# JF PY1T10 Special Relativity

Lecture 13:

Special Relativity, Electricity and Magnetism

### Electric Force

#### Coulomb's Law:

Force which as *stationary* charge  $q_1$  exerts on a *stationary* charge  $q_2$ :



**Source** charge

Force on  $q_2$ :

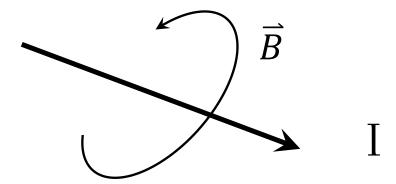
$$F = k \frac{q_1 q_2}{r^2} \hat{r}$$

where 
$$k = \frac{1}{4\pi\epsilon_0}$$
,  $\epsilon_0 = 8.85 \times 10^{-12} \text{Fm}^{-1}$ 

Coulomb's law also holds when  $q_2$  is moving, as long as  $q_1$  is stationary.

# Magnetic Field $(\vec{B})$

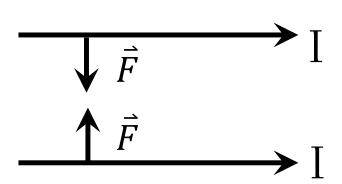
 $\overline{B}$ -field is generated by a moving charge, e.g. a current in a wire



Charge moving with velocity  $\vec{u}$  in  $\vec{B}$  is subject to magnetic force:

$$\vec{F}_{mag} = q\vec{u} \times \vec{B}$$

This equation defines  $\overline{B}$  in terms of a force which is proportional to velocity and charge.



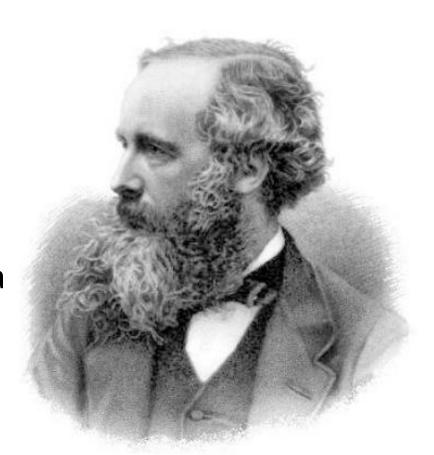
## Maxwell's Equations (1864)

Defined relations between  $\overline{E}$  and  $\overline{B}$  for stationary and moving charges.

#### Developed before SR but still:

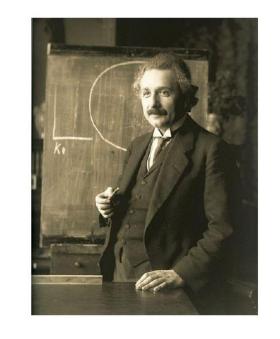
- Correct as  $v \rightarrow c$
- Describes electromagnetic (e.m.) phenomena in *any* inertial frame

But  $\vec{E}$  and  $\vec{B}$  are <u>not</u> the same in different inertial frames.



## Einstein's Insight

" What led me more or less directly to the special theory of relativity was the conviction that the electromotive force acting on a body moving in a magnetic field was nothing else than an electric field"



#### Both at rest:

 $q_1 igordam{\vec{F}_{elec}}{\vec{r}_{elec}}$  only

 $q_2 \bigcirc \vec{E}$  only

#### **Both moving:**

 $q_1 \longrightarrow \vec{F}_{elec} + \vec{F}_{mag}$ 

 $q_2 \longrightarrow \vec{E} \text{ and } \vec{B}$ 

What appears as  $\overrightarrow{B}$  in one frame is nothing else but  $\overline{E}$  when viewed from another frame.

Use S.R. to specify this link between  $\vec{E}$  and  $\vec{B}$ 

## Einstein's Insight

Consider frame in which source charge is at rest.

Coulomb's law tells us the force on test charge  $q_2$ .

Now: Find force on  $q_2$  in another frame using transformation of force.

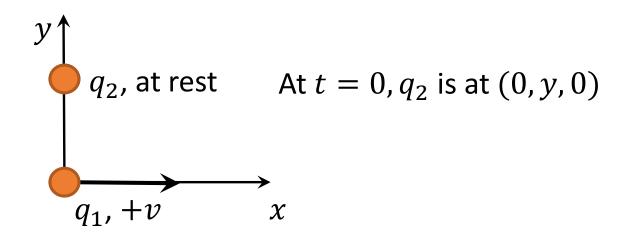
Identify  $\overrightarrow{B}$  from component of force on  $q_2$  which <u>depends</u> on velocity of  $q_2$  Identify  $\overrightarrow{E}$  from component of force on  $q_2$  which is <u>independent</u> of velocity of  $q_2$ 

Invariance of charge: Charge not changed by motion of carrier

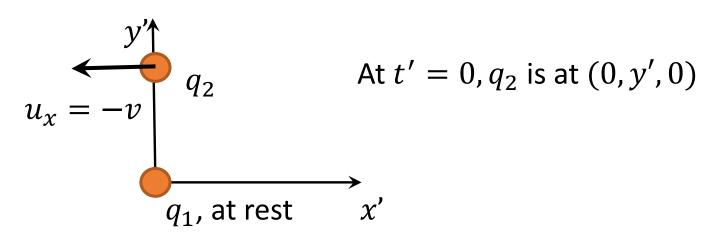
[Evidence: Take neutral atoms where  $\sum q_i = 0$ . Use heat to speed up the  $e^-$ . Still  $\sum q_i = 0$ . Atoms are neutral even though electrons in different orbits have different velocities.]

