

Neutron Spectrum Characterization

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Abstract—This paper outlines a foil activation experiment performed at the Air Force Institute of Technology neutron pile. The purpose of the experiment was to measure the neutron flux spectrum in stringer 2 using foil activation unfolding techniques such as STAYSL. Due to facility limitations, the unfolding analysis was conducted on an experiment with a test snout performed at the National Ignition Facility (NIF). Foil activation experiments are an indirect method of measuring an incident neutron flux and are preferred in high flux environments that will damage electronics or for size constraints. Foil activities produced from threshold and non-threshold reactions were analyzed using a high purity germanium detector that resulted in time-corrected activities post-irradiation in the range of 100 Bq to 10 kBq. The NIF experiment utilized aluminum, gold, indium, and zirconium foils of varying mass at the pinhole, basket, and kinematic base locations of the snout. The irradiation planned for the neutron pile also included tungsten and manganese. The incident neutron fluence was unfolded from the activities and nuclear data with a MCNP simulated NIF initial spectrum and a flat spectrum. The results for the pinhole and basket show varying physical results, while both have p-values under 0.01. However, the results for the basket do not match intuition and likely requires a better starting spectrum. The result for the kinematic base shows a similar result to the pinhole. However, the p-value is near 0.9, which indicates the result is not likely to be representative of the activities produced. The kinematic base is a significant change from the initial source spectrum and requires a better simulated starting spectrum.

Index Terms—neutron flux unfolding, foil activation, high purity germanium, STAYSL

I. INTRODUCTION

CHARACTERIZING the energy dependent neutron environment has many applications to the nuclear sciences community. Determining a neutron flux is important for experiments where the neutron flux requires validation or is not well modeled. Neutrons can be detected using a variety of methods, such as Bonner spheres, long counters, He-3 based detectors, or proton recoil scintillators [1]. Foil activation experiments can also be performed to acquire an indirect measurement of the incident neutron fluence on a set of activation foils. Activation experiments are essential for testing that requires small geometry to fit in the apparatus or in situations where electronics equipment for higher fidelity measuring techniques will be damaged.

This work aims to characterize the neutron spectrum in the Air Force Institute of Technology's (AFIT) Building 470 neutron pile constructed in the 1960s, which contains a plutonium-beryllium (PuBe) neutron source [2]. Previous studies have accomplished results indicating changes to the neutron source term and modeled the neutron flux at various positions [3], [4]. However, the source characterization is still not entirely known.

The neutron source characterization was performed with a foil activation experiment. The foils were analyzed using a

calibrated high purity germanium (HPGe) detector to determine time corrected activities post-irradiation. The HPGe was calibrated for energy and efficiency to back out the initial foil activities post-irradiation.

The pile activation results are not used in the final neutron flux analysis due to issues with the HPGe at AFIT. Instead, this work analyzed an experimental setup at the National Ignition Facility (NIF). The NIF foil activation experiment was performed on a snout for a shot performed in March 2018 [5]. The goal of the NIF experiment was to characterize the NIF source in the aluminum snout which has applications for exploring cross-section uncertainties if there are unexpected results. The incident neutron flux was determined from the foil activation results using Pacific Northwest National Laboratory (PNNL) STAYSL, which uses generalized least-square minimization to determine the flux spectrum for given foil activities and nuclear data.

II. PROBLEM DESCRIPTION

The neutron spectrum characterization requires information from the foil activities. The theory and problem statement contributing the foil activation involves the neutron flux environment, foil activation, HPGe spectroscopy, and neutron flux unfolding. Supporting documentation and work performed is available on Github [13].

A. Building 470 Pile Neutron Environment

The neutron pile was modeled in MCNP version 6.1 based on previous work to gauge expected activation rates and to determine saturation activities [3]. The setup of foil placement and PuBe source is shown in Figure 1. The pile has stringers with fabricated placements for the source and foils to help ensure repeatability of measurements.

The starting PuBe neutron spectrum, shown in Figure 2, ends at 11 MeV and peaks near 3.5 MeV. However, there is significant thermalization of the starting PuBe spectrum by the graphite pile as shown in the simulated stringer 2 spectrum in Figure 3. The simulated spectrum was used to select activation foils based on the modeled foil activation rates, discussed in Section ???.

B. NIF Snout Experiment

A passive foil activation experiment was performed in 2018 at the NIF. A snout composed of aluminum was mounted to the Target and Diagnostic Manipulator (TANDM) 90-348. Three foil activation diagnostic packs of aluminum, zirconium, indium, and gold were placed in the nose cap pinhole (7 cm from source), filter basket (41 cm), and kinematic base (110 cm). The aluminum foil was not used in the pinhole. Figure 4 displays the snout with the activation foil sites [5].

