

Magnetic helicity in the interstellar medium

Niels Oppermann



CITA
ICAT

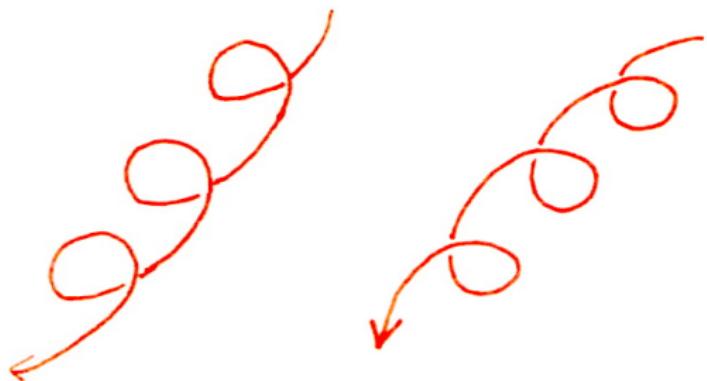
Canadian Institute for
Theoretical Astrophysics
—
L'institut Canadien
d'astrophysique théorique

Workshop on interstellar magnetic fields
IRAP, Toulouse, 2015-04-27

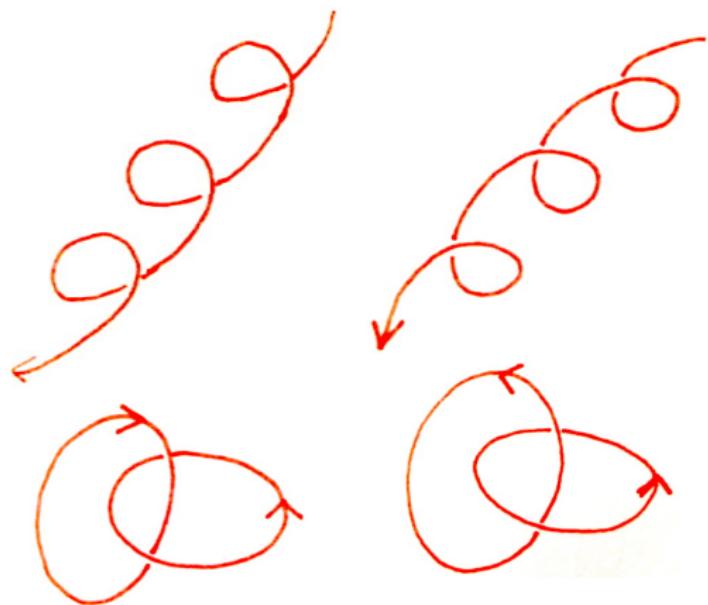
1. What?
2. Why?
3. How?

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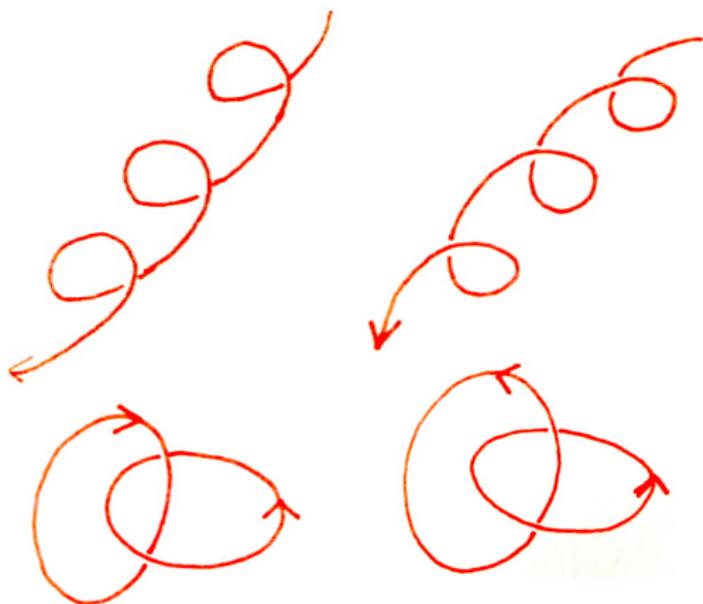
Magnetic helicity – the handedness of the field lines



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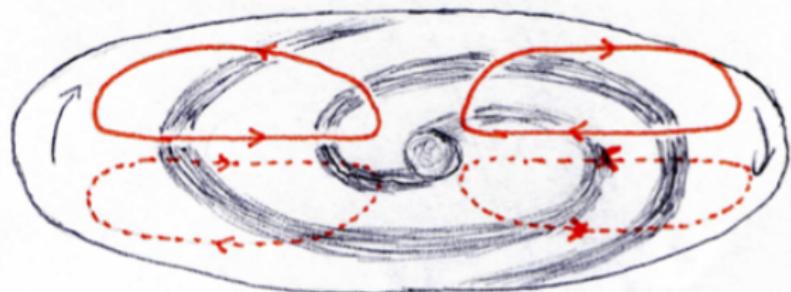
$$H = \int_V d^3x \vec{A} \cdot \vec{B}$$

conserved if:

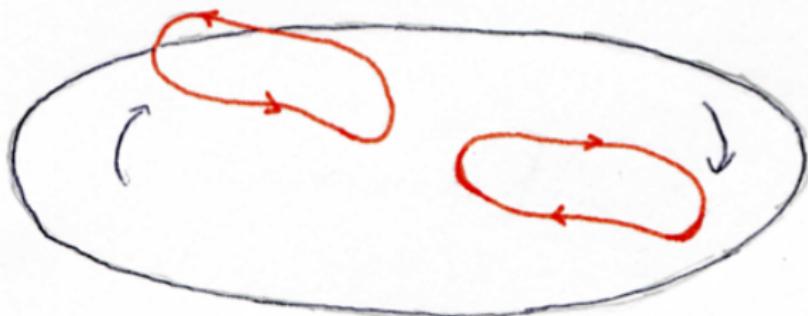
- ▶ conductivity high
- ▶ nothing happens at the surface of V

