



The Galactic Faraday sky

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What it is, how it's done, and why it's useful

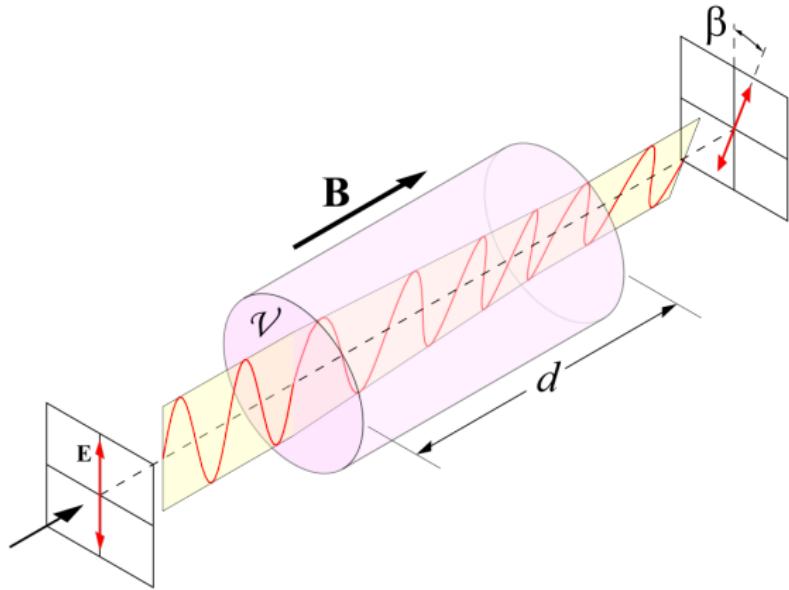
Niels Oppermann

with

G. Robbers, T.A. Enßlin, H. Junklewitz, M.R. Bell, A. Bonafede, R. Braun, J.-A.C. Brown, T.E. Clarke, I.J. Feain, B.M. Gaensler, A. Hammond, L. Harvey-Smith, G. Heald, M. Johnston-Hollitt, U. Klein, P.P. Kronberg, S.A. Mao, N.M. McClure-Griffiths, S.P. O'Sullivan, L. Pratley, T. Robishaw, S. Roy, D.H.F.M. Schnitzeler, C. Sotomayor-Beltran, J. Stevens, J.M. Stil, C. Sunstrum, A. Tanna, A.R. Taylor, and C.L. Van Eck

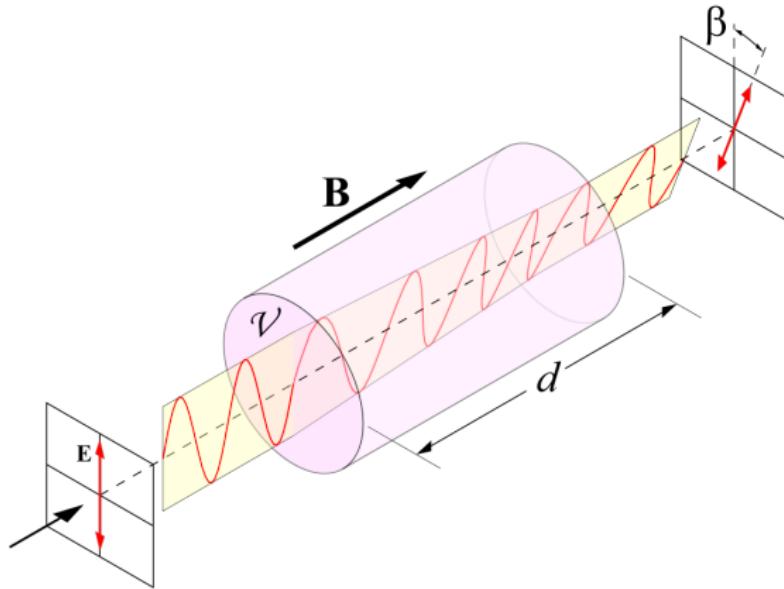
Polarized Foreground Workshop, Garching, 2012-11-26

What it is



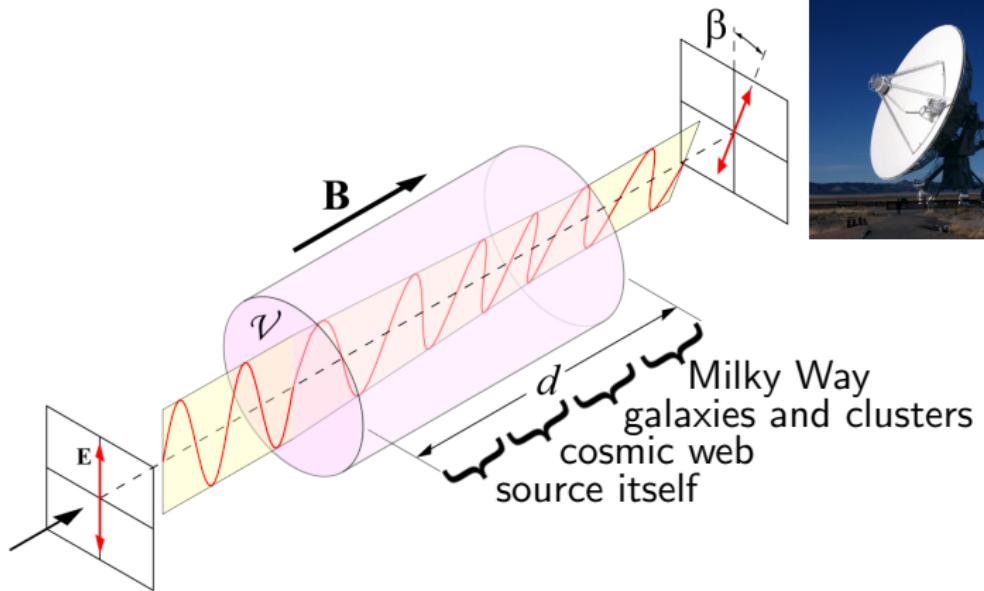
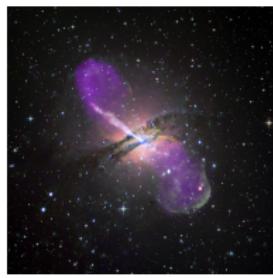
$$d\beta \propto \lambda^2 n_e(\vec{x}) B_r(\vec{x}) dr$$

$$\Rightarrow \beta \propto \lambda^2 \int_{r_{\text{source}}}^0 n_e(\vec{x}) B_r(\vec{x}) dr$$

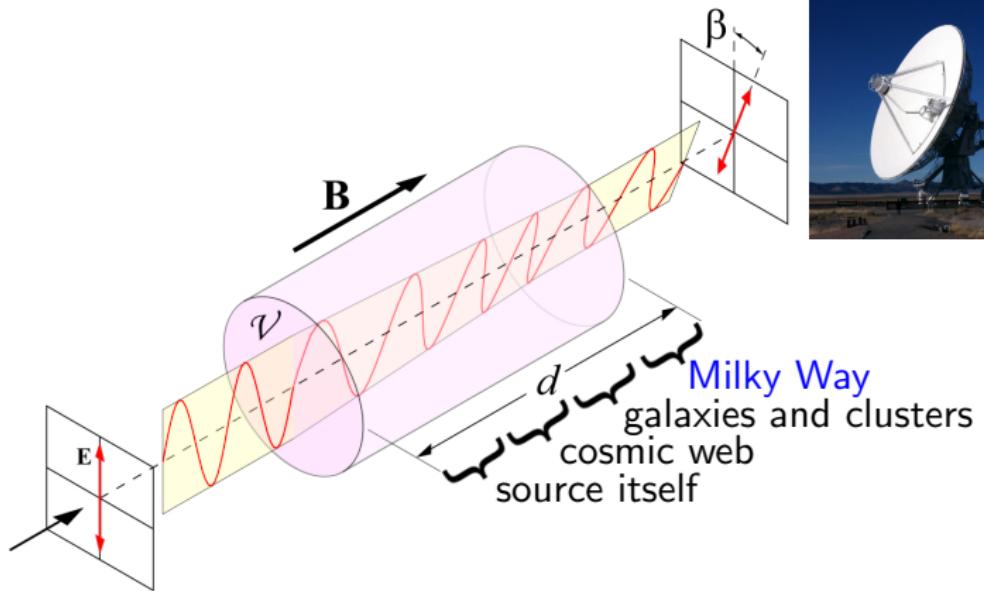
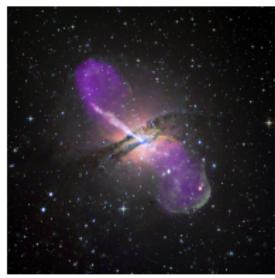