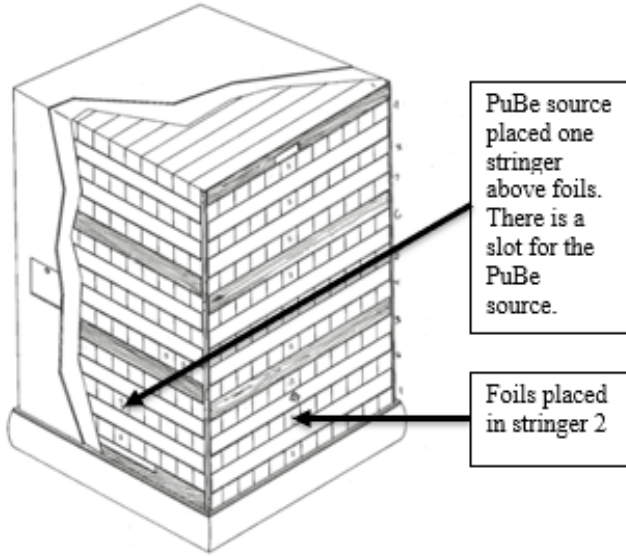


Fig. 1. Activation foil setup for Building 470 graphite pile irradiation.

The NIF source is a deuterium-tritium (DT) capsule which have neutrons at a nominal energy of 14 MeV. The spectrum of neutrons used as a source term is shown in Figure 5.

C. Foil Activation

Foil activation can be used with published nuclear data to provide an estimate of the incident flux on the foils by unfolding techniques. The production rate of radioactive isotopes is negated by radioactive decay processes, which place an upper limit on the radioactivity of a foil [1]. The saturated activity, (A_∞), for a given reaction is given by:

$$A_\infty = R = \int_{E_1}^{E_2} \phi(E) \Sigma(E)_{act} V_{foil} dE \quad (1)$$

the saturated activity is equivalent to the reaction rate (R), which is a function of the energy dependent flux (ϕ), the macroscopic reaction cross-section for the activation reaction channel ($\Sigma(E)_{act}$), and the volume of the foil (V_{foil}). The energy term (E_1) is zero in many cases; however, threshold reactions require the incident neutron to be of higher energy to enable the reaction channel.

A correction needs to be made in cases where the activation is not sufficient to fully saturate the foil. At six half-lives a foil will have reached approximately 98% of its saturation activity, neglecting spatial and energy self-shielding effects [1]. The activation of the foil for a given irradiation time (t_i) is given as a function of the decay constant:

$$A_0 = A_\infty (1 - e^{-\lambda t_i}) \quad (2)$$

Experimental measurements also can be corrected to deduce the original activity of the foil, immediately after irradiation. The measured counts, C , is reduced by the background counts, B . A corrects for the radioactive decay for the time between the end of irradiation and the start of counting (t_{decay}). A similar correction factor based on the count time, t_{count}

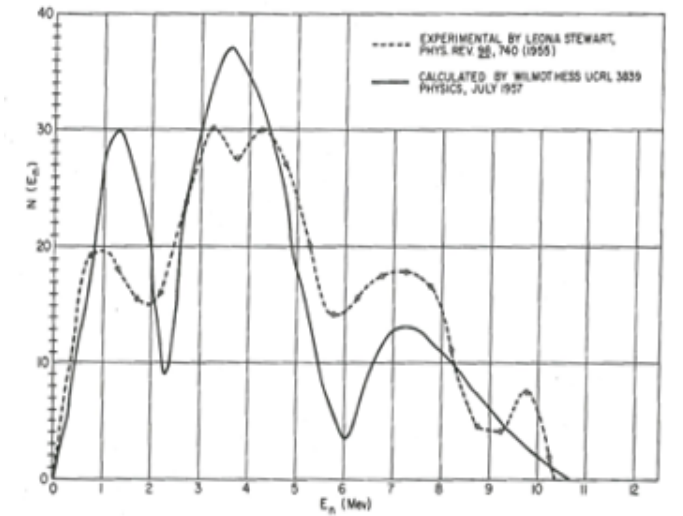


Fig. 2. PuBe neutron emission source spectrum.

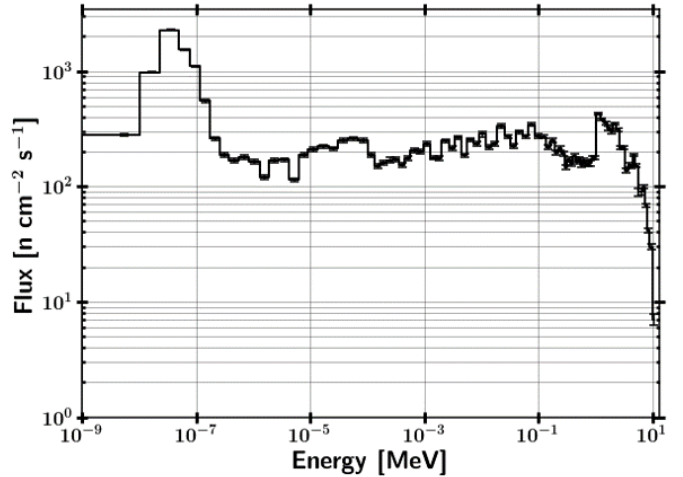


Fig. 3. Foil activation neutron flux spectrum (all statistical errors are under 1.5 percent).

