

Manual of RyMoP

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1 Introduction

RyMoP refers to Rydberg Molecule Potential Programme which is written in C++ and developed by Ke Liao for the 5. Physikalisches Institut, Universität Stuttgart. The programme can calculate the potential in a Rydberg Molecule using the partial wave scattering theory, to be precise only S- and P-wave scattering are included. The programme, though at a primitive stage (no graphic interface or post-process scripts), includes all the core functions needed to calculate the S- and P-wave scattering Rydberg Molecule potential. Currently, the programme has only been tested on Linux (Ubuntu) and Mac OSX with gcc (g++) compiler 4.8 and it works flawlessly. Older version of gcc like 4.2 may not support some features and newer versions like gcc 5 abandons some old features so that the programme cannot be compiled. So choose gcc 4.8 to make life easier.

The programme comes with a compressed file, *RyMoPP.tar.gz*, which when uncompressed will give you 4 folders: *Code*, *Wave*, *Scripts* and *EigenLib*. All of the essential codes are stored in *Code*, *EigenLib* is the Eigen Library, *Scripts* contains some pre- and post-processing scripts and *Wave* contains a selected range of wave functions that you can use. If you want to use more wave functions, you have to copy the corresponding files into this folder. One important thing to note is that the names of the newly copied files have to have the same format as the old ones which are already there. Originally I obtain those wave function files by converting the *.mat files to *.txt files. I wrote a script called *convert.py* to do this job. I will address more about this script later in Scripts 1.

For record, I list the essential files here and explain briefly what they do.

Core Files

1. *Kore.cpp*, which is the core file in this programme and contains all the essential functions;
2. *specfunc.h*, which contains a collection of special functions, including the Spherical Harmonics and another Spherical Harmonics without $\sin\theta$ ¹. Though it is not included directly in *Kore.cpp*, it is contained in 'integral.h';
3. *memalloc.h*, which handles memory allocation. It is used in 'integral.h';
4. *integral.h*, which provides integration of 3D angular integration. In this file, it includes 'memalloc.h' and 'specfunc.h'.
5. *line.h*, which provides the functions used in 'Kore.cpp' to read data files correctly.
6. *compile.sh*, which is used to compile the programme. In this file, it compiles the *Kore.cpp* with the option linked to the 'EigenLib' library, which handles the matrices manipulations.

Scripts Above 5 files are the essential ones. During the process, I also wrote two python scripts, which are

1. *convert.py*, which can convert the wave function files from *.mat format to *.txt format, and the latter is the format that this programme can read. Users have to pay attention to the format of the files' name and the files themselves. The converted wave function files have the eigenenergy of this state at the end of the file.
2. *get_pot.py*, which is the post-processing script to get data from the output files from the main programme.

The implementation of the code involves some theoretical derivations, so the next section is dedicated to the basics of the theory. After that, I will dive into the details of the codes. In the end, some examples will be shown.

¹I will introduce this function in detail. It is crucial to the P-wave scattering calculation.

2 Theoretical Basis

In this section, we go through the basics of the theory. The system under consideration is a Rydberg molecule, which consists of a Rydberg atom and a ground state atom sitting in the range of the former. Due to the scattering between the Rydberg electron and the ground state atom, an effective attraction between these two atoms can form. Our goal is to calculate the effective potential between the Rydberg atom and the ground state atom when the distances between them changes. The distance here refers to the distance between the ground state atom and the nucleus of the Rydberg atom. The general idea is to expand the scattered electronic wave function with the unperturbed eigenfunctions from the isolated Rydberg atom, based on the linear variational principle, we can get the perturbed electronic wave function and energy, in turn, the energy of the electron acts as an effective potential for the ground state atom in a wave similar to the Born-Oppenheimer approximation in solids.

