

Math 135, Spring 2022

Lecture #24: PDEs and boundary value problems

Monday May 23rd

Last time

- We considered the IBVP

$$\begin{cases} \frac{\partial w}{\partial t} = a^2 \frac{\partial^2 w}{\partial x^2} \\ w(t, 0) = 0 = w(t, \pi) \\ w(0, x) = f(x) \end{cases}$$

- We showed that (formally) the solution is given by

$$w(t, x) = \sum_{n=1}^{\infty} b_n e^{-a^2 n^2 t} \sin(nx) \quad \text{where} \quad f(x) = \sum_{n=1}^{\infty} b_n \sin(nx).$$

Learning objectives

Today we will discuss:

- The Dirichlet problem for the unit disc.

The Laplace Equation

2d models

$$w = w(x, y)$$

- We now consider our models in 2d.
- The **wave equation** takes the form

$$\Delta w := \partial_x^2 w + \partial_y^2 w$$

Laplacian of w

$$\frac{\partial^2 w}{\partial t^2} = \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} = \Delta w$$

- The **heat equation** takes the form

$$\frac{\partial w}{\partial t} = \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} = \Delta w$$

- In both cases, steady states are given by solutions of the **Laplace equation**

$$\lim_{t \rightarrow \infty} w(x, t)$$

$$\frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} = 0. \quad [w \text{ is harmonic}]$$

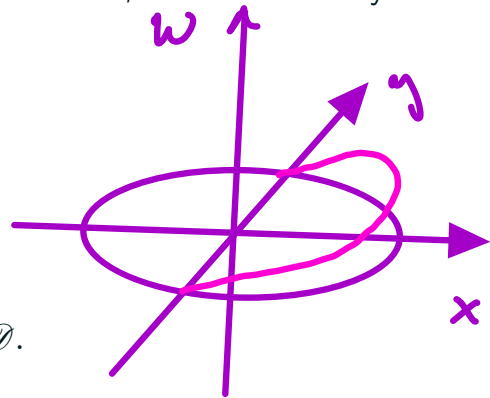
The Dirichlet problem

- Let $\mathcal{D} \subseteq \mathbb{R}^2$ be an open, simply connected region with simple, smooth, closed boundary $\partial\mathcal{D}$.



- The **Dirichlet problem** concerns finding solutions of

$$\begin{cases} \frac{\partial^2 w}{\partial x^2} + \frac{\partial^2 w}{\partial y^2} = 0 & \text{for } (x, y) \in \mathcal{D} \\ w(x, y) = f(x, y) & \text{for } (x, y) \in \partial\mathcal{D}. \end{cases}$$



- We refer to $f(x, y)$ as the **boundary values** of $w(x, y)$.

A special case

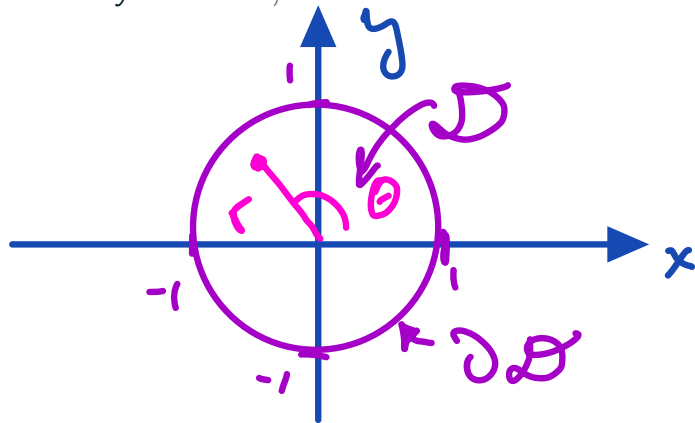
- We will restrict our attention to the case of the unit disc

$$\mathcal{D} = \{x^2 + y^2 < 1\} \quad \text{and} \quad \partial\mathcal{D} = \{x^2 + y^2 = 1\}.$$

- It then makes sense to switch to **polar coordinates**

$$x = r \cos \theta \quad \text{and} \quad y = r \sin \theta,$$

where $r \geq 0$ and $-\pi < \theta \leq \pi$.