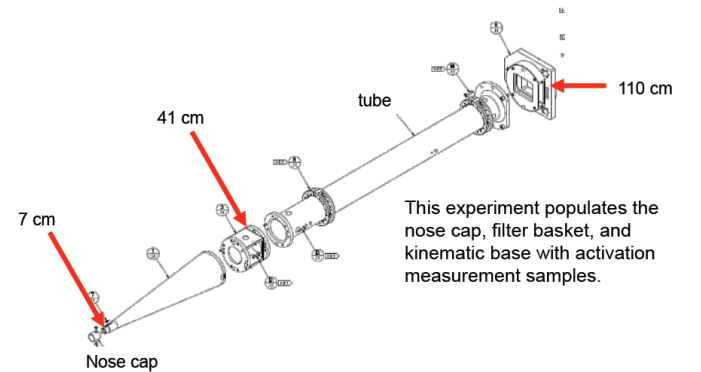


Fig. 4. Passive NIF snout foil activation locations. Foils were placed in the nose cap pinhole (7 cm from source), filter basket (41 cm), and kinematic base (110 cm).

provides a correction for radioactive decay during counting. Additionally, the detector efficiency for the given gamma-ray energy, ϵ , and relative gamma intensity, I_γ , must be taken into account. The gamma intensity may also include a branching

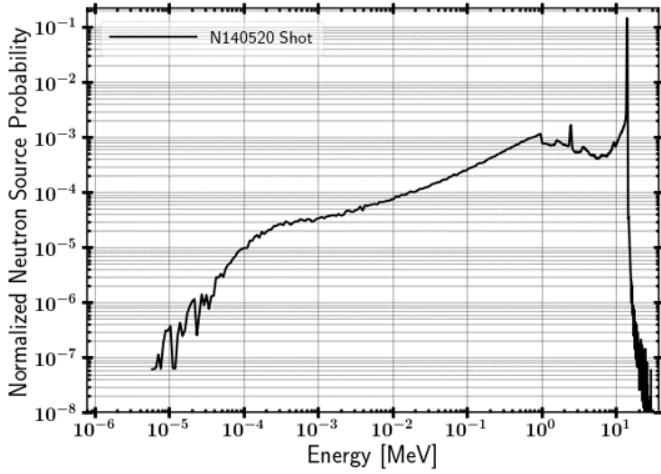


Fig. 5. A typical NIF D-T neutron source spectrum for high-yield indirect drive shots.

ratio if applicable to the decay mechanism. All corrections included, less self-shielding effects, provide a formulation for converting counts to post-irradiation activity as:

$$A_0 = \frac{\lambda(C - B)e^{\lambda t_{decay}}}{\epsilon(1 - e^{-\lambda t_{count}})I_\gamma} \quad (3)$$

The formula can be simplified in the limit of irradiation times much less than the half-life of the activation products. In this case, the reaction rate is much larger than the decay from radiation, so the rate of production of the radioisotope is driven only by the reaction rate. The time integrated flux, or neutron fluence (Φ), can be used to determine the total reactions, R_{total} , over an irradiation period, given by:

$$R_{total} = \int_{E_1}^{E_2} \Phi(E) \Sigma(E)_{act} V_{foil} dE \quad (4)$$

The foils selected must meet several requirements for this experiment. A list of the various requirements that are of importance for a neutron activation foil experiment with energies in the range of thermal to approximately 11 MeV are summarized below [1], [6], [7].

The reaction neutron cross-section is extremely important for foil activation, and there are a few key parameters that should be considered. First, the magnitude of the cross-section determines the reaction rate of the product nuclides. A large cross-section allows for more activation, and therefore better results when analyzing the activation foils. Second, the uniqueness of the cross-section shape is used to unfold the incident neutron energy spectrum. An (n, γ) cross-section may peak in a particular region, which is essential to providing information of the neutron flux in that energy region. Alternatively, a threshold reaction, such as an $(n, 2n)$, is important for providing information of the flux at higher energies. Third, the selected foils for an experiment should cover the entire energy range of the incident neutron flux. Finally, the cross-section must be well characterized with low uncertainty over the neutron energy range of interest.

The decay constant of the product nuclides is important. The half-lives applicable for a particular experiment depend on the time post-irradiation that the foils can be counted. A long lived radioisotope will be available for counting for longer times, but the activity will be reduced due to the lower decay constant. The opposite is true for short half-lives. A half-life on the order of an hour to a few years is in the right direction; however, the half-life must also be balanced with the production of the radioisotope to understand the entire picture.

The elemental and chemical purity of the activation foil should be well known. An unknown composition foil will likely cause erroneous results.

