

0. Motivation

01

02

03

04

05

06

01

02

03

04

05

06





- Tunneling is such an interesting phenomenon. There are still many questions despite the age of this discovery.
- I picked this topic because I want to visualize the dynamics of a wave packet tunneling through a barrier.
- However, there are still leading questions about the tunneling time and such. The technique utilized in my project might not lead to the most accurate results of tunneling and its actual representation, but they show a general picture of the phenomenon.

01

02

03

04

05

06



01

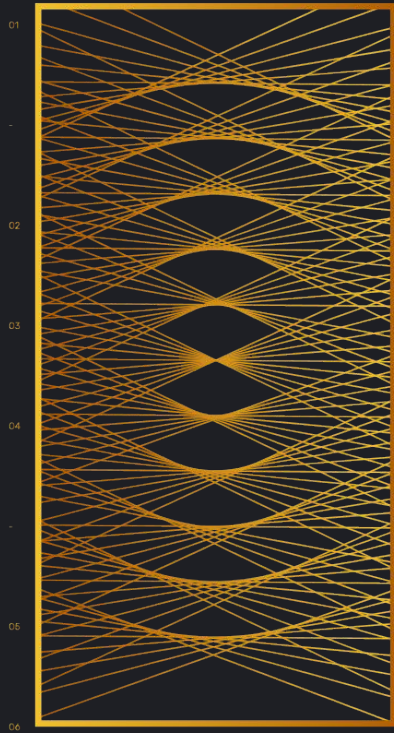
02

03

04

05

06



1. Description of the problem

01

02

03

04

05



01

02

03

04

05

06

06



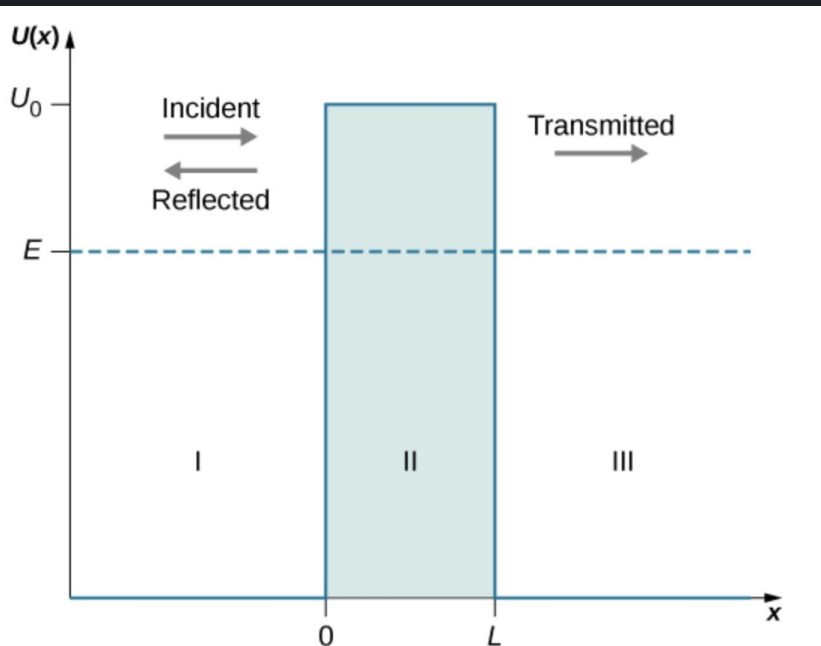
An interesting question occurs when the energy of an incoming wave packet is less than the energy of a barrier.

$$-\frac{\hbar^2}{2m} \frac{d^2\psi(x)}{dx^2} + U(x)\psi(x) = E\psi(x),$$

This situation exhibits the phenomenon of tunneling



General Model



$$-\frac{\hbar^2}{2m} \frac{d^2 \psi_I(x)}{dx^2} = E \psi_I(x),$$

in region *I*: $-\infty < x < 0$,

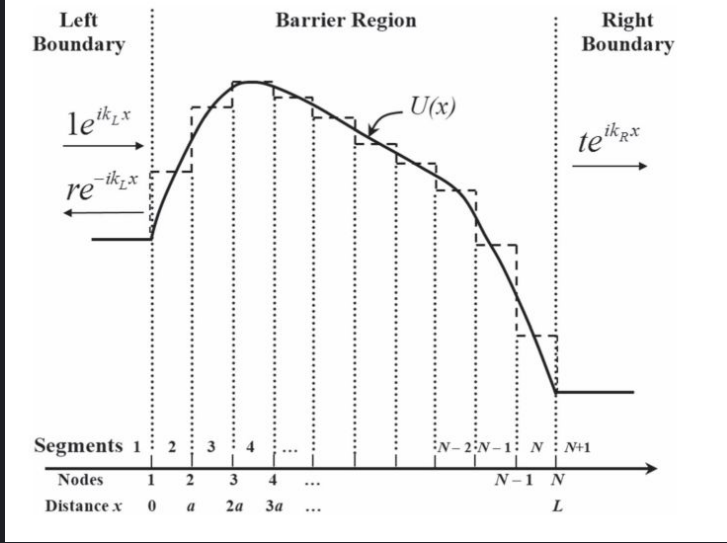
$$-\frac{\hbar^2}{2m} \frac{d^2 \psi_{II}(x)}{dx^2} + U_0 \psi_{II}(x) = E \psi_{II}(x)$$

