A Statistical Dynamical Model to Predict Extreme Events and Anomalous Features in Shallow Water Waves with Abrupt Depth Change

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This manuscript was compiled on November 25, 2018

Understanding and predicting extreme events and their anomalous statistics in complex nonlinear systems is a grand challenge in climate, material, and neuroscience, as well as for engineering design. Recent controlled laboratory experiments in weakly turbulent shallow water with abrupt depth change (ADC) exhibit a remarkable transition from nearly Gaussian statistics in incoming wave trains before the ADC to outgoing waves trains after the ADC with extreme anomalous statistics with large positive skewness of the surface height. Here we develop a statistical dynamical model to explain and quantitatively predict the above anomalous statistical behavior as experimental control parameters are varied. The first step is to use incoming and outgoing truncated Korteweg-de Vries (TKdV) equations matched in time at the ADC. The TKdV equation is a Hamiltonian system which induces incoming and outgoing statistical Gibbs invariant measures which are statistically matched at the ADC. The statistical matching of the known nearly Gaussian incoming Gibbs state at the ADC completely determines the predicted anomalous outgoing Gibbs state, which can be calculated by a simple sampling algorithm, verified by direct numerical simulations, and successfully captures key features of the experiment. There is even an analytic formula for the anomalous outgoing skewness. The strategy here should be useful for predicting extreme anomalous statistical behavior in other dispersive media in different settings.

extreme anomalous event | statistical TKdV model | matching Gibbs measures | surface wave displacement and slope

Understanding and predicting extreme events and their anomalous statistics in complex nonlinear systems is a grand challenge in climate, material and neuroscience as well as for engineering design. This is a very active contemporary topic in applied mathematics with qualitative and quantitative models (1–7) and novel numerical algorithms which overcome the curse of dimension for extreme event prediction in large complex systems (2, 8–11). The occurrence of Rogue waves as extreme events within different physical settings of deep water (12–16) and shallow water (17–19) is an important practical topic.

Recent laboratory experiments in weakly turbulent shallow water reveal a remarkable transition from Gaussian to anomalous behavior as surface waves cross an abrupt depth change (ADC). A normally-distributed incoming wave train, upstream of the ADC, transitions to a highly non-Gaussian outgoing wave train, downstream of the ADC, that exhibits large positive skewness of the surface height and more frequent extreme events (20). Here we develop a statistical dynamical model to explain and quantitatively predict this anomalous behavior as experimental control parameters are varied. The first step is to use incoming and outgoing truncated Korteweg-de Vries (TKdV) equations matched in time at the ADC. The TKdV equation is a Hamiltonian system which induces incoming and outgoing statistical Gibbs invariant measures which are statistically matched at the ADC. The statistical matching of the known nearly Gaussian incoming Gibbs state at the ADC completely determines the predicted anomalous outgoing Gibbs state, which can be calculated by a simple MCMC algorithm, verified by direct numerical simulations, and successfully captures key features of the experiment. There is even an analytic formula for the anomalous outgoing skewness. The strategy here should be useful for predicting extreme anomalous statistical behavior in other dispersive media in different settings.

If you are interested, we could cite papers on 'Optical rogue waves' and/or rogue waves in microwave systems here.

1. Experiments showing anomalous wave statistics induced by an abrupt depth change

Controlled laboratory experiments were carried out in (20) to examine the statistical behavior of surface waves crossing an ADC. In these experiments, nearly unidirectional waves are generated by a paddle wheel and propagate through a long, narrow wave tank. Midway through, the waves encounter a

Significance Statement

Understanding and predicting extreme events and their anomalous statistics in complex nonlinear systems is a grand challenge in applied sciences as well as for engineering design. Recent controlled laboratory experiments in weakly turbulent shallow water with abrupt depth change exhibit a remarkable transition from nearly Gaussian statistics to extreme anomalous statistics with large positive skewness of the surface height. We develop a statistical dynamical model to explain and quantitatively predict the anomalous statistical behavior. The incoming and outgoing wares are modeled by the truncated Korteweg-de Vries equations statistically matched at the depth change. The statistical matching of the known nearly Gaussian incoming Gibbs state completely determines the predicted anomalous outgoing Gibbs state, and successfully captures key features of the experiment.

A.J.M. designed the research. A.J.M., M.N.J.M., and D.Q. performed the research. A.J.M., M.N.J.M., and D.Q. wrote the paper.

The authors declare no conflict of interest.