Faraday depth: $\phi \propto \int_{r_{\text{source}}}^0 n_e(\vec{x}) B_r(\vec{x}) dr$

$$\beta = \phi \lambda^2$$

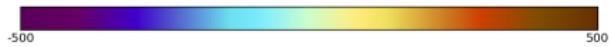
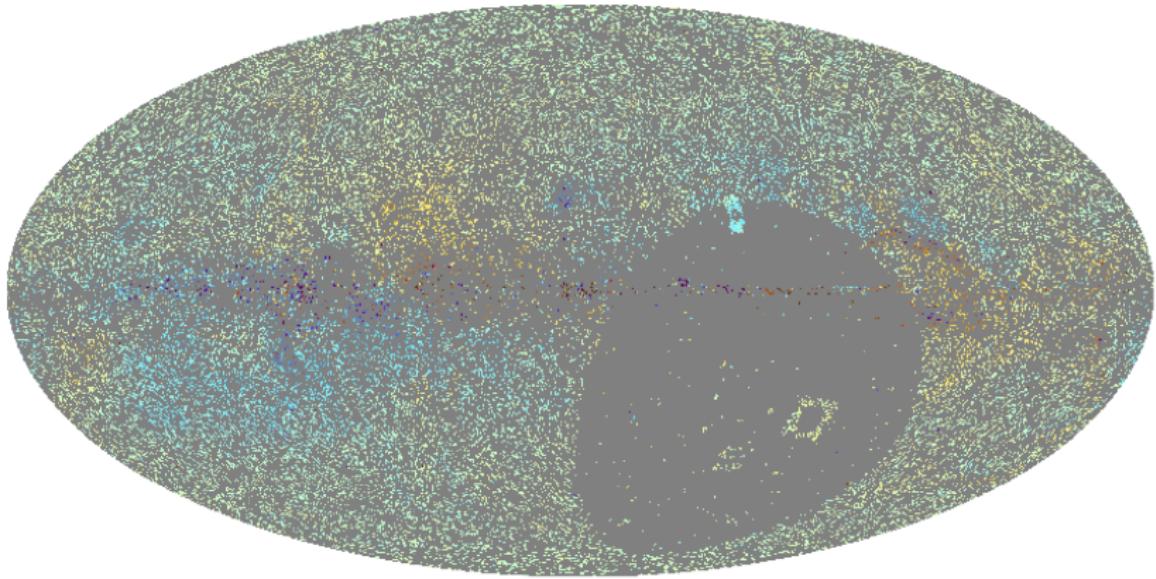


$$\text{Faraday depth: } \phi \propto \int_{r_{\text{source}}}^0 n_e(\vec{x}) B_r(\vec{x}) dr$$
$$\beta = \phi \lambda^2$$

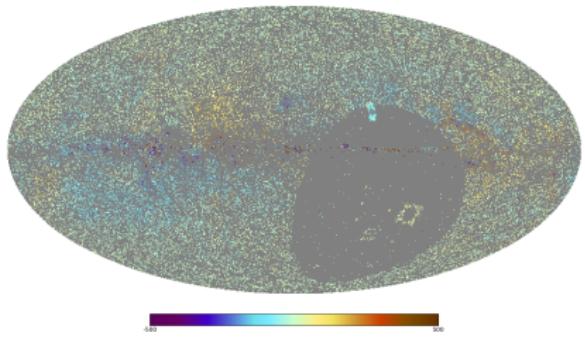


Galactic Faraday depth:

$$\phi \propto \int_{r_{\text{MilkyWay}}}^0 n_e(\vec{x}) B_r(\vec{x}) dr$$



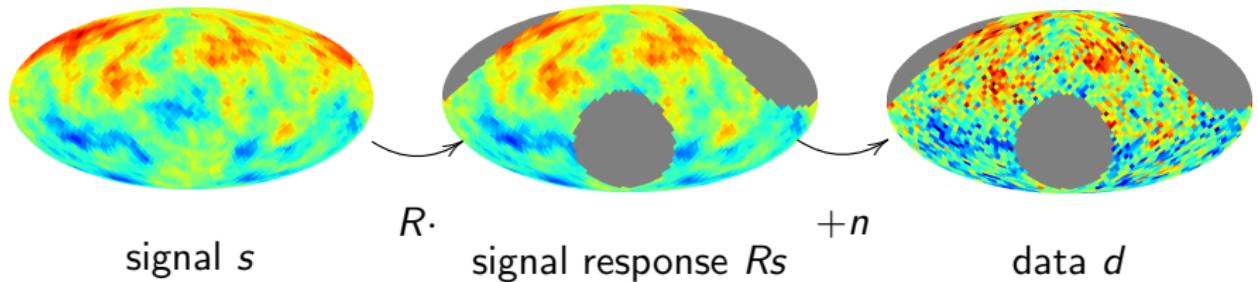
41 330 data points



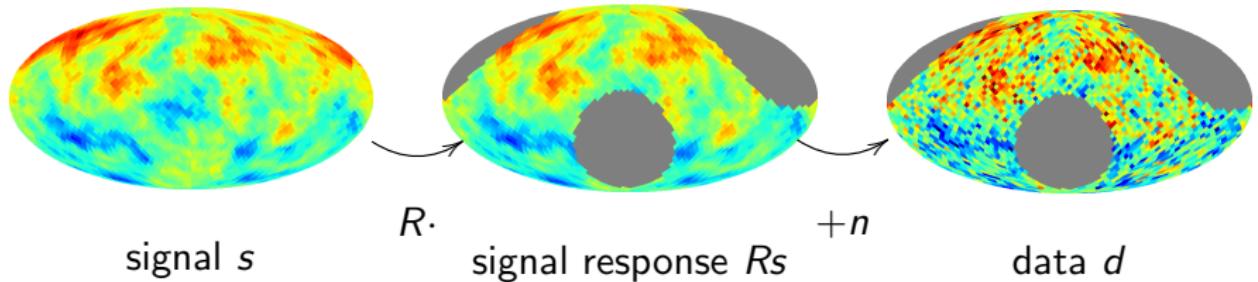
Challenges

- ▶ Regions without data
- ▶ Uncertain error bars:
 - ▶ complicated observations
 - ▶ $n\pi$ -ambiguity
 - ▶ extragalactic contributions unknown

