# Math 135, Spring 2022

Lecture #22: PDEs and boundary value problems

Wednesday May 18<sup>th</sup>

## Recap

• We finished the topic "Fourier Series".

So far we have covered

- Existence and uniqueness theory for ODEs
- Laplace transform
- Fourier series

And we are now going to deal with

Partial Differential Equations (PDEs)

# **Learning objectives**

Today we will discuss:

Fourier series in an arbitrary domain.

ullet Derivation of the 1d wave equation as a model for a vibrating string.

 Solving the initial boundary value problem for the 1d wave equation by separation of variables.

# Last topic on Fourier series

## **Scaling**

Suppose that 
$$f(x)$$
 is defined on  $[-L,L]$ , where  $L>0$ .  
We want to construct the Fourier series for  $f(x)$ .  
Let  $g(t) = f(Lt/T_L)$  be defined on  $[-T,T]$ .  
The Fourier series for  $g(t)$  is
$$\frac{1}{2}a_0 + \sum_{n=1}^{\infty} \{a_n \cos(nt) + b_n \sin(nt) \}$$

shere

$$a_n = \frac{1}{u} \int_{-\pi}^{\pi} g(t) \cos(nt) dt \quad and \quad b_n = \frac{1}{11} \int_{-\pi}^{\pi} g(t) \sin(nt) dt$$

As  $x = Lt/\pi$ , take our Fourier series for f(x) to be  $t = \frac{\pi}{L}x$ 

$$\frac{1}{2}a_0 + \sum_{n=1}^{\infty} \left\{ a_n \cos(n \frac{\pi}{L}) + b_n \sin(n \frac{\pi}{L}) \right\}$$

where t= = = dx  $a_{n} = \frac{1}{n} \int_{-\pi}^{\pi} \int (Lt_{\pi}^{\prime}) \cos(nt) dt = \frac{1}{L} \int_{-L}^{L} \int (x) \cos(n \frac{\pi x}{L}) dx$  $b_{n} = \frac{1}{\pi} \int_{-\pi}^{\pi} \int (Lt/\pi) \sin(nt) dt = \frac{1}{L} \int_{-\pi}^{L} \int (x) \sin(\frac{n\pi x}{L}) dx$  **Definition:** Let f(x) be an integrable function on the interval [-L, L], where L > 0. We define the **Fourier series** for f(x) on [-L, L] to be

$$\frac{1}{2}a_0 + \sum_{n=1}^{\infty} \left[ a_n \cos\left(\frac{n\pi}{L}x\right) + b_n \sin\left(\frac{n\pi}{L}x\right) \right],$$

where

$$a_n = \frac{1}{L} \int_{-L}^{L} f(x) \cos(\frac{n\pi}{L}x) dx$$
 and  $b_n = \frac{1}{L} \int_{-L}^{L} f(x) \sin(\frac{n\pi}{L}x) dx$ .

# An example

Find the Fourier series for f(x) = x on [-1, 1].

**Step 1:** Use any available symmetries to simplify the computation.



Step 2: Compute

$$b_n = 2 \int_0^1 x \sin(n\pi x) dx$$

$$b_n = 2 \int_0^1 x \sin(n\pi x) dx$$

$$dx = 2 \int_0^1 x \sin(n\pi x) dx$$
even

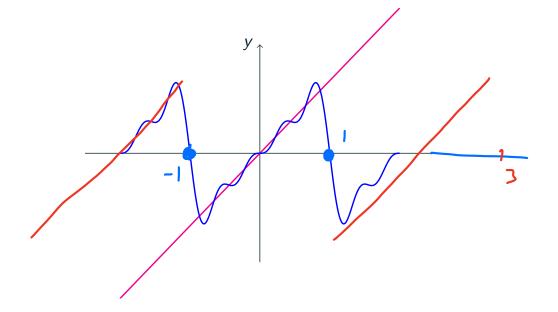
As 
$$u \ge 1$$
,
$$b_n = \left[ -\frac{2}{u_{\overline{u}}} \times coj(u_{\overline{u}} \times 1) \right]_{x=0}^{x=1} + \int_{0}^{1} \frac{2}{u_{\overline{u}}} \cos(u_{\overline{u}} \times 1) dx$$

$$= -\frac{2}{u_{\overline{u}}} (-1)^n$$

$$= -\frac{2}{u_{\overline{u}}} (-1)^{n+1}$$

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So, the Fourier series for 1(x)=x on [-1,Dis



$$f(x) = x$$

$$S_4(x) = \sum_{n=1}^4 (-1)^{n+1} \frac{2}{n\pi} \sin(nx)$$

PDEs and boundary value

problems

at fine t The vibrating string

Let y=y(x,t) be the vertical displacement of a tight string constraint at its ends (y(x,t)

