Star Formation, Black Hole Types, Hawking Radiation

Jeans Mass

The first criterion for gas cloud collapse:

The most enterion for gas cloud collapse: $M > \frac{5R}{m} \frac{k_B T}{m} \equiv M_{\rm jeans}$ $(\overline{m} \text{ is avg. gas particle mass, } R \text{ is cloud radius})$ $R_{\rm jeans} = \left(\frac{15k_B T}{4\pi G m \rho}\right)^{1/2} = \left(\frac{3M_{\rm jeans}}{4\pi \rho}\right)^{1/3}$ Fragmentation

Fragmentation For some gas cloud above $M_{\rm j}$, as cloud collapses, $\rho\uparrow$, so $M_{\rm j}\downarrow$, so now regions of the cloud have $M>M_{\rm j}$ This cycle repeats for smaller and smaller subregions,

Only if T is constant

Once ρ increases enough, $l_{\rm mfp}$ \downarrow : cloud can't cool, then $T\uparrow$ means $M_j\uparrow;$ fragmentation stops

this is one reason why stars form in groups

Star Birth and Getting on the Main Sequence

1) Dark cloud cores, then gravitational collapse 2) Clump of gas becomes protostar, then T-Tauri

3a) Pre MS Star: If massive enough, collapses to begin fusion

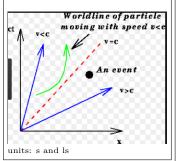
3b) If not massive enough, stays as a brown dwarf the more massive the star, the faster it forms

stars can take 10s of thous to 10s of mil yrs to land on MS

Einstein's Postulates (1915 (GR)) 1) Physics is the same in all ref frames that move at a const vel

2b) An accelerating frame is indist. from a frame at rest in a grav field with g=a

Spacetime Diagrams



(
$$\Delta s$$
)² = $-(c\Delta t)^2 + (\Delta x^2 + \Delta y^2 + \Delta z^2)$

as seen by obs who sees them happen in same place

 $(\Delta s)^2 > 0$: spacelike

Star Clusters

Open Clusters: (stellar nurseries): mostly MS stars,

drift apart as supernovae start popping off Globular Clusters: Dense, bound for life

rare star-star collisions thought to form blue stragglers

Metallicity: What the star formed from

Gas Cooling Mechanisms:
- Dissociation of molecular H, - Ionization of Hydrogen

- Hydrogen recombination (Balmer and lower E γ can escape (not too high E that would just be absorbed by surrounding H)

- Metals: molecular lines, free emission The more metallic, the better at cooling, the more

fragmentation, smaller stars form

Diff cooling mechanisms happen at diff temps: From low T to high, in order:

molecular line emission, low ionization metals,

H & He recomb., highly ionized metals, free-free (deceleration of a charged particle deflected by

other charged particle; loses KE as photon)

Types of Black Holes

Stellar Mass: form by core collapse SN, NS mass transfers

Supermassive: formed by merging, and acreetion, found in centers of galalxies and grow over cosmic time (mils of M_{\odot}) Intermediate: 1000's of M_{\odot} , tentative evidence

Primordial: no clear evidence

Hawking Radiation

in empty space, particle-antiparticle pairs exist for instant before annihilating into energy near event horizon, one can fall in, as the other escapes, but energy debt must be paid, so BH loses mass. Therefore, BHs have temps and BB spectrum:

 $T = 6.2 \times 10^{-8} \text{ K} \times \frac{M_{\odot}}{M}$

 $T = 6.2 \times 10^{-6} \text{ K} \times \frac{1}{M}$ Black holes evaporate if hotter than background CMB T: $(M < 0.008 M_{\text{earth}})$

IMPORTANT: $M(< r) = \int_0^r \rho(r') dV$, $dV = 4\pi r'^2 dr'$ see interstellar gas visibly with:

diffuse cool gas: reflection/scattering dense cool gas: extinction/reddening (objects appear redden

hot gas: emits lines (red: $H\alpha$)

0.0.2Relativity and Accretion Disks

1b) The speed of light in vacuum is same in all ref frames

2) A freefall frame in a grav field is indist. from an iner. frame in no grav. field

Spacetime Interval All observers agree on Δs $(\Delta s)^2 = -(c\Delta t)^2 + (\Delta x^2 + \Delta y^2 + \Delta z^2)$ **Proper Time**: τ : time between 2 events