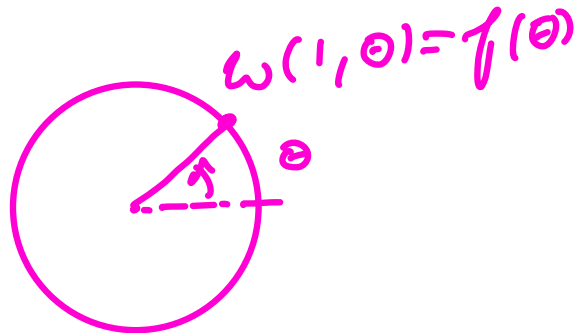
- Suppose that $w = w(r, \theta)$.
- An exercise in the multivariable chain rule (see homework) shows that the Laplace equation can be written as

$$\frac{\partial^2 w}{\partial r^2} + \frac{1}{r} \frac{\partial w}{\partial r} + \frac{1}{r^2} \frac{\partial^2 w}{\partial \theta^2} = 0,$$

for $0 < r < 1$ and $-\pi < \theta \leq \pi$.

- The boundary values can be written as

$$w(1, \theta) = f(\theta).$$



Separation of variables

$$\frac{\partial^2 \omega}{\partial r^2} + \frac{1}{r} \frac{\partial \omega}{\partial r} + \frac{1}{r^2} \frac{\partial^2 \omega}{\partial \theta^2} = 0$$

Look for solutions $\omega(r, \theta) = u(r)v(\theta)$

Then
$$u''(r)v(\theta) + \frac{1}{r} u'(r)v(\theta) + \frac{1}{r^2} u(r)v''(\theta) = 0$$

$$\rightarrow \frac{r^2 u'' + r u'}{u} = -\frac{v''}{v} = \lambda \quad \text{constant}$$

$$\rightarrow \begin{cases} v'' + \lambda v = 0 \\ r^2 u'' + r u' - \lambda u = 0 \end{cases} \quad \leftarrow \text{Euler's equidimensional equation}$$

θ -values

$$v(\theta) = v(\theta + 2\pi)$$

We want to find non-trivial solutions of

$$v'' + \lambda v = 0$$

that are 2π -periodic.

What should we take λ to be?

If $\lambda > 0$:

$$\mu^2 + \lambda = 0$$

$$\mu = \pm \sqrt{-\lambda} = \pm \sqrt{\lambda} i$$

A) n for $n \in \mathbb{N}$

B) n^2 for $n \in \mathbb{N}$

C) $-n^2$ for $n \in \mathbb{N}$

D) $-n$ for $n \in \mathbb{N}$

$$v(\theta) = c_1 \sin(\sqrt{\lambda} \theta) + c_2 \cos(\sqrt{\lambda} \theta)$$

$$\sqrt{\lambda} = n, n \geq 1$$

$$v(\theta) = c_1 e^{\sqrt{-\lambda} \theta} + c_2 e^{-\sqrt{-\lambda} \theta}$$

NOT
periodic

If $\lambda < 0$:

$$\mu^2 + \lambda = 0$$

$$\mu = \pm \sqrt{-\lambda}$$

What is the general solution to

$$v'' + n^2 v = 0$$

when $n \in \mathbb{N}$?

$n \geq 1$

A) $C \sin(n\theta)$

B) $C \cos(n\theta)$

→ C) $C_1 \cos(n\theta) + C_2 \sin(n\theta)$

D) $C_1 e^{n\theta} + C_2 e^{-n\theta}$

$\lambda = 0 \rightarrow$

$v'' = 0 \rightarrow$

$v(\theta) = C_1 + C_2 \theta$

\Rightarrow

$C_2 = 0$

Because v is
2 π periodic

Summary: $v'' + \lambda v = 0$ has a non-trivial, 2π -periodic solution of the form

$$v_n = \begin{cases} \frac{1}{2} a_0 & \text{if } \lambda = 0 \\ a_n \cos(n\theta) + b_n \sin(n\theta) & \text{if } \lambda = n^2 \end{cases}$$