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step in the bottom topography, and thus abruptly transition from one depth to another. The paddle wheel is forced with angle

$$\theta\left(t\right)=\theta_{0}+\Delta\theta\sum_{n=1}^{N}a_{n}\cos\left(\omega_{n}t+\delta_{n}\right),\ E\sim\left(\Delta\theta\right)^{2}\sum_{n}a_{n}^{2}\omega_{n}^{2}.$$

where the weights a_n are Gaussian in spectral space, and the total energy E in the system is determined by the angle amplitude $\Delta\theta$. Optical measurements of the free surface reveal a number of surprising statistical features:

- Distinct statistics are found between the incoming and outgoing wave disturbances: the incoming waves display near-Gaussian statistics, while the outgoing waves show skewness toward the positive displacement.
- The non-Gaussian statistics is related with the total energy contained in the system: larger driving amplitudes $\Delta\theta$ generate stronger skewness in the outgoing PDFs.
- The waves also show different characteristic peak wavelengths in incoming and outgoing flows.

  I am not sure we need to include this last item. It is not really an experimental observation, simply a calculation that can be done using the dispersion relation and the two depths. Could we simply take this last item out?

2. Surface wave turbulence modeled by truncated KdV equation with depth dependence

The surface wave turbulence is modeled by a one-dimensional deterministic dynamical model. The Korteweg-de Vries (KdV) equation (21) is a leading-order approximation of the surface waves that are determined by the balance of nonlinear and dispersive effects in an appropriate far-field limit. Here, the KdV equation is truncated in the first Λ modes (with $J=2\Lambda+1$ grid points) to generate weakly turbulent dynamics. Therefore, the surface disturbance is modeled by the state variable $u_{\Lambda}^{\pm}\left(x,t\right)$ with superscript '$-$' for the incoming waves and '$+$' for the outgoing waves. The Galerkin truncated variable $u_{\Lambda}=\sum_{1\leq|k|\leq\Lambda}\hat{u}_{k}\left(t\right)e^{ikx}$ is normalized with zero mean $\hat{u}_{0}=0$ and unit energy $2\pi\sum_{k=1}^{\Lambda}\left|\hat{u}_{k}\right|^{2}=1$, which are conserved quantities, and $u_{\Lambda}\equiv\mathcal{P}_{\Lambda}u$ denotes the subspace projection. The motion is governed by the truncated KdV equation with depth change $D_{\pm}$ about the state variable $u_{\Lambda}^{\pm}$

$$\frac{\partial u_{\Lambda}^{\pm}}{\partial t}+\frac{D_{\pm}^{-3/2}}{2}E_{0}^{1/2}L_{0}^{-3/2}\frac{\partial}{\partial x}\mathcal{P}_{\Lambda}\left(u_{\Lambda}^{\pm}\right)^{2}+D_{\pm}^{1/2}L_{0}^{-3}\frac{\partial^{3}u_{\Lambda}^{\pm}}{\partial x^{3}}=0, \quad [1]$$

on the normalized periodic domain $x\in[-\pi,\pi]$ with the conserved Hamiltonian decomposed into the difference of two components containing the cubic and quadratic terms

$$\mathcal{H}_{\Lambda}^{\pm}=D_{\pm}^{-3/2}E_{0}^{1/2}L_{0}^{-3/2}H_{3}\left(u_{\Lambda}^{\pm}\right)-D_{\pm}^{1/2}L_{0}^{-3}H_{2}\left(u_{\Lambda}^{\pm}\right),$$
$$H_{3}\left(u\right)=\frac{1}{6}\int_{-\pi}^{\pi}u^{3}dx,\ H_{2}\left(u\right)=\frac{1}{2}\int_{-\pi}^{\pi}\left(\frac{\partial u}{\partial x}\right)^{2}dx.$$

The model [1] is non-dimensionalized in the periodic domain. The depth is assumed to be unit $D_{-}=1$ before the abrupt depth change and becomes $D_{+}<1$ for the flows after the change. We introduce the model parameters $\left(E_{0},L_{0},\Lambda\right)$ based on the following model assumptions:

- The wavenumber truncation Λ is fixed in a moderate value for generating weakly turbulent dynamics;
- The state variable $u_{\Lambda}^{\pm}$ is normalized with zero mean and unit energy, $\mathcal{M}\left(u_{\Lambda}\right)=0,\mathcal{E}\left(u_{\Lambda}\right)=1$, conserved during the evolution, while E_0 characterizes the total energy injected in the system based on the paddle amplitude $\left(\Delta\theta\right)^{2}$;
- The length scale of the system is defined by L_0. The value is chosen so that the resolved scale $\Delta x=2\pi L_{0}/J$ is comparable with the the characteristic wave length λ_c found from the experiments.