How it's done



$$d = Rs + n$$

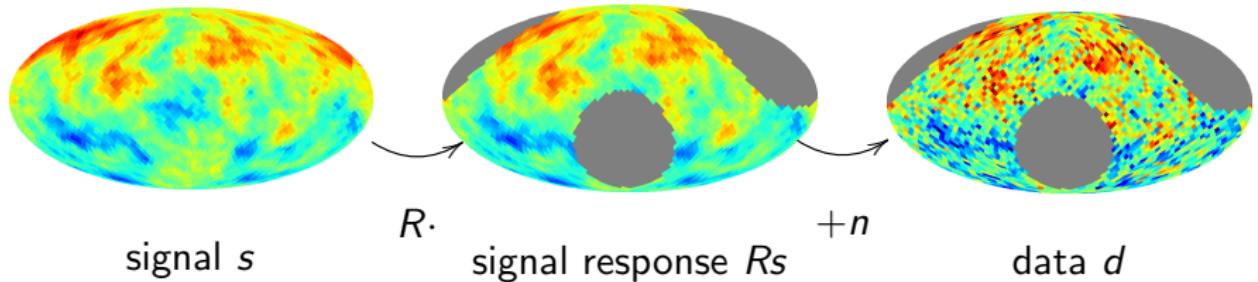


$$\mathcal{P}(s) = \mathcal{G}(s, S)$$

$$\mathcal{P}(n) = \mathcal{G}(n, N)$$

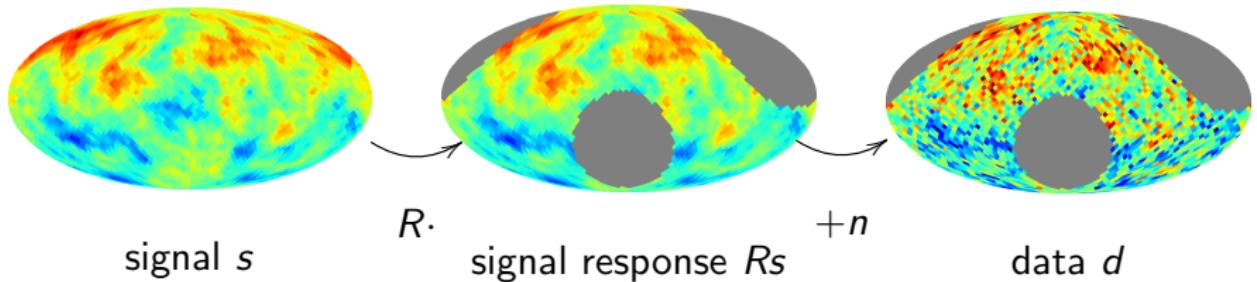
$$d = Rs + n$$

$$\mathcal{G}(s, S) = \frac{1}{|2\pi S|^{1/2}} \exp \left[\frac{1}{2} s^\dagger S^{-1} s \right]$$



$$d = Rs + n$$

$$m = \int \mathcal{D}s \ s \ \mathcal{P}(s|d)$$



Wiener Filter

$$d = Rs + n$$

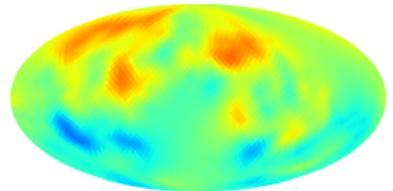
$$m = \int \mathcal{D}s \ s \ \mathcal{P}(s|d)$$

$m = Dj$, where

$$j = R^\dagger N^{-1}d$$

$$D = (S^{-1} + R^\dagger N^{-1} R)^{-1}$$

$$\downarrow DR^\dagger N^{-1}.$$



$$S(\hat{n}, \hat{n}') = \int \mathcal{D}s \ s(\hat{n})s(\hat{n}')\mathcal{P}(s)$$

$$\Rightarrow S_{(\ell m), (\ell' m')} = \int \mathcal{D}s \ s_{\ell m} s_{\ell' m'}^* \mathcal{P}(s)$$

$$\begin{aligned}
S(\hat{n}, \hat{n}') &= \int \mathcal{D}s \ s(\hat{n})s(\hat{n}')\mathcal{P}(s) \\
&= S(\hat{n} \cdot \hat{n}') \\
\Rightarrow S_{(\ell m), (\ell' m')} &= \int \mathcal{D}s \ s_{\ell m} s_{\ell' m'}^* \mathcal{P}(s) \\
&= \delta_{\ell\ell'} \delta_{mm'} C_\ell
\end{aligned}$$

↪ angular power spectrum

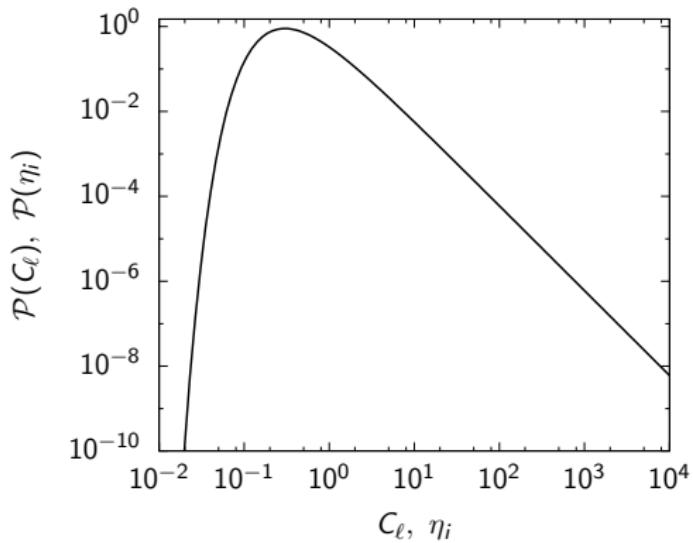
$$\begin{aligned} S(\hat{n}, \hat{n}') &= \int \mathcal{D}s \ s(\hat{n})s(\hat{n}')\mathcal{P}(s) \\ &= S(\hat{n} \cdot \hat{n}') \\ \Rightarrow S_{(\ell m), (\ell' m')} &= \int \mathcal{D}s \ s_{\ell m} s_{\ell' m'}^* \mathcal{P}(s) \\ &= \delta_{\ell \ell'} \delta_{mm'} C_\ell \\ &\hookrightarrow \text{angular power spectrum} \end{aligned}$$

$$N_{ij} = \delta_{ij} \sigma_i^2$$

(uncorrelated noise)

$$S_{(\ell m), (\ell' m')} = \delta_{\ell \ell'} \delta_{mm'} \textcolor{red}{C}_{\ell} \quad N_{ij} = \delta_{ij} \eta_i \sigma_i^2$$

$$S_{(\ell m), (\ell' m')} = \delta_{\ell\ell'} \delta_{mm'} \textcolor{red}{C_\ell} \quad N_{ij} = \delta_{ij} \eta_i \sigma_i^2$$



\Rightarrow marginalize over all possible parameters

Problem: $\mathcal{P}(s|d)$ is non-Gaussian.

Solution: Find Gaussian $\mathcal{G}(s - m, D)$, that best approximates $\mathcal{P}(s|d)$.

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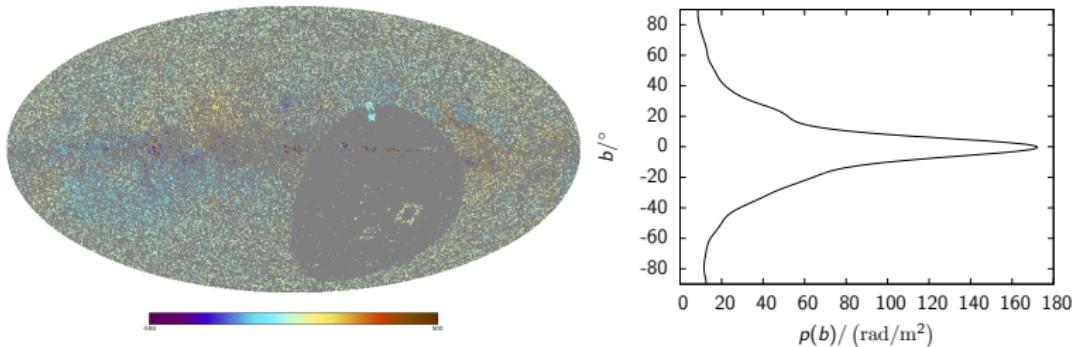
Extended Critical Filter

$$m = Dj, \quad D = \left[\sum_{\ell} C_{\ell}^{-1} S_{\ell}^{-1} + \sum_i \eta_i^{-1} R^{\dagger} N_i^{-1} R \right]^{-1},$$

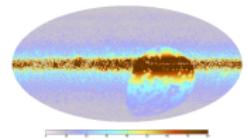
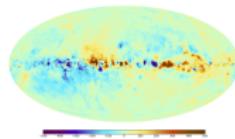
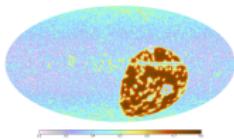
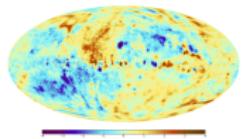
$$j = \sum_i \eta_i^{-1} R^{\dagger} N_i^{-1} d$$

