

# An Introduction to Dark Energy

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## 1 Introduction

In 1998, two independent teams of researchers (the Supernova Cosmology Project and the High-Z Supernova Search Team) made the same exciting discovery: the universe is expanding at an accelerating rate. We've known that the universe has been expanding for a while now, thanks to the work of several astronomers (including Edwin Hubble), who found that light from distant galaxies was redshifted.

### Definition

**Redshift**, often denoted by the letter  $z$ , is the phenomenon of light seeming redder if the source is moving away from us with some high velocity  $v$ . You may have heard of Doppler shifts, which are the basis for redshifts. Just like sirens sound longer as an ambulance drives away from you, the wavelength of light gets stretched as the source moves away from you, causing the light to look redder.

The key idea behind the work of both of these research teams (who shared the Nobel prize in 2011) was to use a specific type of supernova, called a **Type Ia** supernova, to measure distances and redshifts to these objects further out in the universe. More on that later!

The exciting thing is that we still use the same method today, although with a lot more precise data, and a better control on the systematics in making these measurements. That means the error bars get smaller, and our conclusions get stronger! Supernovae also play a very important part in the DESI results that came out last month, which is why this project is super exciting and relevant right now, and we might get a lot of questions at our poster. But have no fear, I'll try my best to prepare us as much as possible with this document :-)

## 2 Cosmology 101

Although cosmology is a very developed and complicated field today, it can mostly be described by a few equations. The first of these equations is thanks

to Albert Einstein, and called the “Einstein Field Equations”. Here it is:

$$G_{\mu\nu} = 8\pi T_{\mu\nu} \quad (1)$$

This is beyond the scope of our project, and rather quite a complicated equation to truly understand and solve (despite how simple it looks!). What it basically means is that the curvature/shape of the universe, which is given by this  $G_{\mu\nu}$  thing is related to the stuff/energy inside the universe itself, represented by this  $T_{\mu\nu}$  thing. There are some deep conclusions that can be further explored in a general relativity course, but we don’t have to worry too much about that.

The other equation that really matters, and is somewhat derivative of this Einstein stuff, is the Friedmann equation. Before we jump to it, here’s another definition.

#### Definition

The **scale factor**  $a$  represents the scale of the expanding universe. You can think of it as a variable that describes the age of the universe, or even the size of the universe. Today, at present day, we set the scale factor to  $a = 1$ . When the universe was just born, we set  $a = 0$ . It is very convenient to think of distances between objects in terms of  $a$ , so we have a constant scale for the actual distance between objects at different times in the universe without having to worry about the expansion that is always happening. For example, if two galaxies are 10 Mpc (big unit of distance) apart today, in the past, when  $a = 0.5$ , they would have really physically been  $10a = 5$  Mpc apart. But it’s more convenient for us to think of them as 10 comoving Mpc across at all points in time, because we don’t have to keep switching between numbers. See Figure 1 for an example. Another useful fact that basically comes from wavelengths is that redshift and scale factor are inherently related by  $a = \frac{1}{1+z}$ , so the two quantities are often used interchangeably to describe the universe when it was older (and smaller).

Here’s the Friedmann equation as promised:

$$H^2 = \left(\frac{\dot{a}}{a}\right)^2 = \frac{8\pi G}{3}\rho - \frac{k}{a^2} + \frac{\Lambda}{3} \quad (2)$$

$\dot{a}$  refers to the time derivative of the scale factor  $a$ .  $\rho$  is the density of the universe,  $k$  is the curvature, and  $\Lambda$  is a mysterious cosmological constant.  $G$  is just Newton’s gravitational constant, and  $H = \frac{\dot{a}}{a}$  is something we refer to as the Hubble parameter. The value of  $H$  at  $a = 1$  (i.e. today) is the Hubble constant  $H_0$ , which is the current expansion rate of the universe, i.e.  $\dot{a}$ , which is the time derivative of the scale of the universe - that makes sense!