in region *II*: $0 < x < L$,

$$-\frac{\hbar^2}{2m} \frac{d^2 \psi_{III}(x)}{dx^2} = E \psi_{III}(x)$$

in region *III*: $L < x < +\infty$,

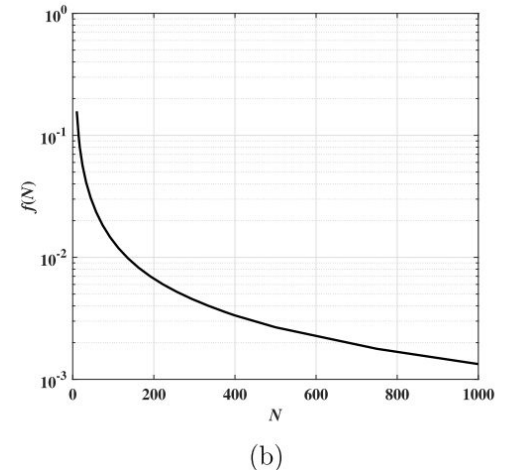
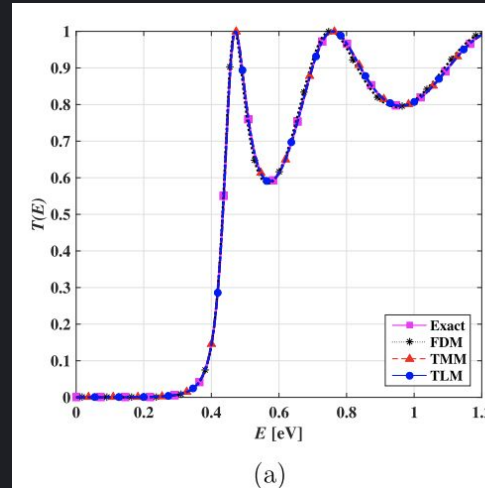
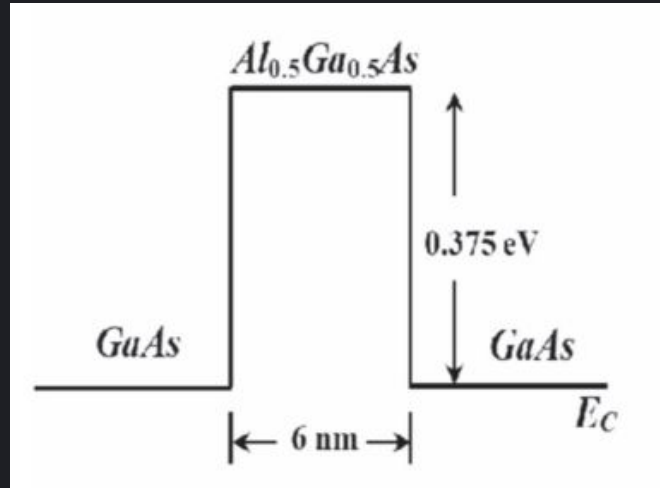
Certain barriers could be solved analytically with the application of PDE when we separate them into multiple equations. Analytically, solving some tunneling problems is impossible and needs the involvement of numerical methods. Therefore, I utilize split operator and fourier methods in this project.



Finite Difference Method (FDM) separates the space into multiple discretizations.

$$\psi_L(x) = \exp(ik_L x) + r \exp(-ik_L x); x < 0,$$

$$\psi_R(x) = t \exp(ik_R x); x > L,$$





Applying central derivative will give us this
discretion of a wave packet.

$$f(x + \Delta x) = f(x) + \Delta x f'(x) + \frac{\Delta x^2}{2!} f''(x) + \dots,$$

$$f(x - \Delta x) = f(x) - \Delta x f'(x) + \frac{\Delta x^2}{2!} f''(x) - \dots,$$

Taylor
Series

01

02

03

$$\left. \frac{d}{dx} \left(\frac{\psi}{m^*} \right) \right|_i = \frac{\psi_{i+1} - \psi_i}{m_i^* a}, \quad \left. \frac{d}{dx} \left(\frac{\psi}{m^*} \right) \right|_{i-1} = \frac{\psi_i - \psi_{i-1}}{m_{i-1}^* a},$$

$$\left. \frac{d^2}{dx^2} \left(\frac{\psi}{m^*} \right) \right|_i = \frac{1}{a} \left(\frac{\psi_{i+1} - \psi_i}{m_i^* a} - \frac{\psi_i - \psi_{i-1}}{m_{i-1}^* a} \right),$$

$$f'(x) = \frac{f(x + \Delta x) - f(x)}{\Delta x},$$

$$f''(x) = \frac{\frac{f(x + \Delta x) - f(x)}{\Delta x} - \frac{f(x) - f(x - \Delta x)}{\Delta x}}{\Delta x} = \frac{f(x + \Delta x) - 2f(x) + f(x - \Delta x)}{(\Delta x)^2}$$

05

06

03

04

05

06



Schrodinger in a discretized form



```
temp1 = (hbar)**2/(2*dx**2*mass(a-dx))
temp2 = (hbar)**2/(2*dx**2*mass(a))
Kl = sqrt(2*mass(a-dx)*(E-potential2(a-dx)))/hbar
psi(2) = (2*Im*temp1*sin(Kl*dx)-(temp2*exp(Im*Kl*dx)+E-potential2(x)-(temp1+temp2))*psi(1))/temp2
do i=2, N-2
```

Find node 1

$$[\eta_1 \exp(ik_L a) + E - U_1 - (\eta_1 + \eta_L)]\psi_1 + \eta_1 \psi_2 = 2i\eta_L \sin(k_L a),$$

```
    x=a+dx*float(i-1)
```

Find node between

$$\eta_{i-1}\psi_{i-1} + [E - U_i - (\eta_i + \eta_{i-1})]\psi_i + \eta_i\psi_{i+1} = 0$$

```
    temp1 = (hbar)**2/(2*dx**2*mass(x))
```

```
    temp2 = (hbar)**2/(2*dx**2*mass(x-dx))
```

```
    psi(i+1) = (temp2*psi(i-1)+(E-potential2(x)-(temp1+temp2))*psi(i))/temp1
```

```
enddo
```

```
!find psiN
```

Find node N

$$\eta_{N-1}\psi_{N-1} + [E - U_N - (\eta_{N-1} + \eta_R) + \eta_R \exp(ik_R a)]\psi_N = 0$$

```
temp1 = (hbar)**2/(2*dx**2*mass(b-dx))
```

```
temp2 = (hbar)**2/(2*dx**2*mass(b))
```

```
KR = sqrt(2*mass(b+dx)*(E-potential2(b+dx)))/hbar
```

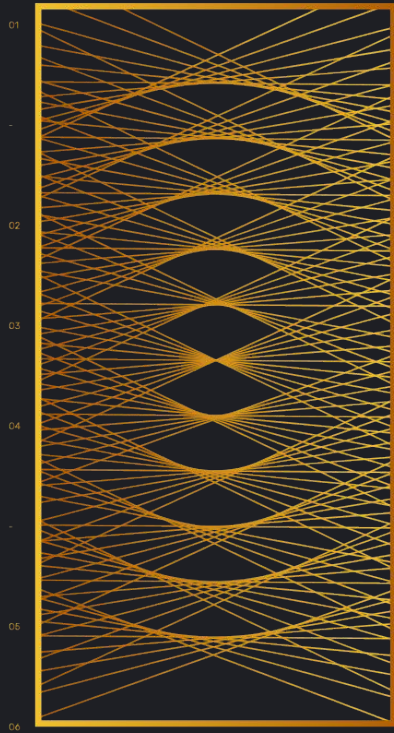
```
psi(N)=-temp1*psi(N-1)/(E-potential2(b)-temp1-temp2+temp2*exp(Im*KR*dx))
```

$$k_{L,R} = \sqrt{2m_{L,R}^*(E - U_{L,R})}/\hbar.$$

$$\tau(E) = |t|^2 k_R / k_L.$$

We'll be able to find all the values of the wave in the barrier





2. Main Numerical Methods Used

01

02

03

04

05

06



01

02

03

04

05

06



Fast Fourier Transform

$$\Delta x \Delta k = \frac{2\pi}{N}$$

01

$$\psi_n = \frac{\Delta k}{\sqrt{2\pi}} \sum_{m=0}^{N-1} \phi_m e^{ik_m x_n} \quad \Leftrightarrow \quad \phi_m = \frac{\Delta x}{\sqrt{2\pi}} \sum_{n=0}^{N-1} \psi_n e^{-ik_m x_n}.$$