### Force on a stationary test charge $q_2$ due to moving source $q_1$

### Source $q_1$ with $v \mid\mid x$ :



### Shift to S', in which $q_1$ is now at rest:



In S', as source is at rest, we can apply Coulomb's Law:

$$F_y' = \frac{kq_1q_2}{y'^2}, F_y' = F_z' = 0$$

### Force on a stationary test charge $q_2$ due to moving source $q_1$

Transform back to inertial frame *S*:

$$F_{y} = \frac{F_{y}'}{\gamma \left(1 + \frac{vu_{x}'}{c^{2}}\right)} = \frac{F_{y}'}{\gamma \left(1 - \frac{v^{2}}{c^{2}}\right)} = \gamma F_{y}'$$

$$F_x = F_z = 0$$
 (as  $u' = u'_x$  and  $F' = F'_y$ ,  $\therefore \overline{F'}.\overline{u'} = 0$ 

Also, y' = y (by construction)

$$\Rightarrow F_{y} = \gamma \frac{kq_{1}q_{2}}{y^{2}}$$

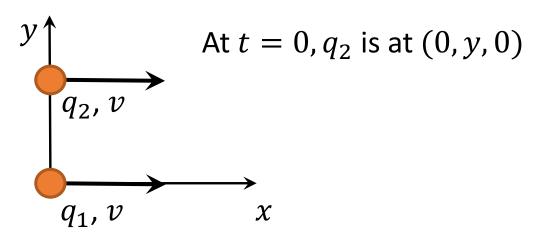
This is Coulomb's law, now modified for moving source charge.

Note: This is for a specific case when the line from the moving source charge to the stationary test charge is  $\perp$  to direction of motion of the source charge. It can be generalized for any point.

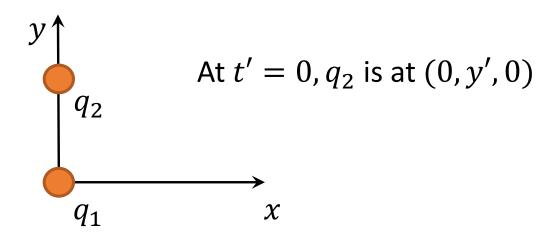
### Force on a moving test charge $q_2$ due to moving source $q_1$

Now assume both source  $q_1$  and test  $q_2$  are moving with velocity  $v \mid\mid x$  at t=0

#### S frame:



S' frame: Both  $q_1$  and  $q_2$  are at rest.



In S', as source is at rest, we can apply Coulomb's Law:

$$F_y' = \frac{kq_1q_2}{y'^2}, F_y' = F_z' = 0, \overrightarrow{u'} = 0$$

Force on a moving test charge  $q_2$  due to moving source  $q_1$ 

Again, transform back into *S*:

$$F_{y} = \frac{F_{y}'}{\gamma \left(1 + \frac{vu_{x}'}{c^{2}}\right)} = \frac{1}{\gamma} F_{y}', \qquad F_{x} = F_{z} = 0$$

$$F_{y} = \frac{1}{\gamma} \frac{kq_{1}q_{2}}{y^{2}} \left(\operatorname{as} \overrightarrow{u'} = 0\right)$$

Compare with previous situation (where  $q_2$  was stationary, and  $q_1$  was moving. We had:

$$F_{y} = \gamma \frac{kq_1q_2}{y^2}$$

### Magnetism

Therefore, we have an <u>extra</u> force on  $q_2$  due to its velocity in S. This is the **magnetic force** on  $q_2$ .

$$F_{mag} = \frac{kq_1q_2}{y^2} \left(\frac{1}{\gamma} - \gamma\right)$$

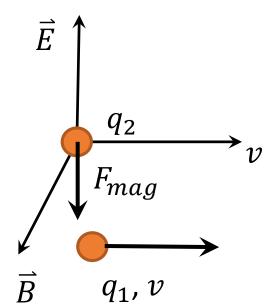
$$= -\frac{v^2}{c^2} \gamma \left(\frac{kq_1q_2}{y^2}\right)$$

$$F_{mag} = -\frac{v^2}{c^2} F_{elec}$$

$$F_{mag}$$
 is given by:  $F_{mag} = q_2 \vec{v} \times \vec{B}$   
If  $\vec{v} \perp \vec{B}$ :  $F_{mag} = q_2 vB$ 

What direction is  $\vec{B}$ ?

