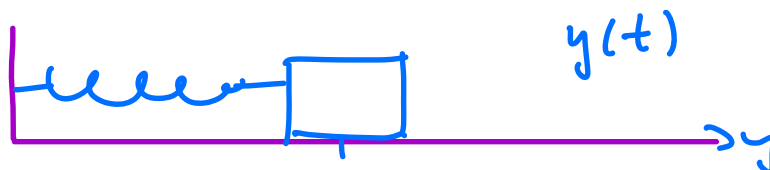
$$y'' + ay' + by = f(t), \quad t \geq 0$$

$$\left. \begin{array}{l} y(0) = 0 \\ y'(0) = 0 \end{array} \right\} \text{at rest I.C.s.}$$



The external force could be a sudden hammer blow in the mechanical system or a lightning stroke on a transmission line.

So, our external force can be a very irregular function.



0

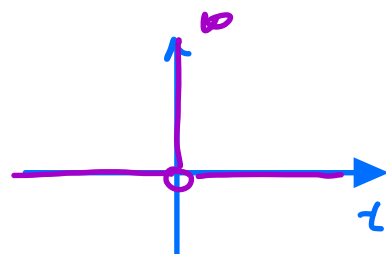
# Dirac Delta "distribution" function

## Def (Dirac Delta function)

- P. Dirac
- L. Schwartz  
1950's

The Dirac Delta function (distribution) is characterized by the following two properties:

$$(1) \quad \delta(t) = \begin{cases} 0, & \text{if } t \neq 0 \\ \infty, & \text{if } t = 0 \end{cases}$$



$$(2) \quad \int_{-\infty}^{\infty} \delta(t) f(t) dt = f(0)$$

for any function that is cont. on an open interval containing 0.

Remark:

a) By property (2), for  $a \geq 0$

$$\begin{aligned} \int_{-\infty}^{\infty} f(t) \delta(t-a) dt &= \int_{-\infty}^{\infty} f(u+a) \delta(u) du \\ &= f(a) \end{aligned}$$

$$u = t - a \rightarrow t = u + a$$

$$du = dt$$

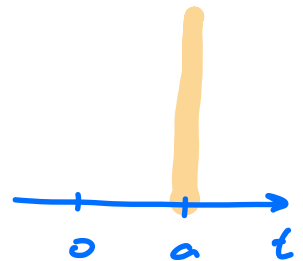
b) let's take  $f(t) \equiv 1$

$$\int_{-\infty}^{\infty} 1 \delta(t) dt = 1$$

$$\text{so } \boxed{\int_{-\infty}^{\infty} \delta(t) dt = 1}$$



c) L.T. of  $\delta(t-a)$  for  $a \geq 0$ ?



$$\mathcal{L}\{\delta(t-a)\}(s) = \int_{-\infty}^{\infty} e^{-st} \delta(t-a) dt$$

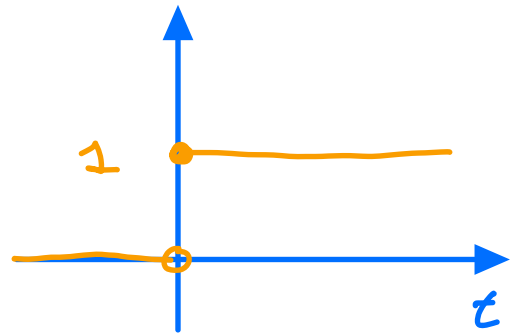
$$= \int_{-\infty}^{\infty} e^{-st} \delta(t-a) dt$$

$$= e^{-sa}$$

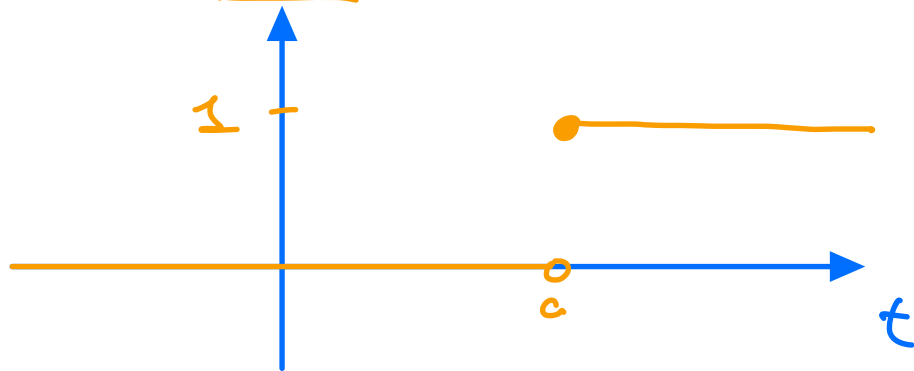
## Unit step function

The unit step function  $u(t)$  is defined by

$$u(t) = \begin{cases} 0, & t < 0 \\ 1, & t \geq 0 \end{cases}$$



Plot  $u(t-a)$  where  $a > 0$ .