The intuition for the distinct model dynamics comes from the balance between the cubic and quadratic terms in the Hamiltonian $\mathcal{H}_{\Lambda}^{\pm}$. After the depth change, $D_{+}<1$, more weight is added in the cubic term, H_3, for stronger nonlinearity and weaker dispersion for the third-order derivative term reflected by the smaller coefficient for H_2 in the Hamiltonian. Since $\frac{\partial u}{\partial x}$ is the slope of the wave height, $H_{2}\left(u\right)$ measures the wave slope energy.

A *deterministic matching condition* is given for the surface displacement $u_{\Lambda}^{\pm}$ agreeing at the locations before and after the abrupt depth change T_{ADC}

$$u_{\Lambda}^{-}\left(x,t\right)|_{t=T_{\mathrm{ADC}}-}=u_{\Lambda}^{+}\left(x,t\right)|_{t=T_{\mathrm{ADC}}+},$$

assuming the abrupt depth change is met at $t=T_{\mathrm{ADC}}$. Equation [1] is not designed to capture the short scale changes in rapid time. On the other hand, we are interested in the model statistical transition before and after the depth change, so it is reasonable to observe the suitable slow-time performance in the large scale structures.

Interpreting experimental parameters in the dynamical model. The model parameters $\left(E_{0},L_{0},\Lambda\right)$ in [1] can be directly linked with the basic scales from the physical problem. The important characterizing parameters measured from the experiments include: $\epsilon=\frac{a}{H_{0}}$ the wave amplitude a to water depth H_0 ratio; $\delta=\frac{H_{0}}{\lambda_{c}}$ the water depth to wavelength scale λ_c ratio; and $D_{0}=\frac{d}{H_{0}}$ the normalized wave depth ratio with incoming flow depth $d=H_{0}$ to the outgoing flow depth $d<H_{0}$. The interpretations and reference values of these model parameters are based on the experimental setup (20). By comparing the characteristic physical scales, the normalized TKdV equation [1] can be linked directly with the measured non-dimensional quantities by

$$L_{0}=6^{\frac{1}{3}}\left(M\epsilon^{\frac{1}{2}}\delta^{-1}\right),\ E_{0}=\frac{27}{2}\gamma^{-2}\left(M\epsilon^{\frac{1}{2}}\delta^{-1}\right), \quad [2]$$

where M defines the computational domain size $L_{d}=M\lambda_{c}$ as M-multiple of the characteristic wavelength λ_c, and $\gamma=\frac{U}{a}$ represents the factor to normalize the total energy in the state variable u_Λ to one.

Consider the spatial discretization $J=2\Lambda+1$ so that the smallest resolved scale is comparable with the characteristic wavelength

$$\Delta x=\frac{2\pi M\lambda_{c}}{J}\lesssim\lambda_{c}\ \Rightarrow\ M=\frac{J}{2\pi}\sim5,\ J=32.$$

Therefore in the practical numerical simulations, we pick $M=5$ and γ varies in the range $[0.5,1]$. Using the reference experimental measurements (20), $\epsilon \in [0.0024, 0.024]\,, \delta \sim 0.22$, and D_0 changes from 1 to 0.24 before and after the depth change. The reference values for the model scales can be estimated in the range $L_0 \in [2, 6]$ and $E_0 \in [50, 200]$. These are the values we will test in the direct numerical simulations. See details about the derivation from scale analysis in *SI Appendix, A*.

3. Equilibrium statistical mechanism for generating the stationary invariant measure

Since the TKdV equation satisfies the Liouville property, the equilibrium invariant measure can be described by an equilibrium statistical formulism (22–24) using a Gibbs measure with the conserved energy $\mathcal{E}_\Lambda$ and Hamiltonian $\mathcal{H}_\Lambda$. The equilibrium invariant measure is dictated by the conservation laws in the TKdV equation. In the case with fixed total energy E_0, this is the *mixed Gibbs measure* in the truncated model with microcanonical energy and canonical Hamiltonian ensembles (22)

$$\mathcal{G}_\theta^\pm\left(u_\Lambda^\pm; E_0\right) = C_\theta^\pm \exp\left(-\theta^\pm \mathcal{H}\left(u_\Lambda^\pm\right)\right)\delta\left(\mathcal{E}\left(u_\Lambda^\pm\right) - E_0\right), \quad [3]$$