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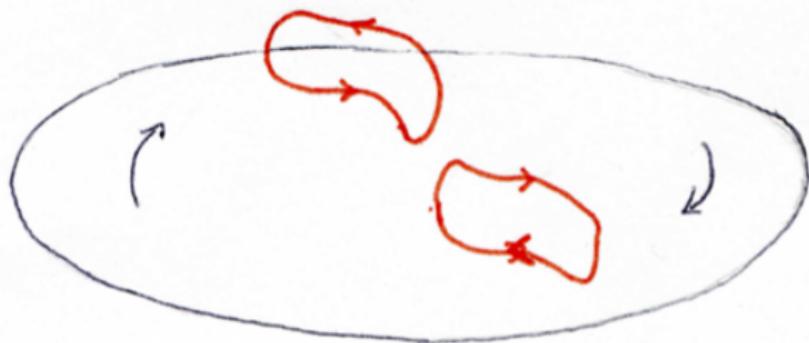
The α - Ω -dynamo



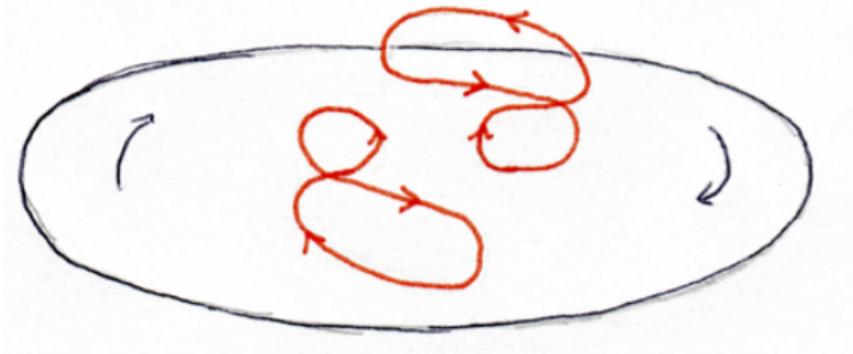
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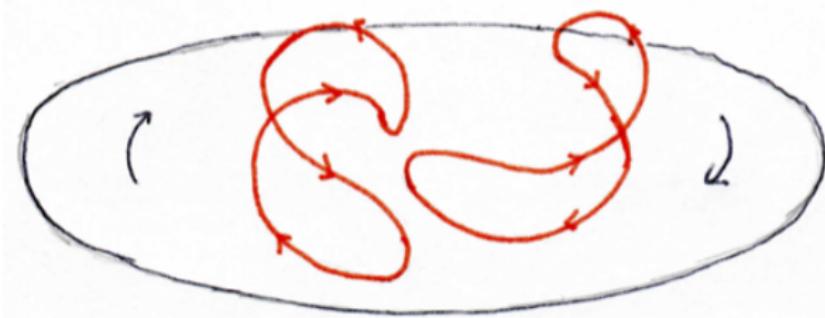
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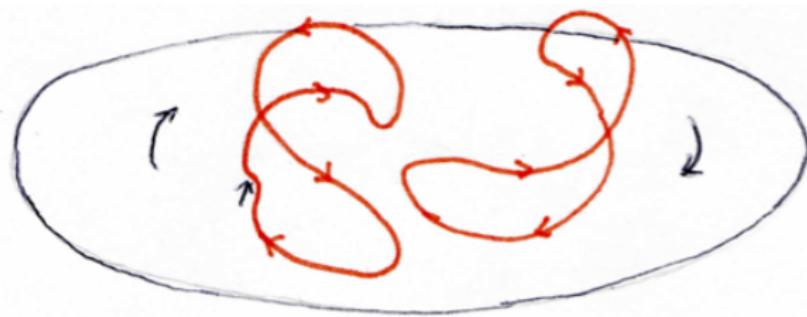
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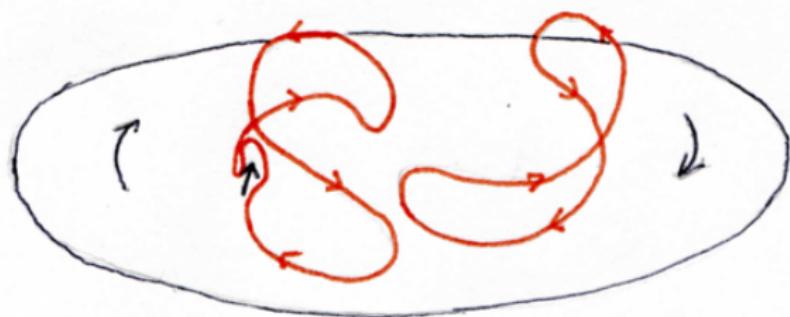
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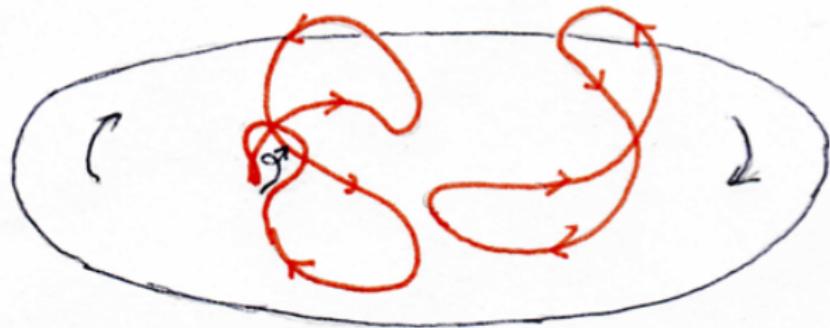
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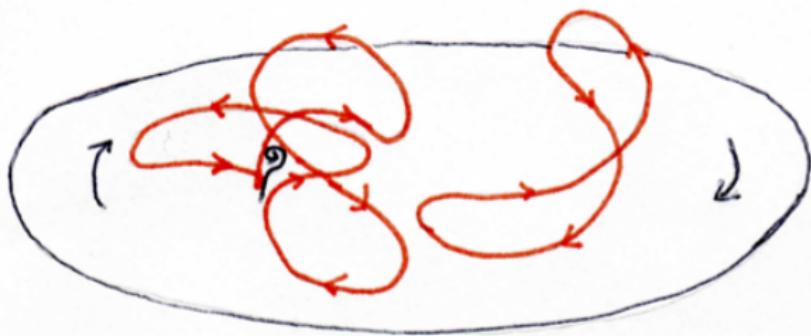
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Ingredients: mean-field theory

induction equation:

$$\frac{\partial \vec{B}}{\partial t} = \vec{\nabla} \times (\vec{v} \times \vec{B}) + \eta \Delta \vec{B}$$

