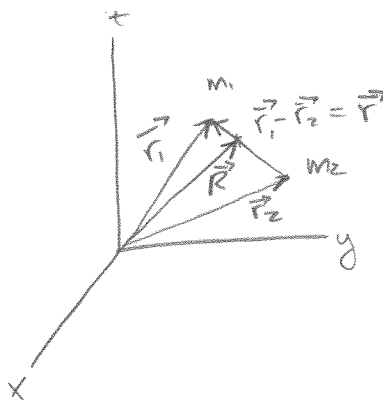


8.2

a)



$$T = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2$$

$$U = U(r) + m_1 g z_1 + m_2 g z_2 \quad \vec{g} = g \hat{z}$$

$$L = T - U$$

$$L = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 - (U(r) + m_1 g z_1 + m_2 g z_2)$$

$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{M}, \quad \vec{r} = \vec{r}_1 - \vec{r}_2$$

Express  $\vec{r}_1$  and  $\vec{r}_2$  in terms of  $\vec{R}$  and  $\vec{r}$ .

$$r_1: \vec{R} = \frac{m_1 \vec{r}_1 + m_2 (\vec{r}_1 - \vec{r})}{M}$$

$$M \vec{R} = m_1 \vec{r}_1 + m_2 \vec{r}_1 - m_2 \vec{r}$$

$$= \vec{r}_1 (m_1 + m_2) - m_2 \vec{r}$$

$$\frac{M \vec{R} + m_2 \vec{r}}{m_1 + m_2} = \vec{r}_1$$

$$\vec{r}_1 = \vec{R} + \frac{m_2}{m_1 + m_2} \vec{r}$$

$$\vec{r}_2: M \vec{R} = m_1 (\vec{r} + \vec{r}_2) + m_2 \vec{r}_2$$

$$= m_1 \vec{r} + m_1 \vec{r}_2 + m_2 \vec{r}_2$$

$$\frac{M \vec{R} - m_1 \vec{r}}{m_1 + m_2} = \vec{r}_2$$

$$\vec{r}_2 = \vec{R} - \frac{m_1}{m_1 + m_2} \vec{r}$$

$$\dot{\vec{r}}_1 = \dot{\vec{R}} + \frac{m_2}{m_1 + m_2} \dot{\vec{r}}$$

$$\dot{\vec{r}}_2 = \dot{\vec{R}} - \frac{m_1}{m_1 + m_2} \dot{\vec{r}}$$

$$L = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 - U(r) - m_1 g z_1 - m_2 g z_2$$

$$\dot{\vec{r}}_1^2 = \left[ \dot{\vec{R}} + \frac{m_2}{m_1+m_2} \dot{\vec{r}} \right]^2$$

$$= \dot{\vec{R}}^2 + 2 \frac{m_2}{m_1+m_2} \dot{\vec{R}} \dot{\vec{r}} + \left( \frac{m_2}{m_1+m_2} \right)^2 \dot{\vec{r}}^2$$

$$\dot{\vec{r}}_2^2 = \dot{\vec{R}}^2 - 2 \frac{m_1}{m_1+m_2} \dot{\vec{R}} \dot{\vec{r}} + \left( \frac{m_1}{m_1+m_2} \right)^2 \dot{\vec{r}}^2$$

$$\frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2 = \frac{1}{2} \left[ (m_1+m_2) \dot{\vec{R}}^2 + \left( \frac{2m_1 m_2}{m_1+m_2} - \frac{2m_1 m_2}{m_1+m_2} \right) \dot{\vec{R}} \dot{\vec{r}} + \left( \frac{m_1 m_2^2}{(m_1+m_2)^2} + \frac{m_2 m_1^2}{(m_1+m_2)^2} \right) \dot{\vec{r}}^2 \right]$$

$$= \frac{1}{2} \left[ M \dot{\vec{R}}^2 + \frac{m_1 m_2}{m_1+m_2} \left( \frac{m_2+m_1}{m_1+m_2} \right) \dot{\vec{r}}^2 \right] \quad \{ M = m_1+m_2 \}$$

$$= \frac{1}{2} \left[ M \dot{\vec{R}}^2 + \frac{m_1 m_2}{M} \dot{\vec{r}}^2 \right]$$

$$= \frac{1}{2} \left[ M \dot{\vec{R}}^2 + \mu \dot{\vec{r}}^2 \right] \quad \{ \mu = \frac{m_1 m_2}{M} \}$$

$$Z = \frac{m_1 z_1 + m_2 z_2}{M}$$

$$M Z = m_1 z_1 + m_2 z_2$$

$$L = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2 - U(r) - g(m_1 z_1 + m_2 z_2)$$

$$L = \frac{1}{2} M \dot{\vec{R}}^2 - M g Z + \frac{1}{2} \mu \dot{\vec{r}}^2 - U(r)$$

$$L_{cm} = \frac{1}{2} M \dot{\vec{R}}^2 - M g Z$$

$$L_{rel} = \frac{1}{2} \mu \dot{\vec{r}}^2 - U(r)$$

$$b) \quad \vec{R} = X \hat{x} + Y \hat{y} + Z \hat{z}$$

$$\dot{\vec{R}}^2 = \dot{X}^2 + \dot{Y}^2 + \dot{Z}^2 \quad \{ \dot{\vec{R}}^2 = \dot{\vec{R}} \cdot \dot{\vec{R}} \}$$

$$L_{cm} = \frac{1}{2} M (\dot{X}^2 + \dot{Y}^2 + \dot{Z}^2) - M g Z$$

$$\frac{\partial L_{cm}}{\partial X} = \frac{\partial L_{cm}}{\partial Y} = 0, \quad \frac{\partial L_{cm}}{\partial Z} = -M g$$

$$\frac{\partial L_{cm}}{\partial \dot{X}} = M \dot{X}, \quad \frac{\partial L_{cm}}{\partial \dot{Y}} = M \dot{Y}, \quad \frac{\partial L_{cm}}{\partial \dot{Z}} = M \dot{Z}$$

$$\frac{\partial L_{cm}}{\partial \dot{q}_i} - \frac{d}{dt} \frac{\partial L_{cm}}{\partial \dot{q}_i} = 0$$

$$X: M\ddot{X} = 0$$

$$Y: M\ddot{Y} = 0$$

$$Z: M\ddot{Z} = -Mg$$

The X and Y components have zero acceleration, and the Z component accelerates at  $-g$ .

$$L_{rel} = \frac{1}{2} \mu \dot{\vec{r}}^2 - U(r)$$

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

$$= (x_1 - x_2) \hat{x} + (y_1 - y_2) \hat{y} + (z_1 - z_2) \hat{z}$$

$$\vec{r} = x_r \hat{x} + y_r \hat{y} + z_r \hat{z} \quad \{r = \text{relative}\}$$

$$\dot{\vec{r}}^2 = \dot{x}_r^2 + \dot{y}_r^2 + \dot{z}_r^2$$

$$L_{rel} = \frac{1}{2} \mu [\dot{x}_r^2 + \dot{y}_r^2 + \dot{z}_r^2] - U(|\vec{r}_1 - \vec{r}_2|) \quad \{r = |\vec{r}_1 - \vec{r}_2|\}$$

$$= \frac{1}{2} \mu [\dot{x}_r^2 + \dot{y}_r^2 + \dot{z}_r^2 - U(x_r) - U(y_r) - U(z_r)]$$

$$\frac{\partial L_{rel}}{\partial x_r} = -\frac{\partial U(x_r)}{\partial x_r}, \quad \frac{\partial L_{rel}}{\partial y_r} = -\frac{\partial U(y_r)}{\partial y_r}, \quad \frac{\partial L_{rel}}{\partial z_r} = -\frac{\partial U(z_r)}{\partial z_r}$$

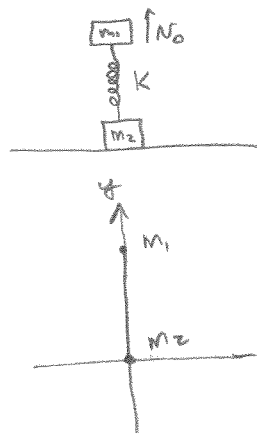
$$\frac{\partial L_{rel}}{\partial \dot{x}_r} = \mu \dot{x}_r, \quad \frac{\partial L_{rel}}{\partial \dot{y}_r} = \mu \dot{y}_r, \quad \frac{\partial L_{rel}}{\partial \dot{z}_r} = \mu \dot{z}_r$$

$$\mu \ddot{x}_r = -\frac{\partial U(x_r)}{\partial x_r}, \quad \mu \ddot{y}_r = -\frac{\partial U(y_r)}{\partial y_r}, \quad \mu \ddot{z}_r = -\frac{\partial U(z_r)}{\partial z_r}$$

This is equivalent to

$$\mu \ddot{\vec{r}} = -\frac{\partial U(r)}{\partial \vec{r}},$$

$$\text{the components of } \mu(\ddot{x}_r + \ddot{y}_r + \ddot{z}_r) = -\left(\frac{\partial}{\partial x_r} U(x_r) + \frac{\partial}{\partial y_r} U(y_r) + \frac{\partial}{\partial z_r} U(z_r)\right).$$



$L = \text{natural length of spring}$

$$\vec{Y} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{m_1 + m_2}$$

$$\vec{Y} = \frac{m_1 \vec{y}_1 + m_2 \vec{y}_2}{m_1 + m_2}$$

$$\vec{Y} = \frac{m_1 \dot{\vec{y}}_1 + m_2 \dot{\vec{y}}_2}{M} \quad \{ M = m_1 + m_2 \}$$

$$\vec{r} = \vec{y}_1 - \vec{y}_2$$

$$\dot{\vec{r}} = \dot{\vec{y}}_1 - \dot{\vec{y}}_2$$

Express  $\vec{y}_1$  and  $\vec{y}_2$  in terms of  $\vec{r}$  and  $\vec{Y}$

$$\begin{aligned} \vec{Y} &= \frac{m_1 \vec{y}_1 + m_2 (\vec{y}_1 - \vec{r})}{M} \\ &= \frac{\vec{y}_1 (m_1 + m_2) - m_2 \vec{r}}{M} \end{aligned}$$

$$\frac{M \vec{Y} + m_2 \vec{r}}{m_1 + m_2} = \vec{y}_1$$

$$\vec{y}_1 = \vec{Y} + \frac{m_2}{m_1 + m_2} \vec{r}$$

$$\boxed{\vec{y}_1 = \vec{Y} + \frac{\mu}{m_1} \vec{r}} \quad \{ \mu = \frac{m_1 m_2}{m_1 + m_2} \}$$

$$\vec{Y} = \frac{m_1 (\vec{y}_2 + \vec{r}) + m_2 \vec{y}_2}{M}$$

$$\frac{M \vec{Y} - m_1 \vec{r}}{m_1 + m_2} = \vec{y}_2$$

$$\vec{y}_2 = \vec{Y} - \frac{m_1}{m_1 + m_2} \vec{r}$$

$$\boxed{\vec{y}_2 = \vec{Y} - \frac{\mu}{m_2} \vec{r}}$$

(2)

$$\dot{\vec{y}}_1 = \dot{\vec{Y}} + \frac{\mu}{m_1} \dot{\vec{r}}$$

$$\dot{\vec{y}}_2 = \dot{\vec{Y}} - \frac{\mu}{m_2} \dot{\vec{r}}$$

$$T = \frac{1}{2} m_1 \dot{\vec{y}}_1^2 + \frac{1}{2} m_2 \dot{\vec{y}}_2^2$$

$$= \frac{1}{2} \left[ m_1 \left( \dot{\vec{Y}} + \frac{\mu}{m_1} \dot{\vec{r}} \right)^2 + m_2 \left( \dot{\vec{Y}} - \frac{\mu}{m_2} \dot{\vec{r}} \right)^2 \right]$$

$$= \frac{1}{2} \left[ m_1 \left( \dot{\vec{Y}}^2 + \frac{2\mu}{m_1} \dot{\vec{Y}} \dot{\vec{r}} + \left( \frac{\mu}{m_1} \right)^2 \dot{\vec{r}}^2 \right) + m_2 \left( \dot{\vec{Y}}^2 - \frac{2\mu}{m_2} \dot{\vec{Y}} \dot{\vec{r}} + \left( \frac{\mu}{m_2} \right)^2 \dot{\vec{r}}^2 \right) \right]$$

$$= \frac{1}{2} \left[ m_1 \dot{\vec{Y}}^2 + 2\mu \dot{\vec{Y}} \dot{\vec{r}} + \frac{\mu^2}{m_1} \dot{\vec{r}}^2 + m_2 \dot{\vec{Y}}^2 - 2\mu \dot{\vec{Y}} \dot{\vec{r}} + \frac{\mu^2}{m_2} \dot{\vec{r}}^2 \right]$$

$$= \frac{1}{2} \left[ \dot{\vec{Y}}^2 (m_1 + m_2) + \dot{\vec{r}}^2 \left( \frac{\mu^2}{m_1} + \frac{\mu^2}{m_2} \right) \right]$$

$$= \frac{1}{2} \left[ M \dot{\vec{Y}}^2 + \mu^2 \left[ \frac{m_1 + m_2}{m_1 m_2} \right] \dot{\vec{r}}^2 \right]$$

$$T = \frac{1}{2} [M \dot{\vec{Y}}^2 + \mu \dot{\vec{r}}^2]$$

$$U = m_1 g y_1 + m_2 g y_2 + \frac{1}{2} k (y_1 - y_2 - L)^2$$

$$= g (m_1 y_1 + m_2 y_2) + \frac{1}{2} k (r - L)^2$$

$$U = M g Y + \frac{1}{2} k (r - L)^2 \quad \left\{ Y = \frac{m_1 y_1 + m_2 y_2}{M} \right\}$$