# Assumptions:

- · Motion is strictly vertical
- · No resistance to bending
- 5 Small vibia Hous
- · Constant density s
- · Constant tarion T



## The 1d wave equation

We are lead to consider the 1d wave equation:

$$\frac{\partial^2 y}{\partial t^2} = c^2 \frac{\partial^2 y}{\partial x^2}$$

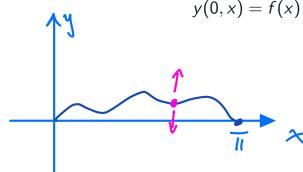
with the **boundary conditions** 

$$y(t,0) = 0 = y(t,\pi)$$

T tension

and the initial conditions

$$y(0,x) = f(x)$$
 and  $\frac{\partial y}{\partial t}(0,x) = g(x)$ .





# Linearity

**Theorem:** If  $y_1(t,x)$  and  $y_2(t,x)$  are solutions of the linear wave equation

$$\frac{\partial^2 y}{\partial t^2} = c^2 \frac{\partial^2 y}{\partial x^2}$$

with the **boundary conditions** 

$$y(t,0)=0=y(t,\pi)$$

and  $C_1$ ,  $C_2$  are constants, then so is

$$y(t,x) = C_1y_1(t,x) + C_2y_2(t,x).$$

**Proof:** Easy computation.

# **Separation of variables**

Solity = 
$$\frac{c}{2} \frac{d^2 x}{d^2 x} \frac{d^2 y}{d^2 y} = \frac{c}{2} \frac{d^2 x}{d^2 x} \frac{d^2 y}{d^2 y} = \frac{c}{2} \frac{d^2 x}{d^2 y} \frac{d^2 y}{d^2 y} = 0$$

Try to find a solution of the form
$$y(x,t) = v(t)u(x) \quad \text{where } u(0) = u(\pi) = 0$$

$$(educated y-ess)$$
Plog this into:
$$v''(t)u(x) = c^2 v(t)u''(x)$$

$$\frac{1}{c^2} \frac{v''(t)}{v(t)} = \frac{u''(x)}{u(x)} = -\lambda \quad \text{a constant}.$$



# **Spatial oscillations**

$$\frac{u''(x)}{u(x)} = -\lambda$$

If  $\lambda < 0$ , what is the general solution of

velve proble

$$u''(x) + \lambda u(x) = 0$$

$$u''(x) + \lambda u(x) = 0$$
 ?  $u(x) = u(x) = 0$ 

A) 
$$C_1 e^{\sqrt{-\lambda}x} + C_2 e^{-\sqrt{-\lambda}x}$$

B) 
$$C_1 \cos(\sqrt{-\lambda}x) + C_2 \sin(\sqrt{-\lambda}x)$$

C) 
$$C_1x + C_2$$

D) None of the above

$$\mu^2 + \lambda =$$

$$\nu' = -\lambda > 0$$

"+ x 4=0 II > < 0: u(x) = c, e + < 2 e Doundary conditions:  $0=u(0)=c_1+c_2 \longrightarrow c_2=-c_1$  $0 = U(\pi) = C_1 e + c_2 e$   $= C_1 (e - e) - \sqrt{-\lambda \pi}$ No non-trival solutions.

-> C1 = 0

If  $\lambda > 0$ , what is the general solution of

$$u''(x) + \lambda u(x) = 0$$
?