I used fast fourier transform package to find the wave packet in momentum and space coordinates

04

```
120 | | call zfftf(N,psi,wsave) !forward Fourier transform
```

05



01

02

03

04

05

06

06



```
122 !~~~~~Momentum Space~~~~~!  
123 !save momentum probability density space  
124 if(mod(nt,n_save)==0)then !--- save psi every n_save steps  
125     nt_tmp=nt_tmp+1  
126     do i=i_min,i_max  
127         psi_save_momentum(i,nt_tmp)=abs(psi(i)*Ns)**2!*sqrt(CMPLX(N,0.0d0))**2  
128         psi_save_momentum_parts(i,nt_tmp)=psi(i)*Ns!*sqrt(CMPLX(N,0.0d0))  
129     enddo  
130 endif  
131 !--- calculate average p and deviation of p ---!  
132 tmp1=0.0  
133 tmp2=0.0  
134 do i=1,N  
135     if(i<=(N/2))then !--- momentum in Fourier space  
136         p=hbar*2.0*pi*(i-1)/(dx*N) !positive  
137     else  
138         p=-hbar*2.0*pi*(N+1-i)/(dx*N) !negative  
139     endif  
140     tmp1=tmp1+p*abs(psi(i)*Ns)**2  
141     tmp2=tmp2+(p**2)*abs(psi(i)*Ns)**2  
142 enddo  
143 av_p(nt)=tmp1*dp !average p  
144 tmp2=tmp2*dp !average p**2  
145 delta_p(nt)=sqrt(abs(tmp2)-av_p(nt)**2)  
146  
147 !psi=psi/Ns !1/sqrtN is needed due to fourth and back fft  
148 call zfftb(N,psi,wsave) !backward Fourier transform
```

Saving the
dynamics of
momentum
space

Calculate the
momentum in Momentum
Space

Record the average momentum and
momentum uncertainty

```

148      call zfftb(N,psi,wsave) !backward Fourier transform
!save the coordinate space
if(mod(nt,n_save)==0)then !<-- save psi every n_save steps
    !nt_tmp=nt_tmp
    do i=i_min,i_max
        psi_save(i,nt_tmp)=abs(psi(i))**2
        psi_save_parts(i,nt_tmp)=psi(i)
    enddo
endif

!--- calculate average x and deviation of x ---!
tmp1=0.0
tmp2=0.0
do i=1,N
    x=a+dx*float(i)
    tmp1=tmp1+x*abs(psi(i))**2
    tmp2=tmp2+(x**2)*abs(psi(i))**2
enddo
av_x(nt)=tmp1*dx !average x
tmp2=tmp2*dx !average x**2
delta_x(nt)=sqrt(abs(tmp2)-av_x(nt)**2) !uncertainty

```

Saving the
dynamics of
coordinate
space

Calculate the
position in
Coordinate Space

Record the average position and
position uncertainty

01

02

03

04

05

06

04

05

06



Split technique for
unitary operator

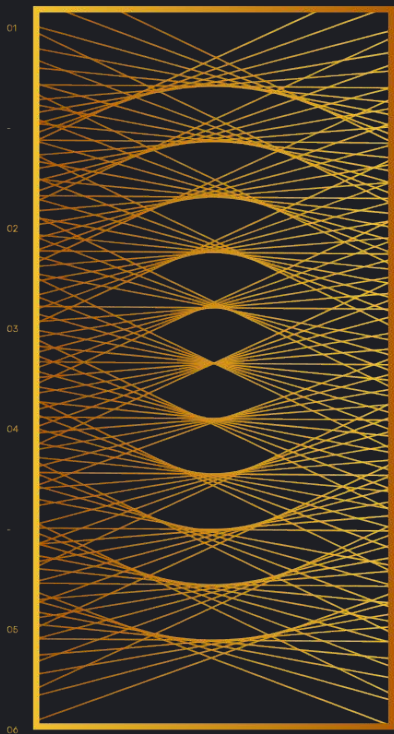
$$\psi(\mathbf{x}, \Delta t) = e^{-i\hat{H}\Delta t/\hbar}\psi(\mathbf{x}, 0)$$

$$\begin{aligned} \psi(\mathbf{x}, t + \Delta t) \\ \approx \mathcal{F}^{-1} \left[e^{-i\mathbf{p}^2\Delta t/4m\hbar} \mathcal{F} \left[e^{-iV(\mathbf{x}, t)\Delta t/\hbar} \mathcal{F}^{-1} \left[e^{-i\mathbf{p}^2\Delta t/4m\hbar} \mathcal{F}[\psi(\mathbf{x}, t)] \right] \right] \right] \end{aligned}$$

```
psi(i)=exp_V*psi(i)  psi=exp_T*psi/float(N)
```

Starting with the
wave packet from
0. I used split
operator to
propagate the wave
packet moving
through a
potential barrier.





3. Tunneling/Results

01

02

03

04

05

06

01

02

03

04

05

06





There are a lot that we can play with tunneling. There are many things we can do such that:

- Interacting with an Arbitrary barrier
- Propagating with an Arbitrary Wave Packet
- Trapping a Wave Packet
- Changing the width/height of the barrier

In this project, I will go through the fourth one mainly to explore the physics of tunneling. Besides, I will add some interesting topics that I could further explore if I have more time.

01

02

03

04

05

06



01

02

03

04

05

06



Simple Potential Barrier

I want to discover what happen to the wave dynamic when I increase the barrier height and the barrier width for a simple potential barrier.

01

02

03

04

05

06



01

02

03

04

05

06

Fixing the width of the barrier at 5, I want to explore how the height affects the Wave Packet.