Interfering reaction channels and decay emissions should be avoided. An example of this is natural copper, which has multiple 511 keV emissions from different reaction channels. It is difficult to distinguish these gamma-rays to determine activation in counting. Similar problems arise in multi-isotope materials that have multiple reactions producing the same nuclide. For example, a Cadmium-106 (n, γ) reaction produces the same isotope as a Cadmium-108 $(n, 2n)$ reaction.

The activation foil should be optically thin so as to not cause perturbations of the neutron flux. An additional benefit of relatively thin foils is that the gamma-ray emissions are not significantly attenuated through self-shielding. The neutron flux should ideally also not be changing substantially over the foil region. In general adding additional foils helps to improve the unfolding results, as long as the entire foil set remains optically thin [8].

The decay nature of the product nuclide should be a gamma-ray emitter. Gamma-ray detection can provide fine energy resolution to determine activation due to specific reaction channels. The discrete gamma-ray emissions provide a means of determining the source and magnitude of the foil activation. The energy of the gamma is also of importance. Semiconductor detection methods have a peak intrinsic efficiency near 100 keV with some variance depending on if it p-type or n-type.

D. HPGe Calibration

The HPGe efficiency as a function of energy was determined using a known multi-nuclide source and measurements of the experimental setup. The absolute efficiency was calibrated at a source distance to obtain the unknown source activity from the activation foils [1]. The absolute efficiency (ϵ_{abs}) was obtained by:

$$\epsilon_{abs} = \frac{\# \text{ Events Recorded}}{\# \text{ Radiation Quanta Emitted}} \quad (5)$$

A calibrated efficiency curve can be used to find the source strength of samples, scaled by the distance to the detector. There are a few types of fits that are appropriate for efficiency curves. Piece-wise curves can work in certain applications. A common efficiency curve is given as a function of N fitting parameters and a reference energy (E_0) was used in this work:

$$\ln(\epsilon_{abs}) = \sum_{i=1}^N a_i \left(\ln\left(\frac{E_\gamma}{E_0}\right) \right)^{i-1} \quad (6)$$

E. Neutron Flux Unfolding

Foil activation experiments are a well established method for determining an incident neutron energy spectrum. The foils are exposed to a nearly equivalent neutron flux, which serves to activate the foil samples through various nuclear reaction channels. Each of the reactions has a unique response function with respect to the neutron flux. The nuclear data and activities of the foils can be used to unfold the incident neutron energy spectrum.

A few examples of studied methods of unfolding matrix inversion, least-square spectral adjustment, and stochastic algorithms [10]. Matrix inversion can lead to non-physical results, such as negative fluxes [10]. Stochastic methods rely on random sampling to derive a best-fit or average over a group of reasonably well-fitting spectra [10].

The least-squares method minimizes the chi-square based on a guess spectrum, activation information, and nuclear data [11]. The least-squares method is also known as spectral adjustment and can incorporate more information, most notably the underlying nuclear data, into the determination of the resultant spectrum [11]. The optimization is performed to minimize varying versions of the chi-square statistic among an energy group structure for the flux and nuclear data.

The general formulation of the least-squares method is derived from minimizing the activation results to the nuclear data and input spectrum [11]. The chi-square, χ^2 , is given as per degrees of freedom (n) as a function of the uncertainty, activation rates, nuclear data, and measured results. The chi-square formulation of the least-squares approach can be reduced if there is no time dependency of the neutron flux as:

$$\frac{\chi^2}{n} = \frac{1}{n} \sum_{i=1}^m \frac{(\sum_{j=1}^N \Sigma_i(E_j) \Phi(E_j) - \frac{A_i}{V_{Foil}})^2}{\sigma_i^2} \quad (7)$$

Providing an initial spectrum is required for many unfolding methods. The activities produced for the foils is often highly degenerate, where an infinite amount of spectra could provide the same end-point. The initial spectrum allows for the insertion of more physics based results to have an impact on the overall result. For neutron spectra, an initial guess spectrum is often created with a particle transport code or a deterministic solution.

The foil activities are used with the underlying nuclear data to unfold the neutron spectrum using Pacific Northwest National Laboratory (PNNL) STAYSL. STAYSL relies on least-squares spectral adjustment based on the chi-square of the measured activities to determine the incident neutron flux [12].

III. DESCRIPTION OF WORK

The selected foils for the experiment are summarized in Table I. The foils were selected based on availability and usability in the nuclear data libraries used by STAYSL.

The irradiation time in the neutron pile was planned to be ten days. This time was chosen to build up the enough activity

to be able to count the foils to 1% statistics. Specifically, Au-197 (n,2n) has a very long half-life, so approximately four half-lives of Au-196 were used for the irradiation time. The expected post-irradiation activities is summarized in Table II.

A. Foil Selection / Irradiation Parameters

The simulated environment in MCNP resulted in no feasible reactions in the high energy region, such as the (n,2n) threshold reactions. The lack of pinning at high energy has a potential issue for the neutron flux unfolding, where there is no reaction pinning the high end of the neutron energy spectrum.