 $(\Delta s)^2 < 0$: timelike $(\Delta s)_2^2 = 0$: null, or lightlike

Gravitational Redshift (less rigorous) $z = \frac{\Delta \lambda}{\lambda_{\rm em}} = \frac{-v}{c} = \frac{-g\Delta t}{c} = \frac{\Delta \lambda}{\lambda_{\rm em}} = \frac{gh}{c^2}$ h is d traveled by γ

according to outside observer redshifted as it exits grav well

SC metric:
$$\mathrm{d}s^2 = -c^2\,\mathrm{d}t^2\left(1 - \frac{2GM}{rc^2}\right) + \mathrm{d}r^2\left(1 - \frac{2GM}{rc^2}\right)^{-1} + r^2\,\mathrm{d}\theta^2 + r^2\sin^2\theta\,\mathrm{d}\phi^2$$
 In flat spacetime:
$$\mathrm{d}s^2 = -c^2\,\mathrm{d}t^2 + \mathrm{d}r^2 + r^2\,\mathrm{d}\theta^2 + r^2\sin^2\theta\,\mathrm{d}\phi^2$$
 Gravitational Time Dilation:
$$\mathrm{d}\tau = \mathrm{d}t\sqrt{1 - \frac{2GM}{rc^2}}$$
 Where dt is time seen by somebody infinitely far away.

$$\begin{aligned} & \text{Grav Redshift: } \lambda_{\text{obs}} = \lambda_{\text{em}} \left(1 - \frac{2GM}{rc^2}\right)^{-1/2} & \frac{\Delta \lambda}{\lambda_{\text{em}}} = \frac{\sqrt{1 - \frac{2GM}{r_2c^2}}}{\sqrt{1 - \frac{2GM}{r_1c^2}}} - 1 \end{aligned}$$

SC Radius (event hori): $R_s = \frac{2GM}{c^2}$, lowest possible STABLE orbit: $r > \frac{3GM}{c^2}$

Accretion Disks

From MFM, as $M(< r) \simeq M_{BH}$: $v_{\rm rot} \simeq \sqrt{GM/r}$

For ring of disk: $T = r^{-3/4} \times \left(\frac{GM\frac{dM}{dt}}{8\pi\sigma_{SB}}\right)^{1/4}$ Where M is BH Mass

 $\begin{array}{l} \mbox{Lotal disk luminosity: } L_{\rm tot} \simeq \frac{1}{2} \frac{GM}{\frac{\rm dM}{\rm dt}} \\ \mbox{Where } r_{\rm inner} \mbox{ is the innermost stable orbit } (3R_S \mbox{ for BH}) \end{array}$

efficiency: how much of M going in is being radiated away: $\eta = \frac{L_{\rm tot}}{\frac{{\rm d}M}{dt}c^2} = \frac{1}{2}\frac{GM}{r_{\rm inner}c^2}$

Eddington Luminosity Limit: $L_{\mathrm{tot,\ max}} = \frac{4\pi cGMm_{\mathrm{proton}}}{\sigma_m}$ $\sigma_T = 6.65 \times 10^{-29} \text{ m}^2, m_p = 1.673 \times 10^{-27} \text{ kg}$

0.0.3**Gravitational Lensing and Finding Planets**

For top right situation (aligned): $\alpha = \frac{4GM}{c^2b}$ (Where b is dist of closest approach)

angular disk radius: $\theta_E = \left(\frac{4GM}{c^2} \frac{D_{\mathrm{ls}}}{D_{ol} D_{os}}\right)^{1/2}$ Maybe: $\theta_E = 10^{-4}$ as $\left(\frac{M}{\mathrm{M}_\odot}\right)^{1/2} \times \left(\frac{D_{ol}}{10 \mathrm{ kpc}}\right)^{1/2}$ For bottom right situation (nonaligned): $D_{os}\beta + D_{is}\alpha = D_{os}\theta \rightarrow \theta^2 - \beta\theta - \theta_E^2 = 0$

$$\theta_{\pm} = \frac{1}{2} \left[\beta \pm (\beta^2 + 4\theta_E^2)^{1/2} \right]$$

brightness†; more light focused on us

Magnification (rightmost pic):