Width=5

01

02

03

04

05

06

01

02

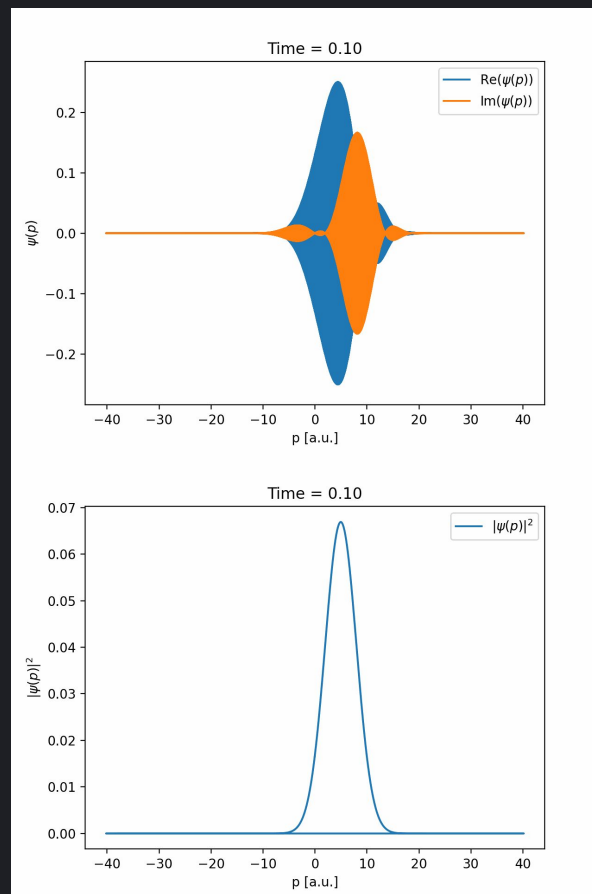
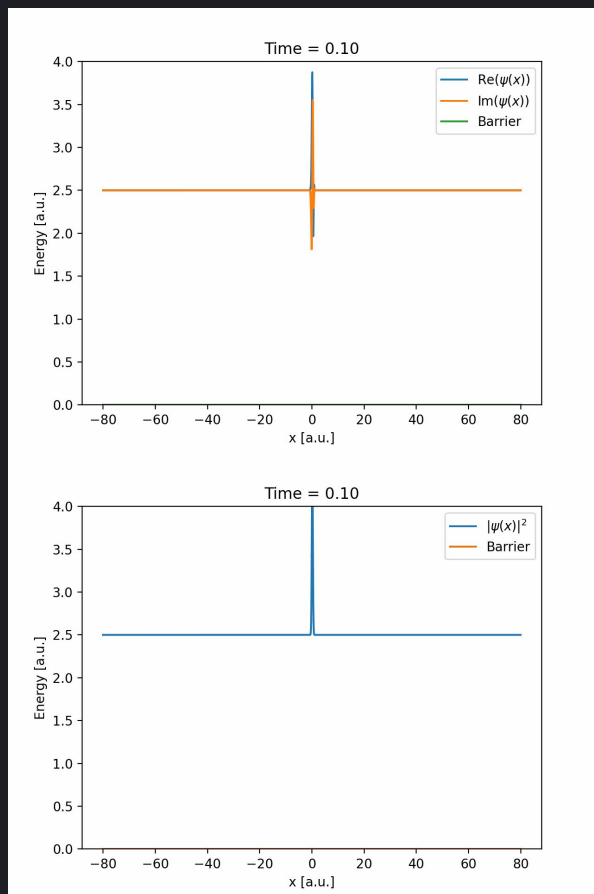
03

04

05

06

⊕ Height=0



01

06

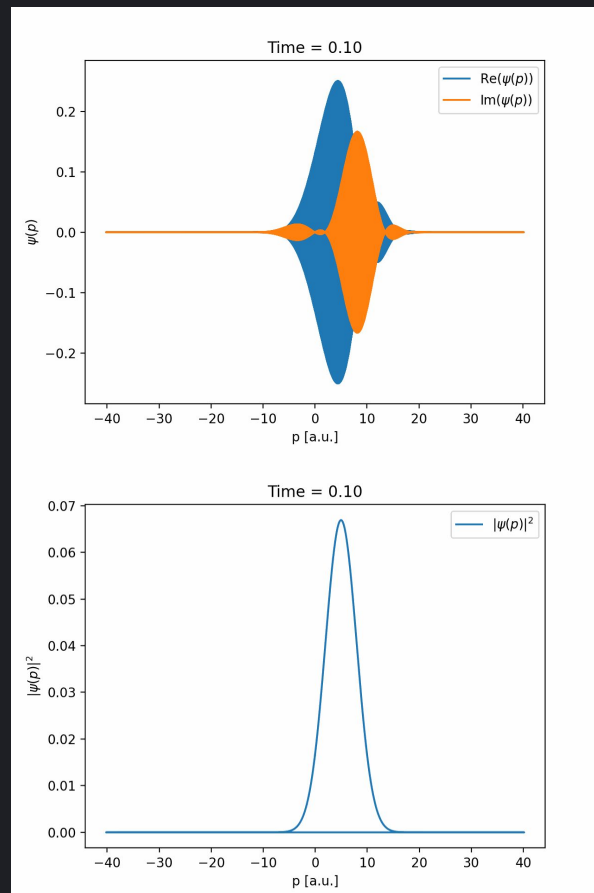
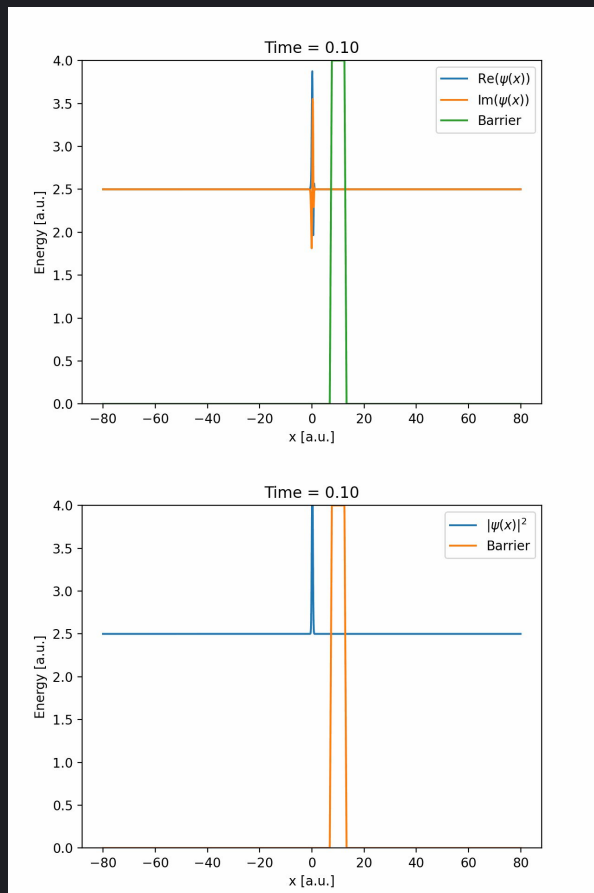
01