B. HPGe Characterization

The HPGe detection system can be streamlined compared to a traditional nuclear counting experiment. An ORTEC HPGe was connected to a DSA1000 which functions to replace many components necessary for a traditional nuclear detection system. The DSA1000 was connected to a computer using the Genie 2000 multi-channel analyzer (MCA) data acquisition software. The bias voltage is set to -4,000 V. The gain was set to 20 and the fine gain to 1.5, this allowed the dynamic range of the MCA to cover the gamma ray energies of interest.

The HPGe was characterized using a multi-nuclide source. The isotopes gamma decay energies in the source range from 60 keV from Am-241 to 1836 keV from Y-88, which covers the energy range of interest for the foil analysis. The resulting energy calibration utilized a linear relationship to map channel to energy as shown in Figure 6.

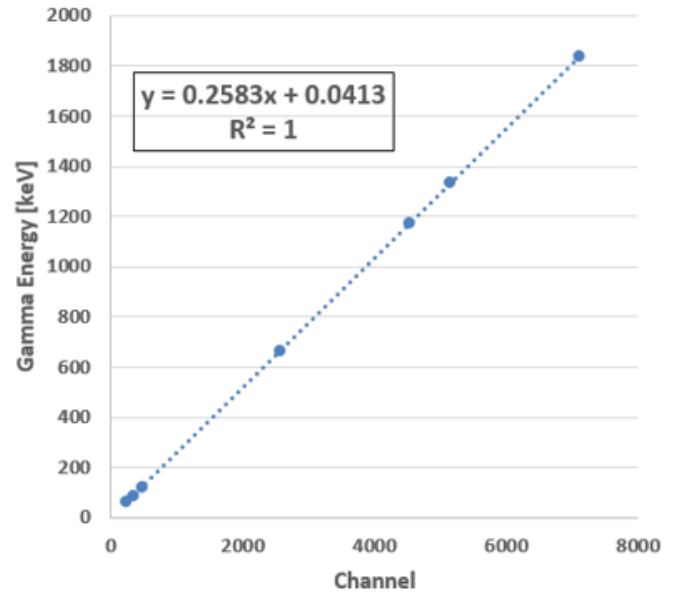


Fig. 6. HPGe energy calibration curve obtained using a multi-nuclide source.

The efficiency calibration was conducted at 10 cm and 5 cm from the base of the HPGe. The solid angle for a quarter inch radius foil at the 10 cm range is 0.66 radians, which is notably less than the idealized point source simplification of 0.79 radians.

A malfunction in the HPGe occurred prior to the completion of collecting data for the 5 cm case. The error prevented

TABLE I
ACTIVATION FOILS SELECTED FOR IRRADIATION IN NEUTRON PILE.

| Foil (thickness) | Reaction | Half-life | Threshold [MeV] | Decay Radiation [keV] (Intensity) | Elemental Purity [%] | Mass [g] |
|----------------------------------|-----------------------|-----------|-----------------|-----------------------------------|----------------------|----------|
| In (0.005" x2) | In-115 (n,g) In-115m | 54.29 min | Thermal | 1293.56 (0.848) | 99.99 | 0.221(7) |
| | In-115 (n,n') In-116m | 4.49 hrs | 0.336 | 336.24 (0.459) | 99.99 | 0.221(7) |
| Al (~2mm x4) | Al-27 (n,a) Na-24 | 15.00 hrs | 3.25 | 1368.63 (0.9999) | 99 | 3.395(3) |
| Au (0.005" x2) | Au-197 (n,2n) Au-196 | 6.17 days | 8.11 | 355.7 (0.87) | 99 | 0.577(2) |
| | Au-197 (n,g) Au-196 | 2.69 days | Thermal | 411.8 (0.9562) | 99 | 0.577(2) |
| W (0.005" x2) | W-186 (n,g) W-187 | 24.00 hrs | Thermal | 685.51 (0.332) | 99.98 | 0.627(2) |
| Mn (0.068 mm to 0.077 mm) x4 | Mn-55 (n,g) Mn-56 | 2.58 hrs | Thermal | 846.8 (0.9885) | 99 | 0.266(0) |
| Zr (Not used in Pile Experiment) | Zr-90 (n,2n) Zr-89 | 78.41 hrs | 12.1 | 909.2 (0.9904) | N/A | N/A |

TABLE II
ACTIVATION INFORMATION FOR SELECTED PILE FOILS.

| Reaction | A _{inf} [Bq] | A ₀ [Bq] | Count Time to 10,000 Counts (at 1% efficiency) [hrs] |
|----------------------|-----------------------|---------------------|--|
| W-186 (n,g) W-187 | 176.90 | 175.52 | 9 |
| In-115 (n,g) In-115m | 1394.07 | 1394.07 | 1 |
| Au-197 (n,g) Au-196 | 820.33 | 684.87 | 7 |
| Mn-55 (n,g) Mn-56 | 256.32 | 256.32 | 3 |
| All others | <1.0 | <1.0 | N/A |

additional data acquisition for the building 470 pile portion of this experiment.