$$L = T - U$$

$$L = \frac{1}{2} M \dot{\vec{Y}}^2 - M g Y + \frac{1}{2} \mu \dot{\vec{r}}^2 - \frac{1}{2} k (r - L)^2$$

$$\frac{\partial L}{\partial Y} = -Mg$$

$$\frac{\partial L}{\partial \dot{\vec{Y}}} = M \dot{\vec{Y}} \rightarrow \frac{d}{dt} M \dot{\vec{Y}} = M \ddot{\vec{Y}}$$

$$\ddot{\vec{Y}} = -g$$

$$\dot{\vec{Y}} = -gt + v_{Y_0}$$

$$Y = -\frac{1}{2} g t^2 + v_{Y_0} t + Y_0$$

$$\frac{\partial L}{\partial r} = -k(r - L)$$

$$\frac{\partial L}{\partial \dot{\vec{r}}} = \mu \dot{\vec{r}} \rightarrow \frac{d}{dt} \mu \dot{\vec{r}} = \mu \ddot{\vec{r}}$$

$$u\ddot{r} + k(r-L) = 0$$

$$u\ddot{r} + kr = kL$$

$$\ddot{r} + \frac{k}{u}r = \frac{k}{u}L$$

$$r_h: r = A_1 e^{\gamma t}$$

$$\gamma^2 + \frac{k}{u} = 0$$

$$\gamma = \pm i\omega, \quad \omega = \sqrt{\frac{k}{u}}$$

$$r_h = A_1 e^{i\omega t} + A_2 e^{-i\omega t}$$

$$r_h = A \cos(\omega t - \delta)$$

$$r_p: r = C_1$$

$$\frac{k}{u}C_1 = \frac{k}{u}L$$

$$C_1 = L$$

$$r_p = L$$

$$r = r_h + r_p$$

$$r = A \cos(\omega t - \delta) + L$$

$$y_1 - y_2 = A [\cos(\omega t) \cos(\delta) + \sin(\omega t) \sin(\delta)] + L$$

$$y_1(0) - y_2(0) = L = A [\cos(\delta)] + L$$

$$A \cos(\delta) = 0$$

$$y_1 - y_2 = A \sin(\omega t) \sin(\delta) + L$$

$$\dot{y}_1 - \dot{y}_2 = A\omega \cos(\omega t) \sin(\delta)$$

$$\dot{y}_1(0) - \dot{y}_2(0) = v_0 = A\omega \sin(\delta)$$

$$\frac{v_0}{\omega} = A \sin(\delta)$$

$$y_1 - y_2 = \frac{v_0}{\omega} \sin(\omega t) + L$$

$$Y = Y_0 + v_{Y_0} t - \frac{1}{2} g t^2$$

$$\frac{m_1 y_1 + m_2 y_2}{M} = Y_0 + v_{Y_0} t - \frac{1}{2} g t^2$$

$$\frac{m_1 y_1(0) + m_2 y_2(0)}{M} = Y_0$$

(4)

$$\frac{m_1 L}{M} = Y_0 \quad \{ y_1(0) = L, y_2(0) = 0 \}$$

$$\frac{m_1 y_1 + m_2 y_2}{M} = \frac{m_1 L}{M} + v_0 t - \frac{1}{2} g t^2$$

$$\frac{m_1 \dot{y}_1 + m_2 \dot{y}_2}{M} = v_0 - g t$$

$$\frac{m_1 \dot{y}_1(0) + m_2 \dot{y}_2(0)}{M} = v_0$$

$$\frac{m_1 v_0}{M} = v_0 \quad \{ \dot{y}_1(0) = v_0, \dot{y}_2(0) = 0 \}$$

$$\frac{m_1 y_1 + m_2 y_2}{M} = \frac{m_1 L}{M} + \frac{m_1 v_0}{M} t - \frac{1}{2} g t^2$$

Now, solve for  $y_1$  and  $y_2$ .

$$m_1 y_1 + m_2 (y_1 - \frac{v_0}{\omega} \sin(\omega t) - L) = m_1 L + m_1 v_0 t - \frac{M}{2} g t^2$$

$$m_1 y_1 + m_2 y_1 - \frac{v_0 m_2}{\omega} \sin(\omega t) - m_2 L = m_1 L + m_1 v_0 t - \frac{M}{2} g t^2$$

$$y_1 M = m_1 L + m_2 L + \frac{v_0 m_2}{\omega} \sin(\omega t) + m_1 v_0 t - \frac{M}{2} g t^2$$

$$y_1 = L + \frac{m_1 v_0}{M} t + \frac{v_0 m_2}{M \omega} \sin(\omega t) - \frac{1}{2} g t^2$$

$$m_1 (y_2 + \frac{v_0}{\omega} \sin(\omega t) + L) + m_2 y_2 = m_1 L + m_1 v_0 t - \frac{M}{2} g t^2$$

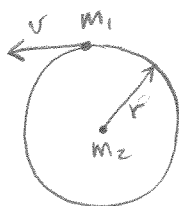
$$m_1 y_2 + m_2 y_2 + \frac{m_1 v_0}{\omega} \sin(\omega t) + m_1 L = m_1 L + m_1 v_0 t - \frac{M}{2} g t^2$$

$$y_2 M = m_1 v_0 t - \frac{m_1 v_0}{\omega} \sin(\omega t) - \frac{M}{2} g t^2$$

$$y_2 = \frac{m_1 v_0}{M} t - \frac{m_1 v_0}{M \omega} \sin(\omega t) - \frac{1}{2} g t^2$$

8.1

a)



$$F = -\frac{G m_1 m_2}{r^2}$$

$$m_1 a = -\frac{G m_1 m_2}{r^2}$$

$$a = -\frac{G m_2}{r^2}$$

$$-\frac{v^2}{r} = -\frac{G m_2}{r^2}$$

$$v^2 = \frac{G m_2}{r}$$

$$v = \left( \frac{G m_2}{r} \right)^{1/2}$$

$$C = 2\pi r$$

$$T = t_p = \frac{C}{v}$$

$$= \frac{2\pi r}{\left( \frac{G m_2}{r} \right)^{1/2}}$$

$$T = 2\pi \sqrt{\frac{r^3}{G m_2}}$$

b)



$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{M}$$

$$\vec{r} = \vec{r}_1 - \vec{r}_2$$

$$T = \frac{1}{2} (M \dot{\vec{R}}^2 + \mu \dot{\vec{r}}^2), \quad \mu = \frac{m_1 m_2}{m_1 + m_2}, \quad M = m_1 + m_2$$

$$U = U(r)$$

$$U(r) = - \int_{r_0}^r \frac{G m_1 m_2}{r'^2} dr'$$

$$= - \frac{G m_1 m_2}{r'} \Big|_{r_0=\infty}^r$$

$$U(r) = - \frac{G m_1 m_2}{r}$$

$$L = T - U$$



Since the C.M. is not accelerating  $L_{cm} = 0$

②

$$L = L_{cm} + L_{rel}$$

$$L = L_{rel}$$

$$L = \frac{1}{2} \mu \dot{r}^2 - U(r)$$

$$\frac{\partial L}{\partial r} = -\frac{\partial U(r)}{\partial r} = -\frac{\partial}{\partial r} \left( \frac{G m_1 m_2}{r} \right)$$

$$\frac{\partial L}{\partial r} = \frac{G m_1 m_2}{r^2}$$

$$\frac{\partial L}{\partial \dot{r}} = \mu \dot{r} \rightarrow \frac{d}{dt} \mu \dot{r} = \mu \ddot{r}$$

$$\mu \ddot{r} = \frac{G m_1 m_2}{r^2} \quad \left\{ \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \frac{\partial L}{\partial r} \right\}$$

$$\frac{m_1 m_2}{M} \ddot{r} = \frac{G m_1 m_2}{r^2}$$

$$\ddot{r} = \frac{GM}{r^2}$$

$$\frac{v^2}{r} = \frac{GM}{r^2}$$

$$v^2 = \frac{GM}{r}$$

$$C = 2\pi r$$

$$T_{rel} = t_p = \frac{C}{v}$$

$$= \frac{2\pi r}{\left(\frac{GM}{r}\right)^{1/2}}$$

$$\boxed{T_{rel} = 2\pi \sqrt{\frac{r^3}{GM}}}$$

As  $m_2 \rightarrow \infty$ ,  $M = m_1 + m_2 \rightarrow m_2 \{m_1 \ll m_2\}$ , so  $T_{rel} \rightarrow T_1$ .

c)  $m_1 = m_2 = M_\odot$

$$r = 1.5 \times 10^8 \text{ km} \quad \{ \text{earth orbital radius} \}$$

$$r = 1.5 \times 10^{11} \text{ m}$$

$$M = 2 M_\odot$$

$$M_\odot = 1.99 \times 10^{30} \text{ kg}$$

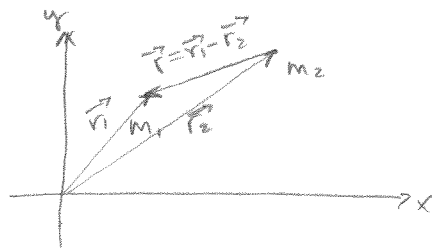
$$G = 6.67 \times 10^{-11} \frac{\text{m}^3}{\text{kg s}^2}$$

$$T_{rel} = 2\pi \sqrt{\frac{(1.5 \times 10^{11} \text{ m})^3}{6.67 \times 10^{-11} \frac{\text{m}^3}{\text{kg s}^2} \cdot [1.99 \times 10^{30} \text{ kg}] \cdot 2}}$$

$$= 2.24 \times 10^7 \text{ s}$$

$$\boxed{T_{rel} = 0.71 \text{ years}}$$

0.8



$$T = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2$$

$$U = \frac{1}{2} K r^2$$

$$L = T - U$$

$$L = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2 - \frac{1}{2} K r^2$$

$$\frac{\partial L}{\partial \vec{R}} = 0$$

$$\frac{\partial L}{\partial \vec{r}} = M \dot{\vec{R}}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \vec{R}} = M \ddot{\vec{R}}$$

$$M \ddot{\vec{R}} = 0$$

$$\dot{\vec{R}} = v_0 \hat{R}$$

$$\boxed{\vec{R} = v_0 t \hat{R} + \vec{R}_0}$$

$$\frac{\partial L}{\partial r} = -kr$$

$$\frac{\partial L}{\partial \dot{r}} = \mu \dot{r}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \mu \ddot{r}$$

$$\mu \ddot{r} + kr = 0 \quad \left\{ \frac{\partial L}{\partial r} - \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = 0 \right\}$$

$$\ddot{r} + \frac{k}{\mu} r = 0$$

$$r = A_1 e^{\gamma t}$$

$$\gamma^2 = -\frac{k}{\mu}$$

$$\gamma = \pm i\omega, \quad \omega = \sqrt{\frac{k}{\mu}}$$

$$r = A_1 e^{i\omega t} + A_2 e^{-i\omega t}$$

$$r = A \cos(\omega t - \delta)$$

$$\vec{R} = \frac{m_1 \vec{r}_1 + m_2 \vec{r}_2}{M}$$

$$\vec{r} = \vec{r}_1 - \vec{r}_2, \quad \vec{r}_1 = x_1 \hat{x} + y_1 \hat{y}, \quad \vec{r}_2 = x_2 \hat{x} + y_2 \hat{y}$$

$$\mu = \frac{m_1 m_2}{m_1 + m_2}$$

Initial conditions:

$$r(0) = r_0$$

$$\dot{r}(0) = v_{0r}$$

$$r = A \cos(\omega t) \cos(\delta) + A \sin(\omega t) \sin(\delta)$$

$$r_0 = A \cos(\delta)$$

$$r = r_0 \cos(\omega t) + A \sin(\omega t) \sin(\delta)$$

$$\dot{r} = -r_0 \omega \sin(\omega t) + A \omega \cos(\omega t) \sin(\delta)$$

$$v_{0r} = A \omega \sin(\delta)$$

$$\frac{v_{0r}}{\omega} = A \sin(\delta)$$

$$r = r_0 \cos(\omega t) + \frac{v_{0r}}{\omega} \sin(\omega t)$$

$$R = v_{0R}(t) + R_0$$

$$\frac{m_1 r_1 + m_2 r_2}{M} = v_{0R}(t) + R_0$$

$$m_1 r_{10} + m_2 r_{20} = M R_0 \quad \{ r_{10} = r_1(0), r_{20} = r_2(0) \}$$

$$m_1 \dot{r}_{10} + m_2 \dot{r}_{20} = M v_{0R} \quad \{ \dot{r}_{10} = \dot{r}_1(0), \dot{r}_{20} = \dot{r}_2(0) \}$$

$$m_1 (\dot{x}_1 \hat{x} + \dot{y}_1 \hat{y}) + m_2 (\dot{x}_2 \hat{x} + \dot{y}_2 \hat{y}) = M v_{0R}$$

$$X \hat{x} + Y \hat{y} = \frac{m_1 \dot{r}_{10} + m_2 \dot{r}_{20}}{M} t + \frac{m_1 \vec{r}_{10} + m_2 \vec{r}_{20}}{M}$$

$$\vec{r}_{10} = x_{10} \hat{x} + y_{10} \hat{y}, \quad \vec{r}_{20} = x_{20} \hat{x} + y_{20} \hat{y}$$

$$\dot{\vec{r}}_{10} = \dot{x}_{10} \hat{x} + \dot{y}_{10} \hat{y}, \quad \dot{\vec{r}}_{20} = \dot{x}_{20} \hat{x} + \dot{y}_{20} \hat{y}$$

$$X \hat{x} + Y \hat{y} = \frac{m_1 (\dot{x}_{10} \hat{x} + \dot{y}_{10} \hat{y}) + m_2 (\dot{x}_{20} \hat{x} + \dot{y}_{20} \hat{y})}{M} t + \frac{m_1 (x_{10} \hat{x} + y_{10} \hat{y}) + m_2 (x_{20} \hat{x} + y_{20} \hat{y})}{M}$$