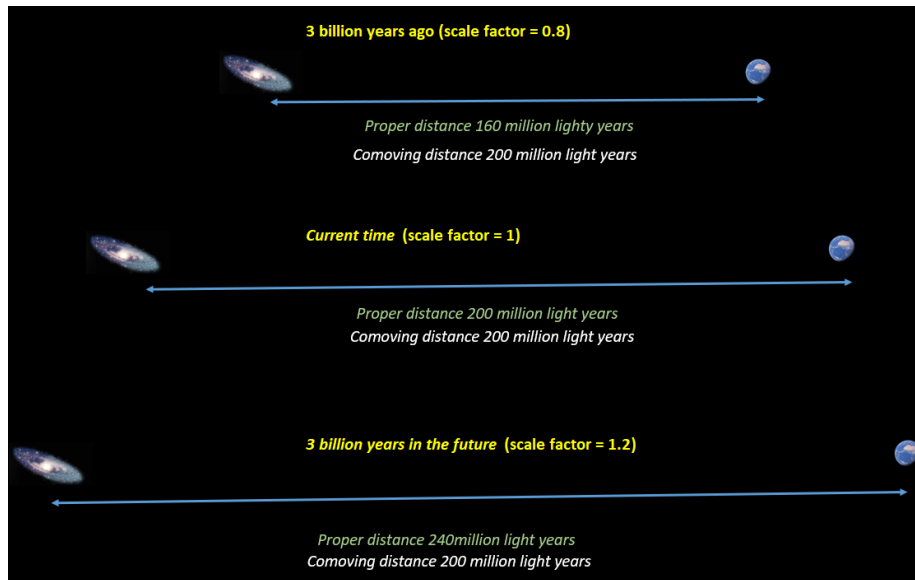


Figure 1: The use of comoving coordinates and scale factors to describe distances

## 2.1 Curvature

As we've discussed, the universe may have some curvature. This is a bit hard to visualize, but imagine how we live on a curved surface, but it looks flat all the time on earth. The same idea can hold in 3D, where the inherent geometry of the universe could also potentially be curved. We parametrize the curvature with this  $k$  number in the Friedmann equation, and it can take on one of three values.

- $k = 0$  : A flat universe, which is nice and easy to work with, like a sheet of paper. The universe is infinite.
- $k = +1$ : A spherical universe, where the universe is closed. Some paths may lead you to the same point again, seemingly parallel lines can intersect, and triangle angles can add to  $> 180^\circ$ . The eventual fate of the universe is to collapse as it closes in on itself with expansion.
- $k = -1$ : A hyperbolic universe, where the universe is open. All paths diverge, triangles can add to  $< 180^\circ$ . The universe does still have infinite space to expand

## 2.2 Cosmological Constant

I mentioned that this  $\Lambda$  term was constant, but mysterious. It actually goes back to the Einstein equations which I told us not to worry about, where he realized (if the universe wasn't expanding, which was the belief at the time)

there needed to be some constant term (on either side of the equation) so that the universe wouldn't just collapse on itself under the force of gravity, and can stay 'static'. However, people found out quickly that the universe is indeed expanding. It turns out that if you keep this constant term with an expanding universe, then the expansion gets faster over time, which could be a nice simple explanation for an accelerating expansion that we measure today.

Here's a nice way to think about it. Let's say that empty space (vacuum) has some energy associated with it. If you expand space (i.e. add more vacuum), you add more vacuum energy, and this leads to even more energy to cause the expansion, so the expansion gets faster and faster as we add more and more space! But why do we call it a constant? The name comes from the role  $\Lambda$  plays in the Einstein Field Equations. But for practical purposes it's constant in terms of energy density. If you look at the units of the Friedmann equation (we can ignore the units of  $G$ ),  $\Lambda$  has essentially the same units as energy density, so the **density** is constant.