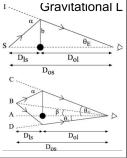
Magnification (rightmost pic): Relative increase in angular size of each image (+ and -):
$$a_{\pm} = \frac{\text{img angular size}}{\text{source angular size}} = \frac{\theta_{\pm} \, \text{d}\theta_{\pm}}{\beta \, \text{d}\beta} = \frac{\theta_{\pm}}{2\beta} \left[1 \pm \frac{\beta}{(\beta^2 + 4\theta_E^2)^{1/2}} \right]$$

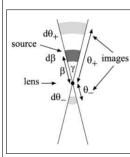
$$a_{\text{tot}} = a_{+} + a_{-} = \frac{u^2 + 2}{u(u^2 + 4)^{1/2}}, \text{ where } u = \beta/\theta_E$$

For small
$$u$$
, $a_{\text{tot}} \simeq 1/u = \theta_E/\beta$ $\frac{\tau_{\text{star}}}{\tau_{\text{planet}}} \simeq \sqrt{\frac{M_{\text{star}}}{M_{\text{planet}}}}$

Duration of lensing event: $\tau = \frac{\theta_E D_{ol}}{v} \propto M_{\rm planet}$ where v is v b/w source and lens (planet) in the plane of the lens

 $\sigma_{\text{LOS}}^2 = \sum_{N} \frac{1}{N} (v_i - \overline{v})^2$





- 1) Microlensing: find rockey planets
- w/long period: one time only
- 2) Radial Velocity of star: best finding massive plans in close orbits (hot jupiters) (didn't form that close)



 $M_{\mathrm{planet}} \sin i \approx \left(\frac{\tau}{2\pi G}\right)^{1/3} \times |v_{1, \mathrm{\ obs}}| M_{\mathrm{star}}$

Where τ is orbital period, i is shown in pic 3) Transit Detection: $\frac{\Delta f}{f} = \left(\frac{r_{\text{planet}}}{r_{\text{star}}}\right)^2$, $P = \frac{r_* + r_p}{d(p,*)}$

4) Direct Imaging: Only possible for big, young planets, far away from star

0.0.4 Galaxies

General Galaxy Stuff
$$\rho \propto 1/r^2$$

From Virial Thm, for ell. gals and gal clusters:
$$M = \frac{^{5\sigma_{\text{LOS}}^2 R}}{^{6\sigma}} \qquad \sigma_{\text{LOS}}^2 = \sum_{i,N} \frac{1}{N} (v_i - \overline{v})^2$$

Active Galactic Nuclei Size of the solar system

gas moving relativistically center has Supermassive Black Hole Emit mostly x-ray and radio (unlike stars)

Eddington Luminosity applies AGN Variability: $R < c\Delta t \simeq 37.5 \text{ AU}$ AGN Jets: extend up to one Mpc emits at all wavelengths, peaks in IR, UV

Types: Seyfert Galaxies: $10^8 - 10^{11} L_{\odot}$ Quasars: $3 \times 10^{10} - 4.3 \times 10^{14} L_{\odot}$ Some Non-AGN galaxies have SMBHs (non-accreting) But not all gals has BHs

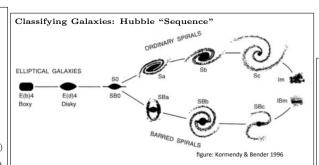
SMBH mass correlates to central bulge mass (but SMBH does not effect gal structure) ↑ due to merging galaxies merging SMBHs too

Isolated Galaxy Formation

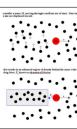
(Galaxies don't really form in isolation) 1) 1st stars form spheroid: Es, S0s tend to have old stellar populations 2) Gas comes in, forms disk as it is the config of lowest energy that maintains \vec{L} . New stars form from disk gas

Galactic Mergers
Steps: 1) random dispersion of stars sphereoids develop (violent relaxation) Dark Matter increases "target size" Dynamical Friction helps galaxies merge More dyn fric if: slower relative motion, greater satellite mass, larger DM halos (mergers start when halos touch) Gas flows inward due to tidal forces, new stars form in huge burst due to replenished

gas, now a lot of stars in center Most star form. was 10-11 Gyr ago Galaxy Clusters: filled w hot gas, gravitational lensing implies mass, all indicates lots of DM. mostly elliptical gal because of many mergers