C. NIF Snout Experiment Activities

The HPGe analyzed dataset used for the NIF experiment was provided from Lawrence Livermore National Lab (LLNL). The counts per channel data were also provided. A validation test was performed with Los Alamos National Laboratory's Peak Easy on the pinhole results for the indium foil. The main In-115m peak at 336 keV is shown in Figure 7. The foils were 1 mm thick with the exception of the gold foils, which were 0.1 mm thick.

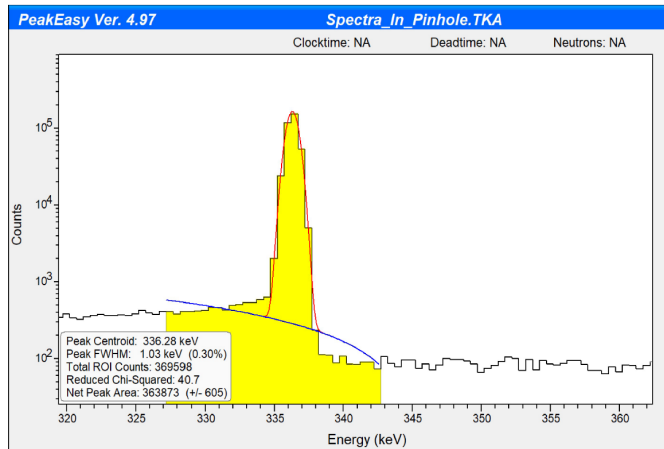


Fig. 7. 336 keV In-115m peak counts determined using Peak Easy.

There was general agreement between the counts for In-115m and In-116m. The counts identified by LLNL resulted in

TABLE III
ACTIVATION INFORMATION FOR THE NIF EXPERIMENT FOILS.

| Kinematic Base | | | | | | |
|----------------|----------|----------|-------------|---------------|----------------------|---------------------|
| isotope | mass (g) | A0 (Bq) | N0 (nuclei) | Percent Error | Gamma Self Shielding | "sig-phi" (at/at-s) |
| Au196g | 3.733 | 5.74E+02 | 4.42E+08 | 1.60 | 0.974 | 3.98E-14 |
| Au198 | 3.733 | 6.63E+02 | 2.23E+08 | 1.20 | 0.980 | 1.99E-14 |
| In115m | 14.35 | 5.57E+03 | 1.30E+08 | 1.20 | 0.950 | 1.82E-15 |
| In116m | 14.35 | 2.12E+05 | 9.97E+08 | 1.60 | 0.998 | 1.33E-14 |
| Zr89 | 12.555 | 1.13E+03 | 4.61E+08 | 1.50 | 0.980 | 5.59E-15 |
| Na-24 | 12.56 | 3.63E+03 | 2.83E+08 | 7.90 | 0.993 | 1.02E-15 |
| Basket | | | | | | |
| Au196g | 0.9393 | 9.47E+02 | 7.29E+08 | 1.20 | 0.974 | 2.61E-13 |
| Au198 | 0.9393 | 2.00E+02 | 6.72E+07 | 1.20 | 0.980 | 2.39E-14 |
| In115m | 0.4189 | 9.30E+02 | 2.17E+07 | 1.20 | 0.950 | 1.04E-14 |
| In116m | 0.4189 | 1.69E+04 | 7.91E+07 | 2.80 | 0.998 | 3.61E-14 |
| Zr89 | 0.2626 | 1.76E+02 | 7.15E+07 | 1.10 | 0.980 | 4.15E-14 |
| Na-24 | 0.0962 | 4.11E+02 | 3.20E+07 | 1.20 | 0.993 | 1.50E-14 |
| Pinhole | | | | | | |
| Au196g | 0.148 | 6.97E+03 | 5.37E+09 | 1.30 | 0.974 | 1.22E-11 |
| Au198 | 0.148 | 1.07E+02 | 3.58E+07 | 5.30 | 0.980 | 8.08E-14 |
| In115m | 1.182 | 1.23E+05 | 2.87E+09 | 0.70 | 0.950 | 4.88E-13 |
| In116m | 1.182 | 1.53E+05 | 7.18E+08 | 2.00 | 0.998 | 1.16E-13 |
| Zr89 | 1.008 | 2.77E+04 | 1.13E+10 | 1.00 | 0.980 | 1.71E-12 |

361,000 \pm 0.23% for In-115m and 1,383 \pm 2.89% for In-116m. The calculated results from Peak Easy determined 364,000 \pm 0.38% for In-115m and 1,402 \pm 3.8% for In-116m. It is expected that multiple peaks were used in the LLNL results thereby lowering the overall uncertainty. A summary of the foils and reactions used for the snout experiment is given in Table III. The value of "sig-phi" was a STAYSL calculated reaction rate used for the unfolding of the source spectrum.