with θ representing the "inverse temperature". The distinct statistics in the upstream and downstream waves can be controlled by the parameter value of the inverse temperature. The negative temperature regime, $\theta^\pm < 0$, is the appropriate regime to predict the experiments as shown below. In the incoming flow field, the inverse temperature θ^- is chosen so that $\mathcal{G}_\theta^-$ has Gaussian statistics. Using the above invariant measures [3], the expectation of any functional $F(u)$ can be computed based on the Gibbs measure

$$\langle F\rangle_{\mathcal{G}_\theta} \equiv \int F(u)\,\mathcal{G}_\theta(u)\,du.$$

The value of θ in the invariant measure is specified from $\langle H_\Lambda\rangle_{\mathcal{G}_\theta}$ (22, 24). The invariant measure also predicts an equilibrium energy spectrum without running the TKdV equation directly. On the other hand, the time autocorrelation and transient statistics about the state variable u_Λ cannot be recovered from the statistical theory.

Statistical matching condition in invariant measures before and after the abrupt depth change. The Gibbs measures $\mathcal{G}_\theta^\pm$ are defined based on the different inverse temperatures $\theta^\pm$ on the two sides of the solutions

$$\mu_t^-\left(u_\Lambda^-; D_-\right), \quad u_\Lambda\,|_{t=T_{\mathrm{ADC}-}} = u_0,\ t < T_{\mathrm{ADC}};$$
$$\mu_t^+\left(u_\Lambda^+; D_+\right), \quad u_\Lambda\,|_{t=T_{\mathrm{ADC}+}} = u_0,\ t > T_{\mathrm{ADC}},$$

where u_0 represents the deterministic matching condition between the incoming and outgoing waves. The two distributions, μ_t^-, μ_t^+ should also be matched at the depth change location T_{ADC}, so that

$$\mu_\infty^-(u_\Lambda) = \mu_{t=T_{\mathrm{ADC}}}^-(u_\Lambda) = \mu_{t=T_{\mathrm{ADC}}}^+(u_\Lambda).$$

In matching the flow statistics before and after the abrupt depth change, first we use the conservation of the deterministic Hamiltonian H_Λ^+ after the depth change. Then assuming ergodicity (22, 23), the statistical expectation for the Hamiltonian $\left\langle H_\Lambda^+\right\rangle$ is conserved in time after the depth change at $t = T_{\mathrm{ADC}}$ and should stay in the same value as the system approaches equilibrium as $t \to \infty$. The final statistical matching condition to get the outgoing flow statistics with parameter θ^+ can be found by

$$\left\langle H_\Lambda^+\right\rangle_{\mathcal{G}_\theta^+} = \left\langle H_\Lambda^+\right\rangle_{\mathcal{G}_\theta^-}, \quad [4]$$

with the outgoing flow Hamiltonian H_Λ^+ and the Gibbs measures $\mathcal{G}_\theta^\pm$ before and after the abrupt depth change.

4. The nearly Gaussian incoming statistical state

The incoming flow is always characterized by a near-Gaussian distribution in the wave displacement. It is found that a physically consistent Gibbs measure should take negative values in the inverse temperature parameter $\theta < 0$, where a proper distribution function and a decaying energy spectrum are generated (see (24) and *SI Appendix, B.1* for the explicit simulation results). The upstream Gibbs measure $\mathcal{G}_\theta^-$ with $D_- = 1$ displays a wide parameter regime in $\left(\theta^-, E_0\right)$ with near-Gaussian statistics. In the left panel of Figure 1 (a), the inflow skewness κ_3^- varies only within small amplitudes among changing values of E_0 and θ^-. The incoming flow PDF then can be determined by picking the proper parameter value θ^- in the near Gaussian regime with small skewness. In contrast, the downstream Gibbs measure $\mathcal{G}_\theta^+$ with $D_+ = 0.24$ shown in the right panel of Figure 1 (a) generates much larger skewness κ_3^+ with highly skewed PDFs as the absolute value of θ^+ and the total energy level E_0 increase in amplitude. The solid lines in Figure 1 (c) offer a further confirmation of the transition from near-Gaussian statistics with tiny κ_3^- to strongly skewed distribution κ_3^+ after the depth change.