Left to Right: Gas infall & Outer disk growth, less spheroid, more thin disk, more rotation, more cold gas & young stars, Right to Left:Galaxy interactions (bulge/inner disk growth) Ellipticals: Red and Dead; Irr: Most common, most gas vs. stars Spiral: Most star formation (blue light means young stars) Ellipticals: "early type" (misnomer)



More on Quasars

1) Quasars are SMBH perpetually at Edd. Lum. 2) Quasars are not near MW b/c they were only common during first few Gyr after the BB because they require a lot of gas & merger activity.
3) Now, gas is mostly used up and

mergers are less common
4) Not all are strong radio sources

Star Cosmology Players:

Hubble (extra-gal distances) Leavitt (Cephied Period-L relation) Slipher (Redshifts of "spiral nebulae"

The Evolving Universe

Hubble's Law: $v \simeq H_0 d$ $H_0 \simeq 70 \text{ km/(s Mpc)}$ $z = \frac{\text{dist today}}{\text{dist at emission}} - 1 = \frac{\Delta \lambda}{\lambda} \simeq \frac{v}{c}$ ABOVE ARE ONLY VIABLE AT SMALL z

d = a(t)x Where d is actual distance b/w two things \boldsymbol{x} is their constant comoving distance H_0 is evaluated at t_0 , today; $a(t_0) = 1$

x is their coinstant comoving distance $H = \frac{\dot{a}(t)}{a(t)}$ H_0 is evaluated at t_0 , today; $a(t_0) = 1$ So, as H_0 is always increasing, so is $\dot{a}(t_0)$. (a-DOT) Cosmol. Redshift:

was 1/5 the distance it is today back then most star formation was 10-11 Gyr ago

The Geometry of the Universe

Observations & Cosmo Principle imply the universe is homogenous and isotropic Observations & Cosmo Principle imply the universe is nomogenous and isotropic (no special positions are directions) FRW Metric: $ds^2 = -c^2 dt^2 + a^2(t)$ [geometry of space] $ds^2 = -c^2 dt^2 + a^2(t) [dx^2 + dy^2 + dz^2]$ Where dx, dy, dz are all comoving dist. Spherical Coords: $ds^2 = -c^2 dt^2 + a^2(t) [dr^2 + r^2 \sin^2 \theta d\phi^2 + r^2 d\theta^2]$

Possible Universe Geometries

 ρ_{crit} is the total energy density needed for the universe to be flat

Closed (Spherical): $\rho_{\rm tot} > \rho_{\rm crit}$, parallel lines cross, triangle angles $> 180^{\circ}$, will collapse: big crunch (this could be altered by DE though)

Open (Hyperbolic): $\rho_{\rm tot} < \rho_{\rm crit}$, parallel lines diverge, triangle angle $< 180^{\circ}$, will expand forerver (not altered by DE)

Flat: $\rho_{\rm tot} = \rho_{\rm crit}$, parallel lines stay parallel, triangle angles = 180° , will expand forever, but at slower and slower rate (altered by DE) Evidence points to this geometry

Today: $\rho_{\text{matter}} = 0.3 \rho_{\text{crit}}$

The Friedmann Equation (In a flat universe)

$$\left(\frac{\dot{a}}{a}\right)^2 \equiv H^2 = \frac{8\pi G}{3} \rho_{\rm tot}$$
, $\rho_{\rm tot} \equiv \rho_{\rm crit}$ (either/both

are not fixed with time, evaluated at a chosen H) $\begin{array}{l} \rho_{\rm crit,0} \text{ is simply evaluated at today } (t_0,) \ H_0. \\ \rho_{\rm crit,0} = 9 \times 10^{-27} kg/m^3 \ \Omega_i \equiv \frac{\rho_{\rm i,0}}{\rho_{\rm crit}} \\ \text{for } i \in \{\text{matter, radiation, dark energy}(\Lambda)\} \end{array}$

So Ω_i is ALWAYS evaluated at today

$$\frac{G}{a} = -\frac{4\pi G}{3} \left(\rho_{\text{tot}} + \frac{3P}{c^2} \right)$$
 Where P is pressure