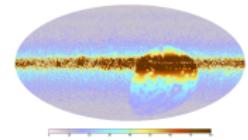
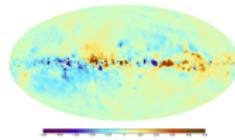
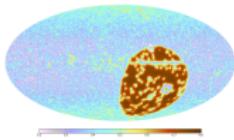
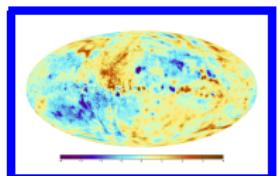
$$C_{\ell} = \frac{1}{\alpha_{\ell} + \ell - 1/2} \left[q_{\ell} + \frac{1}{2} \text{tr} \left(\left(mm^{\dagger} + D \right) S_{\ell}^{-1} \right) \right]$$

$$\eta_i = \frac{1}{\alpha_i} \left[q_i + \frac{1}{2} \text{tr} \left(\left((d - Rm)(d - Rm)^{\dagger} + RDR^{\dagger} \right) N_i^{-1} \right) \right]$$

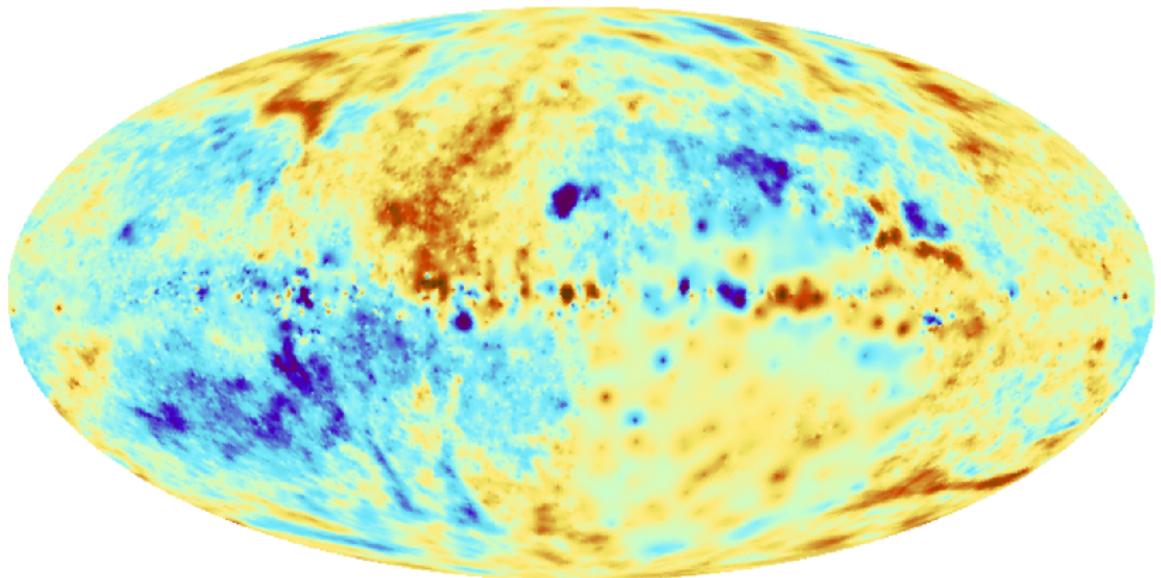


- ▶ Approximate $s(b, l) := \frac{\phi(b, l)}{p(b)}$ as a statistically isotropic Gaussian field
- ▶ R : multiplication with $p(b)$ and projection on directions of sources
- ▶ $N_{ij} = \delta_{ij}\eta_i\sigma_i^2$

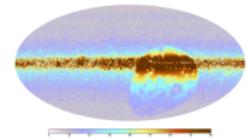
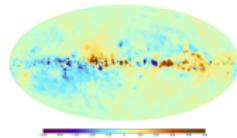
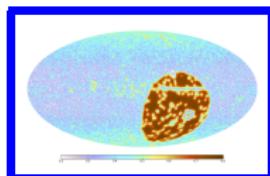
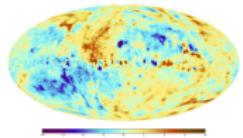




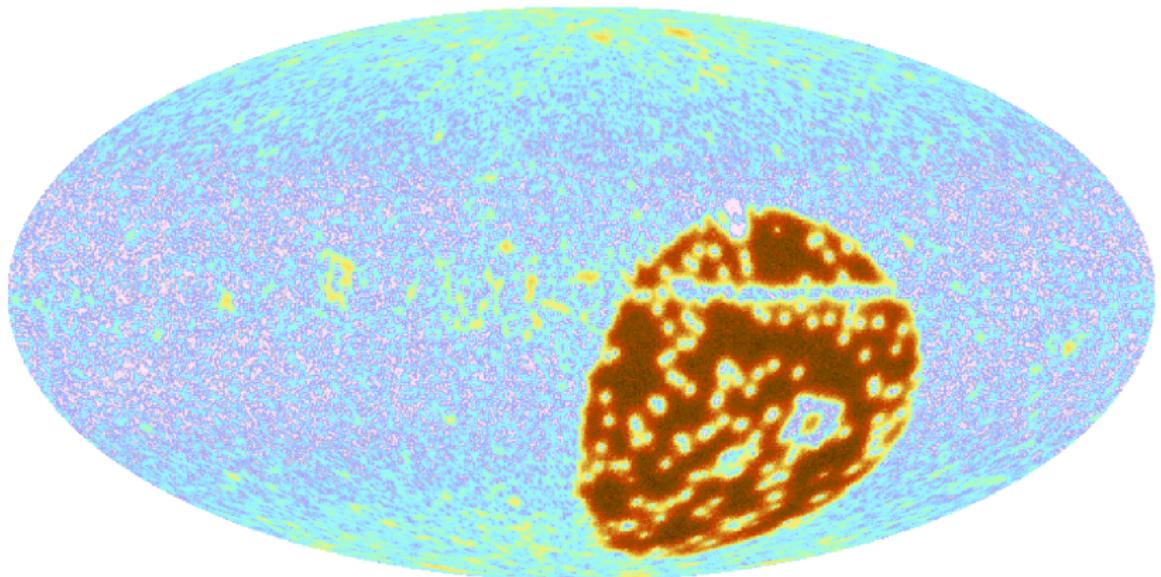
posterior mean of the signal



m

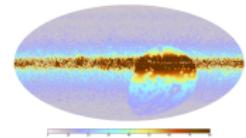
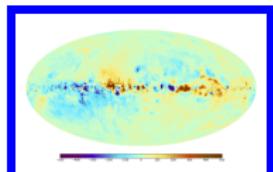
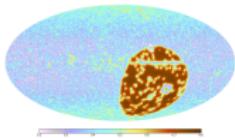
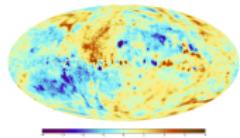