Equating components:

$$X = \frac{m_1 \dot{x}_{10} + m_2 \dot{x}_{20}}{M} t + \frac{m_1 x_{10} + m_2 x_{20}}{M}$$

$$Y = \frac{m_1 \dot{y}_{10} + m_2 \dot{y}_{20}}{M} t + \frac{m_1 y_{10} + m_2 y_{20}}{M}$$

$$r = r_0 \cos(\omega t) + \frac{v_{0r}}{\omega} \sin(\omega t)$$

$$r_1 - r_2 = r_0 \cos(\omega t) + \frac{v_{0r}}{\omega} \sin(\omega t)$$

$$(x_1 - x_2) \hat{x} + (y_1 - y_2) \hat{y} = [(x_{10} - x_{20}) \hat{x} + (y_{10} - y_{20}) \hat{y}] \cos(\omega t) + \frac{(\dot{x}_{10} - \dot{x}_{20}) \hat{x} + (\dot{y}_{10} - \dot{y}_{20}) \hat{y}}{\omega} \sin(\omega t)$$

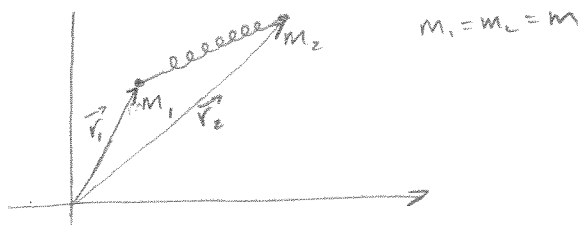
Equating components:

$$x_1 - x_2 = (x_{10} - x_{20}) \cos(\omega t) + \frac{\dot{x}_{10} - \dot{x}_{20}}{\omega} \sin(\omega t)$$

$$y_1 - y_2 = (y_{10} - y_{20}) \cos(\omega t) + \frac{\dot{y}_{10} - \dot{y}_{20}}{\omega} \sin(\omega t)$$

The  $X$  and  $Y$  motions are not accelerated. The  $x_1 - x_2$  and  $y_1 - y_2$  motion is periodic with angular frequency  $\sqrt{\frac{k}{\mu}}$ .

8.9



a)

$$T = \frac{1}{2} m_1 \dot{\vec{r}}_1^2 + \frac{1}{2} m_2 \dot{\vec{r}}_2^2$$

$$T = \frac{1}{2} m (\dot{\vec{r}}_1^2 + \dot{\vec{r}}_2^2)$$

$$U = \frac{1}{2} k (r - L)^2, \quad r = |\vec{r}_1 - \vec{r}_2|$$

$$L = T - U$$

$$L = \frac{1}{2} m (\dot{\vec{r}}_1^2 + \dot{\vec{r}}_2^2) - \frac{1}{2} k (|\vec{r}_1 - \vec{r}_2| - L)^2$$

In terms of CM:

$$L_{cm} = \frac{1}{2} M \dot{\vec{R}}^2$$

$$L_{rel} = \frac{1}{2} \mu \dot{\vec{r}}^2 - \frac{1}{2} k (r - L)^2$$

 $\dot{\vec{r}}$  in polar coordinates

$$\dot{\vec{r}} = \dot{r} \hat{r} + r \dot{\phi} \hat{\phi}$$

$$\dot{r}^2 = \dot{\vec{r}} \cdot \dot{\vec{r}}$$

$$\dot{r}^2 = \dot{r}^2 + r^2 \dot{\phi}^2$$

$$L = L_{cm} + L_{rel}$$

$$L = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu (\dot{r}^2 + r^2 \dot{\phi}^2) - \frac{1}{2} k (r - L)^2$$

$$b) \quad \frac{\partial L}{\partial R} = 0$$

$$\frac{\partial L}{\partial \dot{R}} = M \dot{R}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{R}} = M \ddot{R}$$

$$M\ddot{R} = 0$$

$$\dot{R} = \dot{R}_0$$

$$R = \dot{R}_0 t + R_0$$

$$\vec{R} = X\hat{x} + Y\hat{y}$$

$$\dot{\vec{R}} = \dot{X}\hat{x} + \dot{Y}\hat{y}$$

$$X = \dot{X}_0 t + X_0$$

$$Y = \dot{Y}_0 t + Y_0$$

$$c) \frac{\partial L}{\partial r} = \mu r \dot{\phi}^2 - k(r-L)$$

$$\frac{\partial L}{\partial \dot{r}} = \mu \dot{r}$$

$$\frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = \mu \ddot{r}$$

$$\mu \ddot{r} - \mu r \dot{\phi}^2 + k(r-L) = 0 \quad \frac{\partial L}{\partial r} - \frac{d}{dt} \frac{\partial L}{\partial \dot{r}} = 0$$

$$\frac{\partial L}{\partial \phi} = 0$$

$$\frac{\partial L}{\partial \dot{\phi}} = \mu r^2 \dot{\phi}$$

$$\frac{d}{dt} (\mu r^2 \dot{\phi}) = 0 \quad \left\{ \frac{\partial L}{\partial \phi} - \frac{d}{dt} \frac{\partial L}{\partial \dot{\phi}} = 0 \right\}$$

$$\mu r^2 \dot{\phi} = \text{const.}$$

i)  $r$  is constant

$$\ddot{r} = 0$$

$$\mu r^2 \dot{\phi} = \text{const.}$$

$$\dot{\phi} = \frac{\text{const.}}{\mu r^2} = \text{const.}$$

$$\mu r \dot{\phi}^2 - k(r-L) = 0$$

$$\mu r \dot{\phi}^2 - kr + kL = 0$$

$$r(\mu \dot{\phi}^2 - k) = -kL$$

$$r = \frac{kL}{(k - \mu \dot{\phi}^2)}$$

relative uniform circular motion.

ii)  $\phi$  is constant

$$\dot{\phi} = 0$$

$$\mu \ddot{r} + kr = kL$$

$$\ddot{r} + \frac{k}{\mu} r = \frac{k}{\mu} L$$

$$r_h = A \cos(\omega t - \beta), \quad \omega^2 = \frac{k}{\mu}$$

$$r_P: r = C$$

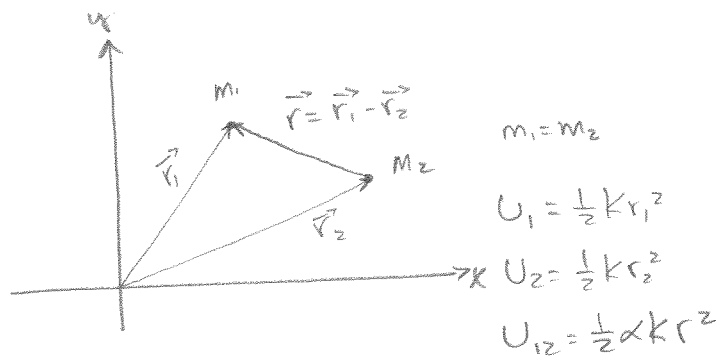
$$C = L$$

$$r_P = L$$

$$r = A \cos(\omega t - \beta) + L, \quad \omega^2 = \frac{k}{\mu} = \frac{k}{\frac{m^2}{2m_1}} = \frac{2k}{m_1}$$

Oscillation relative to each other,

8.10



$$a) T = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2$$

$$M = 2m_1$$

$$\mu = \frac{m_1^2}{m_1 + m_1} = \frac{m_1}{2}$$

$$T = m_1 \dot{\vec{R}}^2 + \frac{1}{4} m_1 \dot{\vec{r}}^2$$

$$\frac{m_1 \vec{r}_1 + m_1 \vec{r}_2}{2m_1} = \vec{R}$$

$$\vec{r}_1 + \vec{r}_2 = 2\vec{R}$$

$$\vec{r}_1 - \vec{r}_2 = \vec{r}$$

$$2\vec{r}_1 = 2\vec{R} + \vec{r}$$

$$\vec{r}_1 = \vec{R} + \frac{\vec{r}}{2}$$

$$2\vec{r}_2 = 2\vec{R} - \vec{r}$$

$$\vec{r}_2 = \vec{R} - \frac{\vec{r}}{2}$$

$$U_1 + U_2 = \frac{1}{2} k (r_1^2 + r_2^2)$$

$$= \frac{1}{2} k \left[ \vec{R}^2 + \vec{R} \cdot \vec{r} + \frac{\vec{r}^2}{4} + \vec{R}^2 - \vec{R} \cdot \vec{r} + \frac{\vec{r}^2}{4} \right]$$

$$= \frac{1}{2} k \left[ 2\vec{R}^2 + \frac{\vec{r}^2}{2} \right] = \frac{1}{2} k \left[ 2R^2 + \frac{r^2}{2} \right] \quad \{ \vec{R}^2 = \vec{R} \cdot \vec{R} = R^2, \vec{r}^2 = \vec{r} \cdot \vec{r} = r^2 \}$$

$$U = U_1 + U_2 + U_{12}$$

$$= \frac{1}{2} k \left[ 2R^2 + \frac{r^2}{2} \right] + \frac{1}{2} \alpha k r^2$$

$$= \frac{1}{2} k \left[ 2R^2 + \frac{r^2}{2} + \alpha r^2 \right]$$

$$L = T - U$$

$$L = m_1 \dot{\vec{R}}^2 + \frac{1}{4} m_1 \dot{\vec{r}}^2 - \frac{1}{2} k \left[ 2R^2 + \frac{r^2}{2} + \alpha r^2 \right]$$

$$b) L = L_{cm} + L_{rel}$$

$$L_{cm} = m_1 \dot{R}^2 - \frac{1}{2} k 2 R^2$$

$$L_{rel} = \frac{1}{4} m_1 \dot{r}^2 - \frac{1}{2} k \left[ \frac{1}{2} + \alpha r^2 \right]$$

CM:

$$\frac{\partial L_{cm}}{\partial R} = -2KR$$

$$\frac{\partial L_{cm}}{\partial \dot{R}} = 2m_1 \dot{R}$$

$$\frac{d}{dt} \frac{\partial L_{cm}}{\partial \dot{R}} = 2m_1 \ddot{R}$$

$$\frac{\partial L_{cm}}{\partial R} - \frac{d}{dt} \frac{\partial L_{cm}}{\partial \dot{R}} = 0$$

$$2m_1 \ddot{R} + 2KR = 0$$

$$\ddot{R} + \frac{K}{m_1} R = 0$$

$$R = A_1 e^{\gamma t}$$

$$\gamma^2 + \frac{K}{m_1} = 0$$

$$\gamma = \pm i\omega, \omega_1 = \sqrt{\frac{K}{m_1}}$$

$$R = A_1 e^{i\omega t} + A_2 e^{-i\omega t}$$

$$\boxed{R = A \cos(\omega_1 t - \phi_1)} \rightarrow \text{C.M. oscillates}$$

Rel:

$$\frac{\partial L_{rel}}{\partial r} = -k \left[ \frac{r}{2} + \alpha r \right]$$

$$\frac{\partial L_{rel}}{\partial \dot{r}} = \frac{1}{2} m_1 \dot{r}$$

$$\frac{d}{dt} \frac{\partial L_{rel}}{\partial \dot{r}} = \frac{1}{2} m_1 \ddot{r}$$

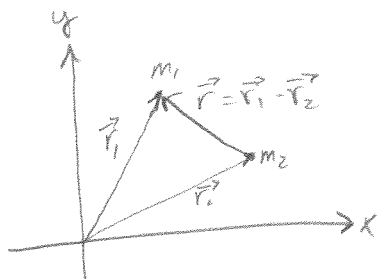
$$\frac{\partial L_{rel}}{\partial r} - \frac{d}{dt} \frac{\partial L_{rel}}{\partial \dot{r}} = 0$$

$$\frac{1}{2} m_1 \ddot{r} + k \left[ \frac{1}{2} + \alpha \right] r = 0$$

$$\ddot{r} + \frac{K}{m} [1 + 2\alpha] r = 0$$

$$\boxed{r = A \cos(\omega_2 t - \phi_2)}, \omega_2 = \sqrt{\frac{K[1+2\alpha]}{m}} \rightarrow \text{masses oscillate relative to each other}$$





$$U = \frac{1}{2} K r^2$$

$$T = \frac{1}{2} M \dot{\vec{R}}^2 + \frac{1}{2} \mu \dot{\vec{r}}^2$$

$$U = \frac{1}{2} K r^2$$

$$L_{cm} = \frac{1}{2} \mu \dot{\vec{r}}^2 - \frac{1}{2} K r^2$$

$$\frac{\partial L_{cm}}{\partial \vec{r}} = -K \vec{r}$$

$$\frac{\partial L_{cm}}{\partial \dot{\vec{r}}} = \mu \dot{\vec{r}}$$

$$\frac{d}{dt} \frac{\partial L_{cm}}{\partial \dot{\vec{r}}} = \mu \ddot{\vec{r}}$$

$$\mu \ddot{\vec{r}} + K \vec{r} = 0$$

$$\ddot{\vec{r}} + \frac{K}{\mu} \vec{r} = 0$$

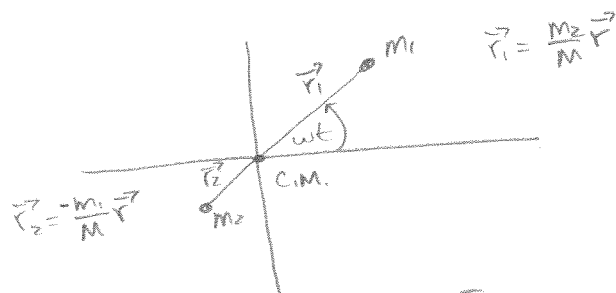
$$r = A \cos(\omega t - \delta), \quad \omega = \sqrt{\frac{K}{\mu}}$$

$$\vec{r} = A \cos(\omega t) \cos(\delta) + A \sin(\omega t) \sin(\delta)$$

$$A \cos(\delta) = C_1$$

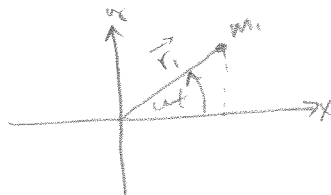
$$A \sin(\delta) = C_2$$

$$\vec{r} = C_1 \cos(\omega t) + C_2 \sin(\omega t)$$



Center of Mass Frame

Motion of  $m_1$  in CM frame



$$\vec{r}_1 = \frac{m_2}{M} (C_1 \cos(\omega t) \hat{x} + C_2 \sin(\omega t) \hat{y})$$

$$x = \frac{C_1 m_2}{M} \cos(\omega t)$$

$$y = \frac{C_2 m_2}{M} \sin(\omega t)$$

$$\frac{xM}{C_1 m_2} = \cos(\omega t)$$

$$\frac{yM}{C_2 m_2} = \sin(\omega t)$$

$$\cos^2(\omega t) + \sin^2(\omega t) = 1$$

$$\left(\frac{M}{C_1 m_2}\right)^2 x^2 + \frac{M}{C_2 m_2} y^2 = 1 \rightarrow \text{Equation of an ellipse with CM center.}$$