2.1 Atomic Wave Functions

Firt of all, we need the unperturbed wave functions of a isolated atom to be the basis functions. Consider the Schrödinger equation for a Rydberg atom:

$$\left[\frac{\mathbf{p}^2}{2m_e} + V_{mod}(r) \right] |\Psi\rangle = E |\Psi\rangle, \quad (2.1)$$

where $V_{mod}(r)$ is a effective potential and we split the wave function into radial and angular part

$$\Psi(r, \theta, \phi) = R(\mathbf{r}) \tilde{Y}(\theta, \phi). \quad (2.2)$$

In the internal wiki of Pi5, there is a detailed explanation about how to get the raidal part of the wave function, so we will skip this part here. But we have to point out that the final calculated result in the wave function file is only the radial part and it is not directly stored as r and $R(r)$, instead a transform of the variables is used:

$$\begin{aligned} x &= \sqrt{r}, \\ X(x) &= x^{3/2} R(r). \end{aligned} \quad (2.3)$$

As for the angular part, since we take the spin orbital interaction into account, spinors, which are the combinations of two spherical harmonics, should be used here.

$$\begin{aligned} \tilde{Y}_{(j-\frac{1}{2}, \frac{1}{2})jm} &= \left(\sqrt{\frac{j+m}{2j}} Y_{j-\frac{1}{2}, m-\frac{1}{2}}, \sqrt{\frac{j-m}{2j}} Y_{j-\frac{1}{2}, m+\frac{1}{2}} \right), \\ \tilde{Y}_{(j+\frac{1}{2}, \frac{1}{2})jm} &= \left(-\sqrt{\frac{j-m+1}{2j+2}} Y_{j+\frac{1}{2}, m-\frac{1}{2}}, \sqrt{\frac{j+m+1}{2j+2}} Y_{j+\frac{1}{2}, m+\frac{1}{2}} \right). \end{aligned} \quad (2.4)$$

2.2 Pseudopotential in Molecular Hamiltonian

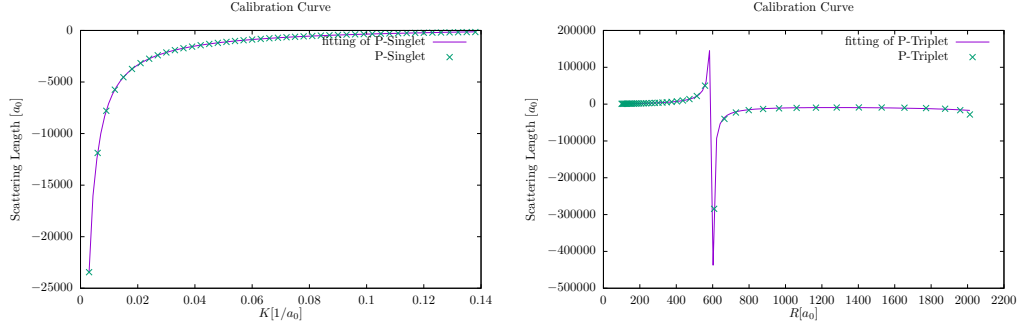
For a highly excited Rydberg atom interacting with a ground state atom, the Hamiltonian is

$$H = \frac{\mathbf{P}^2}{M} + H_{el} + V_{n,e}(\mathbf{r}, \mathbf{R}) \quad (2.5)$$

where H_{el} is the atomic Hamiltonian we mentioned in last subsection, M is the mass of the atom, \mathbf{P} and \mathbf{R} are the relative momentum and position of the ground state atom with respect to the ionic core of the Rydberg atom respectively and \mathbf{r} is the relative position of the Rydberg electron to the ionic core. $V_{n,e}(\mathbf{r}, \mathbf{R})$ is the famous Fermi-Pseudopotential:

$$V_{n,e}(\mathbf{r}, \mathbf{R}) = 2\pi A_s[k(R)] \delta(\mathbf{r} - \mathbf{R}) + 6\pi A_p^3[k(R)] \overleftarrow{\nabla} \delta(\mathbf{r} - \mathbf{R}) \overrightarrow{\nabla} \quad (2.6)$$

where $A_s(k) = -\tan[\delta_0(k)]/k$ and $A_p^3(k) = -\tan[\delta_1(k)]/k^3$ denote the energy-dependent griplet s- and p-wave scattering lengths. It is worth pointing out that $\delta_{l=0,1}(k)$ are the energy dependent phase shifts and $\frac{k(R)^2}{2} = E_{kin} = -\frac{1}{2}n^{*2} + \frac{1}{R}$ where n^* is the effective principal quantum number.



Unfortunately, we do not have analytical expressions for $\delta_{l=0,1}(k)$, so we have to do a fitting based on the data we have got. See figure () for an example of $N = 35$.