Note that

$$\mathcal{L}\{u(t)\}(s) = \int_0^{\infty} e^{-st} u(t) dt = \frac{1}{s}, \quad s > 0$$

Remark:

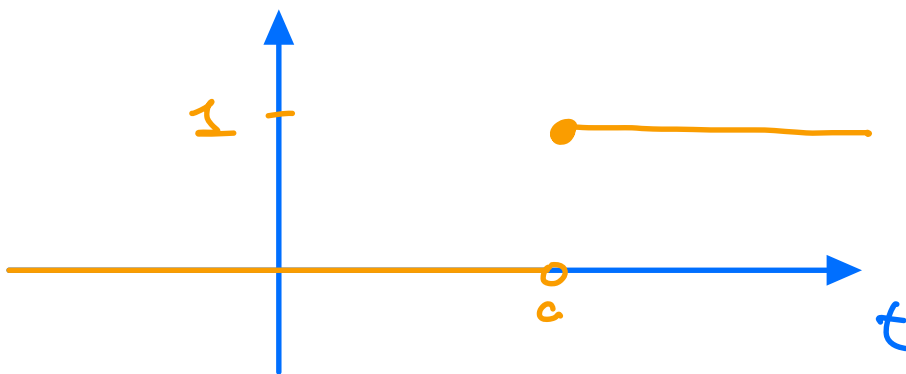


$$\int_{-a}^t \delta(x-a) dx = \begin{cases} 0, & \text{if } t < a \\ 1, & \text{if } t \geq a \end{cases}$$

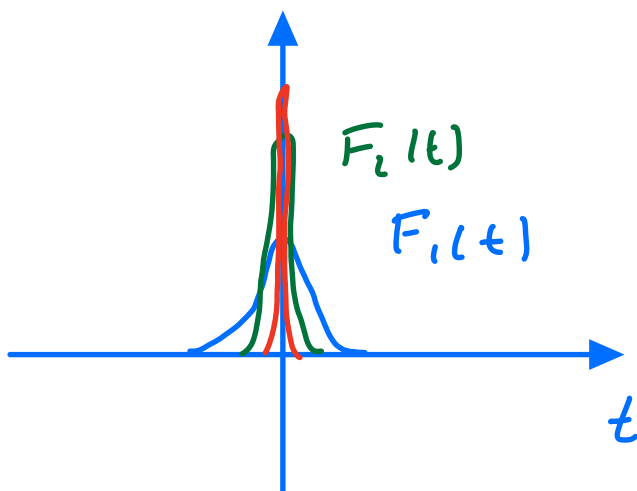
$$= u(t-a)$$

der. w.r.t. time

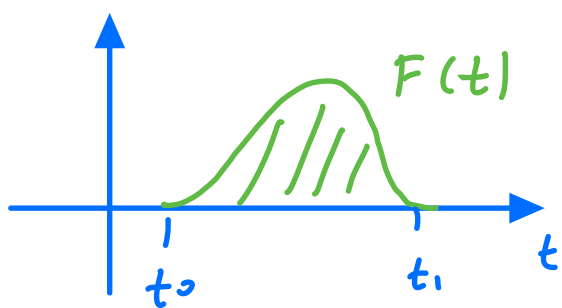
$$\delta(t-a) = u'(t-a)$$



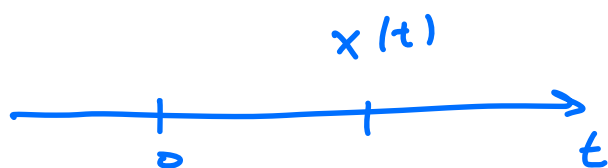
So, the derivative of the step. func.  
is the Dirac Delta distribution !!



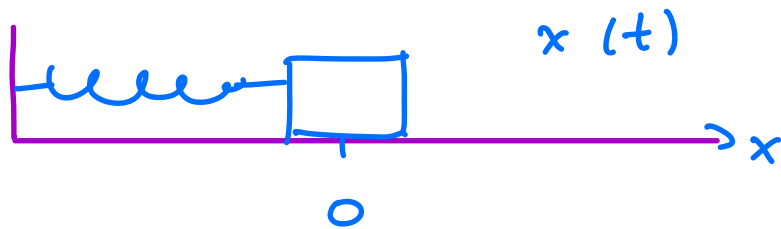
$$\int_{-\infty}^{\infty} S(t) dt = \lim_{n \rightarrow \infty} \int_{-\infty}^{\infty} F_n(t) dt$$



$$\begin{aligned} \text{Impulse} &= \int_{t_0}^{t_1} F(t) dt \\ &= \int_{t_0}^{t_1} m \dot{v}(t) dt \end{aligned}$$



$$= m v(t_1) - m v(t_0)$$



$$\begin{cases} x''(t) + x(t) = \delta(t) & x'' = -x \\ x(0) = 0 \\ x'(0) = 0 \end{cases}$$