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with $\vec{B} = \langle \vec{B} \rangle + \delta \vec{B}$ and $\vec{v} = \langle \vec{v} \rangle + \delta \vec{v}$:

$$\frac{\partial \langle \vec{B} \rangle}{\partial t} = \underbrace{- \left(\langle \vec{v} \rangle \cdot \vec{\nabla} \right) \langle \vec{B} \rangle}_{\text{advection}} + \underbrace{\left(\langle \vec{B} \rangle \cdot \vec{\nabla} \right) \langle \vec{v} \rangle}_{\text{stretching}} + \vec{\nabla} \times \underbrace{\langle \delta \vec{v} \times \delta \vec{B} \rangle}_{\vec{\mathcal{E}}} + \eta \Delta \langle \vec{B} \rangle$$

large-scale motions

Ingredients: mean-field theory

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“electromotive force”: $\vec{\mathcal{E}} \approx \left[\underbrace{\left\langle \delta \vec{j} \cdot \delta \vec{B} \right\rangle}_{\text{current helicity}} - \underbrace{\left\langle \delta \vec{v} \cdot (\vec{\nabla} \times \delta \vec{v}) \right\rangle}_{\alpha \text{ kinetic helicity}} + \dots \right] \langle \vec{B} \rangle + \dots$

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Ingredients: mean-field theory + helicity conservation

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$$H = \int_V d^3x \vec{A} \cdot \vec{B}$$

$$\frac{\partial H}{\partial t} = -2\eta \int_V d^3x \vec{j} \cdot \vec{B} \approx 0$$

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$$\langle H \rangle = \int_V d^3x \langle \vec{A} \rangle \cdot \langle \vec{B} \rangle$$

$$\frac{\partial \langle H \rangle}{\partial t} \approx +2 \int_V d^3x \vec{\mathcal{E}} \cdot \langle \vec{B} \rangle$$

$$\langle \delta H \rangle = \int_V d^3x \left\langle \delta \vec{A} \cdot \delta \vec{B} \right\rangle$$

$$\frac{\partial \langle \delta H \rangle}{\partial t} \approx -2 \int_V d^3x \vec{\mathcal{E}} \cdot \langle \vec{B} \rangle$$

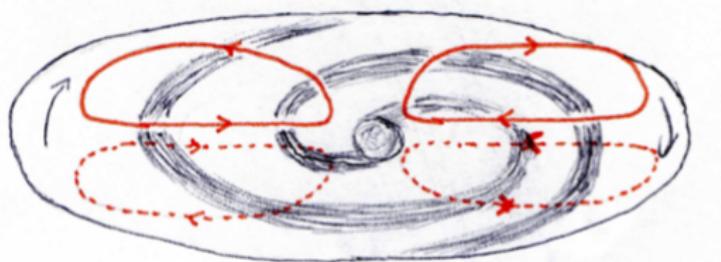
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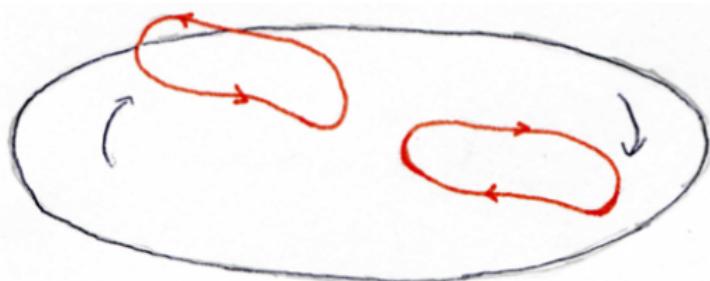
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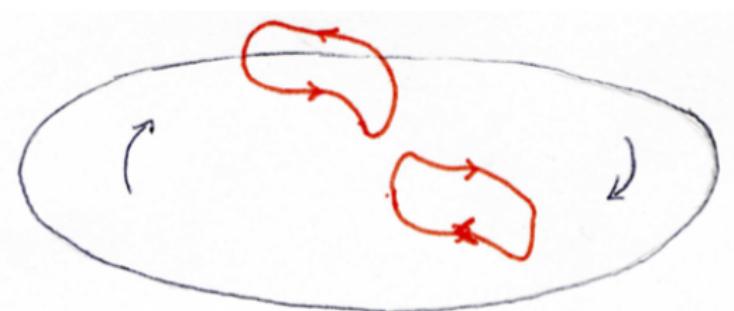
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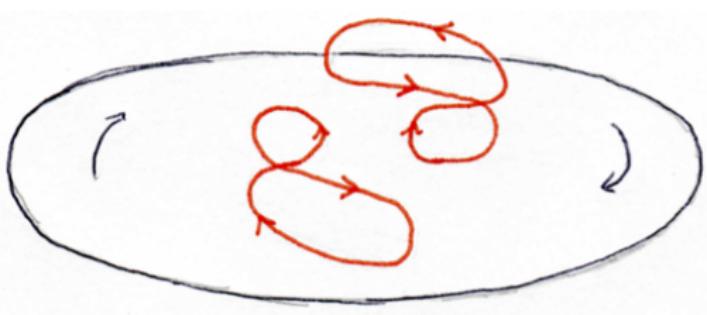
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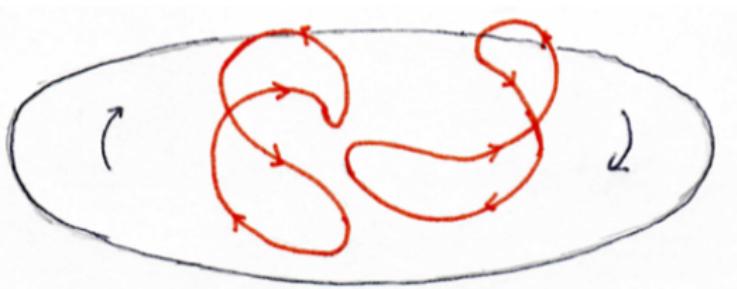
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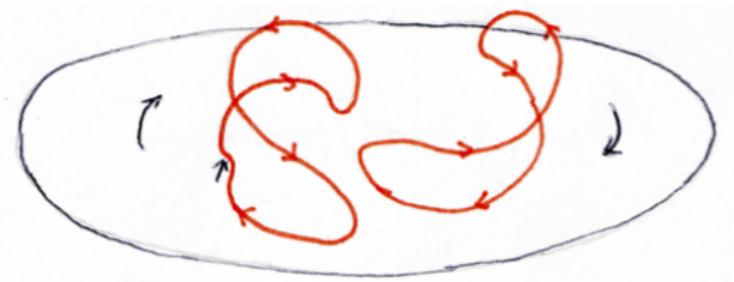
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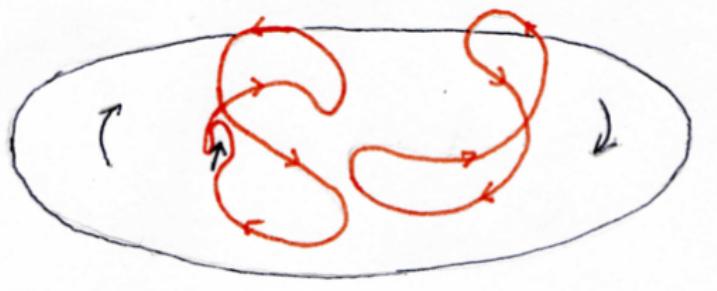
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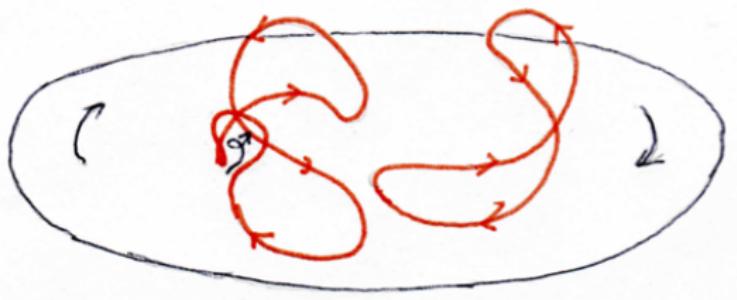
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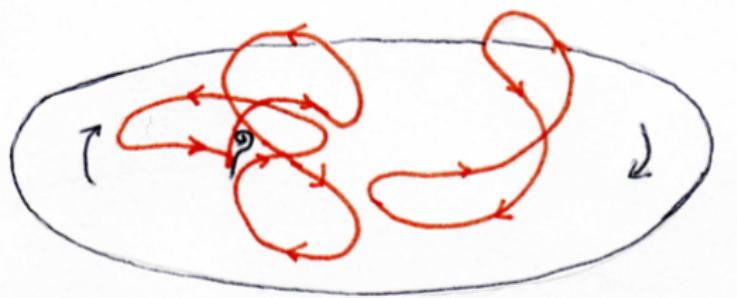
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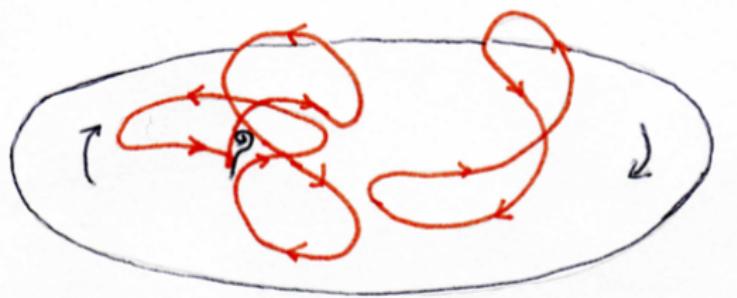
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$$\frac{\partial \langle H \rangle}{\partial t} \approx + 2 \int_V d^3x \vec{\mathcal{E}} \cdot \langle \vec{B} \rangle$$