In the next step, the value of the downstream θ^+ is determined based on the matching condition [4]. The expectation $\left\langle H_\Lambda^+\right\rangle_{\mathcal{G}_\theta^-}$ about the incoming flow Gibbs measure can be calculated according to the predetermined parameter values of θ^- as well as E_0 from the previous step. For the direct numerical experiments shown later in Figure 2, we pick proper choices of test parameter values as $L_0 = 6, E_0 = 100$ and $\theta^- = -0.1, -0.3, -0.5$. More test cases with different system energy E_0 can be found in *SI Appendix, B.2* where similar transition from near Gaussian symmetric PDFs to skewed PDFs in the flow state $u_\Lambda^\pm$ can always be observed.

Direct numerical model simulations. Besides the prediction of equilibrium statistical measures from the equilibrium statistical approach, another way to predict the downstream model statistics is through running the dynamical model [1] directly. The TKdV equation is found to be ergodic with proper mixing property as measured by the decay of autocorrelations as long as the system starts from a negative inverse temperature state as described before. For direct numerical simulations of the TKdV equations, a proper symplectic integrator is required to guarantee the Hamiltonian and energy are conserved in time. It is crucial to use the symplectic scheme to guarantee the exact conservation of the energy and Hamiltonian since they are playing the central role in generating the invariant measure and the statistical matching. The symplectic schemes used here for the time integration of the equation is the 4th-order midpoint method (25). Details about the mixing properties from different initial states and the numerical algorithm are described in *SI Appendix, C*.

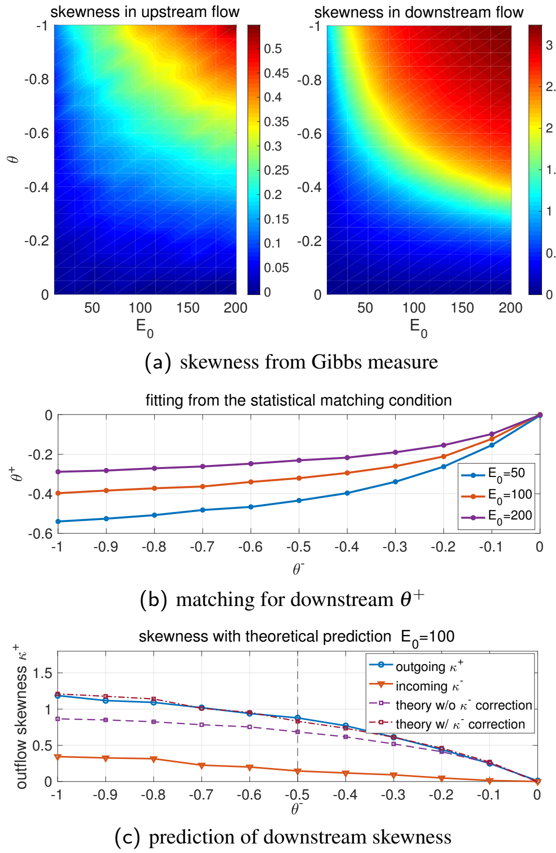


Fig. 1. Fist row: skewness from the Gibbs measures in incoming and outgoing flow states with different values of total energy E_0 and inverse temperature θ (notice the different scales in the incoming and outgoing flows); Second row: outgoing flow parameter θ^+ as a function of the incoming flow θ^- computed from the statistical matching condition with three energy level E_0; Last row: skewness in the outgoing flow with the matched value of θ^+ as a function of the inflow parameter θ^- (the theoretical predictions using [5] are compared).

5. Predicting extreme anomalous behavior after the ADC by statistical matching

With the inflow statistics well described and the numerical scheme set up, we are able to predict the downstream anomalous statistics starting from the near-Gaussian incoming flow going through the abrupt depth change from $D_- = 1$ to $D_+ = 0.24$. First, we consider the statistical prediction in the downstream equilibrium measure directly from the matching condition. The downstream parameter value θ^+ is determined by solving the nonlinear equation [4] as a function of θ^+, $F\left(\theta^+\right) = \left\langle H_\Lambda^+\right\rangle_{\mathcal{G}_\theta^+}\left(\theta^+\right) - \left\langle H_\Lambda^+\right\rangle_{\mathcal{G}_\theta^-} = 0$. In the numerical approach, we adopt a modified secant method avoiding the stiffness in the parameter regime (see the *SI Appendix, B.2* for the algorithm). The fitted solution is plotted in Figure 1 (b) as a function of the proposed inflow θ^-. A nonlinear θ^-–θ^+ relation is discovered from the matching condition. The downstream inverse temperature θ^+ will finally saturate at some level. The corresponding downstream skewness of the wave displacement u_Λ predicted from the statistical matching of Gibbs measures is plotted in Figure 1 (c). In general, a large positive skewness for outgoing flow κ_3^+ is predicted from the theory, while the incoming flow skewness κ_3^- is kept in a small value in a wide range of θ^-. Note that with $\theta^- \sim 0$ (that is, using the microcanonical ensemble only with energy conservation), the outflow statistics are also near Gaussian with weak skewness. The skewness in the outflow statistics grows as the inflow parameter value θ^- increases in amplitude.