Sol. Taking the L.T.

$$(s^2 + 1)\underline{X}(s) = 1$$

$$\underline{X}(s) = \frac{1}{s^2 + 1}$$

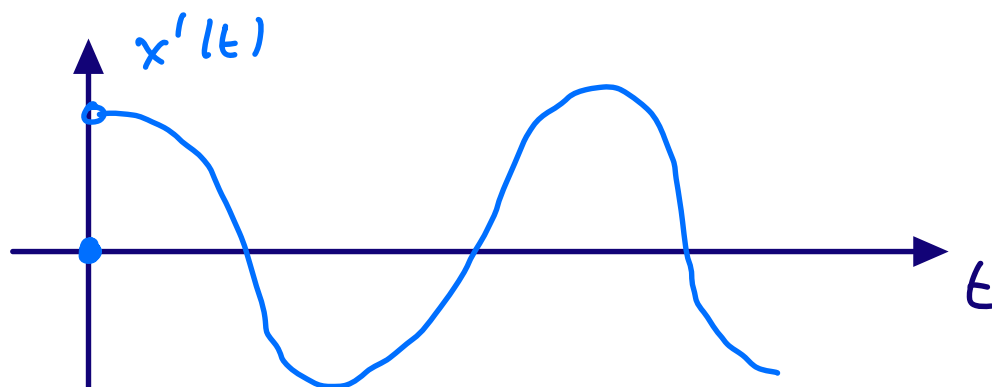
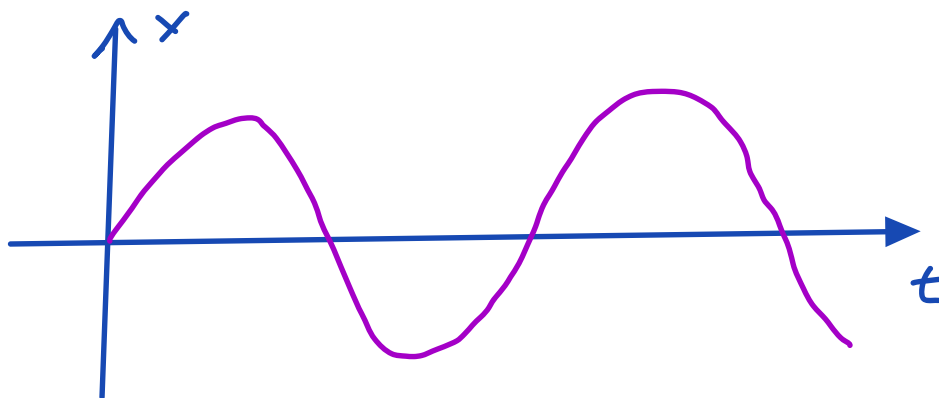
$$\rightarrow x(t) = \mathcal{I}^{-1}\left\{\frac{1}{s^2 + 1}\right\}(t) = \sin(t)$$

Remark:  $x(0) = 0$

$$x'(0) = \cos(t)|_{t=0} = 1 \neq 0$$

momentum goes from  $m\dot{x}(0) = 0$

$$\downarrow$$
$$\lim_{t \rightarrow \infty} m\dot{x}(t) = 1$$

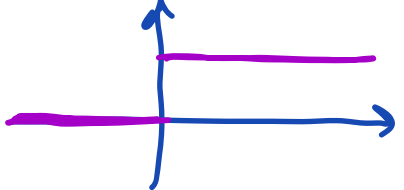




→ Now let's consider the 2nd order  
Lin D.E

$$\begin{cases} A''(t) + aA'(t) + bA(t) = u(t) \\ A(0) = 0 \\ A'(0) = 0 \end{cases}$$

↑  
step function



$a, b \in \mathbb{R}.$

Applying the L.T. we get

$$\begin{aligned} s^2 \mathcal{L}\{A\}(s) + a s \mathcal{L}\{A\}(s) + b \mathcal{L}\{A\}(s) \\ = \mathcal{L}\{u(t)\}(s) = \frac{1}{s} \end{aligned}$$

$$\rightarrow \mathcal{L}\{A\}(s) = \frac{1}{s} \frac{1}{s^2 + as + b}$$

The function  $A$  is called "indicial response" and it will help us to solve the general problem

$$\begin{cases} y''(t) + ay'(t) + by(t) = f(t) \\ y(0) = 0 \\ y'(0) = 0 \end{cases}$$