D. Unfolding NIF Experiment with STAYSL

STAYSL has several modules that are used to unfold the neutron spectrum from the calculated activities. The main components used in this analysis are SHIELD, SIG-PHI Calculator, and PNNL STAYSL. The Beam Correction factor was not used because the NIF irradiation time is much less than the half-lives of the reaction products.

A complete walk-through of the analysis is available on Github [13]. SHIELD was used to generate energy dependent neutron self-shielding factors for non-threshold reactions. SHIELD is not used on high energy threshold reactions because there is negligible shielding.

The SIG-PHI Calculator was used to consolidate all of the reaction information. Additionally, this module generates

gamma-ray shielding factors. The input was generated from this information and the nuclei production information.

STAYSL requires a starting guess spectrum. For this work, the NIF source was binned into the 140 group STAYSL structure, and the flux was scaled by the spherical divergence to the foil locations.

STAYSL was run for the pinhole, basket, and kinematic base for two cases. First, a guess spectrum of the NIF source was used with 100 percent flux uncertainty up to 13 MeV. The flux uncertainty from 13-15 MeV was set to five percent for the pinhole based on the reaction cross-section and threshold reactions expected results. Less was known about the resultant spectra for the kinematic base and basket, so the 13-15 MeV flux uncertainty was set to 100 percent. STAYSL was run iteratively using STAYSL.py until the χ^2 changed by less than 0.1. The product spectrum is used as the guess spectrum for subsequent iterations. The flux uncertainty was not updated until the solution converged based on the χ^2 . The degrees of freedom (ν) for the pinhole is 4, because there are 5 foils. The other sets have 5 degrees of freedom.

The second method of iteratively solving started with a flat guess spectrum. A flat guess spectrum utilizes no a priori information about the flux, so it is a good indicator if the result matches up with the NIF guess spectrum iterations. The flat guess spectrum started with 100 percent uncertainty in all energy groups.

IV. RESULTS

The unfolded results from STAYSL for the pinhole, basket, and kinematic base are shown in Figures 8 - 10. A summary of the χ^2 values and associated p-values are provided in Table IV. The basket results are truncated to a minimum fluence of $1 \text{ n} - \text{cm}^{-2}$.

TABLE IV
SUMMARY OF UNFOLDED RESULTS FOR PINHOLE, BASKET, AND KINEMATIC BASE.

| Location | χ^2 / ν | ν | p-value |
|---------------------------|----------------|-------|---------|
| Pin NIF Guess | 4.4 | 4 | 0.0015 |
| Pin Flat Guess | 6.6 | 4 | 3e-5 |
| Basket NIF Guess | 18.7 | 5 | 1e-18 |
| Basket Flat Guess | 17.7 | 5 | 1e-17 |
| Kinematic Base NIF Guess | 0.35 | 5 | 0.88 |
| Kinematic Base Flat Guess | 0.53 | 5 | 0.78 |

The results show good generally good agreement between the NIF guess and flat spectrum. No modeling information was available to use as a guess spectrum. The χ^2 values can be iterated to shrink the value; however, care must be exercised in how many iterations to perform as non-physical results might be created.

The pinhole results in Figure 8 show that there is a steep decrease in flux as energy decreases, which was not expected. The high energy bins at 13-15 MeV have very good agreement with expected values; however, the rest of the spectra was unexpected to some extent. The large epithermal peak energy is approximately 800 keV; which is approximately the first excited state in Al-27 (843 keV). There is no physical basis for the near complete removal of neutrons between 1 and 10 MeV,

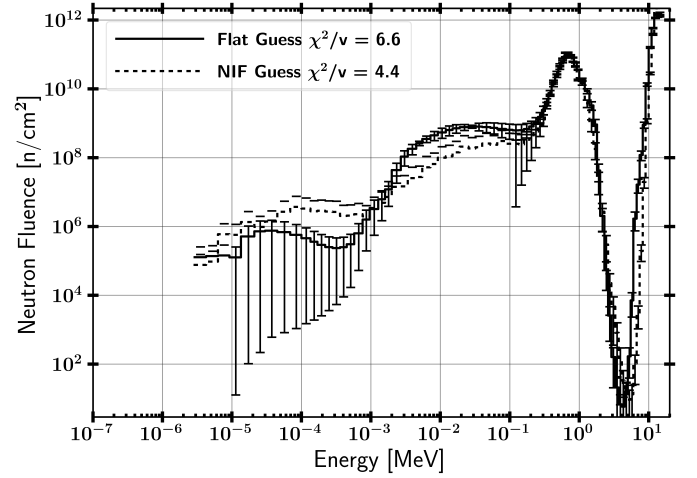


Fig. 8. Unfolded Spectrum for pinhole starting with NIF and flat guess.