$n \in \mathbb{Z}_+$
" $\{1, 2, 3, \dots\}$

r -values

If we write

$$u(r) = \phi(\rho) \quad \text{for} \quad \rho = \ln r,$$

what is

$$\underline{r^2 u'' + ru' \quad ?}$$

A) $e^{-2\rho} \phi'' - e^{-\rho} \phi' + \phi$

B) ϕ''

C) $e^{2\rho} \phi'' + e^{\rho} \phi'$

D) None of the above

Hint: We can write $r^2 u'' + ru' = r(ru')'$.

$$u(r) = \phi(\rho) \quad \rho = \ln r$$

$$r^2 u'' + r u' = r (r u')'$$

$$r u' = r \frac{du}{dr} = r \frac{d\phi}{d\rho} \frac{d\rho}{dr} = \cancel{r} \phi' \frac{1}{\cancel{r}} = \phi'(\rho)$$

$$\begin{aligned} r (r u')' &= r \frac{d}{dr} (r u') = r \frac{d}{dr} (\phi'(\rho)) = r \phi''(\rho) \frac{d\rho}{dr} \\ &= \cancel{r} \phi''(\rho) \frac{1}{\cancel{r}} \\ &= \phi''(\rho) \end{aligned}$$

So

$$r^2 u'' + r u' = \phi''(\rho)$$

What is the general solution of

$$r^2 u'' + ru' = 0 \quad ?$$

$$r^2 u'' + ru' - \lambda u = 0$$

$$\phi''(\xi) = 0$$

→

$$\begin{aligned} \phi(\xi) &= c_1 + c_2 \xi \\ &= c_1 + c_2 \ln r \end{aligned}$$

→

- A) C_1
- B) $C_1 r + C_2 r^{-1}$
- C) $C_1 \cos(r) + C_2 \sin(r)$
- D) $C_1 + C_2 \ln r$

What is the general solution of

where $n \geq 1$?

$$\underbrace{r^2 u'' + r u'}_{\phi''(\rho)} - n^2 u = 0,$$

$$r^2 u'' + r u' - \underline{\underline{n^2}} u = 0$$

A) $C_1 r^n$

→ B) $C_1 r^n + C_2 r^{-n}$

C) $C_1 \cos(nr) + C_2 \sin(nr)$

D) $C_1 + C_2 \ln(nr)$

$$\begin{aligned} \rightarrow \phi'' - n^2 \phi &= 0 & \rho &= \ln r \\ \rightarrow \phi(\rho) &= C_1 e^{n\rho} + C_2 e^{-n\rho} \\ &= C_1 r^n + C_2 r^{-n} \end{aligned}$$

Solving our PDE

Trying to solve $\partial_r^2 w + \frac{1}{r} \partial_r w + \frac{1}{r^2} \partial_{\theta}^2 w = 0$

If we set $w(r, \theta) = u(r)v(\theta)$ then we have non-trivial when

$$\frac{r^2 u'' + r u'}{u} = -\frac{v''}{v} = n^2$$

Then:
$$v_n(\theta) = \begin{cases} \frac{1}{2} a_0 & \text{if } n=0 \\ a_n \cos(n\theta) + b_n \sin(n\theta) & \text{if } n \geq 1 \end{cases}$$

$$u_n(r) = r^n$$

$C_1 r^n + C_2 r^{-n}$
imposing cond. at $r=0$

$$n=1 \quad (a_1 \cos(\theta) + b_1 \sin(\theta))(r)$$

is a sol. $\partial_r^2 w + \frac{1}{r} \partial_r w + \frac{1}{r^2} \partial_\theta^2 w = 0$

$$n=2 \quad (a_2 \cos(2\theta) + b_2 \sin(2\theta))(r^2)$$

⋮

more generally

$$w(r, \theta) = \frac{1}{2} a_0 + \sum_{n=1}^{\infty} (a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta))$$