uncertainty of the signal map

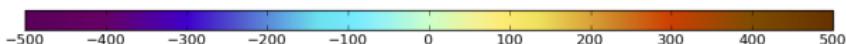
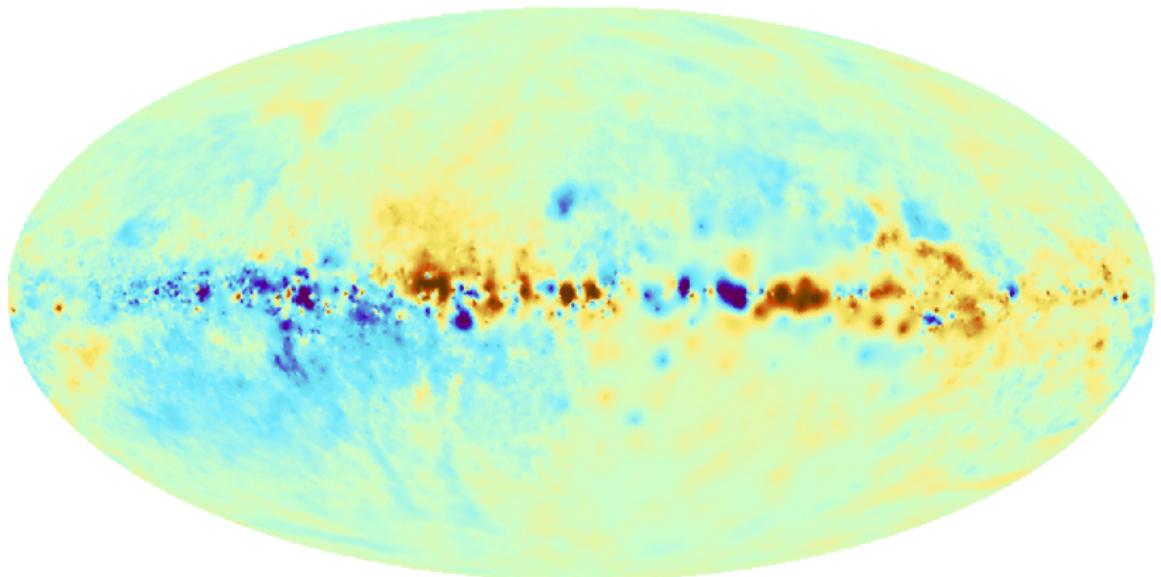


0.2 0.3 0.4 0.5 0.6 0.7 0.8

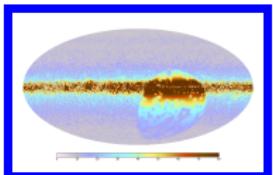
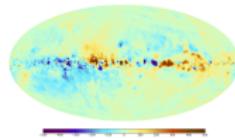
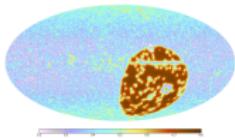
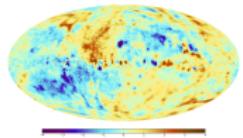
$$\sqrt{\text{diag}(D)}$$



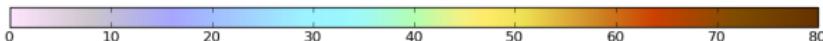
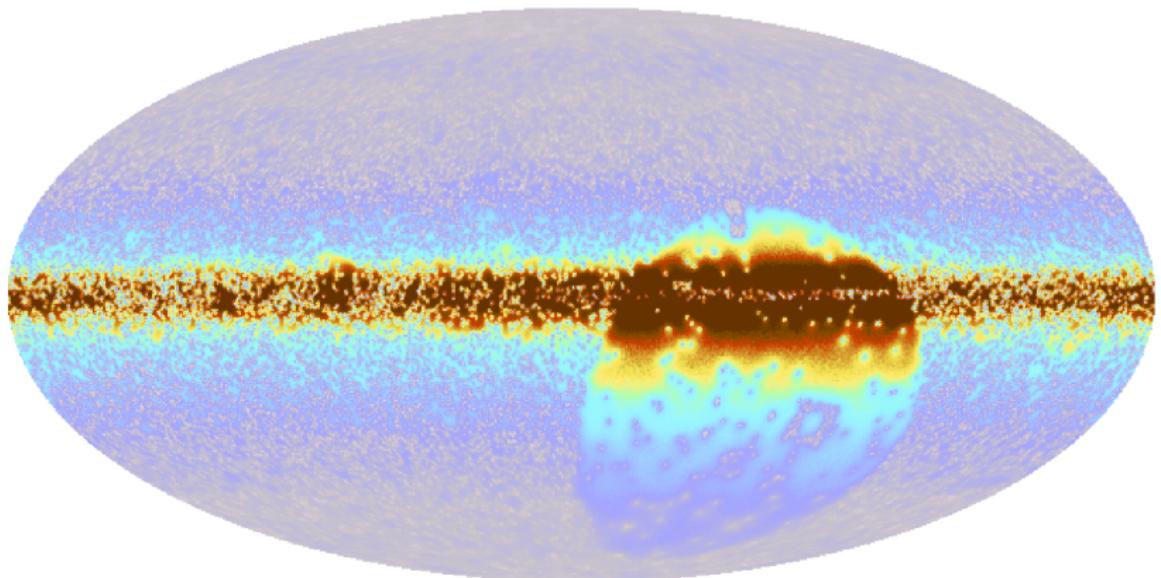
posterior mean of the Faraday depth



$\text{pm}/(\text{rad}/\text{m}^2)$

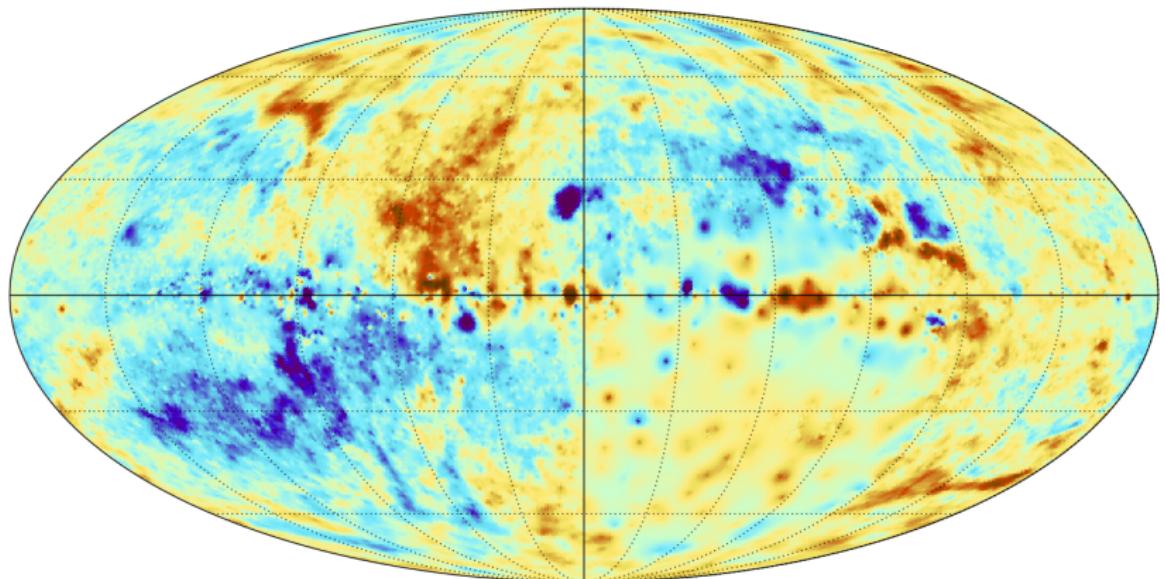


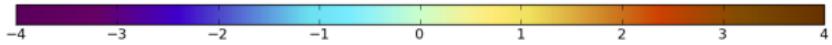
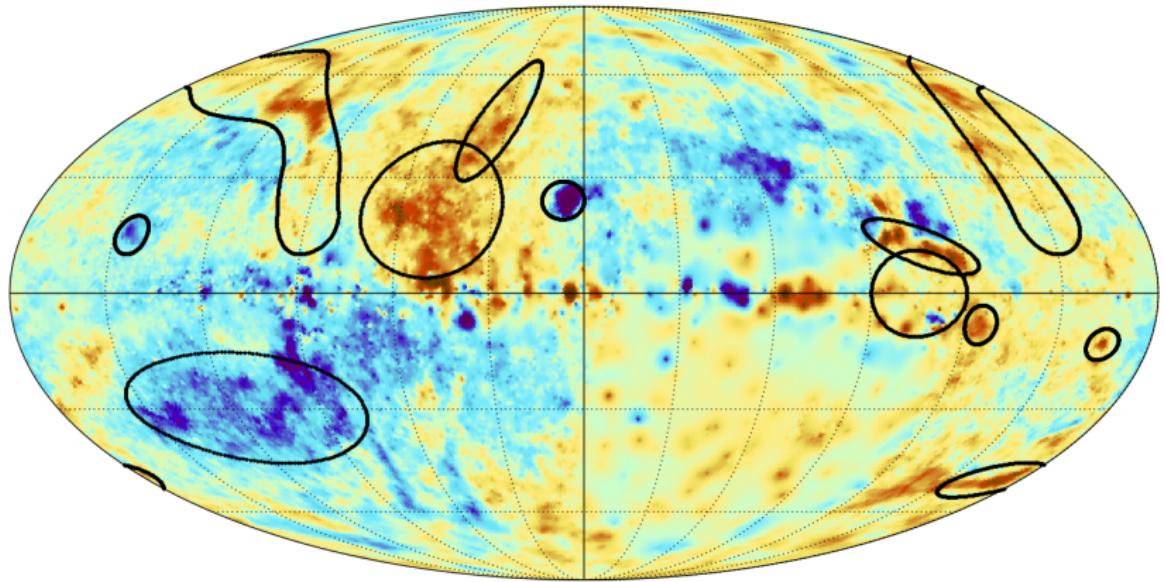
uncertainty of the Faraday depth