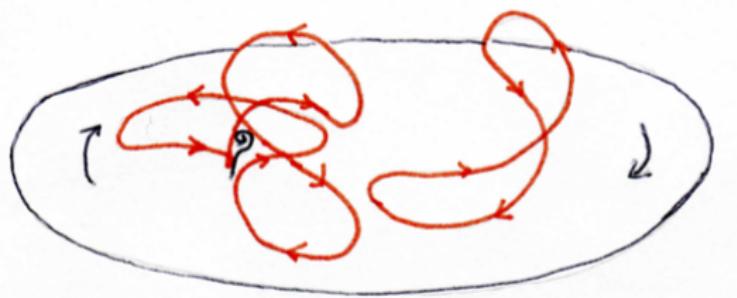
$$\frac{\partial \langle \delta H \rangle}{\partial t} \approx - 2 \int_V d^3x \vec{\mathcal{E}} \cdot \langle \vec{B} \rangle$$

Ingredients: mean-field theory + helicity conservation

$$\frac{\partial \langle \vec{B} \rangle}{\partial t} = \underbrace{- \left(\langle \vec{v} \rangle \cdot \vec{\nabla} \right) \langle \vec{B} \rangle}_{\text{advection}} + \underbrace{\left(\langle \vec{B} \rangle \cdot \vec{\nabla} \right) \langle \vec{v} \rangle}_{\text{stretching}} + \vec{\nabla} \times \underbrace{\left\langle \delta \vec{v} \times \delta \vec{B} \right\rangle}_{\vec{\mathcal{E}}} + \eta \Delta \langle \vec{B} \rangle$$

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problem: Helicity quenches the dynamo.

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problem: Helicity quenches the dynamo.

$$\frac{\partial \langle \delta H \rangle}{\partial t} \approx - 2 \int_V d^3x \vec{\mathcal{E}} \cdot \langle \vec{B} \rangle$$

solution: Move helicity around.

Predictions

- ▶ If kinetic helicity due to cyclonic turbulence:
⇒ opposite signs **above and below the plane**
e.g.: Ferrière (1998), Brandenburg et al. (2014)
- ▶ Due to helicity conservation:
⇒ opposite signs on **small and large scales**
e.g.: Subramanian (2002)
- ▶ If helicity flux dominates kinetic helicity:
⇒ same sign on all scales, but opposite signs above and
below the plane
e.g.: Vishniac et al. (in prep.)