Motion of  $m_2$  in CM frame

$$\vec{r}_2 = \frac{-m_1}{M} (C_1 \cos(\omega t) \hat{x} + C_2 \sin(\omega t) \hat{y})$$

$$x = -\frac{m_1 C_1}{M} \cos(\omega t)$$

$$y = -\frac{m_1 C_2}{M} \sin(\omega t)$$

$$-\frac{xM}{C_1 m_1} = \cos(\omega t)$$

$$-\frac{yM}{C_2 m_1} = \sin(\omega t)$$

$$\cos^2(\omega t) + \sin^2(\omega t) = 1$$

$$\left(\frac{M}{C_1 m_1}\right)^2 x^2 + \left(\frac{M}{C_2 m_1}\right) y^2 = 1 \rightarrow \text{Equation of an ellipse with CM center}$$

$$0.1 \times U_{\text{eff}}(r) = -\frac{Gm_1m_2}{r} + \frac{l^2}{2\mu r^2}$$

$$a) \frac{dU}{dr} = \frac{Gm_1m_2}{r^2} - \frac{l^2}{\mu r^3}$$

$$-\frac{dU}{dr} = \mu \ddot{r}$$

$$\mu \ddot{r} = -\frac{Gm_1m_2}{r^2} + \frac{l^2}{\mu r^3}$$

For a fixed radius  $\ddot{r} = 0$

Let  $m_1 = m_{\text{planet}} = m_p$

$m_2 = m_{\text{sun}} = m_s$

$$\frac{Gm_p m_s}{r^2} = \frac{l^2}{\mu r^3}$$

$$r_0 = \frac{l^2}{Gm_p m_s \mu}, \quad \mu = \frac{m_p m_s}{m_p + m_s}$$

$$r_0 = \frac{l^2(m_p + m_s)}{G(m_p m_s)^2}$$

b) Taylor expansion of  $U_{\text{eff}}$

$$U_{\text{eff}} = U_{\text{eff}0} + \frac{dU_{\text{eff}}}{dr}\bigg|_{r_0} (r-r_0) + \frac{1}{2} \frac{d^2 U_{\text{eff}}}{dr^2}\bigg|_{r_0} (r-r_0)^2 + \dots$$

The constant  $U_{\text{eff}0}$  can be dropped. At  $r=r_0$ ,  $\frac{dU_{\text{eff}}}{dr} = 0$ . It is recognized that the quadratic term is of the form  $U = \frac{1}{2} k x^2$ . Therefore,

$$\omega_r^2 = \frac{k}{\mu} = \frac{1}{\mu} \frac{d^2 U_{\text{eff}}}{dr^2}\bigg|_{r_0}$$

$$\frac{d}{dr} \frac{dU_{\text{eff}}}{dr} = \frac{d}{dr} \left( \frac{Gm_1m_2}{r^2} - \frac{l^2}{\mu r^3} \right)$$

$$= -\frac{2Gm_1m_2}{r^3} + \frac{3l^2}{\mu r^4}$$

$$k = \frac{d^2 U_{\text{eff}}}{dr^2}\bigg|_{r_0} = -\frac{2Gm_1m_2}{r_0^3} + \frac{3l^2}{\mu r_0^4}$$

$$= -\frac{2Gm_p m_s G^3 (m_p m_s)^6}{l^6 (m_p + m_s)^3} + \frac{3l^2 G^4 (m_p m_s)^8}{\mu l^8 (m_p + m_s)^4}$$

$$= -\frac{2G^4 (m_p m_s)^7}{l^6 (m_p + m_s)^3} + \frac{3G^4 (m_p m_s)^7}{l^6 (m_p + m_s)^3}$$

$$= \frac{G^4 (m_p m_s)^7}{l^6 (m_p + m_s)^3}$$

(2)

$$\omega_r^2 = \frac{k}{\mu}$$

$$\omega_r^2 = \frac{G^4 (m_p m_s)^6}{l^6 (m_p + m_s)^2}$$

the radial equation is

$$\mu \ddot{r} = -\frac{G m_p m_s}{r^2} + \frac{l^2}{\mu r^3}$$

and with  $r = r_0 + \epsilon$ ,  $\epsilon \ll 1$

$$\begin{aligned} \mu \ddot{\epsilon} &= -G m_p m_s (r_0 + \epsilon)^{-2} + \frac{l^2}{\mu} (r_0 + \epsilon)^{-3} \quad \{ (a+b)^n = a^n + n a^{n-1} b + \frac{n(n-1)}{2} a^{n-2} b^2 \} \\ &\approx -G m_p m_s (r_0^{-2} - 2 r_0^{-3} \epsilon + \dots) + \frac{l^2}{\mu} (r_0^{-3} - 3 r_0^{-4} \epsilon + \dots) \quad \{ \text{terms of } \epsilon^2 \text{ and higher are ignored} \} \\ &\approx -\frac{G m_p m_s}{r_0^2} + \frac{2 G m_p m_s}{r_0^3} \epsilon + \frac{l^2}{\mu r_0^3} - \frac{3 l^2}{r_0^4} \epsilon \\ &\approx \left( \frac{2 G m_p m_s}{r_0^3} - \frac{3 l^2}{r_0^4} \right) \epsilon + \frac{l^2}{\mu r_0^3} - \frac{G m_p m_s}{r_0^2} \\ &\approx \frac{-G^4 (m_p m_s)^7}{l^6 (m_p + m_s)^3} \epsilon + \frac{l^2}{\mu r_0^3} - \frac{G m_p m_s}{r_0^2} \end{aligned}$$

The homogeneous equation is of the form  $F = -Kx$ , a restoring force, and is therefore stable. Another view is that the Taylor expansion about  $r_0$  gives an effective potential of  $\frac{1}{2} K(r-r_0)^2$  which is the potential of a harmonic oscillator.

Period of oscillation:

$$\begin{aligned} T_r &= \frac{2\pi}{\omega_r} \\ &= \frac{2\pi (m_p + m_s) l^3}{G^2 (m_p m_s)^3} \end{aligned}$$

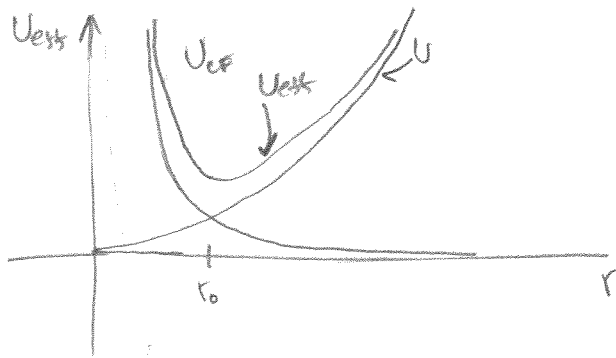
Period of orbit

$$\dot{\phi} = \frac{l}{\mu r^2}$$

$$\begin{aligned} T_0 &= \frac{2\pi}{\dot{\phi}} \\ &= \frac{2\pi \mu r_0^2}{l} \\ &= \frac{2\pi \frac{m_p m_s}{l} \frac{(m_p + m_s)^2 l^4}{G^2 (m_p m_s)^4}}{l} \\ &= \frac{2\pi (m_p + m_s) l^3}{G^2 (m_p m_s)^3} \rightarrow T_r = T_0 \end{aligned}$$

8.13

$$a) U = \frac{1}{2}Kr^2, U_{cf} = \frac{l^2}{2\mu r^2}, U_{eff} = U + U_{cf}$$



$$b) U_{eff} = \frac{1}{2}Kr^2 + \frac{l^2}{2\mu r^2}$$

$$\frac{dU_{eff}}{dr} = Kr - \frac{l^2}{\mu r^3}$$

$$\left. \frac{dU_{eff}}{dr} \right|_{r_0} = 0$$

$$Kr_0 = \frac{l^2}{\mu r_0^3}$$

$$r_0^4 = \frac{l^2}{\mu K}$$

$$r_0 = \left( \frac{l^2}{\mu K} \right)^{1/4}$$

$$c) U_{eff} = U_{eff_0} + \left. \frac{dU_{eff}}{dr} \right|_{r_0} (r - r_0) + \frac{1}{2} \left. \frac{d^2 U_{eff}}{dr^2} \right|_{r_0} (r - r_0)^2 + \dots$$

$$\frac{d^2 U_{eff}}{dr^2} = \frac{d}{dr} \left( Kr - \frac{l^2}{\mu r^3} \right)$$

$$= K + \frac{3l^2}{\mu r^4}$$

$$\left. \frac{d^2 U_{eff}}{dr^2} \right|_{r_0} = K + \frac{3l^2}{\mu r_0^4}$$

$$= K + \frac{3l^2}{\mu} \frac{\mu K}{l^2}$$

$$= K + 3K$$

$$= 4K$$

$$\omega_r^2 = \frac{4K}{\mu} \rightarrow \omega_r = \sqrt{\frac{4K}{\mu}}$$

8.17

$$U = Kr^n, \quad Kn > 0$$

$$U_{\text{eff}} = U + U_{\text{cf}}$$

$$= Kr^n + \frac{l^2}{2\mu r^2}$$

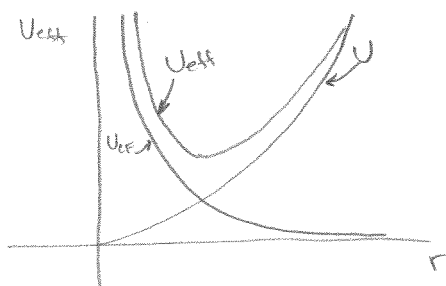
$$F = -\frac{dU_{\text{eff}}}{dr}$$

$$= -(Kn r^{n-1} - \frac{l^2}{\mu r^3})$$

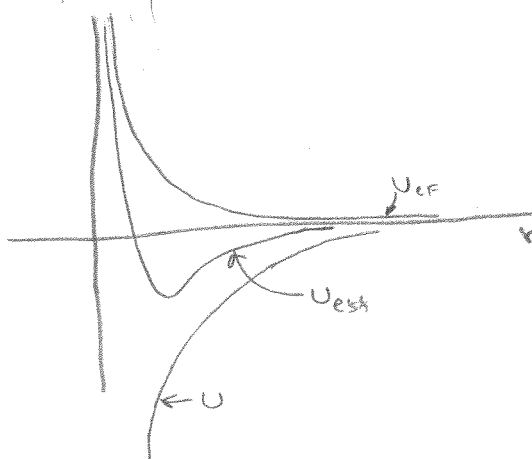
$$F = -Kn r^{n-1} + \frac{l^2}{\mu r^3}$$

The condition  $Kn > 0$  indicates the force is a central force and a centrifugal force.