02

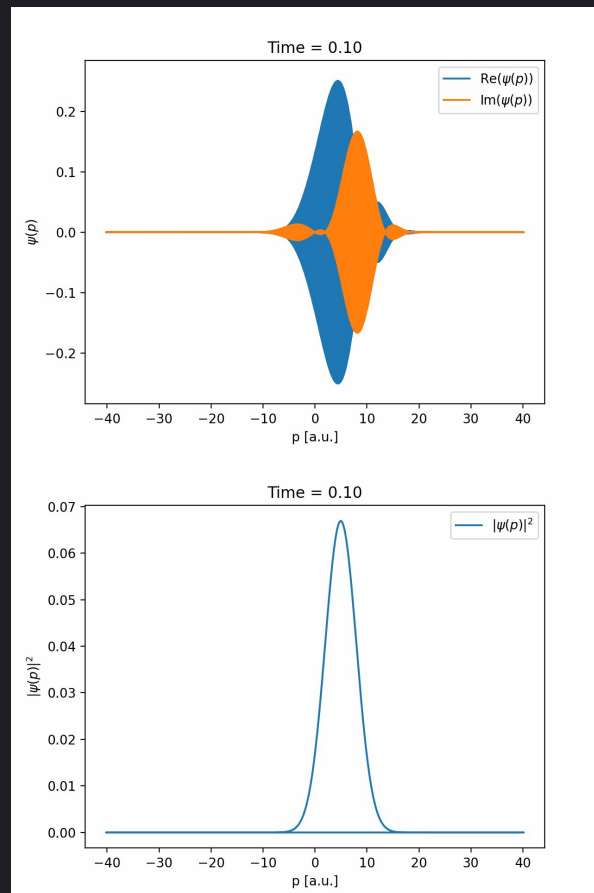
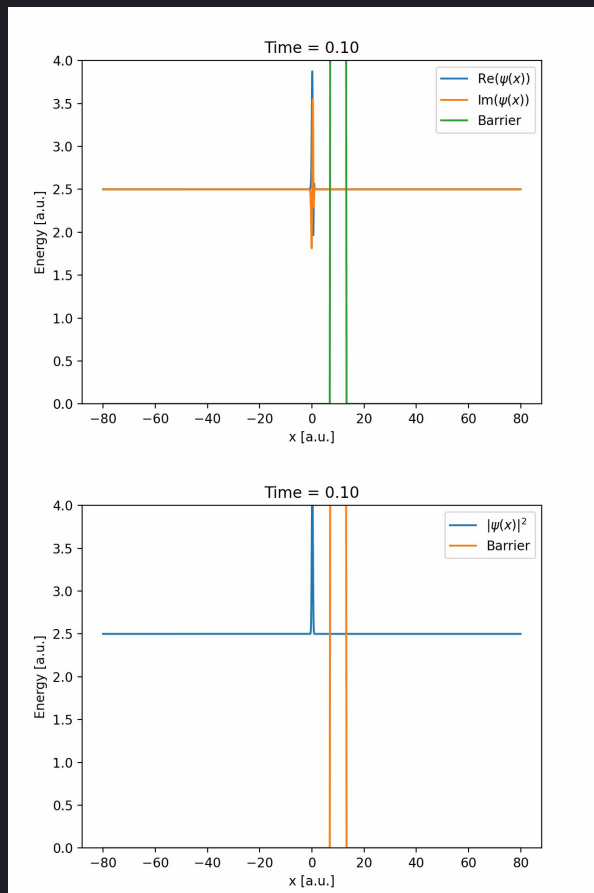
03

04

05



⊕ Height=20



01

01

02

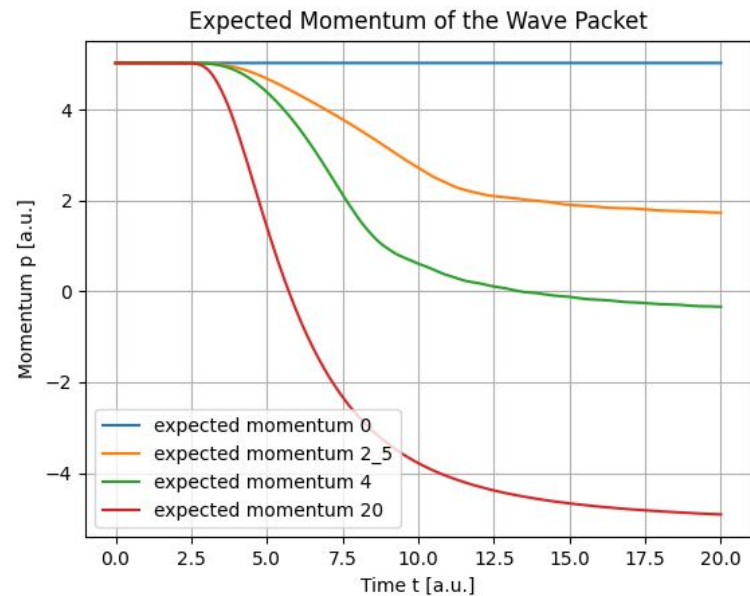
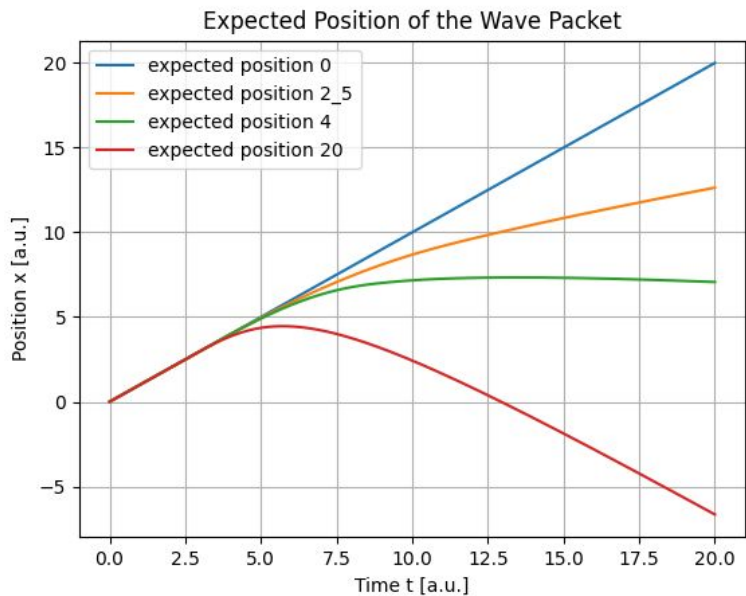
03

04

05

06

06



01

02

03

04

05

06

01

02

03

04

05

06

Now, we fix the height at 4 to investigate how the packet change with the width of the barrier