1. What?
2. Why?
3. How?

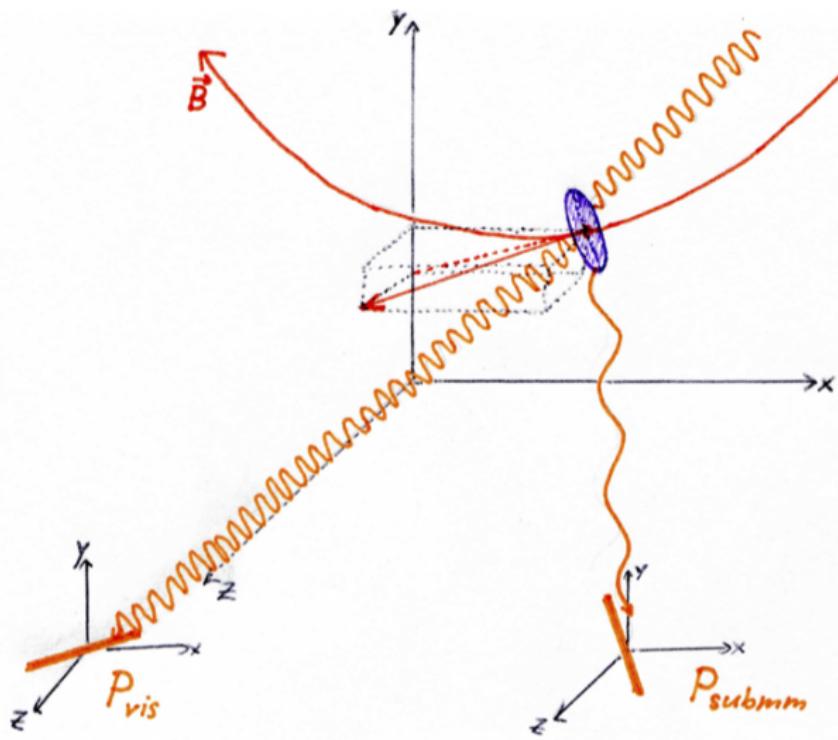
Synchrotron



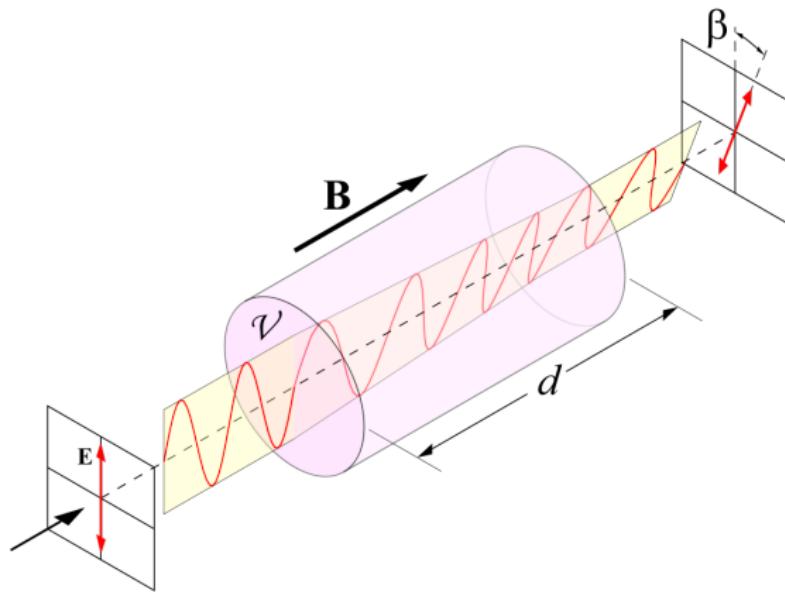
for $n_{\text{CRE}}(E) \propto E^{-\gamma}$:

$$P(\lambda) = Q(\lambda) + iU(\lambda) \propto \lambda^{\frac{\gamma-1}{2}} \int dz n_{\text{CRE}} B_{\perp}^{\frac{\gamma+1}{2}} e^{2i\chi}$$

Dust

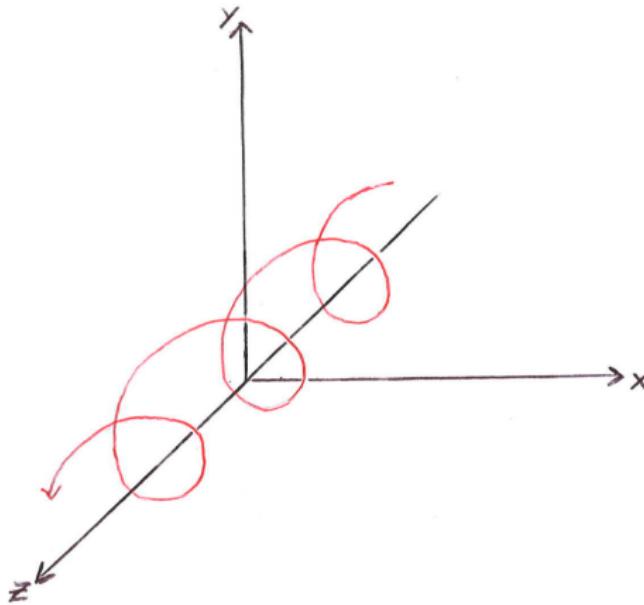


Faraday rotation



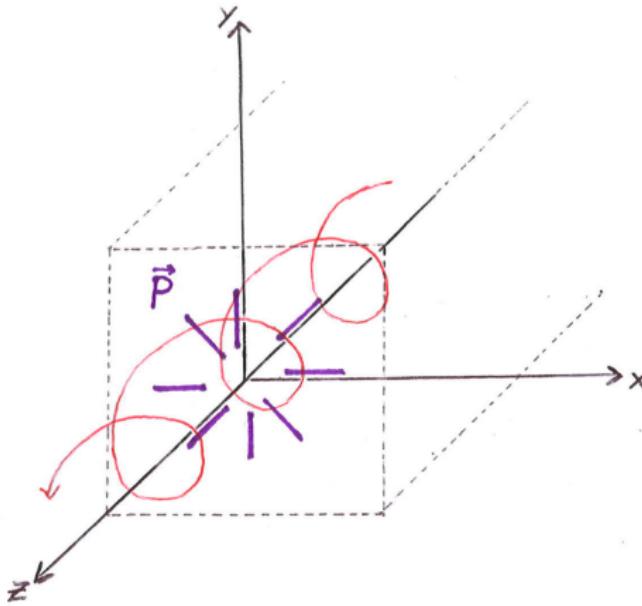
$$\beta \propto \lambda^2 \phi(r_{\text{source}})$$

$$\phi(r_{\text{source}}) \propto \int_{r_{\text{source}}}^0 n_e(\vec{x}) B_r(\vec{x}) \, dr$$