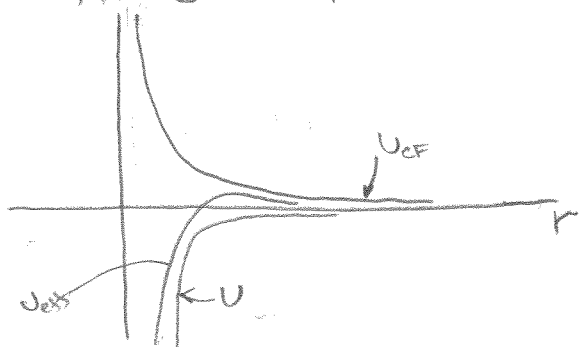
i)  $n=2 \rightarrow U = Kr^2$



ii)  $n=-1 \rightarrow U = -\frac{K}{r}$



iii)  $n=-3 \rightarrow U = -\frac{K}{r^3}$



b) Fixed radius  $\rightarrow \ddot{r} = 0 \rightarrow F = 0$

$$Kn r_0^{n-1} = \frac{l^2}{\mu r_0^3}$$

$$r_0^{n+2} = \frac{l^2}{Kn\mu}$$

$$r_0 = \left( \frac{l^2}{Kn\mu} \right)^{\frac{1}{n+2}}$$

(2)

For stability  $\left. \frac{d^2 U_{\text{eff}}}{dr^2} \right|_{r_0} > 0$  {concave up}

$$\frac{d^2 U_{\text{eff}}}{dr^2} = \frac{d}{dr} \left( K n r^{n-1} - \frac{l^2}{u r^3} \right)$$

$$= K n (n-1) r^{n-2} + \frac{3l^2}{u r^4}$$

$$K n (n-1) r_0^{n-2} + \frac{3l^2}{u r_0^4} > 0$$

$$K n (n-1) r_0^{n+2} + \frac{3l^2}{u} > 0$$

$$r_0 = \left( \frac{l^2}{K n u} \right)^{\frac{1}{n+2}}$$

$$K n (n-1) \left( \frac{l^2}{K n u} \right)^{\frac{n+2}{n+2}} + \frac{3l^2}{u} > 0$$

$$(n-1) \frac{l^2}{u} + 3 \frac{l^2}{u} > 0$$

$$n+2 > 0$$

$$n > -2$$

The sketches for  $n=2$  and  $n=-1$  show stability. The  $n=-3$  sketch shows no stable equilibrium points. So

$n > -2$  is in agreement with the sketches.

c) Small radial oscillations:

$$K = \left. \frac{d^2 U_{\text{eff}}}{dr^2} \right|_{r_0} = K n (n-1) r_0^{n-2} + \frac{3l^2}{u r_0^4}$$

$$K r_0^4 = K n (n-1) r_0^{n+2} + \frac{3l^2}{u}$$

$$K r_0^4 = K n (n-1) \left( \frac{l^2}{K n u} \right)^{\frac{n+2}{n+2}} + \frac{3l^2}{u}$$

$$= (n-1) \frac{l^2}{u} + 3 \frac{l^2}{u}$$

$$= \frac{l^2}{u} (n+2)$$

$$K = \frac{l^2}{u r_0^4} (n+2)$$

$$\omega_r^2 = \frac{K}{u} = \frac{l^2}{u^2 r_0^4} (n+2)$$

$$\omega_r = \frac{l}{\mu r_0^2} (n+z)^{1/2}$$

$$T_{osc} = \frac{2\pi}{\omega_r}$$

$$T_{osc} = \frac{2\pi \mu r_0^2}{l \sqrt{n+z}}$$

$$\dot{\phi} = \frac{l}{\mu r_0^2}$$

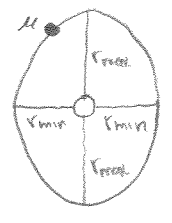
$$T_{orb} = \frac{2\pi}{\dot{\phi}}$$

$$T_{orb} = \frac{2\pi \mu r_0^2}{l}$$

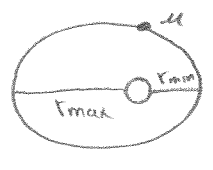
$$T_{osc} = \frac{T_{orb}}{\sqrt{n+z}}$$

If  $\sqrt{n+z}$  is a rational number then  $p\omega_{osc} = q\omega_{orb}$  and produces a Lissajous figure which is a closed orbit.

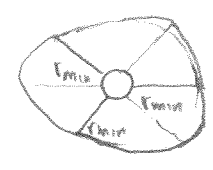
i)  $n=2 \rightarrow T_{osc} = \frac{T_{orb}}{2}$



ii)  $n=-1 \rightarrow T_{osc} = T_{orb}$



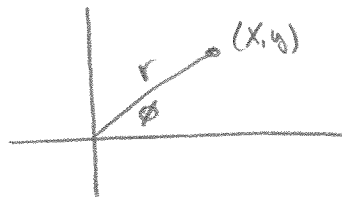
iii)  $n=7 \rightarrow T_{osc} = \frac{T_{orb}}{3}$





8.16

$$r = \frac{c}{1 + e \cos \phi}$$



$$x = r \cos \phi$$

$$y = r \sin \phi$$

$$r = (x^2 + y^2)^{1/2}$$

$$r + e r \cos \phi = c$$

$$(x^2 + y^2)^{1/2} + e x = c$$

$$x^2 + y^2 = (c - e x)^2$$

$$= c^2 - 2 c e x + e^2 x^2$$

$$x^2 - e^2 x^2 + 2 c e x + y^2 = c^2$$

$$(1 - e^2) x^2 + 2 c e x + y^2 = c^2$$

$$x^2 + \frac{2 c e x}{1 - e^2} + \frac{y^2}{1 - e^2} = \frac{c^2}{1 - e^2}$$

$$x^2 + \frac{2 c e x}{1 - e^2} + \left( \frac{c e}{1 - e^2} \right)^2 + \frac{y^2}{1 - e^2} = \frac{c^2}{1 - e^2} + \left( \frac{c e}{1 - e^2} \right)^2$$

$$\left( x + \frac{c e}{1 - e^2} \right)^2 + \frac{y^2}{1 - e^2} = \frac{c^2 - c^2 e^2 + c^2 e^2}{(1 - e^2)^2} = \frac{c^2}{(1 - e^2)^2}$$

$$\frac{(x + d)^2}{a^2} + \frac{y^2}{b^2} = 1$$

$$a^2 = \frac{c^2}{(1 - e^2)^2} \rightarrow a = \frac{c}{1 - e^2}$$

$$b^2 = \frac{c^2}{1 - e^2} \rightarrow b = \frac{c}{\sqrt{1 - e^2}}$$

$$d = \frac{c e}{1 - e^2} \rightarrow d = a e$$

8.17

$$a) G = \vec{r} \cdot \vec{p}$$

$$\frac{dG}{dt} = \frac{d}{dt}(\vec{r} \cdot \vec{p})$$

$$= \frac{d\vec{r}}{dt} \cdot \vec{p} + \frac{d\vec{p}}{dt} \cdot \vec{r}$$

$$\frac{dG}{dt} = \vec{v} \cdot \vec{p} + \vec{F} \cdot \vec{r}$$

$$dG = (\vec{v} \cdot \vec{p} + \vec{F} \cdot \vec{r}) dt$$

$$dG = (\vec{v} \cdot m\vec{v} + \vec{F} \cdot \vec{r}) dt$$

$$\int_0^t dG = \int_0^t m v^2 dt' + \int_0^t \vec{F} \cdot \vec{r} dt' \quad \{ \vec{v} \cdot \vec{v} = v^2 \}$$

$$\frac{\int dG}{t} = \frac{1}{t} \int_0^t m v^2 dt + \frac{1}{t} \int_0^t \vec{F} \cdot \vec{r} dt$$

$$\frac{G(t) - G(0)}{t} = 2\langle T \rangle + \langle \vec{F} \cdot \vec{r} \rangle, \quad \langle f \rangle = \frac{1}{t} \int_0^t f dt'$$

b) If the orbit is periodic then  $t$  can be made as large as desired because the righthand side of the equation is the time average. Therefore, the left side of the equation approaches zero as  $t \rightarrow \infty$ .

c) Let  $t \rightarrow \infty$

$$0 = 2\langle T \rangle + \langle \vec{F} \cdot \vec{r} \rangle$$

$$\vec{F} = -\frac{d}{dr}(U) \hat{r}$$

$$= -\frac{d}{dr}(kr^n) \hat{r}$$

$$\vec{F} = -nkr^{n-1} \hat{r}$$

$$\vec{F} \cdot \vec{r} = -nkr^{n-1} \hat{r} \cdot r \hat{r}$$

$$= -nkr^n$$

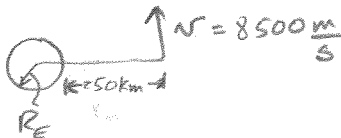
$$\vec{F} \cdot \vec{r} = -nU$$

$$2\langle T \rangle = -\langle \vec{F} \cdot \vec{r} \rangle$$

$$= n\langle U \rangle$$

$$\boxed{\langle T \rangle = \frac{n}{2} \langle U \rangle}$$

8.18



$$r_{\min} = \frac{C}{1+E}$$

$$C = \frac{l^2}{\gamma \mu} \quad , \quad \gamma = G m_1 m_2 = G M_e m_s, \quad \mu = \frac{M_e m_s}{M_e + m_s}$$

$$l = m_s r_{\min}^2 \dot{\phi}$$

$$v = r_{\min} \dot{\phi}$$

$$l = m_s r_{\min} v \quad , \quad r_{\min} = R_E + 250 \text{ km} = 6.4 \times 10^6 \text{ m} + 2.5 \times 10^5 \text{ m} = 6.65 \times 10^6 \text{ m}$$

$$C = \frac{m_s^2 r_{\min}^2 v^2}{G M_e m_s \left( \frac{M_e m_s}{M_e + m_s} \right)}$$

$$C = \frac{r_{\min}^2 v^2}{G \frac{M_e^2}{M_e + m_s}}$$

$$r_{\min} + r_{\min} E = C$$

$$E = \frac{C - r_{\min}}{r_{\min}}$$

$$= \frac{C}{r_{\min}} - 1$$

$$= \frac{r_{\min} v^2}{G \frac{M_e^2}{M_e + m_s}} - 1$$

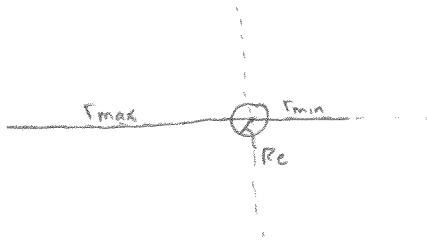
$$E \approx \frac{r_{\min} v^2}{G M_e} - 1 \quad \{ M_e + m_s \approx M_e \}$$

$$\frac{G M_e}{R_E^2} = g$$

$$G M_e = R_E^2 g$$

$$E = \frac{r_{\min} v^2}{R_E^2 g} - 1 = \frac{(6.65 \times 10^6)^2 (8500)^2}{(6.4 \times 10^6)^2 (9.81)} - 1 = 0.1957 = \boxed{0.20}$$

8.19



$$r_{\min} = R_E + 300 \text{ km}$$

$$r_{\max} = R_E + 3000 \text{ km}$$

$$r_{\min} = \frac{c}{1+E}$$

$$r_{\max} = \frac{c}{1-E}$$

$$\frac{r_{\min}}{r_{\max}} = \frac{1-E}{1+E}$$

$$r_{\min} + r_{\min} E = r_{\max} - r_{\max} E$$

$$E(r_{\min} + r_{\max}) = r_{\max} - r_{\min}$$

$$E = \frac{r_{\max} - r_{\min}}{r_{\max} + r_{\min}}$$

$$= \frac{R_E + 3000 \text{ km} - R_E - 300 \text{ km}}{R_E + 3000 \text{ km} + R_E + 300 \text{ km}}$$

$$= \frac{2700 \text{ km}}{2R_E + 3300 \text{ km}}$$

$$= \frac{2700 \text{ km}}{2(6.4 \times 10^3 \text{ km}) + 3300 \text{ km}}$$

$$= 0.1677$$

$$\boxed{E = 0.1677}$$

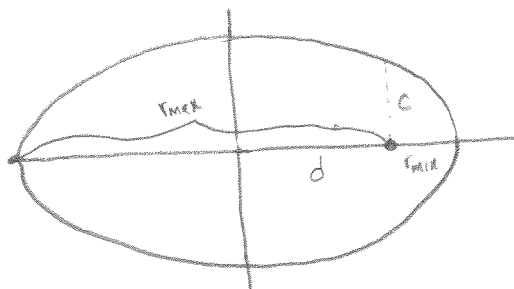
$$r = \frac{c}{1+E \cos \phi} \quad \phi = 90^\circ$$

$$r = c$$

$$h = c - R_E$$

$$= 7824 - 6400$$

$$\boxed{h = 1424 \text{ km}}$$



$a$  = semi-major axis

$$a = \frac{r_{\min} + r_{\max}}{2}$$

$$a = \frac{c}{1-E^2}$$

$$c = a(1-E^2)$$

$$= \frac{r_{\min} + r_{\max}}{2} (1-E^2)$$

$$= \frac{[2(6.4 \times 10^3) + 3300]}{2} (1-(0.1677)^2)$$

$$= 7824 \text{ km}$$

$$r_{\max} = \frac{C}{1-\epsilon}$$

$$C = \frac{l^2}{\gamma \mu}$$

$$r_{\max} = \frac{l^2}{\gamma \mu (1-\epsilon)}$$

$$r_{\max} \gamma \mu (1-\epsilon) = l^2$$

$$\lim_{l \rightarrow 0} l^2 = \lim_{\epsilon \rightarrow 1} C_1 (1-\epsilon) \quad \{ C_1 = r_{\max} \gamma \mu \}$$

$$0 = \lim_{\epsilon \rightarrow 1} (1-\epsilon)$$

$$\boxed{\epsilon \rightarrow 1}$$

$$r_{\min} = \frac{C}{1+\epsilon}$$

$$r_{\min} = \frac{l^2}{\gamma \mu (1+\epsilon)}$$

$$\begin{aligned} \lim_{l \rightarrow 0} r_{\min} &= \lim_{l \rightarrow 0} \frac{l^2}{\gamma \mu (1+\epsilon)} \\ &= \frac{0}{\gamma \mu (2)} \end{aligned}$$

$$\lim_{l \rightarrow 0} r_{\min} \rightarrow 0$$

When  $l \rightarrow 0$  with fixed  $r_{\max}$ ,  $C \rightarrow 0$ ,  $r_{\min} \rightarrow 0$ , and  $\epsilon \rightarrow 1$  {parabolic orbit}. So, the elliptical orbit becomes very thin and long. Since  $r_{\min} \rightarrow 0$  then  $a \rightarrow \frac{r_{\max}}{2}$ .