$$s^2 \mathcal{L}\{y\}(s) + as \mathcal{L}\{y\}(s) + b \mathcal{L}\{y\}(s) = \mathcal{L}\{f\}(s)$$

$$\rightarrow \mathcal{L}\{y\}(s) = \frac{\mathcal{L}\{f\}(s)}{s^2 + as + b}$$

$$\rightarrow \frac{1}{s} \mathcal{L}\{y\}(s) = \frac{1}{s} \frac{1}{s^2 + as + b} \mathcal{L}\{f\}(s)$$

$$= \mathcal{L}\{A\}(s) \mathcal{L}\{f\}(s)$$

$$= \mathcal{L}\{(A * f)(t)\}(s)$$

$$= \mathcal{L}\{(f * A)(t)\}(s)$$

$$= \mathcal{L}\left\{\int_0^t f(\tau) A(t-\tau) d\tau\right\}(s)$$

$$\rightarrow \mathcal{L}\{y\}(s) = s \mathcal{L}\left\{\int_0^t f(\tau) A(t-\tau) d\tau\right\}(s)$$

$$[\mathcal{L}(h'(t))(s) = s \mathcal{L}(h(t))(s) - h(0)]$$

$$\rightarrow \mathcal{L}(y)(s) = \mathcal{L}\left(\frac{d}{dt} \int_0^t f(\tau) A(t-\tau) d\tau\right)(s)$$

$$\rightarrow y(t) = \frac{d}{dt} \int_0^t f(\tau) A(t-\tau) d\tau$$

Leibniz's rule

$$F(t) = \int_{u(t)}^{v(t)} G(t, x) dx$$

$$\frac{d}{dt} F(t) = G(t, v(t)) v'(t) - G(t, u(t)) u'(t) + \int_{u(t)}^{v(t)} \partial_t G(t, x) dx$$

$$y(t) = f(t) A(t-t) + \int_0^t \partial_t (f(\tau) A(t-\tau)) d\tau$$

$$= f(t) A(0) + \int_0^t f(\tau) A'(t-\tau) d\tau$$

or

$$y(t) = \int_0^t A(t-\tau) f'(\tau) d\tau + f(0) A(t)$$

~~type~~

Example: Solve

$$\begin{cases} y'' - y' - 6y = \underline{2e^{3t}} & f(t) \\ y(0) = 0 \\ y'(0) = 0 \end{cases}$$

Sol. First step: get the "indicial response" function  $A(t)$

$$\mathcal{L}\{A\}(s) = \frac{1}{s(s^2 + s - 6)} \quad \begin{array}{l} 1^2 - 4(-6) \\ \text{"} \\ 25 > 0 \end{array}$$

Applying partial fraction decomp.

$$A(t) = -\frac{1}{6} + \frac{1}{15}e^{-3t} + \frac{1}{10}e^{2t}$$

$$y(t) = \int_0^t A(t-\tau) f'(\tau) d\tau + f(0) A(t)$$

$$y(t) = \int_0^t \left[ -\frac{1}{6} + \frac{1}{15}e^{-3(t-\tau)} + \frac{1}{10}e^{2(t-\tau)} \right] \times \underline{6e^{3\tau}} d\tau$$

$$+ 2 \left[ -\frac{1}{6} + \frac{1}{15} e^{-3t} + \frac{1}{10} e^{2t} \right]$$

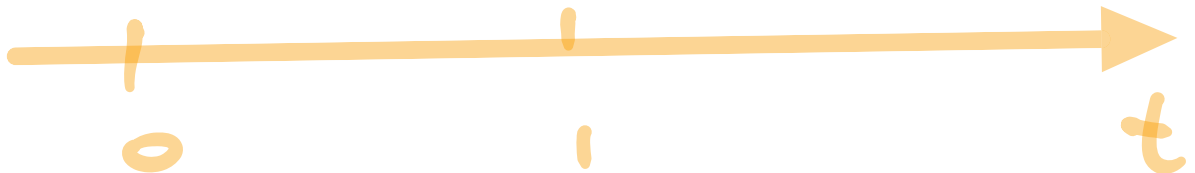
$$= \frac{1}{3} e^{3t} + \frac{1}{15} e^{-3t} - \frac{2}{5} e^{2t}$$


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Finally:

$$f(t) = \begin{cases} 0, & 0 < t \leq 1 \\ e^t, & t > 1 \end{cases}$$

$$\mathcal{L}\{f(t)\}(s) = \int_0^{\infty} e^{-st} f(t) dt$$



## Final Remark :

There are entire books just devoted to the Laplace Transform. Here we have just scratched the surface.

- If you would like to see more examples of partial fraction decomposition consult the book

Fundamentals of DE and boundary value problems by

Nagle, Saff and Snider