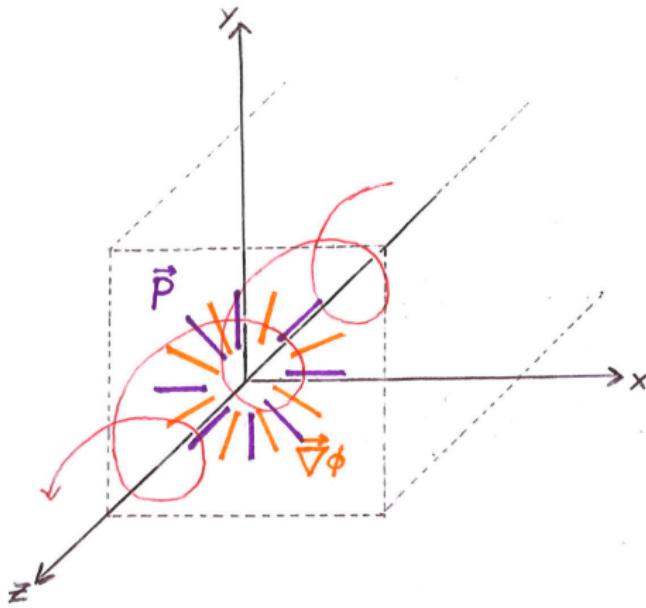
LITMUS procedure

Junklewitz et al. (2011)
Oppermann et al. (2011)



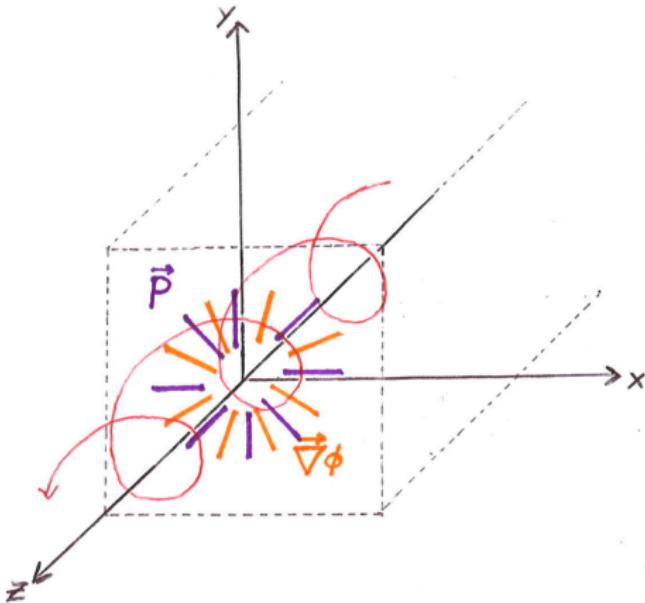
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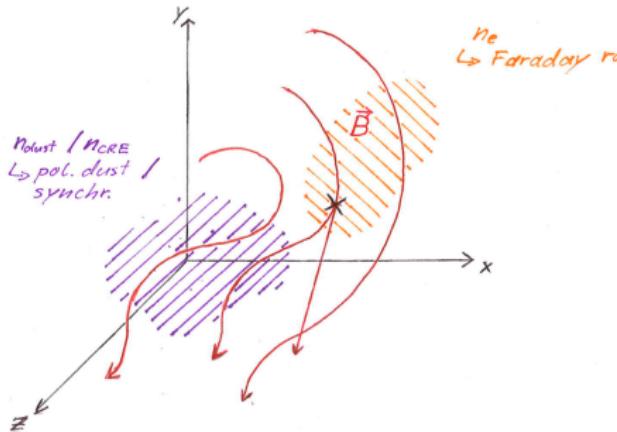


LITMUS procedure

Junklewitz et al. (2011)
Oppermann et al. (2011)



but:

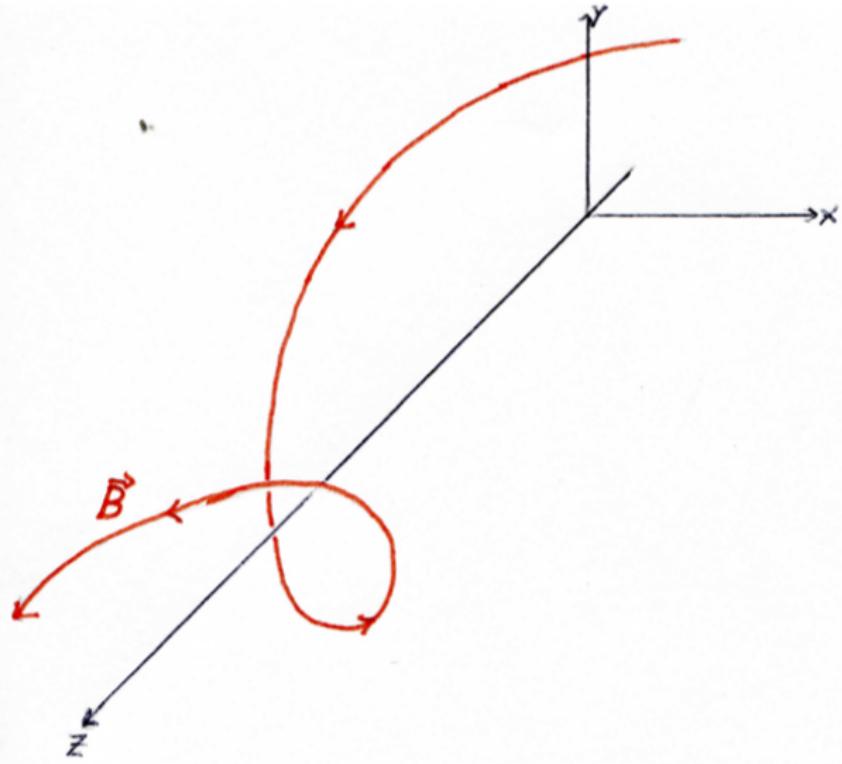


Junklewitz et al. (2011)
Oppermann et al. (2011)