8.21

a) See Solution 8.20

$$b) \tau^2 = 4\pi^2 \frac{a^3 \mu}{\gamma}$$

$$a = \frac{r_{\max}}{2}$$

$$\mu = \frac{M_s m_c}{M_s + m_c} \approx m_c$$

$$\gamma = G M_s m_c$$

$$\tau^2 = 4\pi^2 \frac{(r_{\max}/2)^3 m_c}{G M_s m_c}$$

$$\tau^2 = \frac{\pi^2 (r_{\max})^3}{2 G M_s}$$

$$\tau = \frac{\pi (r_{\max})^{3/2}}{\sqrt{2 G M_s}}$$

$$c) t = \sqrt{\frac{m_c}{2}} \int_{x_0}^x \frac{dx'}{\sqrt{E - U(x')}}$$

$$E = T_0 + U_0$$

$$E = -\frac{G M_s m_c}{r_{\max}}$$

$$U(r) = -\frac{G M_s m_c}{r}$$

$$m = m_c$$

$$r_0 = x_0 = r_{\max}$$

$$r = x = 0$$

$$t = \sqrt{\frac{m_c}{2}} \int_{r_{\max}}^0 \frac{dr}{\sqrt{E + \frac{G M_s m_c}{r}}}$$

$$\text{Let } \frac{G M_s m_c}{r} = -E \sec^2 \theta \quad \{E < 0\}$$

$$\frac{G M_s m_c \cos^2 \theta}{-E} = r$$

$$\frac{2 G M_s m_c \cos \theta \sin \theta d\theta}{E} = dr$$

$$U = - \int -\frac{G M_s m_c}{r^2} dr$$

$$U = -\frac{G M_s m_c}{r}$$

$$\begin{aligned}
 t &= \sqrt{\frac{m_c}{2}} \int_{\theta_0}^{\theta} \frac{\frac{2GM_s m_c}{E} \cos\theta \sin\theta d\theta}{\sqrt{E - E \sec^2\theta}} \\
 &= 2GM_s m_c \sqrt{\frac{m_c}{2}} \frac{1}{(-E)^{3/2}} \int_{\theta_0}^{\theta} \frac{\cos\theta \sin\theta d\theta}{\sqrt{-1 + \sec^2\theta}} \\
 &= 2GM_s m_c \sqrt{\frac{m_c}{2}} \left(\frac{GM_s m_c}{r_{\max}}\right)^{3/2} \int_{\theta_0}^{\theta} \frac{\cos\theta \sin\theta d\theta}{\tan\theta}
 \end{aligned}$$

$$= \frac{2(r_{\max})^{3/2}}{\sqrt{2GM_s}} \int_{\theta_0}^{\theta} \cos^2\theta d\theta$$

$$\begin{aligned}
 \frac{GM_s m_c \cos^2\theta}{-E} &= r \\
 r_{\max} \cos^2\theta &= r
 \end{aligned}$$

$$\begin{aligned}
 r=0 &\rightarrow \theta = \frac{\pi}{2} \\
 r=r_{\max} &\rightarrow \theta_0 = 0
 \end{aligned}$$

$$= \frac{2(r_{\max})^{3/2}}{\sqrt{2GM_s}} \frac{1}{2} \int_0^{\pi/2} (1 + \cos 2\theta) d\theta$$

$$= \frac{(r_{\max})^{3/2}}{\sqrt{2GM_s}} \left[ \theta \Big|_0^{\pi/2} + \frac{1}{2} \sin 2\theta \Big|_0^{\pi/2} \right]$$

$$= \frac{(r_{\max})^{3/2}}{\sqrt{2GM_s}} \left[ \frac{\pi}{2} \right]$$

$$t = \frac{\pi (r_{\max})^{3/2}}{2\sqrt{2GM_s}}$$

d) The motion would be that of a one-dimensional oscillator.

$$T_{\text{osc}} = 4t \quad \{ \text{time to complete one cycle} \}$$

$$T_{\text{osc}} = \frac{2\pi (r_{\max})^{3/2}}{\sqrt{2GM_s}}$$

e) Comparing with b)

$$T_{\text{osc}} = 2\tau$$

8.22

$$a) F(r) = \frac{K}{r^3}$$

$$U_{\text{eff}} = U + U_{\text{cf}}$$

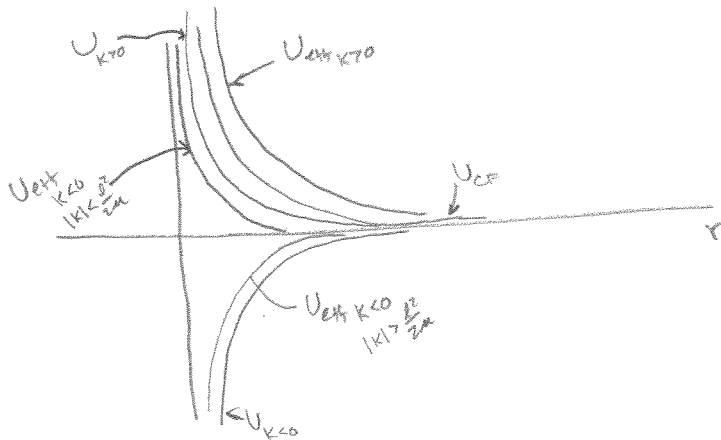
$$U = - \int \vec{F} \cdot d\vec{r} \quad \{ \vec{F} = \frac{K}{r^3} \hat{r} \}$$

$$= -K \int r^{-3} dr$$

$$U = \frac{K}{2} r^{-2} = \frac{K}{2r^2}$$

$$U_{\text{cf}} = \frac{l^2}{2\mu r^2}$$

$$U_{\text{eff}} = \frac{K}{2r^2} + \frac{l^2}{2\mu r^2}$$



$E > 0$ :  $U_{\text{eff}}, K > 0$  and  $U_{\text{eff}}, K < 0, |K| < \frac{l^2}{2\mu}$  produce open orbits that approach infinity.

$E < 0$ :  $U_{\text{eff}}, K < 0, |K| > \frac{l^2}{2\mu}$  produces a well such that the object is attracted to the force center,

$$b) \mu''(\phi) = -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} F$$

$$= -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} \frac{K}{r^3}$$

$$= -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} K \mu(\phi)^3 \quad \{ \mu = \frac{1}{r} \}$$

$$= -\mu(\phi) - \frac{\mu K}{r^2} \mu(\phi)$$

$$= -\mu(\phi) \left[ 1 + \frac{\mu K}{r^2} \right]$$

$$\mu''(\phi) + \mu(\phi) \left[ 1 + \frac{\mu K}{r^2} \right] = 0$$

$$r^2 + \left[ 1 + \frac{\mu K}{r^2} \right] = 0 \quad \{ \mu(\phi) = A e^{\gamma \phi} \}$$

$$r^2 = - \left[ 1 + \frac{\mu K}{r^2} \right]$$



$$\gamma = \pm i\omega$$

$$\omega = \sqrt{1 + \frac{\mu k}{L^2}}$$

i)  $k > 0 \rightarrow \omega$  is real

$$u = A \cos(\omega\phi - d)$$

$$\frac{1}{r} = A \cos(\omega\phi - d)$$

$$r = \frac{r_0}{\cos(\omega\phi - d)}$$

This is the equation for a straight line in agreement with  $U_{eff}$ ; unbounded orbit,  $k > 0$

ii)  $-1 < \frac{\mu k}{L^2} \leq 0$

$\omega = \sqrt{1 + \frac{\mu k}{L^2}} > 0 \rightarrow$  same condition as  $k > 0$  and agrees with the sketch; unbounded orbit.

iii)  $\frac{\mu k}{L^2} = -1$

$$\omega = 0$$

$r = r_0 \rightarrow$  no acceleration

iv)  $\frac{\mu k}{L^2} < -1$

$$\omega < 0 \rightarrow \gamma^2 = \omega^2$$

$$\gamma = \pm \omega$$

$$u = A_1 e^{\omega\phi} + A_2 e^{-\omega\phi}$$

$$= A_1 [e^{\omega\phi} + e^{-\omega\phi}]$$

$$\frac{1}{r} = A_2 \cosh(\phi) \quad \{ A_2 = \frac{A_1}{2} \}$$

$$r = \frac{1}{A_2 \cosh(\phi)}$$

$$r = A_3 \operatorname{sech}(\phi) \quad \{ A_3 = \frac{1}{A_2} \}$$

This is a hyperbolic trajectory and agrees with the sketch; unbounded orbit.

Therefore, an inverse cube central force does not produce bounded orbits.

8.23

$$F(r) = -\frac{K}{r^2} + \frac{\gamma}{r^3}, \quad K > 0, \gamma > 0$$

$$a) \mu''(\phi) = -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} F$$

$$\mu'(\phi) = -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} \left[ -\frac{K}{r^2} + \frac{\gamma}{r^3} \right], \quad r = \frac{1}{\mu}$$

$$\mu''(\phi) = -\mu(\phi) - \frac{\mu}{r^2 \mu(\phi)^2} \left[ -K\mu(\phi)^2 + \gamma\mu(\phi)^3 \right]$$

$$r^2 = \frac{1}{\mu^2}, \quad r^3 = \frac{1}{\mu^3}$$

$$\mu''(\phi) = -\mu(\phi) - \frac{\gamma\mu}{r^2} \mu(\phi) + \frac{\mu K}{r^2}$$

$$= -\mu(\phi) \left[ 1 + \frac{\gamma\mu}{r^2} \right] + \frac{\mu K}{r^2}$$

$$\mu''(\phi) + \mu(\phi) \left[ 1 + \frac{\gamma\mu}{r^2} \right] = \frac{\mu K}{r^2}$$

$$\mu_h: \mu(\phi) = A_1 e^{\gamma\phi}$$

$$\gamma^2 + \left[ 1 + \frac{\gamma\mu}{r^2} \right] = 0$$

$$\gamma^2 = -\left[ 1 + \frac{\gamma\mu}{r^2} \right]$$

$$\gamma = \pm i\omega, \quad \omega = \sqrt{1 + \frac{\gamma\mu}{r^2}}$$

$$\mu_h(\phi) = A_2 \cos(\omega\phi - \delta)$$

$$\mu_p: \mu = C_1$$

$$C_1 \left[ 1 + \frac{\gamma\mu}{r^2} \right] = \frac{\mu K}{r^2}$$

$$C_1 \left[ \frac{r^2 + \gamma\mu}{r^2} \right] = \frac{\mu K}{r^2}$$

$$C_1 = \frac{\mu K}{r^2} \left[ \frac{r^2}{r^2 + \gamma\mu} \right]$$

$$= \frac{\mu K}{r^2 + \gamma\mu}$$

$$\mu_p = \frac{\mu K}{r^2 + \gamma\mu}$$

$$\mu = A_2 \cos(\omega\phi - \delta) + \frac{\mu K}{r^2 + \gamma\mu}$$

$$\frac{1}{r} = A_2 \cos(\omega\phi - \delta) + \frac{\mu K}{r^2 + \gamma\mu}$$

$$= A_2 \cos(\omega\phi) + A_3 \quad \sum \delta = 0 \text{ for convenience, } A_3 = \frac{\mu K}{r^2 + \gamma\mu}$$

$$= A_3 \left[ \frac{A_2}{A_3} \cos(\omega\phi) + 1 \right]$$

$$= A_3 \left[ A_4 \cos(\omega\phi) + 1 \right] \quad \sum A_4 = \frac{A_2}{A_3}$$

$$r = \frac{\frac{1}{A_3}}{1 + A_4 \cos(\omega \phi)}$$

$$r = \frac{c}{1 + \epsilon \cos(\beta \phi)}, \quad c = \frac{1}{A_3}, \quad A_4 = \epsilon, \quad \omega = \beta$$

b)  $c = \frac{1}{A_3}$

$$c = \frac{l^2 + \gamma \mu}{\mu k}$$

$$\beta = \omega$$

$$\beta = \sqrt{1 + \frac{\gamma \mu}{l^2}}$$

$$r = \frac{\frac{l^2 + \gamma \mu}{\mu k}}{1 + \epsilon \cos\left(\left[1 + \frac{\gamma \mu}{l^2}\right]\phi\right)}$$

For  $0 < \epsilon < 1$ ,  $r$  is never  $\infty$  ( $1 + \epsilon \cos(\beta \phi) > 0$ ) and is therefore a bounded orbit.

c) For the orbit to be closed it has to return to its initial conditions.

$$\begin{aligned} \beta &= \sqrt{1 + \frac{\gamma \mu}{l^2}} \\ &= \sqrt{\frac{l^2 + \gamma \mu}{l^2}} = \omega \end{aligned}$$

$$\begin{aligned} T_{osc} &= \frac{2\pi}{\omega} \\ &= \frac{2\pi}{\beta} \end{aligned}$$

$$T_{orb} = \frac{2\pi}{\dot{\phi}}$$

For closed orbits

$$n T_{orb} = m T_{osc}, \quad n \text{ and } m \text{ are rational numbers}$$

$$n \frac{2\pi}{\dot{\phi}} = m \frac{2\pi}{\beta}$$

$$\boxed{\beta = \frac{m}{n} \dot{\phi}}$$

When  $\gamma \rightarrow 0$ :  $r = \frac{\frac{l^2}{\mu k}}{1 + \epsilon \cos(\phi)}$  which is a Kepler orbit.