For a second approach, we can use direct numerical simulations starting from the initial state sampled from the incoming flow Gibbs measure $\mathcal{G}_\theta^-$ and check the transient changes in the model statistics. Figure 2 illustrates the change of statistics as the flow goes through the abrupt depth change. The first row plots the changes in the skewness and kurtosis for the state variable u_Λ after the depth change at $t=0$. The PDFs in the incoming and outgoing flow states are compared with three different initial inverse temperatures θ^-. After the depth changes to $D_0 = 0.24$ abruptly at $t=0$, both the skewness and kurtosis jump to a much larger value in a short time, implying the rapid transition to a highly skewed non-Gaussian statistical regime after the depth change. Further from Figure 2, different initial skewness (but all relatively small) is set due to the various values of θ^-. With small $\theta^- = -0.1$, the change in the skewness is not very obvious (see the second row of Figure 2 for the incoming and outgoing PDFs of u_Λ). In comparison, if the incoming flow starts from the initial parameter $\theta^- = -0.3$ and $\theta^- = -0.5$, much larger increase in the skewness is induced from the abrupt depth change. Furthermore, in the detailed plots in the third row of Figure 2 for the downstream PDFs under logarithmic scale, fat tails towards the positive direction can be observed, which represent the extreme events in the downstream flow (see also Figure 3 for the time-series of u_Λ).

As a result, the downstream statistics in final equilibrium predicted from the direct numerical simulations here agree with the equilibrium statistical mechanism prediction illustrated in Figure 1. The prediction from these two different approaches confirm each other.

6. Analytic formula for the upstream skewness after the ADC

A statistical link between the upstream and downstream energy spectra can be found for an analytical prediction of the skewness in the flow state u after the ADC. The skewness of the state variable u_j at one spatial grid point is defined as the ratio between the third and second moments

$$\kappa_3 = \left\langle u_j^3\right\rangle_\mu / \left\langle u_j^2\right\rangle_\mu^{\frac{3}{2}}.$$

Now we introduce mild assumptions on the distribution functions:

- The upstream equilibrium measure μ_- has a relatively small skewness so that

$$\left\langle H_3\right\rangle_{\mu_-} = \frac{1}{6}\int_{-\pi}^{\pi}\left\langle u^3\right\rangle_{\mu_-} dx \equiv \epsilon;$$

- The downstream equilibrium measure μ_+ is homogeneous at each physical grid point, so that the second and third moments are invariant at each grid point

$$\left\langle u_j^2\right\rangle_{\mu_+} = \sigma^2 = \pi^{-1}, \quad \left\langle u_j^3\right\rangle_{\mu_+} = \sigma^3\kappa_3 = \pi^{-\frac{3}{2}}\kappa_3^+.$$

Then the skewness of the downstream state variable u_Λ^+ after the ADC is given by the difference between the inflow and