$w(1, \theta) = f(\theta)$ so the series is the Fourier series of $f(\theta)$

An example

Use the expression

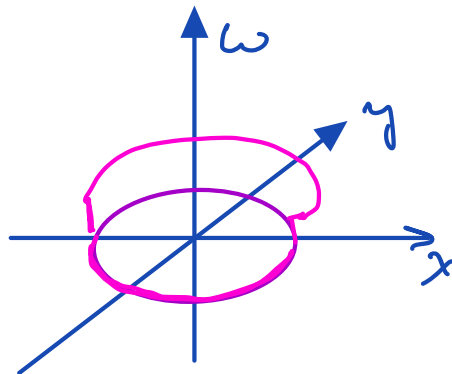
$$w(r, \theta) = \frac{1}{2}a_0 + \sum_{n=1}^{\infty} \left\{ a_n r^n \cos(n\theta) + b_n r^n \sin(n\theta) \right\}$$

to solve the Dirichlet problem for the unit circle with boundary values

$$w(1, \theta) = f(\theta) = \begin{cases} 1 & \text{if } 0 \leq \theta \leq \pi \\ 0 & \text{otherwise.} \end{cases}$$

$$\partial_{r^2}^2 w + \frac{1}{r} \partial_r w + \frac{1}{r^2} \partial_{\theta^2}^2 w = 0$$

$$D = B(0, 1)$$



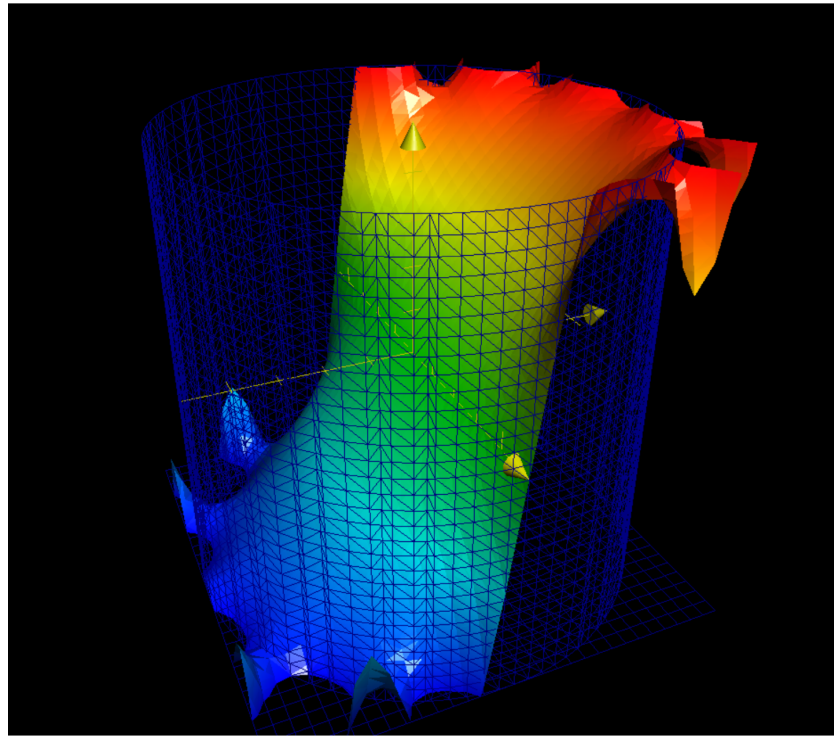
$f(\theta) = \mathbb{1}_{[0, \pi]}(\theta)$ so the Fourier coefficients are

$$\begin{aligned} a_n &= \frac{1}{\pi} \int_{-\pi}^{\pi} \cos(n\theta) \mathbb{1}_{[0, \pi]}(\theta) d\theta = \frac{1}{\pi} \int_0^{\pi} \cos(n\theta) d\theta \\ &= \begin{cases} 1 & \text{if } n=0 \\ 0 & \text{otherwise} \end{cases} \end{aligned}$$

$$b_n = \frac{1}{\pi} \int_0^{\pi} \sin(n\theta) d\theta = \frac{1}{n\pi} [1 - (-1)^n]$$

plug in \rightarrow

$$w(r, \theta) = \frac{1}{2} + \sum_{n=1}^{\infty} \frac{2}{\pi(2n-1)} r^{2n-1} \sin((2n-1)\theta)$$



$$w(r, \theta) \approx \frac{1}{2} + \sum_{n=1}^{11} \frac{2}{2\pi(2n-1)} \sin((2n-1)\theta)$$

See you next time!