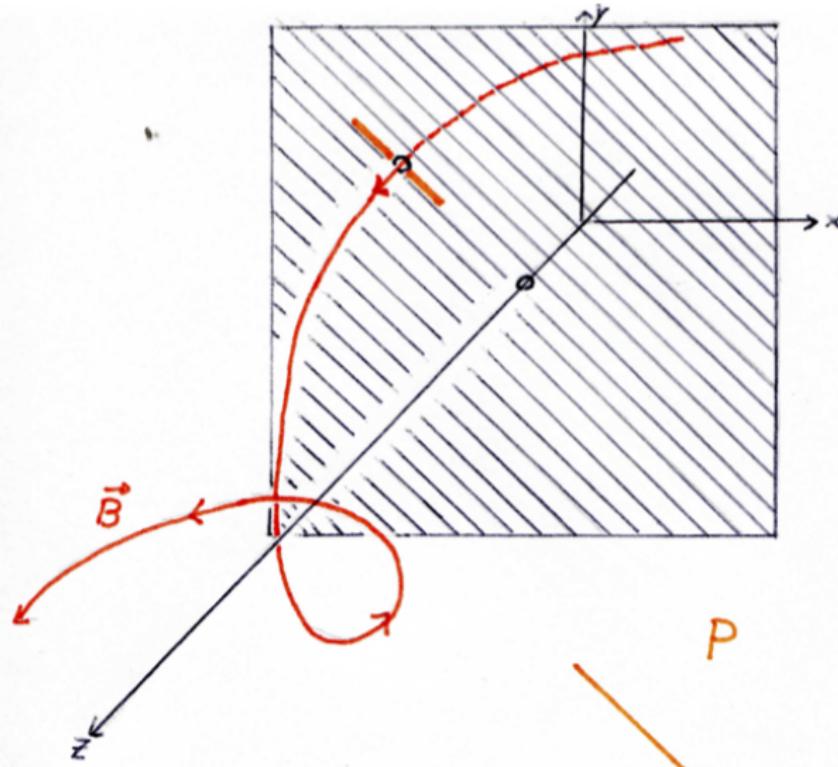
Faraday rotated synchrotron radiation

$$P(\lambda) \propto \int_0^{\infty} dz p(z) e^{2i \lambda^2 \phi(z)}$$

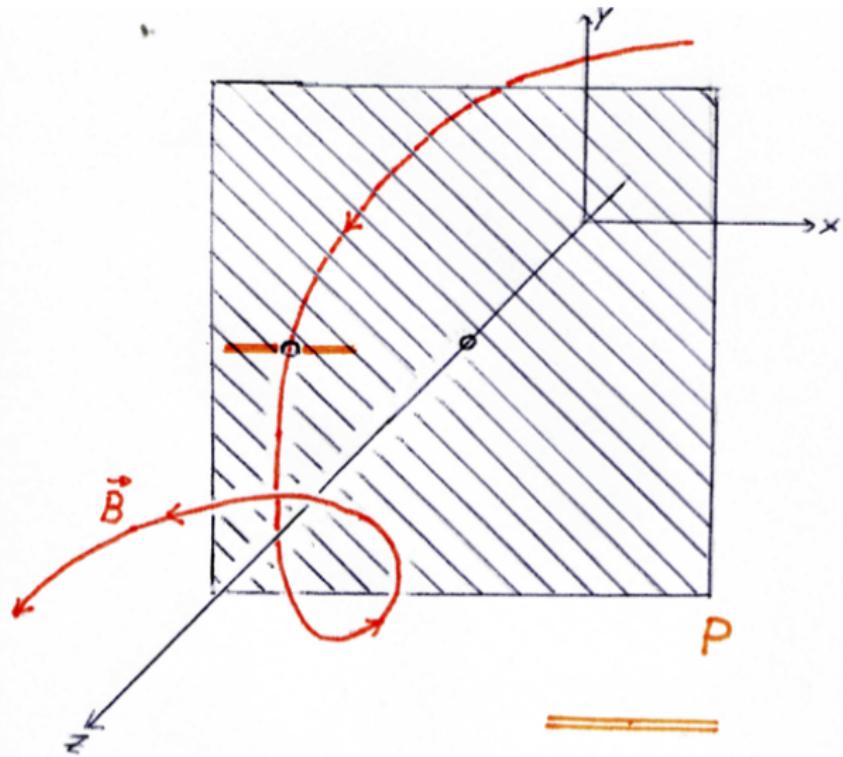
in general: Faraday depolarization



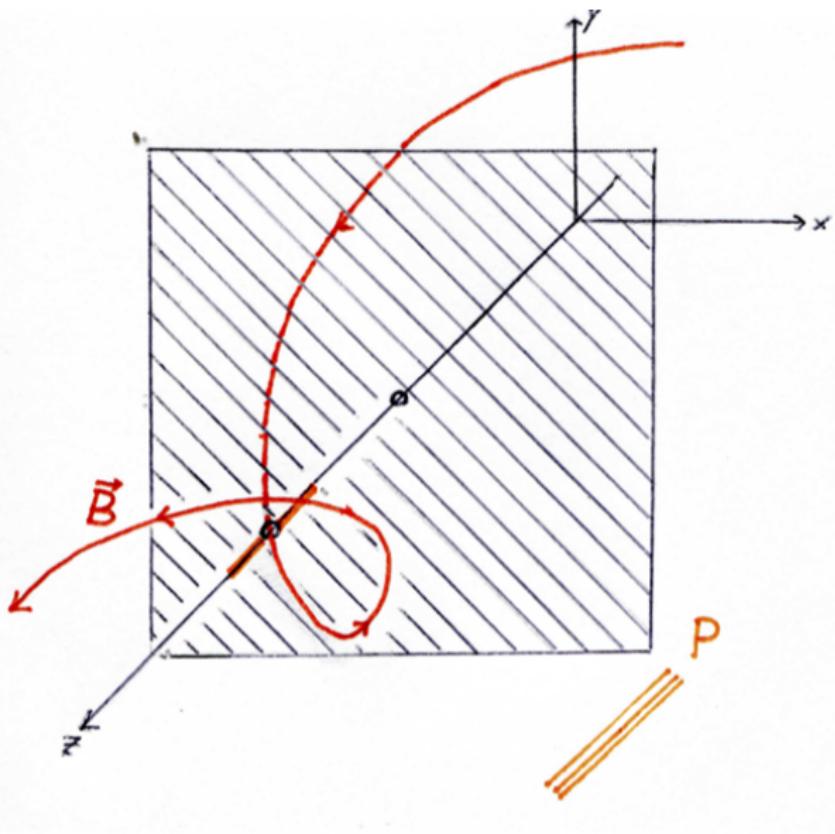
Brandenburg et al. (2014)