8.24

$$F(r) = -\frac{k}{r^2} - \frac{\gamma}{r^3}$$

$$u''(\phi) = -u(\phi) - \frac{\mu}{r^2 u(\phi)^2} F$$

$$u''(\phi) = -u(\phi) - \frac{\mu}{r^2 u(\phi)^2} \left[ -\frac{k}{r^2} - \frac{\gamma}{r^3} \right]$$

$$u''(\phi) = -u(\phi) - \frac{\mu}{r^2 u(\phi)^2} \left[ -k u(\phi)^2 - \gamma u(\phi)^3 \right], \quad r = \frac{1}{u}$$

$$u''(\phi) = -u(\phi) + \frac{\mu \gamma}{r^2} u(\phi) + \frac{\mu k}{r^2}$$

$$u''(\phi) = -u(\phi) \left[ 1 - \frac{\mu \gamma}{r^2} \right] + \frac{\mu k}{r^2}$$

$$u''(\phi) + u(\phi) \left[ 1 - \frac{\mu \gamma}{r^2} \right] = \frac{\mu k}{r^2}$$

$$u_h: u(\phi) = A_1 e^{\gamma \phi}$$

$$\gamma^2 + \left[ 1 - \frac{\mu \gamma}{r^2} \right] = 0$$

$$\gamma^2 = - \left[ 1 - \frac{\mu \gamma}{r^2} \right]$$

$$\gamma_1 = \pm i \omega, \quad \omega = \sqrt{1 - \frac{\mu \gamma}{r^2}}$$

$$u_h(\phi) = A_2 \cos(\omega \phi - \ell)$$

$$u_p: u = C_1$$

$$C_1 \left[ 1 - \frac{\mu \gamma}{r^2} \right] = \frac{\mu k}{r^2}$$

$$C_1 \left[ \frac{r^2 - \mu \gamma}{r^2} \right] = \frac{\mu k}{r^2}$$

$$C_1 = \frac{\mu k}{r^2 - \mu \gamma}$$

$$u_p = \frac{\mu k}{r^2 - \mu \gamma}$$

$$u = A_2 \cos(\omega \phi) + \frac{\mu k}{r^2 - \mu \gamma}, \quad \ell = 0 \text{ for convenience}$$

$$\frac{1}{r} = A_3 \left[ \frac{A_2}{A_3} \cos(\omega \phi) + 1 \right], \quad A_3 = \frac{\mu k}{r^2 - \mu \gamma}$$

$$r = \frac{\frac{1}{A_3}}{1 + \epsilon \cos(\omega \phi)}$$

$$r = \frac{\frac{r^2 - \mu \gamma}{\mu k}}{(1 + \epsilon \cos(\omega \phi))}$$

For  $l^2 > a^2$

$$\omega = \sqrt{\frac{l^2 - a^2}{l^2}} > 0$$

This is a bounded orbit since

$$\gamma_1^2 = \pm i\omega$$

resulting in a cosine solution.

For  $l^2 < -a^2$

$$\omega = \sqrt{\frac{l^2 - a^2}{l^2}}$$

$$l^2 - a^2 < 0$$

Then

$$\gamma_1^2 - \omega_1^2 = 0, \quad \omega_1^2 = \frac{l^2 - a^2}{l^2}$$

$$\gamma_1^2 = \omega_1^2$$

$$\gamma_1 = \pm \omega_1$$

$$\begin{aligned} u_h(\phi) &= A_1 e^{\pm \omega \phi} \\ &= A_1 e^{\omega \phi} + A_2 e^{-\omega \phi} \\ &= \frac{A_2}{2} [e^{\omega \phi} + e^{-\omega \phi}] \end{aligned}$$

$$u_h(\phi) = A_2 \cosh(\omega \phi)$$

This is an unbounded orbit.

Y.26

$$r = \frac{c}{1 + \epsilon \cos \phi}, \quad 0 < \epsilon < 1 \quad \{1^{\text{st}} \text{ Law}\}$$

$$\frac{dA}{dt} = \frac{1}{2} r^2 \dot{\phi} = \text{constant} \quad \{2^{\text{nd}} \text{ Law}\}$$

$$\frac{1}{2} r^2 \dot{\phi} = C_1$$

$$\vec{r} = r \hat{r}(\phi)$$

$$\dot{\vec{r}} = \dot{r} \hat{r} + r \frac{d\hat{r}}{d\phi} \frac{d\phi}{dt}$$

$$\dot{\vec{r}} = \dot{r} \hat{r} + r \dot{\phi} \hat{\phi} \quad \left\{ \frac{d\hat{r}}{d\phi} = \hat{\phi} \right\}$$

$$\begin{aligned} \ddot{\vec{r}} &= \ddot{r} \hat{r} + \dot{r} \frac{d\hat{r}}{d\phi} \frac{d\phi}{dt} + \dot{r} \dot{\phi} \hat{\phi} + r (\ddot{\phi} \hat{\phi} + \dot{\phi} \frac{d\hat{\phi}}{d\phi} \frac{d\phi}{dt}) \\ &= \ddot{r} \hat{r} + \dot{r} \dot{\phi} \hat{\phi} + \dot{r} \dot{\phi} \hat{\phi} + r \ddot{\phi} \hat{\phi} - r \dot{\phi}^2 \hat{r} \quad \left\{ \frac{d\hat{\phi}}{d\phi} = -\hat{r} \right\} \end{aligned}$$

$$\ddot{\vec{r}} = (\ddot{r} - r \dot{\phi}^2) \hat{r} + (2 \dot{r} \dot{\phi} + r \ddot{\phi}) \hat{\phi}$$

2<sup>nd</sup> Law

$$\dot{\phi} = \frac{2C_1}{r^2}$$

$$\ddot{\phi} = -\frac{4C_1}{r^3} \dot{r}$$

$$\begin{aligned} \ddot{\vec{r}} &= (\ddot{r} - r \dot{\phi}^2) \hat{r} + \left( 2 \dot{r} \left( \frac{2C_1}{r^2} \right) + r \left( -\frac{4C_1}{r^3} \dot{r} \right) \right) \hat{\phi} \\ &= (\ddot{r} - r \dot{\phi}^2) \hat{r} + \left( \frac{4C_1}{r^2} \dot{r} - \frac{4C_1}{r^2} \dot{r} \right) \hat{\phi} \end{aligned}$$

$$\ddot{\vec{r}} = (\ddot{r} - r \dot{\phi}^2) \hat{r} \quad \{ \text{establishes that a force directed through the center exists} \}$$

$$\dot{r} = -c(1 + \epsilon \cos \phi)^{-2} (-\epsilon \sin \phi) \dot{\phi}$$

$$= \frac{c}{1 + \epsilon \cos \phi} \frac{\epsilon \sin \phi}{1 + \epsilon \cos \phi} \frac{2C_1}{r^2}$$

$$= \frac{r^2}{c} \epsilon \sin \phi \frac{2C_1}{r^2}$$

$$\dot{r} = \frac{2\epsilon C_1}{c} \sin \phi$$

$$\ddot{r} = \frac{2\epsilon C_1}{c} \cos \phi \dot{\phi}$$

$$= \frac{2\epsilon C_1}{c} \cos \phi \frac{2C_1}{r^2}$$

$$\ddot{r} = 4C_1^2 \frac{\epsilon}{c} \cos \phi \frac{1}{r^2}$$

$$\frac{c}{r} = 1 + \epsilon \cos \phi$$

$$\epsilon \cos \phi = -1 + \frac{c}{r}$$

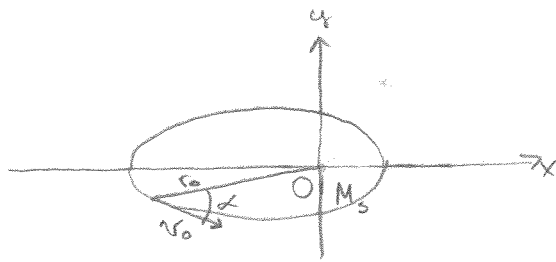
$$\ddot{r} = \frac{-4C_1^2}{cr^2} + \frac{4C_1^2}{r^3}$$

$$\ddot{\vec{r}} = \left( -\frac{4C_1^2}{cr^2} + \frac{4C_1^2}{r^3} - r \left( \frac{2C_1}{r^2} \right)^2 \right) \hat{r}$$

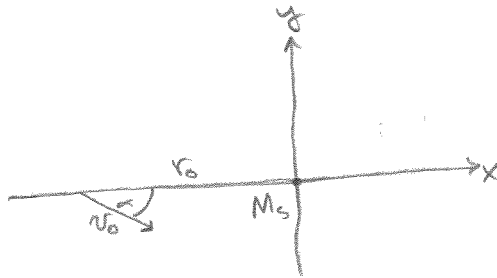
$$= \left( -\frac{4C_1^2}{cr^2} + \frac{4C_1^2}{r^3} - \frac{4C_1^2}{r^3} \right) \hat{r}$$

$$\boxed{\ddot{\vec{r}} = -\frac{4C_1^2}{c} \frac{1}{r^2} \hat{r}}$$

8.27



Orient coordinate system such that the comet is on the x-axis at  $t_0$ .



Angular momentum is conserved

$$l = m r^2 \dot{\phi}$$

$$= m r_0 v_{0 \tan} \text{ at } t_0$$

$$l = m r_0 v_0 \sin \alpha$$

Now,

$$F(r) = -\frac{\gamma}{r^2} \rightarrow \gamma = G m_c M_s$$

$$c = \frac{l^2}{\gamma \mu}$$

$$= \frac{(m_c r_0 v_0 \sin \alpha)^2}{G m_c M_s \left( \frac{M_s m_c}{M_s + m_c} \right)}$$

$$c \approx \frac{r_0^2 v_0^2 \sin^2 \alpha}{G M_s} \left\{ \frac{M_s m_c}{M_s + m_c} \approx m_c, M_s \gg m_c \right\}$$

Energy - eccentricity relationship

$$E = \frac{\gamma^2 \mu}{2 l^2} (\epsilon^2 - 1)$$

Energy of comet

$$E = \frac{1}{2} \mu \dot{r}^2 + U_{\text{eff}}(r)$$

$$= \frac{1}{2} \mu \dot{r}^2 + \frac{l^2}{2 \mu r^2} - \frac{G M_s m_c}{r}$$

$\dot{r}$  = radial velocity

$$\dot{r} = v_0 \cos \alpha$$

Equating energy expressions

$$\frac{1}{2} \mu \dot{r}^2 + \frac{l^2}{2\mu r^2} - \frac{GM_s m_c}{r} = \frac{v^2 \mu}{2l^2} (\epsilon^2 - 1)$$

$$\epsilon^2 = \frac{2l^2}{v^2 \mu} \left[ \frac{1}{2} \mu \dot{r}^2 + \frac{l^2}{2\mu r^2} - \frac{GM_s m_c}{r} \right] + 1$$

$$= \frac{2(m_c r_0 v_0 \sin \alpha)^2}{(GM_s m_c)^2 m_c} \left[ \frac{1}{2} m_c (v_0 \cos \alpha)^2 + \frac{(m_c r_0 v_0 \sin \alpha)^2}{2 m_c r_0^2} - \frac{GM_s m_c}{r_0} \right] + 1$$

$$= \frac{2(r_0 v_0 \sin \alpha)^2}{(GM_s)^2} \left[ \frac{(v_0 \cos \alpha)^2}{2} + \frac{(v_0 \sin \alpha)^2}{2} - \frac{GM_s}{r_0} \right] + 1$$

$$\epsilon^2 = \frac{2(r_0 v_0 \sin \alpha)^2}{(GM_s)^2} \left[ \frac{v_0^2}{2} - \frac{GM_s}{r_0} \right] + 1$$

$$r_0 = \frac{c}{1 + \epsilon \cos(\phi_0)}$$

$$\frac{c}{r_0} = 1 + \epsilon \cos(\phi_0)$$

$$\frac{c}{r_0} - 1 = \epsilon \cos(\phi_0)$$

$$\frac{c - r_0}{\epsilon r_0} = \cos(\phi_0)$$

Given:  $r_0 = 1 \times 10^{11} \text{ m}$

$$v_0 = 45 \times 10^3 \frac{\text{m}}{\text{s}}$$

$$\alpha = 50^\circ$$

$$M_s = 2.0 \times 10^{30} \text{ kg}$$

$$G = 6.7 \times 10^{-11} \frac{\text{N} \cdot \text{m}^2}{\text{kg}^2}$$

$$\epsilon^2 = -0.433 + 1$$

$$= 0.566$$

$$\epsilon = 0.753$$



$$c = 8.868 \times 10^{10} \text{ m}$$

$$\cos \phi_0 = \cos^{-1} \left[ \frac{c - r_0}{\epsilon r_0} \right]$$

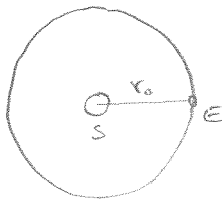
$$\delta = \cos^{-1} [-0.150] \quad \{ \phi_0 = \delta \}$$

$$\delta = 98.646^\circ$$

$$\delta = 1.72 \text{ radians}$$

8.29

Initial conditions:



$$U = -\frac{GM_s M_E}{r_0}$$

Virial Theorem:  $U = Kr^n \rightarrow T = n \frac{U}{2}$

$$U_0 = -\frac{GM_s M_E}{r_0} \rightarrow T_0 = \frac{GM_s M_E}{2r_0} \quad \{n = -1\}$$

For a circular orbit velocity is always tangential to orbit path; no radial velocity. Also,  $\epsilon = 0$ .

At the instant of sun mass change  $\{M_s = \frac{M_{s0}}{2}\}$

$$U = -\frac{GM_s M_E}{2r_0} = \frac{U_0}{2}$$

$$T = T_0 = \frac{GM_s M_E}{2r_0} = \frac{1}{2} M_E v_0^2$$

$$v_0^2 = \frac{GM_s}{r_0}$$

$$E = U_{eff}(r_{min})$$

$$= U_{eff}(r_0)$$

$$= -\frac{\gamma}{r_0} + \frac{l^2}{2\mu r_0^2} \quad \left\{ \mu = \frac{M_s M_E}{M_s + M_E} \approx M_E \right\}$$

$$E = -\frac{GM_s M_E}{2r_0} + \frac{l^2}{2M_E r_0^2} \quad \left\{ \gamma = \frac{GM_s M_E}{2} \right\}$$

$$l = M_E r_0 v_0$$

$$l^2 = M_E^2 r_0^2 v_0^2$$

$$= M_E^2 r_0^2 \frac{GM_s}{r_0}$$

$$l^2 = GM_s M_E^2 r_0$$

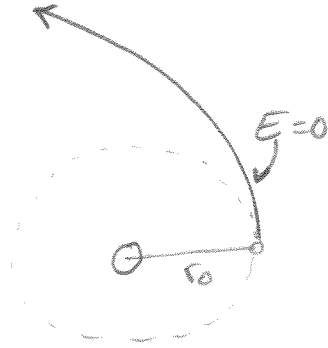
(2)

$$E = -\frac{GM_s M_E}{2r_0} + \frac{GM_s M_E^2 r_0}{2M_E r_0^2}$$

$$= -\frac{GM_s M_E}{2r_0} + \frac{GM_s M_E}{2r_0}$$

$$E = 0$$

When  $E=0$  the path is a parabola.