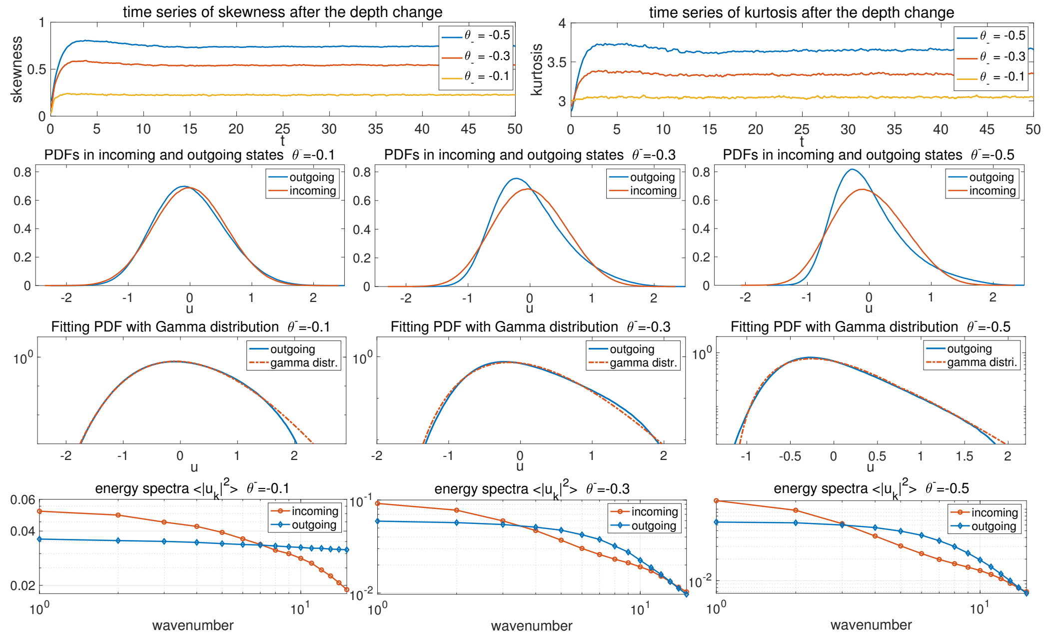


Fig. 2. Changes in the statistics of the flow state going through the abrupt depth change. The initial ensemble is set with the incoming flow Gibbs measure with different inverse temperature θ^-. First row: time evolution of the skewness and kurtosis. The abrupt depth change is taking place at $t = 0$; Second row: inflow and outflow PDFs of u_Λ; Third row: the downstream PDFs fitted with Gamma distributions with consistent variance and skewness (in log coordinate in y); Last row: energy spectra in the incoming and outgoing flows.

outflow wave slope energy of u_x

$$\kappa_3^+ = \frac{3}{2}\pi^{\frac{1}{2}} L_0^{-\frac{3}{2}} E_0^{-\frac{1}{2}} D_+^2 \int_{-\pi}^{\pi} \left[\left\langle u_x^2 \right\rangle_{\mu+} - \left\langle u_x^2 \right\rangle_{\mu-} \right] dx + 3\pi^{\frac{1}{2}}\epsilon. \quad [5]$$

The detailed derivation is shown in *SI Appendix, B.2*. In particular, the downstream skewness with near-Gaussian inflow statistics $\epsilon \ll 1$ is positive if and only if the difference of the incoming and outgoing wave slope energy is positive. This means that there is more small scale wave slope energy in the outgoing state. As an evidence, in the last row of Figure 2 in all the weak and strong skewness cases, the outflow energy spectrum always has a slower decay rate than the inflow energy spectrum which possesses stronger energy in larger scales and weaker energy in the smaller scales.

In Figure 1 (c), we compare the accuracy of the theoretical estimation [5] with numerical tests. In the regime with small incoming inverse temperature θ^-, the theoretical formula offers a quite accurate approximation of the third-order skewness using only information from the second-order moments of the wave-slope spectrum.

7. Key features from experiments captured by the statistical dynamical model

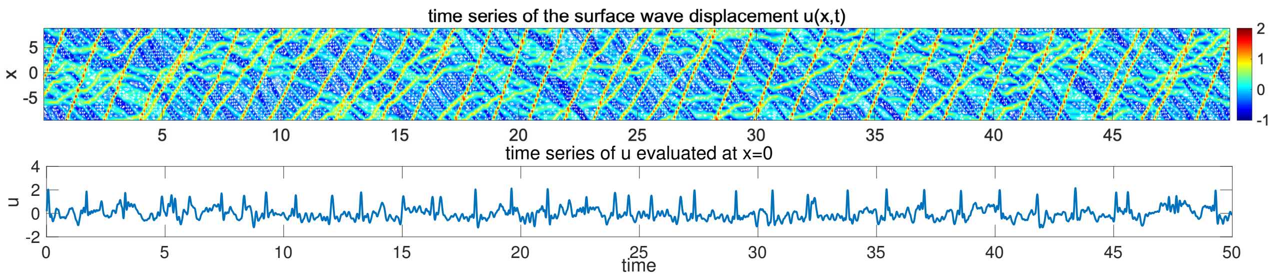
In this final section, we emphasize the crucial features generated by the statistical dynamical model [1] by making comparison with the experimental observations in (20). As from the scale analysis displayed in Section 2, the theory is set in the same parameter regime as the experimental setup.