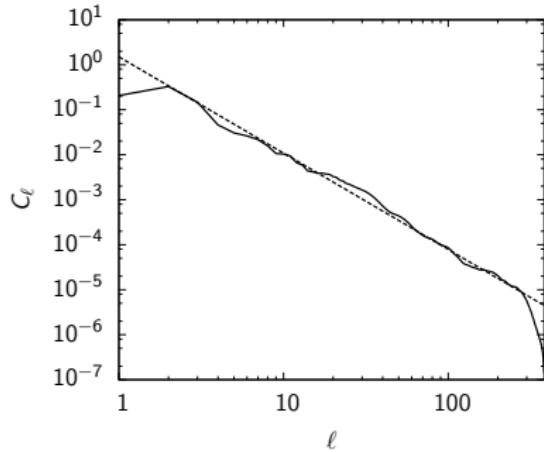


$$p\sqrt{\text{diag}(D)} / (\text{rad/m}^2)$$

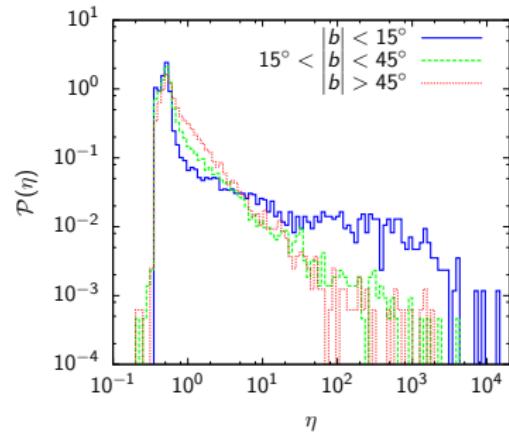
Why it's useful



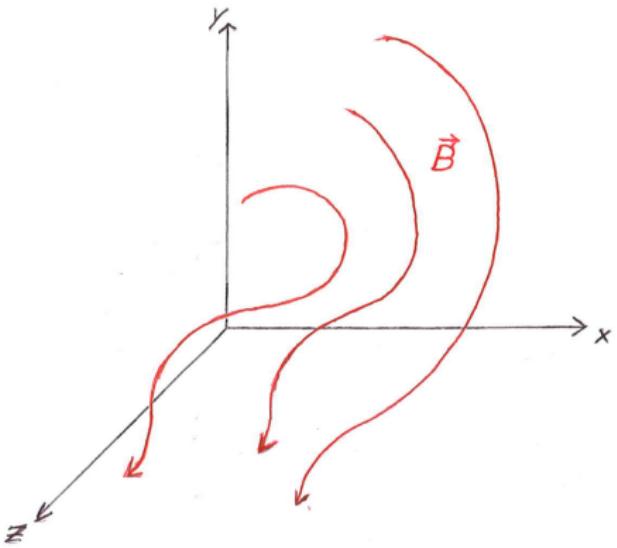


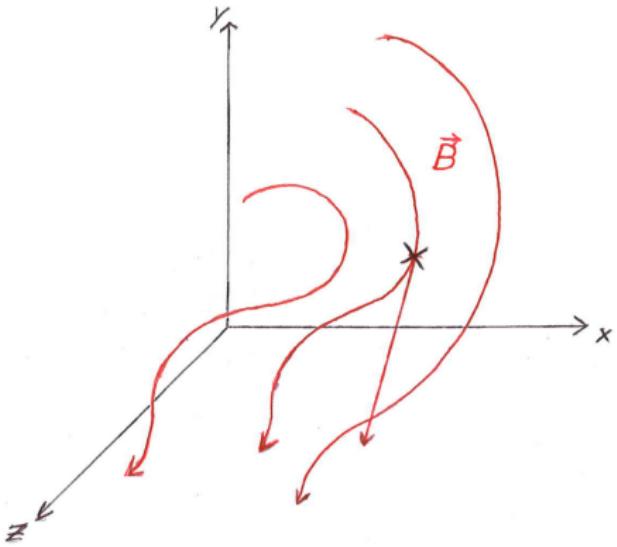


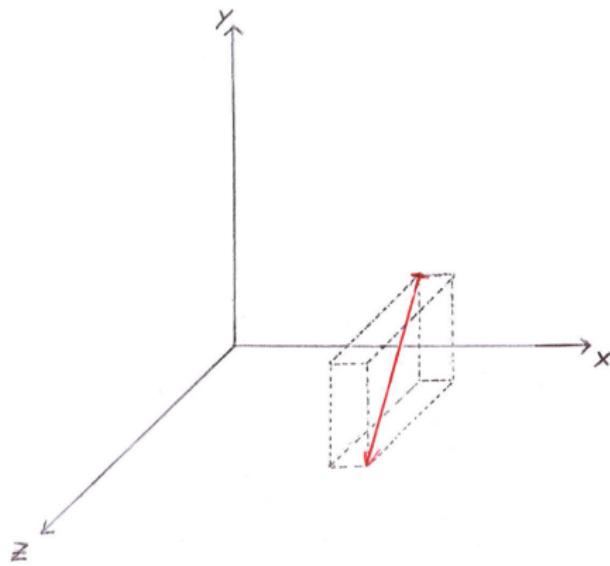
$$C_\ell \propto \ell^{-2.17}$$

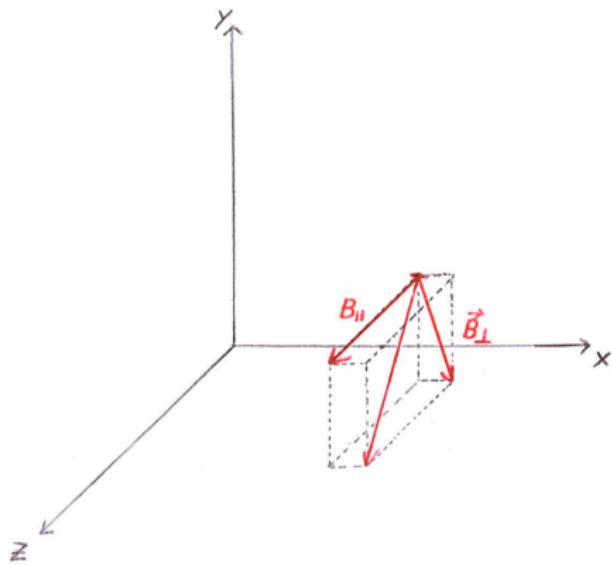


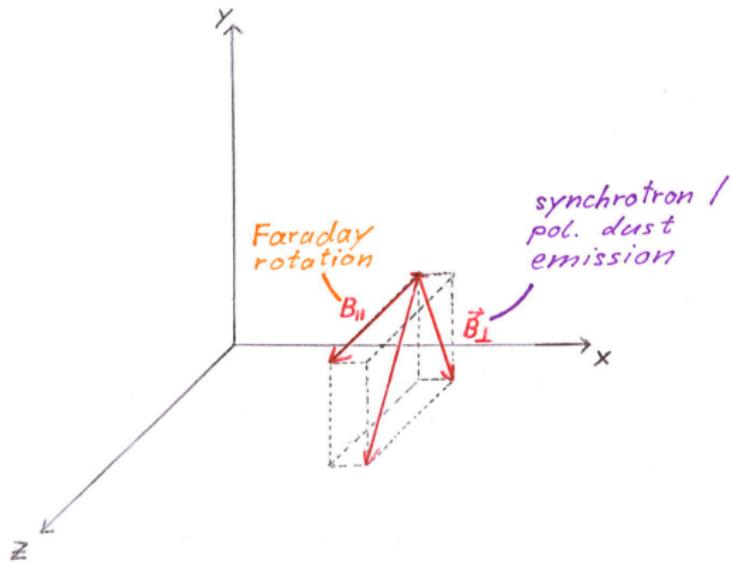
$$N_{ij} = \langle n_i n_j \rangle = \delta_{ij} \eta_i \sigma_i^2$$