Height=4

01

02

03

04

05

06

01

02

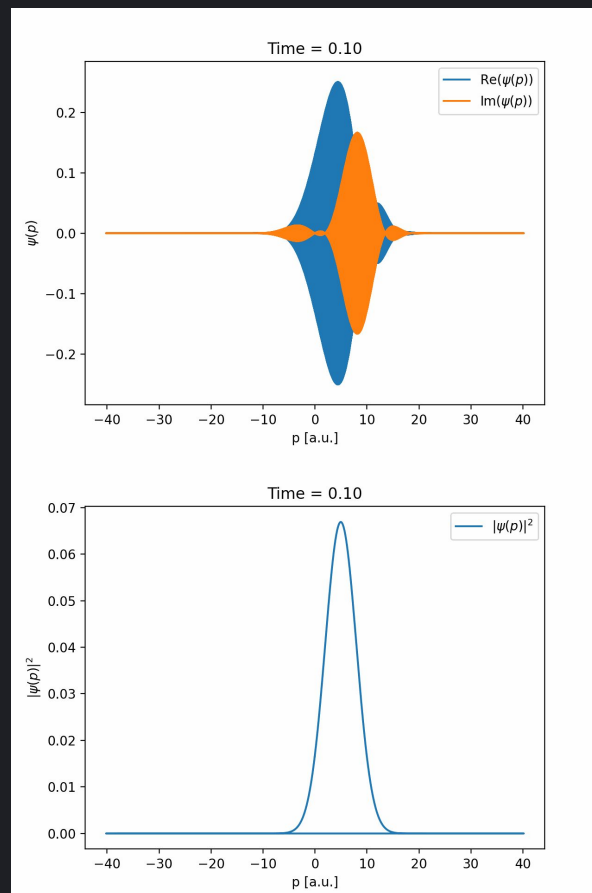
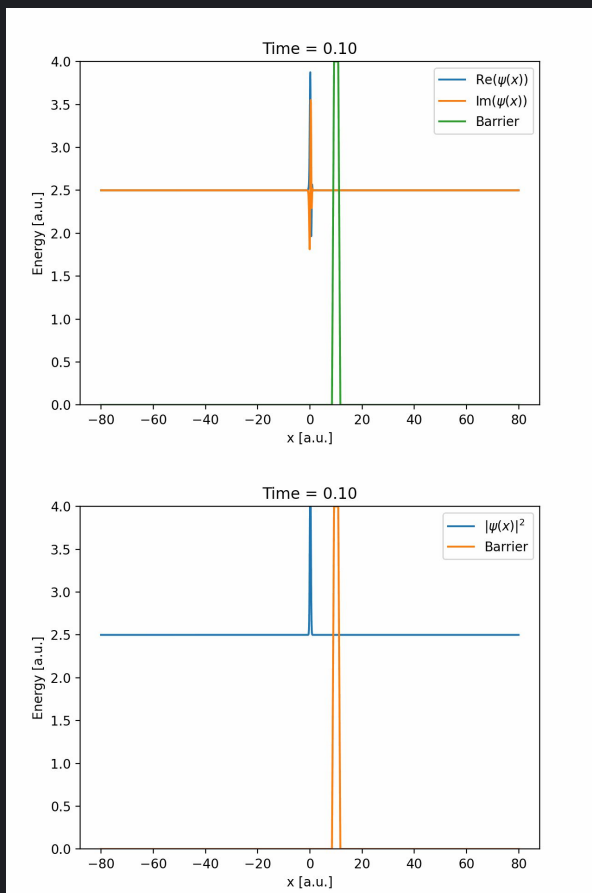
03

04

05

06

⊕ Width=2.5



01

01

02

03

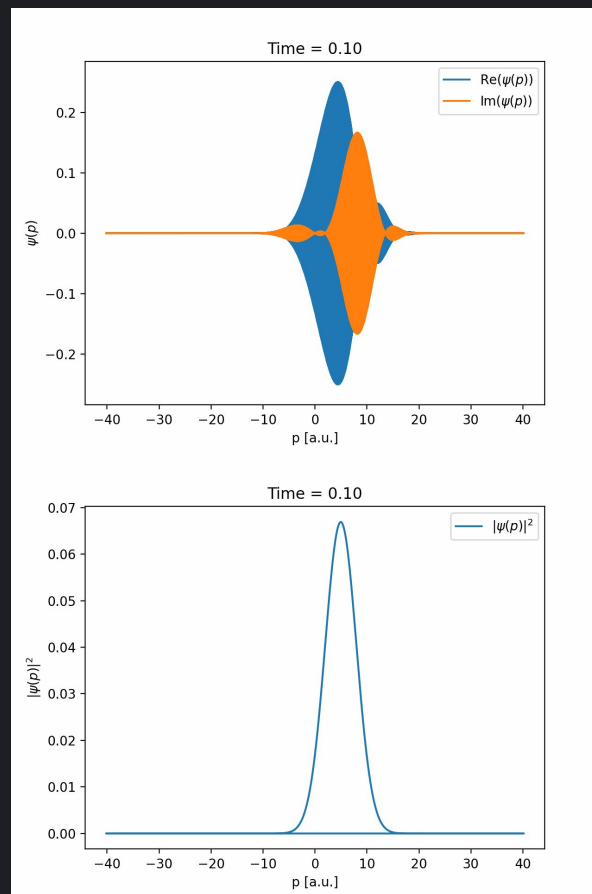
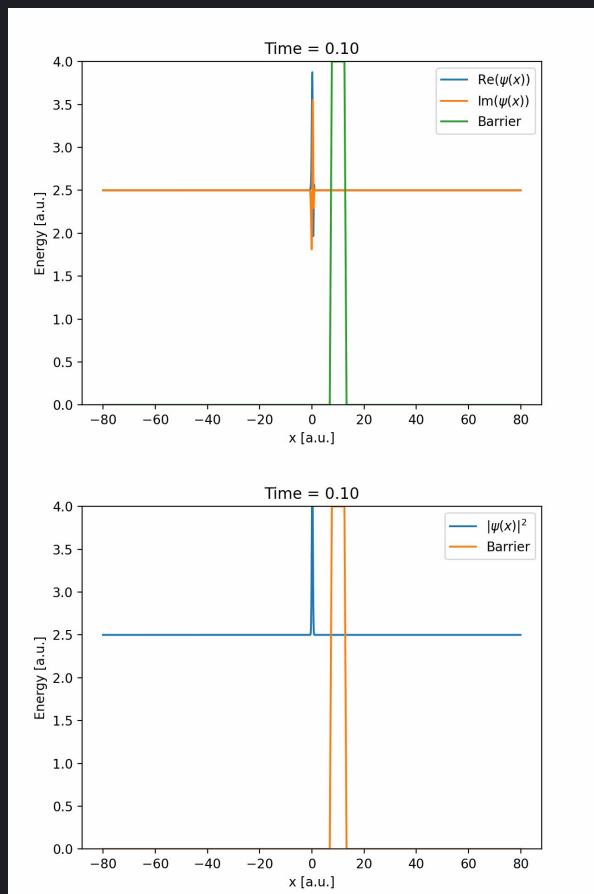
04