A) 
$$C_1 e^{\sqrt{\lambda}x} + C_2 e^{-\sqrt{\lambda}x}$$

B) 
$$C_1 \cos(\sqrt{\lambda}x) + C_2 \sin(\sqrt{\lambda}x)$$

C) 
$$C_1x + C_2$$

D) None of the above

 $\begin{cases} u'' + \lambda u = 0 & \lambda > 0 \\ u(0) = u(\pi) = 0 \end{cases}$ 4(x)= C1 cos (JJXx) +C2 Sm(JJXx)  $Q = u(0) = C_1$ 

0= 4(T)= (2 Sin (N) T) 1 Tf M∈ Z+={1,2,...}, i.e. λ=n2 /or n21

so we have non-trivial solution, u(x)= C2 sin(nx)

## **Eigenfunctions**

• We have shown that we only have non-trivial solutions of the **boundary value problem** 

$$\begin{cases} u'' + \lambda u = 0 & -\frac{d}{dx} u = \lambda u \\ u(0) = 0 = u(\pi) \end{cases}$$

when  $\lambda=n^2$  for  $n=1,2,\ldots$  [exercise: Check what happens when  $\lambda=0$ ]

- We can think of  $-\frac{d^2}{dx^2}$ , together with the boundary conditions  $u(0) = 0 = u(\pi)$  as a **linear operator**, i.e.  $-\frac{d^2}{dx^2}: \mathcal{X} \to \mathcal{Y}$  where  $\mathcal{X}, \mathcal{Y}$  are certain space functions.
- We call the functions

$$u_n(x) = \sin(nx)$$

**eigenfunctions** of this operator and we call the values  $\lambda_n = n^2$  the corresponding eigenvalues. L un = nun

We are trying to solve  $\frac{v''}{c'v} = \frac{u''}{u} = -\lambda$   $\lambda = u^2, \quad \frac{u''}{u} = -u^2 \text{ has a non-trivial col.}$ 

Weed to I dos

$$\frac{U^{11}}{L^2N} = -N^2$$

v"+c2420=0

### **Time oscillation**

What is the general solution of

$$v''(t) + c^2 n^2 v(t) = 0 ?$$

A) 
$$C_1e^{nct} + C_2e^{-nct}$$

B) 
$$C_1 \cos(nct) + C_2 \sin(nct)$$

C) 
$$C_1 nct + C_2$$

D) None of the above

• Suppose we take 
$$\frac{\partial y}{\partial t}(0,x)=g(x)=0$$
. (The case where  $\int (x)=0$   $g(x)\neq 0$  is an the  $H\omega$ )

ullet Consequently, for each  $n\geq 1$  we have a corresponding solution of

$$\begin{cases} v'' + n^2c^2v = 0\\ v'(0) = 0 \end{cases}$$

given by

$$v_n(t) = \cos(nct).$$

$$v_n(t) = \cos(nct) + \cos(nct) + \cos(nct)$$

 $M(x_{it}) = v(t)u(x)$ ,  $\frac{\partial y}{\partial t}(x_{io}) = 0 = v'(t)u(x) = v'(t) = 0$   $v'(t) = -C_{1}uc_{1}s_{1}s_{1}(uct) + C_{2}uc_{1}c_{2}s_{1}(uct)$  $So, v'(o) = C_{2}uc_{1} = v(c_{2})c_{2}$ 

# Solving our PDE

Trying to solve 
$$\frac{\partial^2 y}{\partial t^2} = c^2 \frac{\partial y}{\partial x^2}$$
$$y(o,t) = 0 = y(\pi,t)$$
$$\frac{\partial y}{\partial t}(x,t) = 0$$
$$y(x,0) = 1(x)$$

For each 
$$n \ge 1$$
,  $u_n(x) = \sin(nx)$  and  $u_n(t) = \cos(nct)$   
 $\sin(nx) = \frac{u_n}{u_n} = -u^2$ 

$$\frac{\partial u}{\partial v} = \frac{uv}{uv} = -u^{2}$$

$$y_{n}(x_{1}t) = v_{n}(t) u_{n}(t) i_{1} a_{3} sol$$

$$\int \frac{\partial^{2}y}{\partial t} = c^{2} \frac{\partial^{2}y}{\partial x^{2}}$$

$$y(0|t| = y(|t|) = 0$$

$$\int \frac{\partial^{2}y}{\partial t} = c^{2} \frac{\partial^{2}y}{\partial x^{2}}$$

But then,  $y(x_1 t) = \sum_{k=1}^{\infty} b_k y_k(x_1 t) = \sum_{k=1}^{\infty} b_k (os(nct)sin(nx))$ 

and it solves

Inital condition: y(x,t)= = bn cos(nct) sn(ux) 1(x)= y(x,0) = = = busin(hx) The sine series for flx