### Magnetism



If  $\vec{B} \perp \vec{v}$  (as expected for the magnetic field due to current in direction of  $\vec{v}$  of  $q_1$ ) then:

 $F_{mag}$  is in the opposite direction to  $F_{elec}$ 

This is consistent with  $F_{mag} = -\frac{v^2}{c^2} F_{elec}$ 

$$|qvB| = \left| \frac{v^2}{c^2} \gamma \left( \frac{kq_1 q_2}{y^2} \right) \right|$$

But  $F_{elec} = q_2 E$  where  $E = \gamma \frac{kq_1}{v^2}$ 

$$\Rightarrow B = \frac{1}{c^2} vE$$

Also, the figure suggests  $\vec{B}$  is parallel to  $\vec{v} \times \vec{E}$ :

Thus we can write:  $\vec{B} = \frac{1}{c^2} \vec{v} \times \vec{E}$ , where  $\vec{E} = \gamma \frac{kq_1}{y^2} \hat{n}$  and  $k = \frac{1}{4\pi\epsilon_0}$ 

 $\overline{B}$  depends on the velocity of  $q_1$ 

$$\left| F_{mag} \right| = \frac{v^2}{c^2} \left| F_{elec} \right|$$

Usually  $F_{elec} \gg F_{mag}$  so we cannot easily observe  $F_{mag}$ .

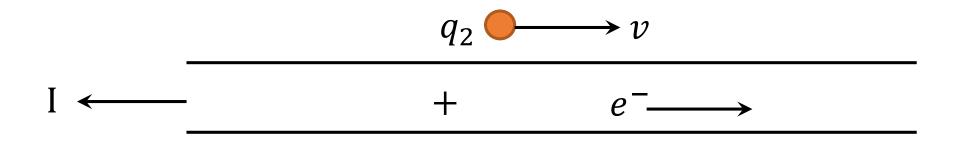
If we get rid of  $F_{elec}$  we should be able to observe  $F_{mag}$ . We can do this by neutralizing with positive charge.

Consider arrangement where  $F_{elec}=0$ , e.g. electron current in metal wire where the positive ions are at rest. Overall the wire is neutral – the positive and negative charges cancel.

$$e^{-} + e^{-} + e^{-} + e^{-} + e^{-} + e^{-}$$

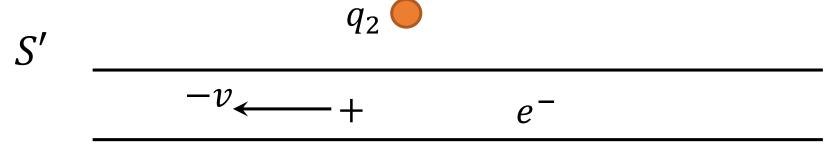
 $e^-$  move with drift velocity, v

Consider a moving test charge, outside the wire, with charge  $q_2$ , and moving || to wire the same velocity, v:



The net  $\overline{E}$  at  $q_2=0$  (the electric fields from the electrons and positive ions cancel).

If the test charge was stationary, it would therefore feel no force. As it is moving, it will feel a magnetic field due to the moving electrons.



Jump to a reference frame S' moving with velocity v w.r.t S. Now  $q_2$  is stationary, the  $e^-$  are stationary, and the positive ions are moving with velocity -v.

In S', the distance between electrons will have increased [by factor  $^{1}/_{\gamma}$ ], but the distance between ions will have decreased [by factor  $\gamma$ ] (using Lorentz contraction).

They are no longer equal and opposite! Have a net charge density of  $D' = D(\gamma - 1/\gamma)$ .

In S', the test charge is stationary, and the force on it, F', depends only on the electric field of both positive and negative charges.

The net force on  $q_2$  is:

$$F_y' = \frac{2kD'}{r}q_2 = \frac{2kD}{r}\left(\gamma - \frac{1}{\gamma}\right)q_2$$

Transform back to *S*:

$$F_{y} = \frac{F_{y}'}{\gamma} = \frac{2kD}{r} \left( 1 - \frac{1}{\gamma^{2}} \right) q_{2} = \frac{v^{2}}{c^{2}} \left( \frac{2kD}{r} \right) q_{2} = \frac{v^{2}}{c^{2}} E q_{2}$$

Since there is no net electric force in S, this force  $F_y$  must be entirely the magnetic force.

So the magnetic field observed in S is related to the electric field observed in S.

The electric field in S' arises from Lorentz contraction effects!

In a typical copper wire, the drift velocity is:

$$v = 6 \times 10^{-2} \text{cm/sec} = 0.6 \text{mm/sec}$$

i.e., 
$$\frac{v}{c} = 2 \times 10^{-12}$$

The magnetic field can be ascribed to the effect of seemingly negligible relativistic contraction effects at low speeds!