05

06

06

⊕ Width=5



01

01

02

03

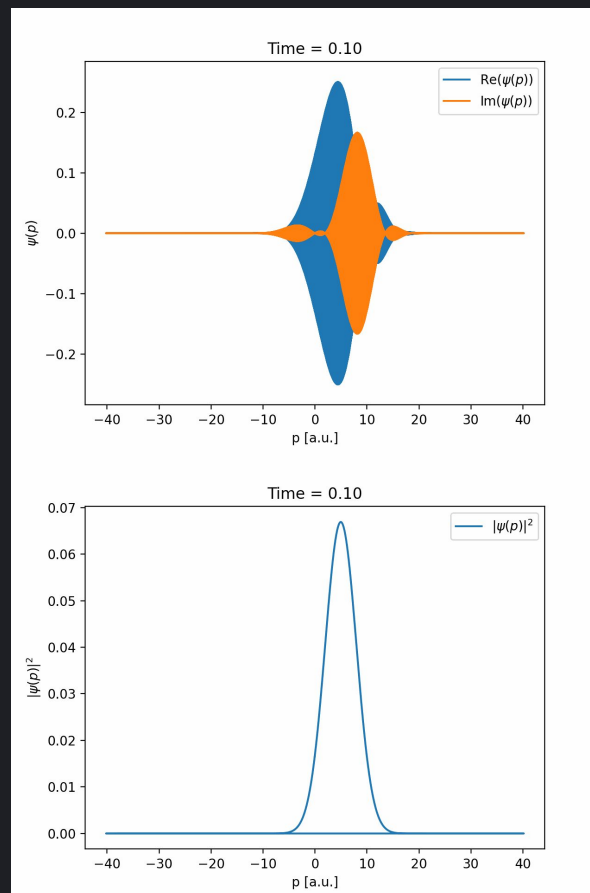
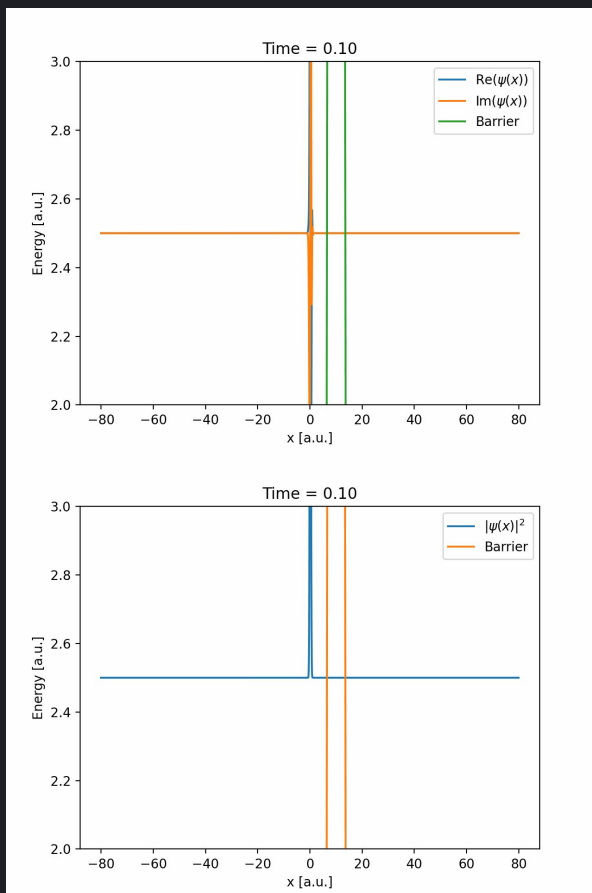
04

05

06

06

⊕ Width=7.5



01

06

01

02

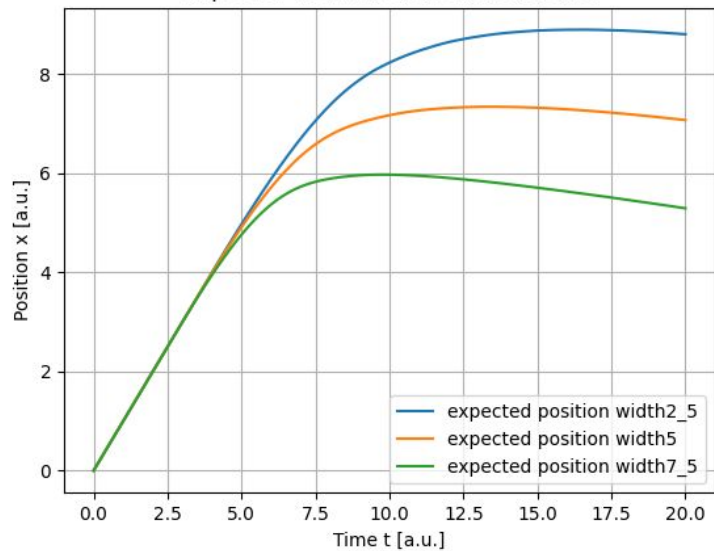
03

04

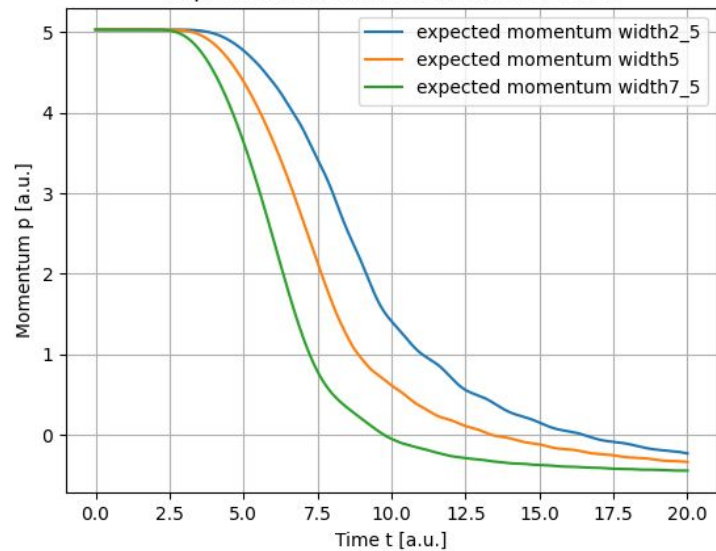
05



Expected Position of the Wave Packet



Expected Momentum of the Wave Packet



Conclusion:

- We have shown the effects that the width and height of a barrier to an incoming wave packet.
- Firstly, increasing the height of a barrier will limit more parts of the wave packet to tunnel through.
- Secondly, increasing the the width of a barrier will result in the wave packet taking a shorter time to 'feel' the barrier significantly.
- Note: These results are derived from the operator technique while these might not represent the actual physical phenomena.