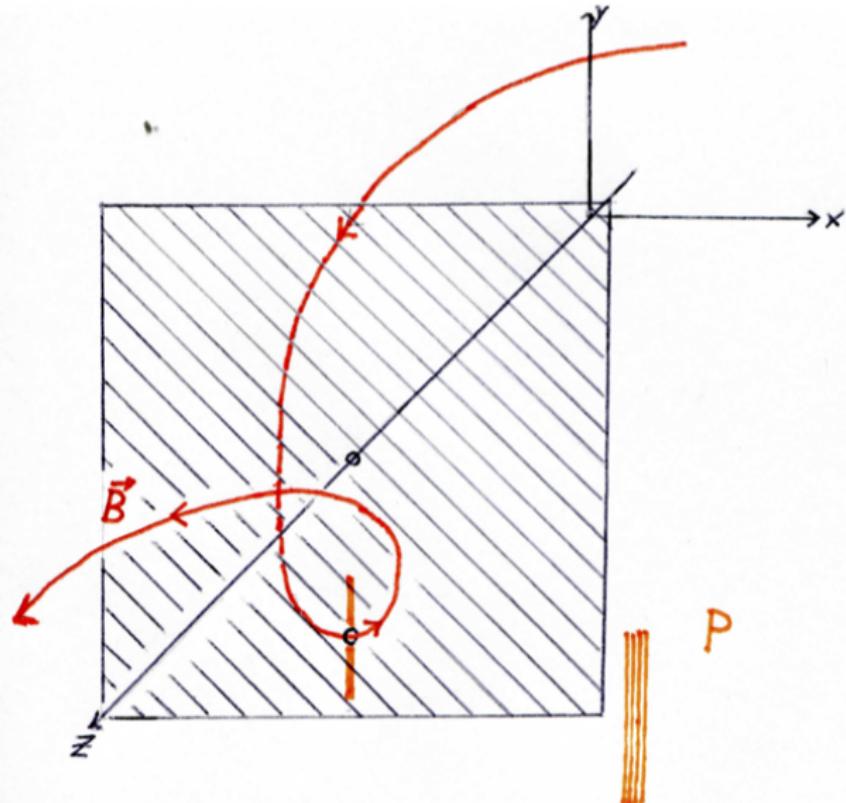
Brandenburg et al. (2014)



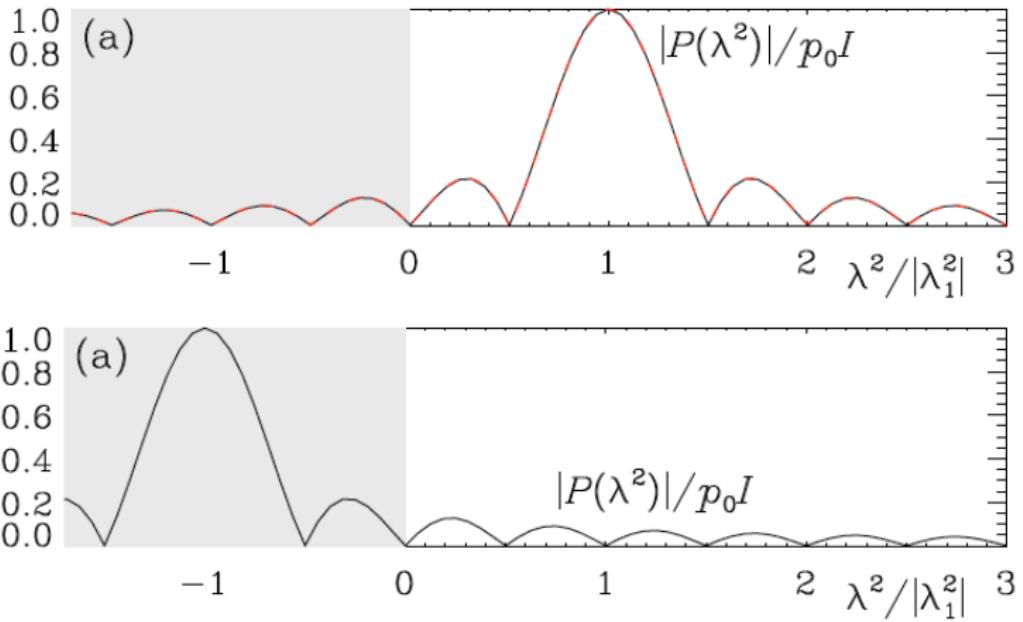
Brandenburg et al. (2014)



Brandenburg et al. (2014)

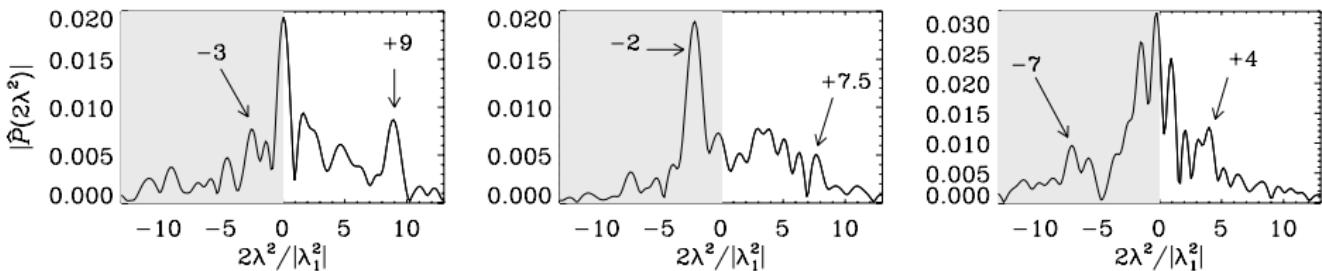


Brandenburg et al. (2014)



Brandenburg et al. (2014)

simulations of helical turbulence



(anti-)correlation between Faraday rotation and polarization degree

Brandenburg et al. (2014)
Volegova et al. (2010)

Magnetic helicity – summary

1. What?

- ▶ twistiness or handedness of magnetic field

2. Why?

- ▶ may tell us if, why, and how the Galactic dynamo works

3. How?

- ▶ need all three B -field components in 3D
- ▶ combine observables
- ▶ details not entirely clear



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