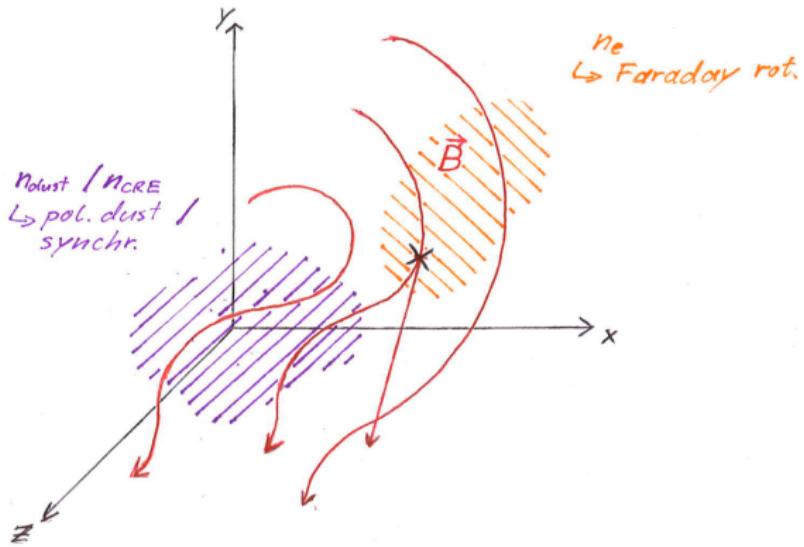












Summary

- ▶ New map of the **Galactic** contribution to Faraday depth
- ▶ Extragalactic contributions filtered out via spatial correlation structure
- ▶ Potential for studies of
 - ▶ Interstellar medium
 - ▶ Galactic magnetic field
 - ▶ Extragalactic sources

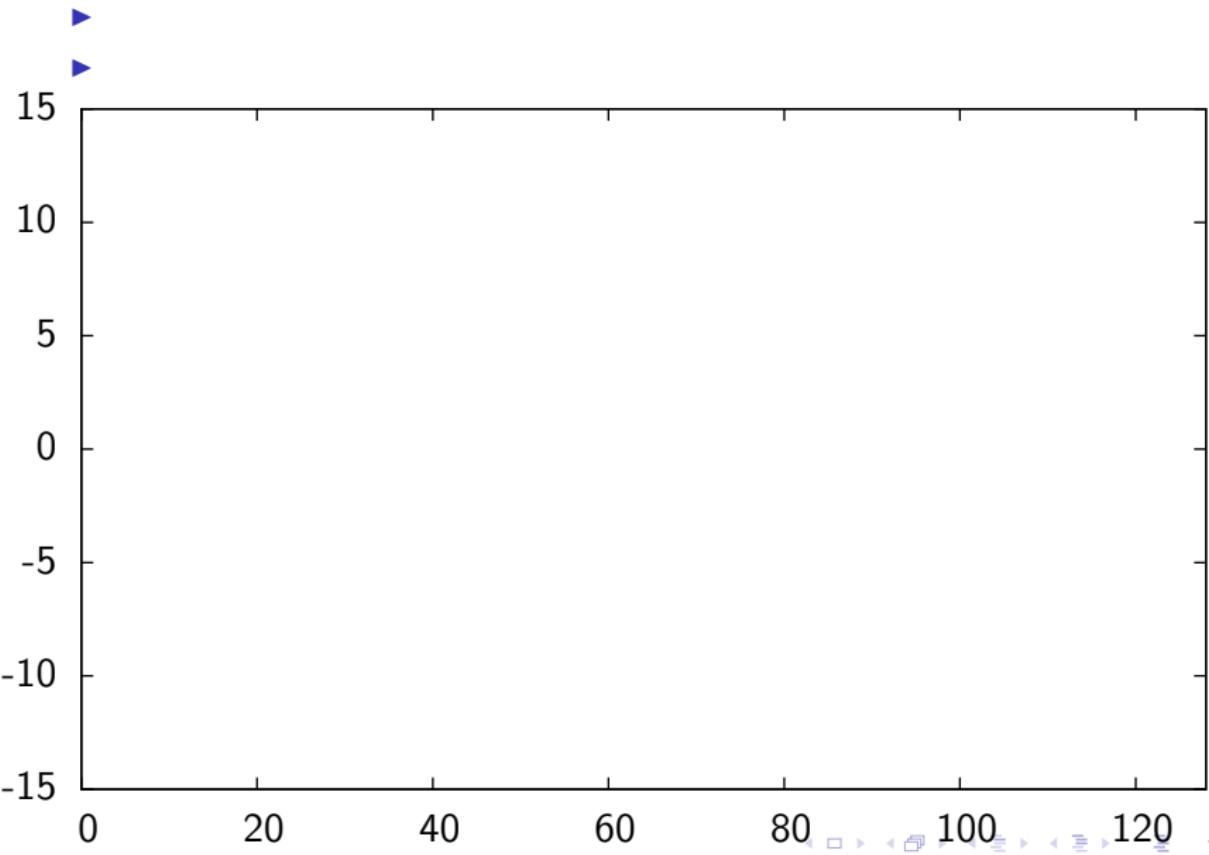
All results available at

<http://www.mpa-garching.mpg.de/ift/faraday/>

Backup

1D test case

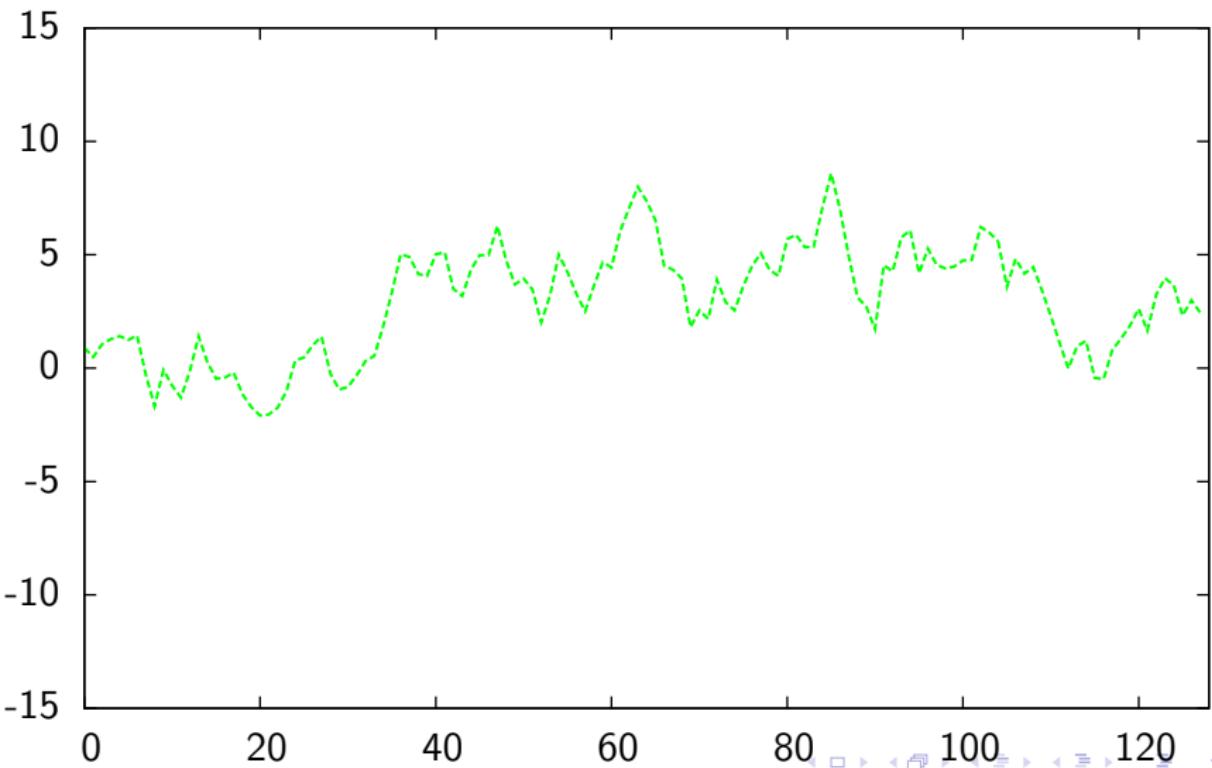
Assumptions:



1D test case

Assumptions:

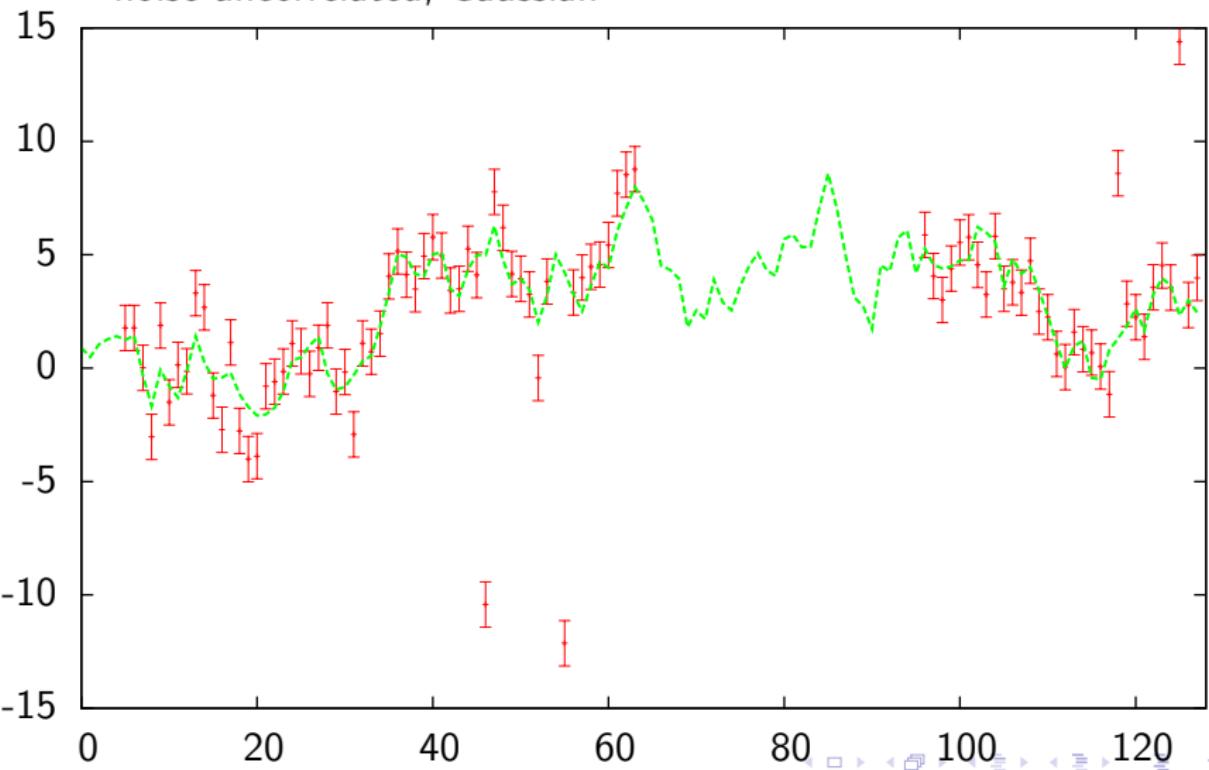
- ▶ signal field statistically homogeneous Gaussian random field
- ▶



1D test case

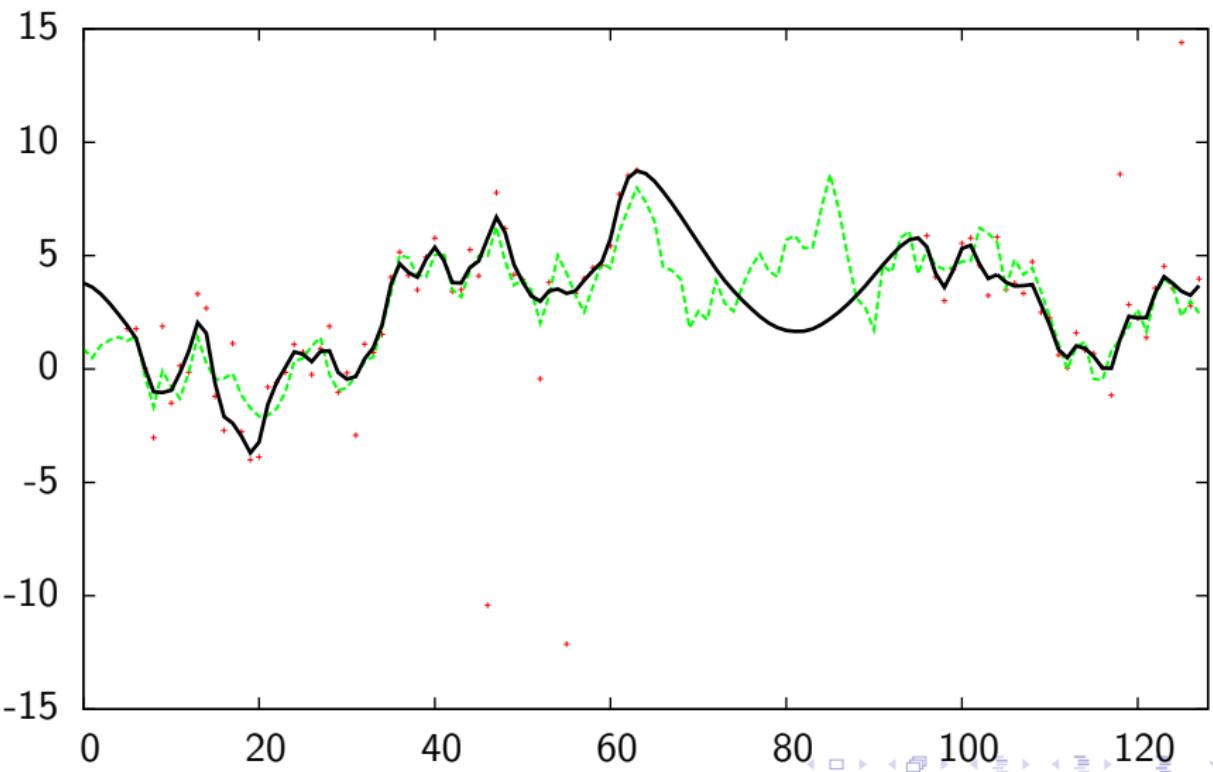
Assumptions:

- ▶ signal field statistically homogeneous Gaussian random field
- ▶ noise uncorrelated, Gaussian



1D test case

- ▶ Reconstruct (iteratively):
signal, power spectrum, noise variance



1D test case

- ▶ Reconstruct (iteratively):
signal, power spectrum, noise variance