So Ω_i is ALWAYS evaluated at today today, $\Omega_{\rm m} \simeq 0.3$, $\Omega_{\rm de} \simeq 0.7$, $\Omega_{\rm r} \simeq 9 \times 10^{-5}$ Second Friedmann Equation: (double dot a in numer.) $\frac{\ddot{a}}{a} = -\frac{4\pi G}{3} \left(\rho_{\rm tot} + \frac{3P}{c^2} \right)$ Where P is pressure $\rho_{\rm m} \propto \frac{1}{a^3}$, $\rho_{\rm r} \propto T^4$, $T \propto \frac{1}{a} \propto (1+z)$ Density and More General Case: $\rho \propto a^{-3(1+\omega)}$, $\omega_m = 0$, $\omega_r = \frac{1}{3}$, $\omega_{de} = -1$ ω is the eqn. of state parameter: $\omega = \frac{P}{c^2\rho}$

From this, we find:
$$\rho_{\rm de} \simeq {\rm constant},$$
 $\rho_m = \rho_{(m,0)} a^{-3} \rightarrow \rho_m \propto a^{-3}$ $\rho_r = \rho_{(r,0)} a^{-4} \rightarrow \rho_r \propto a^{-4}$ From homework, we found:

 $H(t) = H_0 \sqrt{\Omega_m a(t)^{-3} + \Omega_r a(t)^{-4} + \Omega_{de}}$

Calculating Age of Universe:

$$t_0 = \int_0^{a(t_0)} \frac{\mathrm{d}a'}{\dot{a}'} = \int_0^1 \frac{\mathrm{d}a'}{a'H(a')}$$

For radiation only universe: $H = H_0 a^{-2}$

For radiation only dinverse:
$$H = H_0 a$$
 Calculating Size of (Observable) Universe: (derived from FRW metric for some γ traveling to now) radius= c $\int \frac{\mathrm{d}t'}{a} = c$ $\int \frac{\mathrm{d}a'}{a'^2 H(a')}$

 t_{em} a_{em} Where ' just indicates dummy variables

The Origins of Matter:

$$\gamma$$
's \leftrightarrow e^+ , e^- if $E_{\gamma}=h\nu>E_{e^+/e^-}=mc^2$ e^+ande^- continuously annihilate Equilibrium stops when radiation too cold but annihilation contines, and fast **Freeze-out:** annihilations stop b/c the partics. are dilute enough; b/c so few partics. left, and b/c the expans. of the universe Matter-Antimatter Asymmetry

Most likely because of some unexplained matter-antimatter asymmetry (like 1,000,000

Dark Matter Origins:

The freeze-out process predict the observed DM density if DM particles interact via Weak F. IF DM was produced by freeze-out, (only possible if it has antiparticles,) then there should be equal amts of DM and DantiM If both DM and M freeze out, they won't necessarily have same density today, as freezeout is dependent on MFP of particles annihilating

Big Bang Nucleosynthesis

NOT full atoms, just nuclei $T = 10^9 \text{ K} 2p^+ + 2n^0 \to 2D^2 \to \text{He}^4$ $t_0 = 10 \text{ s, lasted about } 200 \text{ s}$ Universe was hot and dense enough to trigger this fusion, but cool enough: $\overline{E}_{\gamma} = k_b T = 0.08$ MeV, $E_{
m deuterium} = -2.22 eV$ Universe had to be cool enough for so not even one in a billion photons (due to freeze-out) would have enough energy to destroy deuterium D That fusion only possible with free neutrons, which decay quickly, so doesn't last long; this is why no elements heavier than Be

formed; window was too short before too cool observations of primordial He and D show that universe is only 4% atoms

The Cosmic Microwave Background

became atoms

Surface of Last Scatter: 380,000 yrs after BB, when T = 3000 K, hydrogen atoms formed, after which photons could travel freely:

$$l_{mfp} = \left(\frac{n_e}{n_{e \text{ free}}}\right) \times \frac{7 \times 10^{28} \text{ m}}{(1+z)^3}$$

 $l_{mfp} = \left(\frac{n_e}{n_e \text{ free}}\right) \times \frac{7 \times 10^{28} \text{ m}}{(1+z)^3}$ Because more of the free electrons joined nuclei (low enough energy free γ to not ionize atoms) so now universe was transparent from then on

 $\begin{array}{l} l_{universe} \approx \frac{c}{H} \\ \text{Our motion actually redshifts \& blueshifts the CMB} \\ \text{CMB is the perfect BB radiator, T=2.725} \\ \text{confirms hot big bang. there are very tine mK flucs.} \end{array}$ for in this