- The transition from near-Gaussian to skewed non-Gaussian distribution as well as the jump in both skewness and kurtosis observed in the experiment observations (Fig. 1 of (20)) can be characterized by the statistical model simulation results (see the first and second row of Figure 2). Notice that the difference in the decay of third and fourth moments in the far end of the downstream regime from the experimental data is due to the dissipation effect in the flow from the wave absorbers that is not modeled in the statistical model here. The model simulation time-series plotted in Figure 3 can be compared with the observed time sequences from experiments (Fig. 1 of (20)). The downstream simulation generates waves with strong and frequent intermittency towards the positive displacement, while the upstream waves show symmetric displacements in two directions with at most small peaks in slow time. Even in the time-series at a single location $x = 0$, the long-time variation displays similar structures.

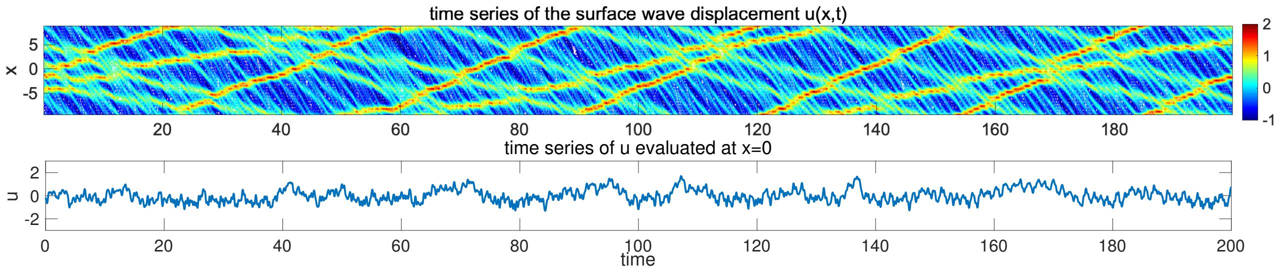
- The downstream PDFs in experimental data are estimated with a Gamma distribution in Fig. 2 of (20). Here in the same way, we can fit the highly skewed outgoing flow PDFs from the numerical results with the Gamma distribution

$$\rho(u; k, \alpha) = \frac{e^{-k}\alpha^{-1}}{\Gamma(k)} \left(k + \alpha^{-1}u\right)^{k-1} e^{-\alpha^{-1}u}.$$

The parameters (k, α) in the Gamma distribution are fitted according to the measured statistics in skewness and variance, that is, $\sigma^2 = k\alpha^2, \kappa_3 = 2/\sqrt{k}$. And the excess kurtosis of the Gamma distribution can be recovered as $\kappa_4 = 6/k$. As shown in the third row of Figure 2, excellent agreement in the PDFs with the Gamma distributions is reached in consistency with the experimental data observations. The accuracy with this approximation increases



(a) downstream state $D_0 = 0.24$



(b) upstream state $D_0 = 1$

Fig. 3. Realization of the downstream and upstream flow solutions $u_\Lambda^\pm$. Note the larger vertical scale in the downstream time-series plot.

as the initial inverse temperature θ^- increases in value to generate more skewed distribution functions.

- The experiments also have the up and down stream power spectra in time (Fig. 4 of (20)), which shows more energy at small time scales, i.e., a relatively slower decay rate in the downstream compared with the upstream case. This is also observed in the direct numerical simulations here (detailed results shown in *SI Appendix, C.2*). The downstream state contains more energetic high frequencies. The peak frequency illustrates the occurrence of the transporting waves along the water tank.

I do not understand the last sentence above. Why does the peak frequency illustrate the occurrence of transporting water waves?

8. Concluding discussion

We have developed a statistical dynamical model to explain and predict extreme events and anomalous features in shallow water waves with abrupt depth change. The theory is based on the dynamical modeling strategy consisting of the TKdV equation matched at the abrupt depth change with conservation of energy and Hamiltonian. Predictions can be made of the extreme events and anomalous features by matching incoming and outgoing statistical Gibbs measures before and after the abrupt depth transition. The statistical matching of the known nearly Gaussian incoming Gibbs state completely determines the predicted anomalous outgoing Gibbs state, which can be calculated by a simple sampling algorithm, verified by direct numerical simulations, and successfully captures key features of the experiment. An analytic formula for the anomalous outgoing skewness is also derived. The strategy here should be useful for predicting extreme statistical events in other dispersive media in different settings.

ACKNOWLEDGMENTS. This research of A. J. M. is partially supported by the Office of Naval Research through MURI N00014-16-1-2161. D. Q. is supported as a postdoctoral fellow on the second grant. M.N.J.M would like to acknowledge support from Simons grant 524259.

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