01

02

03

04

05

06

01

02

03

04

05

06

Ongoing/future projects

01

02

03

04

05

06

06

05

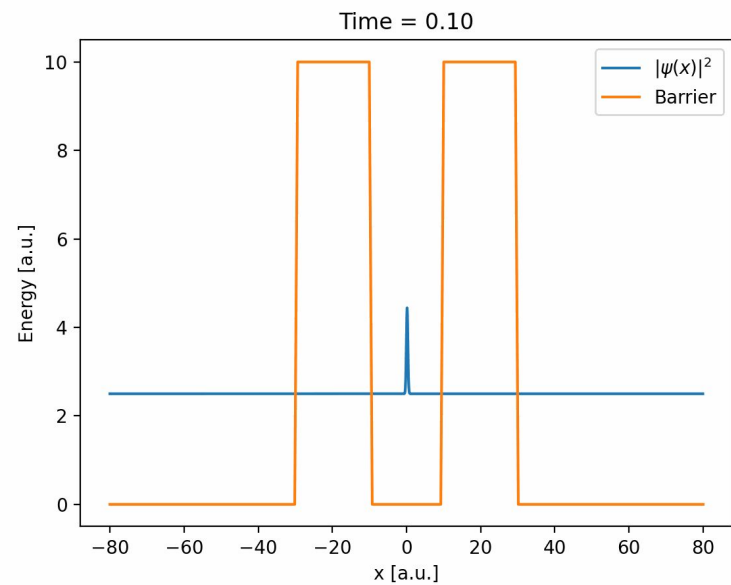
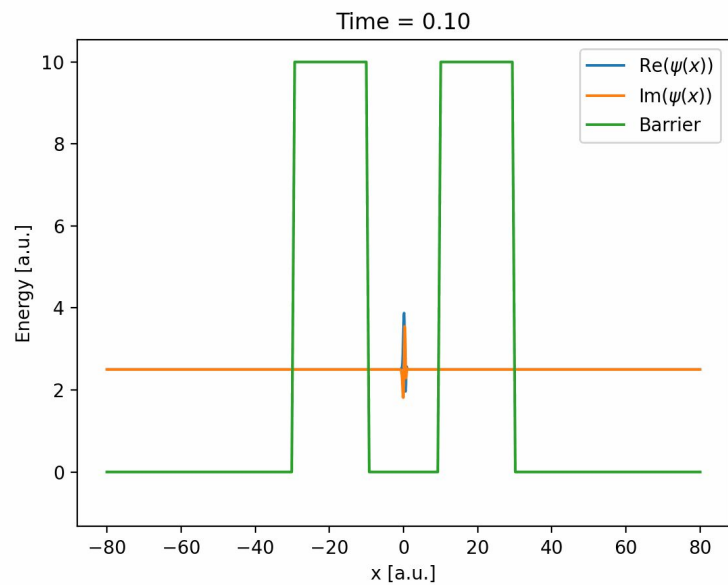
04

03

02

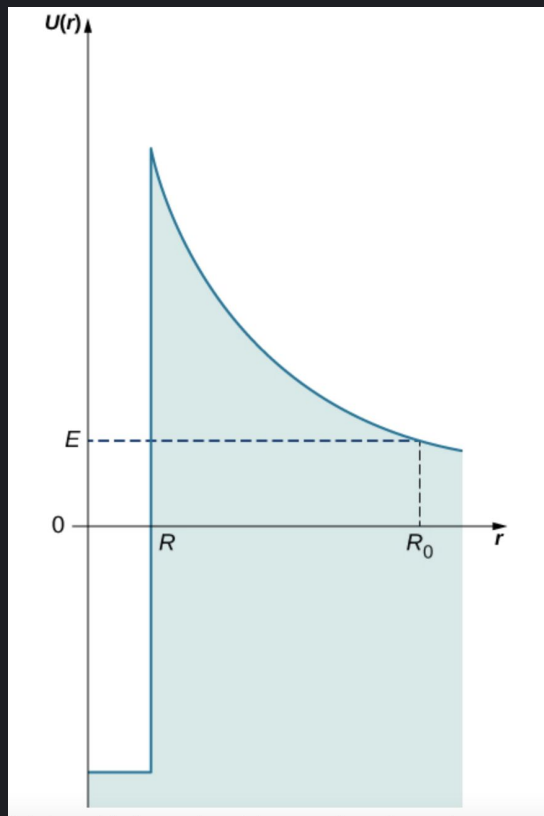
01

Trapped Wave Packet

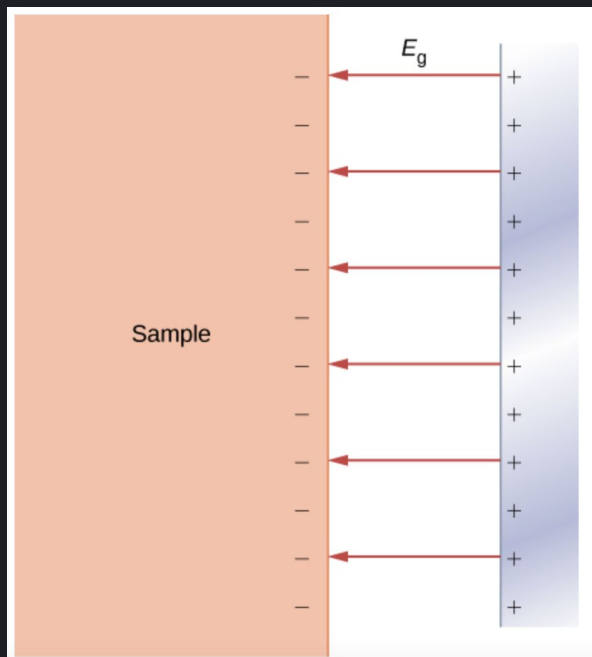




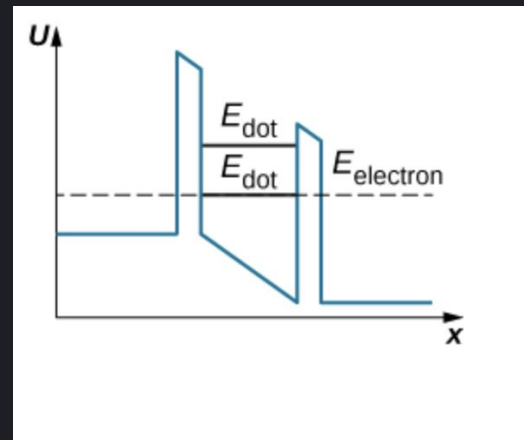
Radioactive Decay



Field Emission

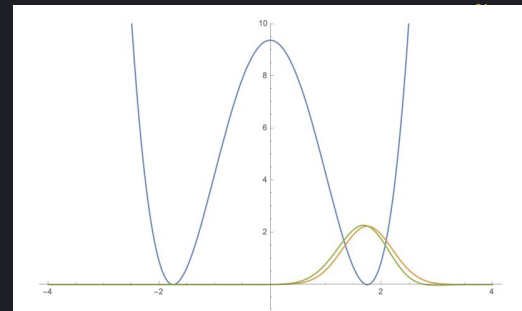


Explore Quantum Dot



03

Double Well





Resources:

- Numerical simulation of tunneling through arbitrary potential barriers applied on MIM and MIIM rectenna diodes
- Quantum Tunneling of Particles through Potential Barriers

01

02

03

04

05

06



01

02

03

04

05

06