Dark Matter Candidates

WIMPS are leading candidate; leading candidate for WIMP: SUSY fermion partners of Weak force carriers (partners of W, Z, Higgs bosons). Other proposals MACHOS: They cant be DM (at least more than 10% of it). Looking to the DM halo, MACHOS would produce much more gravitational microlensing than is actually seen, if it was all MACHOS. Neutrinos: ruled out because they're too hot: the large scale structure that

neutrionis would make is not what is observed. Ultra-cold H2 globules: DM can't be made out of regular matter

MOND: Can explain some things,

but not everything. Still continues to not die. 1° disks in CMB, casually separated by photon path since BB; yet, still all have almost exactly same temp

0.0.6Misc Info, Derivations

Misc Info

HII Regions: when O, B stars form, they emit radiation that can ionize surrounding gas, making HII region $\rho/n = \overline{m} \qquad g_{surf} \simeq \frac{GM}{R^2}$

Press. broad comparing 2 stars: $\frac{\Delta\lambda_1}{\Delta\lambda_2}=\frac{n_1}{n_2}=\frac{P_1/kT_1}{P_2/kT_2}$ $N=M/\overline{m},$ Metallicity is frac. that's not H or He

$$\begin{array}{l} \text{Roche Lobe, Tidal:} \\ \frac{M_1}{(a-R_L)^2} - \frac{M_2}{R_L^2} = \frac{M_1}{a^2} - \frac{R_L(M_1 + M_2)}{a^3} \\ R_L \simeq a \left(\frac{M_2}{3M_1}\right)^{1/3} & F_{\text{tide}} \simeq \frac{2GMmh}{r^3} \end{array}$$

From V thm: collapse if U > 2K: $\frac{3GM^2}{5R} > 2 \times \frac{3Mk_bT}{2\overline{m}}$

or falling dm:
$$\Delta U_{\rm grav} = \frac{-GM \, \mathrm{d}m}{r^2} \, \mathrm{d}r$$

 $\frac{\Delta E_{\rm thermal}}{\Delta E_{\rm thermal}} = \frac{1}{2} \frac{GM \, {\rm d}m \, {\rm d}r}{r^2 \, {\rm d}t}$ $\frac{\Delta E_{\rm thermal}}{{\rm d}t} = L_{\rm ring} = \sigma_{SB} T^4 (2\pi r \, {\rm d}r) \times 2$ Solve these for each other to get answer.

Total Disk Luminosity:
$$L_{\rm tot} = \int_{r_{\rm inner}}^{r_{\rm outer}} 2(2\pi r\,{\rm d}r)\sigma_{SB}T^4$$
 Eddington Luminosity:

Consider some e^- trying to fall into object.

There's pt where rad. press overcomes grav. attrac on e^-

There's pt where rad, press overcomes grav, average on $n_{\gamma} = \frac{L_{\nu}}{4\pi r^2 ch\nu}$ Scattering rate of $e^- = n_{\gamma} \sigma_T c$ Each scatter: $\Delta p_e = \Delta p_{\gamma} = \frac{h\nu}{c}$ $F_{\gamma} = \frac{\mathrm{d}p}{\mathrm{d}t}$, $F_{\gamma_{tot}} = \frac{L_{tot} \sigma_T}{2\pi r^2 c}$; $F_{\gamma} > \frac{GMm_p}{r^2}$ $(m_p$ because electrons bound to protons)

Cosmo Redshift Deriv:
For photon going towards Earth (FRW Metric):
$$c^2 dt^2 = a^2 dr^2 \rightarrow dr = \pm c \frac{dt}{a}$$

$$\int_{rgal}^{r=0} dr = \pm c \int_{tem}^{tobs} \frac{dt}{a} = c \int_{tem+\delta tem}^{tobs+\delta t} Double taylor expand on t_{em} and t_{obs}

$$r_{gal} = c \int_{tem}^{tobs} \frac{dt}{a} + \left(\frac{-c}{a}\right)_{tem}^{tem+\delta tem}$$

$$0 = \frac{-c}{a_{em}} \delta t_{em} + \frac{c}{a_{obs}} \delta t_{obs}$$

$$densities:$$

$$\frac{d\rho}{dt} = \frac{3\pi}{a_{em}} G \left(2 \frac{\tilde{a}}{a^2} - 2 \frac{\tilde{a}^2}{a^3} \right) = -3H(\rho + \frac{P}{c^2}) = -3H(1 + \omega)\rho$$$$