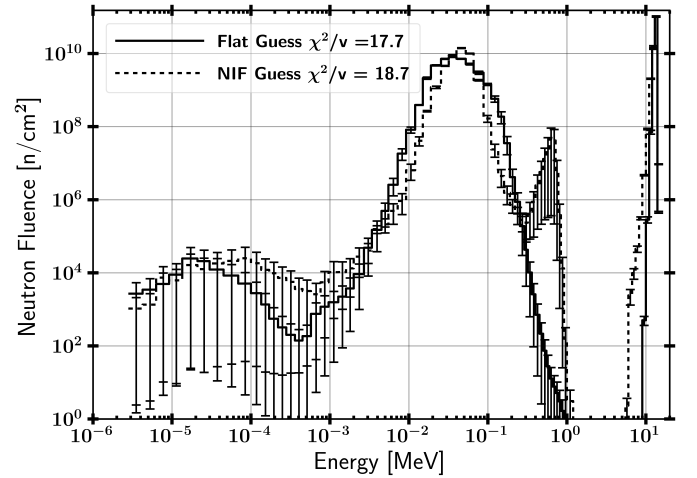


Fig. 9. Unfolded Spectrum for the basket starting with NIF and flat guess (Truncated at $1 \text{ n} - \text{cm}^{-2}$).

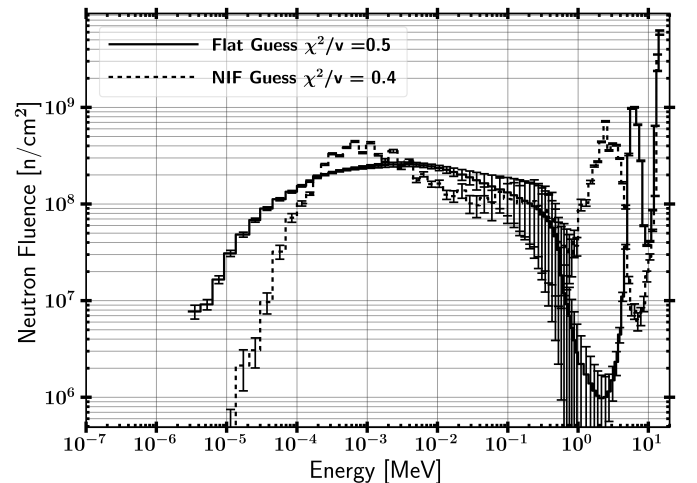


Fig. 10. Unfolded Spectrum for the kinematic base starting with NIF and flat guess.

while the source distribution indicates there should be a large portion present. The peak is shaped like an neutron evaporation spectrum. Past the epithermal peak, there is a decaying thermal tail. Starting from a flat spectrum produced nearly identical results, so it would have been ideal to have a threshold reaction near 1 MeV. The p-values for the pinhole indicate the result is very likely in both cases. Again, there is some degeneracy in the solution.

The basket results are similar to the pinhole with some key differences. First, the main 13-15 MeV peak has dropped by approximately six orders of magnitude, which is not representative of the linear attenuation of 34 cm of aluminum ($\Sigma \approx 0.1 \text{ cm}^{-1}$). There is good agreement between the flat and NIF guess spectrum, so the threshold reactions are indicating the dropoff. Second, the neutron flux between approximately one to ten MeV is nearly removed. There is not a process that would remove all of the source neutrons in this region, but the result appears for the flat starting spectrum as well. The last difference is that the 800 keV peak only is present in the flat spectrum, and both guess spectra result in a peak at approximately 40 keV. The p-values for both spectra indicate that the result matches the activation. The result may have been improved with a better initial guess.

The kinematic base has the largest variance in results. At 13-15 MeV, the neutron flux is at approximately 10^9 , which is more in line with expectations. The NIF guess and flat guess both have an epithermal energy peak; however, the location is different. In both spectra there is a large thermal tail from scattering. The results for the kinematic base are likely the farthest from the true distribution as the initial guess spectrum was farthest from reality. The p-values for the kinematic base indicate that the resultant spectra are not indicative of the neutron environment.

V. CONCLUSIONS

Characterizing the neutron energy spectrum in the snout experiment at the NIF provided mixed results. The foil activities and underlying nuclear data were applied in STAYSL to unfold the spectrum with STAYSL for the pinhole, basket, and kinematic base locations in the snout. The results in each case were not expected. An epithermal peak was produced and higher energy neutrons were removed, which does not align with anticipated results. The reduced χ^2 result indicated that the pinhole and basket unfolds matched the activation results.

The unfolded neutron spectrum results highlight and reinforce necessity of having an initial guess spectrum for the generalized least-squares minimization. Each unfold, completed with the NIF starting spectrum and a flat spectrum, produced similar results; however, there are some differences that require attention. The kinematic base in particular did not converge to a similar solution over a large energy range. The results presented would be improved with a better starting guess for the initial spectrum.

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