8.30

$$r(\phi) = \frac{c}{1 + e \cos \phi}$$

i)  $e = 1$ 

$$r = \frac{c}{1 + \cos \phi}$$

$$x = r \cos \phi$$

$$r^2 = x^2 + y^2$$

$$r = \frac{c}{1 + \frac{x}{r}}$$

$$= \frac{c}{\frac{r+x}{r}}$$

$$r = \frac{rc}{r+x}$$

$$\frac{c}{r+x} = 1$$

$$c = r + x$$

$$c = (x^2 + y^2)^{1/2} + x$$

$$x^2 + y^2 = (c - x)^2$$

$$= c^2 - 2cx + x^2$$

$$y^2 = c(c - 2x)$$

$$y = [c(c - 2x)]^{1/2}$$

ii)  $e > 1$ 

$$r = \frac{c}{1 + e \cos \phi}$$

$$r = \frac{c}{1 + \frac{ex}{r}}$$

$$= \frac{c}{\frac{r+ex}{r}}$$

$$r = \frac{rc}{r+ex}$$

$$1 = \frac{C}{r+Ex}$$

$$r+Ex = C$$

$$r = C - Ex$$

$$x^2 + y^2 = (C - Ex)^2$$

$$= C^2 - 2ECx + E^2x^2$$

$$x^2(1-E^2) + 2ECx + y^2 = C^2$$

$$x^2 + \frac{2EC}{1-E^2} + \frac{y^2}{1-E^2} = \frac{C^2}{1-E^2}$$

$$x^2 + \frac{2EC}{1-E^2} + \left(\frac{EC}{1-E^2}\right)^2 + \frac{y^2}{1-E^2} = \frac{C^2}{1-E^2} + \left(\frac{EC}{1-E^2}\right)^2$$

$$\begin{aligned} \left(x + \frac{EC}{1-E^2}\right)^2 + \frac{y^2}{1-E^2} &= \frac{C^2}{1-E^2} + \frac{E^2C^2}{(1-E^2)^2} \\ &= \frac{C^2(1-E^2) + E^2C^2}{(1-E^2)^2} \\ &= \frac{C^2 - C^2E^2 + E^2C^2}{(1-E^2)^2} \\ &= \frac{C^2}{(1-E^2)^2} \end{aligned}$$

$$\frac{\left(x + \frac{EC}{1-E^2}\right)^2}{\frac{C^2}{(1-E^2)^2}} + \frac{\frac{y^2}{1-E^2}}{\frac{(1-E^2)C^2}{(1-E^2)^2}} = 1$$

$$\frac{\left(x - \frac{EC}{E^2-1}\right)^2}{\frac{C^2}{(1-E^2)^2}} + \frac{\frac{y^2}{1-E^2}}{\frac{C^2}{1-E^2}} = 1$$

$$\boxed{\frac{\left(x - \frac{EC}{E^2-1}\right)^2}{\frac{C^2}{(1-E^2)^2}} - \frac{\frac{y^2}{1-E^2}}{\frac{C^2}{1-E^2}} = 1, \quad \gamma = \frac{EC}{E^2-1}, \quad \alpha^2 = \frac{C^2}{(1-E^2)^2}, \quad \beta^2 = \frac{C^2}{E^2-1}}$$

8.31

$$F(r) = \frac{Y}{r^2} = Y u^2$$

$$u''(\phi) = -u(\phi) - \frac{\mu}{\ell^2 u(\phi)^2} F$$

$$= -u(\phi) - \frac{\mu}{\ell^2 u(\phi)^2} Y u^2$$

$$u''(\phi) = -u(\phi) - \frac{\mu Y}{\ell^2}$$

 $u_h$ :

$$u''(\phi) + u(\phi) = 0$$

Let  $u(\phi) = A_1 e^{\gamma \phi}$

$$\gamma^2 + 1 = 0$$

$$\gamma = \pm i$$

$$u_h(\phi) = A_1 e^{\pm i \phi}$$

$$= A_1 e^{i \phi} + A_1 e^{-i \phi}$$

$$= \frac{A}{2} e^{-i(\phi-\delta)} + \frac{A}{2} e^{i(\phi-\delta)} \quad \{ A_1 = \frac{A}{2} e^{-i\delta} \quad A_1^* = \frac{A}{2} e^{i\delta} \}$$

$$u_h(\phi) = A \cos(\phi - \delta)$$

 $u_p$ :

$$u_p(\phi) = C_1$$

$$C_1 = -\frac{\mu Y}{\ell^2}$$

$$u_p(\phi) = -\frac{\mu Y}{\ell^2}$$

$$u(\phi) = u_h(\phi) + u_p(\phi)$$

$$= A \cos(\phi - \delta) - \frac{\mu Y}{\ell^2}$$

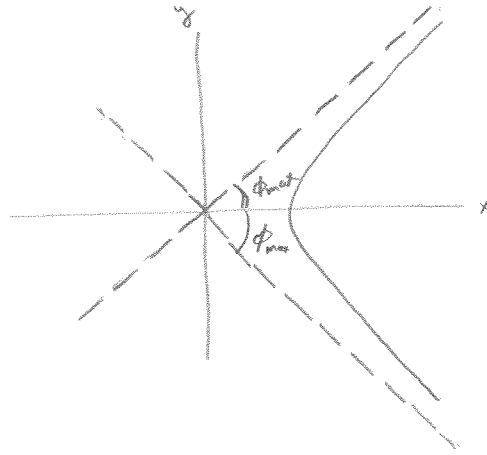
$$u(\phi) = \frac{Y \mu}{\ell^2} (-1 + \epsilon \cos \phi), \quad \delta = 0, \quad \epsilon = \frac{A \ell^2}{Y \mu}$$

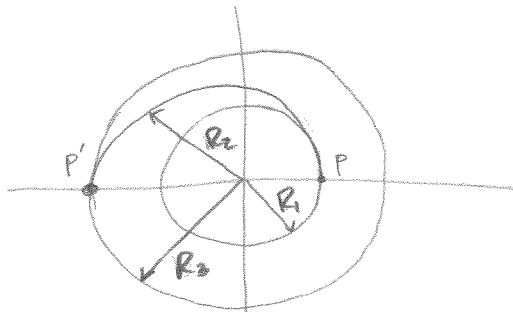
$$\frac{1}{r} = \frac{Y \mu}{\ell^2} (-1 + \epsilon \cos \phi)$$

$$r = \frac{\frac{\ell^2}{Y \mu}}{-1 + \epsilon \cos \phi} = \frac{c}{-1 + \epsilon \cos \phi}$$

(3)

If  $\epsilon < 1$  then  $r < 0$  which is not a possible condition. Therefore,  $\epsilon \geq 1$  is only allowed. This corresponds to  $E \geq 0$ . When  $E = 0$  ( $\epsilon = 1$ ) the only value for  $\phi$  is  $\phi = 0$  giving  $r = \infty$ . For  $E > 0$  ( $\epsilon > 1$ ) then when  $\epsilon \cos \phi_{\max} = 1$   $r \rightarrow \infty$ . Therefore,  $-\phi_{\max} < \phi < \phi_{\max}$  is the allowable range and defines the asymptotes of the hyperbola.





$$R_2 = \frac{C_2}{1 + e_2 \cos \phi}$$

$$R_3 = R_{\text{zapogee}}$$

$$R_{\text{zapogee}} = \frac{C_2}{1 - e_2} \quad \{ \phi = 180^\circ \}$$

$$\frac{C_2}{1 - e_2} = \frac{C_3}{1 + e_3}$$

$$C_3 = \gamma^2 C_2$$

$$e_3 = 0$$

$$\frac{C_2}{1 - e_2} = \gamma^2 C_2$$

$$\frac{1}{1 - e_2} = \gamma^2$$

$$e_2 = \gamma^2 e_1 + (\gamma^2 - 1)$$

$$= \gamma^2 - 1 \quad \{ e_1 = 0 \}$$

$$\gamma^2 = \frac{1}{1 - (\gamma^2 - 1)}$$

$$= \frac{1}{2 - \gamma^2} \quad \gamma^2 = \frac{2R_3}{R_1 + R_3} \quad \{ e.g. 8.72 \}$$

$$= \frac{1}{2 - \frac{2R_3}{R_1 + R_3}}$$

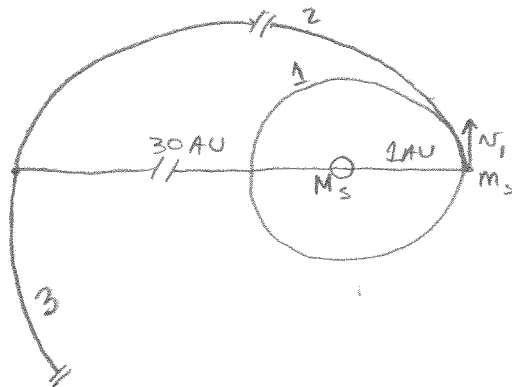
$$= \frac{1}{\frac{2R_1 + 2R_3 - 2R_3}{R_1 + R_3}}$$

$$= \frac{R_1 + R_3}{2R_1}$$

$$\boxed{\gamma' = \sqrt{\frac{R_1 + R_3}{2R_1}}}$$



8.34



Path 2 is an ellipse

$$a = \frac{30+1}{2} \text{ AU}$$

$$a = 15.5 \text{ AU}$$

$$\tau_2^2 = 4\pi^2 \frac{a^3 \mu}{Y}$$

$$\mu = \frac{M_s m_s}{M_s + m_s} = m_s$$

$$Y = G M_s m_s$$

$$\tau_2^2 = \frac{4\pi^2 (15.5)^3 \text{ AU}^3}{G M_s}$$

$$G = 6.7 \times 10^{-11} \frac{\text{N} \cdot \text{m}^2}{\text{kg}^2}$$

$$M_s = 2.0 \times 10^{30} \text{ kg}$$

$$1 \text{ AU} = 1.5 \times 10^{11} \text{ m}$$

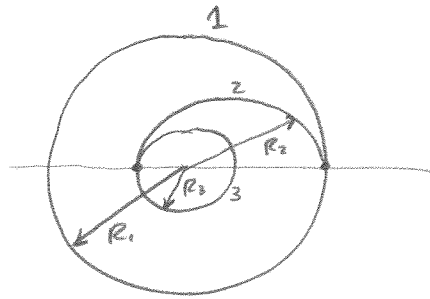
$$\tau_2^2 = \frac{4\pi^2 (23.25 \times 10^{11})^3}{(6.7 \times 10^{-11})(2.0 \times 10^{30})}$$

$$\tau_2 = 1.92 \times 10^9 \text{ s}$$

$$\tau_2 = 61.02 \text{ years}$$

$$\text{Time of transfer} = \frac{\tau_2}{2} = 30.5 = 31 \text{ years}$$

8.35



$$R_3 = \frac{R_1}{4}$$

1<sup>st</sup> stage:

$$\frac{c_1}{1+e_1} = \frac{c_2}{1-e_2} \quad \leftarrow \text{apogee}$$

$$e_1 = 0$$

$$c_2 = \gamma^2 c_1$$

$$c_1 = \frac{\gamma^2 c_1}{1-e_2}$$

$$1-e_2 = \gamma^2$$

$$e_2 = 1-\gamma^2$$

$$R_3 = \frac{c_2}{1+e_2} \leftarrow \text{perigee}$$

$$= \frac{\gamma^2 c_1}{1+(1-\gamma^2)}$$

$$= \frac{\gamma^2 R_1}{2-\gamma^2}$$

$$2R_3 - \gamma^2 R_3 = \gamma^2 R_1$$

$$2R_3 = \gamma^2 (R_1 + R_3)$$

$$\gamma^2 = \frac{2R_3}{R_1 + R_3}$$

$$= \frac{\frac{2R_1}{4}}{R_1 + \frac{R_1}{4}}$$

$$= \frac{\frac{R_1}{2}}{\frac{5R_1}{4}}$$

$$\boxed{\gamma = \sqrt{\frac{2}{5}}}$$

2<sup>nd</sup> stage:

$$\frac{c_2}{1+e_2} = \frac{c_3}{1+e_3}$$

$$e_3 = 0$$

$$c_3 = \gamma'^2 c_2$$

$$\frac{c_2}{1+(1-\gamma'^2)} = \gamma'^2 c_2$$

$$\frac{1}{2-\gamma'^2} = \gamma'^2 \quad \left\{ \gamma'^2 = \frac{1}{2 - \frac{2R_3}{R_1 + R_3}} = \frac{1}{\frac{2R_1 + 2R_3 - 2R_3}{R_1 + R_3}} \right.$$

$$\gamma'^2 = \frac{1}{2 - \frac{2}{5}}$$

$$= \frac{1}{\frac{10-2}{5}}$$

$$\boxed{\gamma' = \sqrt{\frac{5}{8}}}$$

$$v_3 = \gamma' \cdot \frac{v_{2,apo}}{v_{2,per}} \cdot \gamma \cdot v_1$$

$$= \sqrt{\frac{R_1 + R_3}{2R_1}} \cdot \frac{R_1}{R_3} \cdot \sqrt{\frac{2R_3}{R_1 + R_3}} \cdot v_1$$

$$= \sqrt{\frac{R_1}{R_3}} \cdot v_1$$

$$= \sqrt{\frac{R_1}{\frac{R_1}{4}}} \cdot v_1$$

$$= \sqrt{4} \cdot v_1$$

$$\boxed{v_3 = 2v_1}$$