## 2.3 Components of the Universe

We can do some reframing of the Friedmann equation by defining a critical density  $\rho_c$

$$\rho_c = \frac{3H^2}{8\pi G} \quad (3)$$

If the universe is more dense than  $\rho_c$ , and we assume  $\Lambda = 0$ , then the universe would be too dense, and collapse in on itself, like a closed universe. Likewise, if  $\rho < \rho_c$ , it would expand forever. Another way of thinking of critical density is solving for exactly when  $k = 0$  in the Friedmann equation.

We can define different density parameters  $\Omega$  for objects by dividing their densities by the critical density.

$$\Omega_i = \frac{\rho_i}{\rho_c} = \frac{8\pi G \rho_i}{3H^2} \quad (4)$$

With some work, we can rewrite the Friedmann equation as

$$\frac{H^2}{H_0^2} = \Omega_R a^{-4} + \Omega_M a^{-3} + \Omega_k a^{-2} + \Omega_\Lambda \quad (5)$$

where  $\Omega_R$  is the radiation energy density parameter,  $\Omega_M$  is the matter density parameter,  $\Omega_k$  is the curvature density parameter, and  $\Omega_\Lambda$  is the dark energy density parameter today. All quantities are their values today. In practice, we can let  $\Omega_R$  and  $\Omega_k$  go to 0 because they are just super tiny.

The matter density should make sense, as matter density goes like mass/volume, and volume is exactly like  $a^{-3}$ . As space expands, you expect the matter density to drop like  $a^{-3}$ . We once again see that  $\Omega_\Lambda$  is constant; it quite literally stays the same density throughout all time, and across all scales, if it is a constant. But is it??

## 2.4 Evolving Dark Energy?

Let's suppose that  $\Lambda$  isn't actually a constant, and maybe it changes how it behaves over time. What would that look like?

### Definition

The **equation of state** parameter  $w$  describes simply the ratio of an object's density to its pressure. That is  $w = \frac{p}{\rho}$ . A cosmological constant, by definition has  $w = -1$ , or a negative pressure associated with density. If you think about space expanding as you add more of it, that kinda makes sense. Matter, on the other hand has a  $w = 0$ , because it doesn't really apply much pressure in comparison to its mass. Radiation has a  $w = \frac{1}{3}$ , which you will learn in future physics/astronomy classes. It turns out that these exponents of the scale factors are all related to  $w$ ; specifically component  $\Omega_i$  goes as  $a^{-3(1+w)}$ . Double check this for matter and radiation, and the cosmological constant!

The idea is to change the equation of state for dark energy. If it was evolving, then we would expect its equation of state to also evolve with time. We can Taylor expand  $w$  in terms of the scale factor  $a$ . But rather than expanding around  $a = 0$ , which would have been at the birth of the universe, we expand around today,  $a = 1$ . Specifically we let

$$w = w_0 + w_a(1 - a) \quad (6)$$

so the equation of state today is given by  $w_0$ , and it changes over time with slope  $w_a$ . A cosmological constant corresponds to  $w_0 = -1$ , and  $w_a = 0$ .

## 3 Type Ia Supernovae

Type Ia Supernovae (read as “type one a”) are supernovae involving a white dwarf star and a companion star. While both stars are stable individually, when they orbit each other, the white dwarf can exceed its critical mass as it draws stuff from its companion, leading to a very bright and beautiful explosion. To be honest, people aren't exactly sure about the processes that cause these explosions - the prominent ideas have changed even over my (relatively short) lifetime, but the fact remains that these types of explosions are excellent ‘standard candles’.

### Definition

A **standard candle** is an object that can be used as a distance indicator because it has a fixed luminosity/brightness. Objects that are far away are less bright, and this scaling with distance  $d$  is very well measure to be  $L \propto \frac{1}{d^2}$ . If you know the inherent or true luminosity of an object, you can easily measure the distance to it. Astronomers are weird and use something called magnitude  $M$  to measure the brightness of an object. It is a logarithmic quantity, and also higher numbers mean fainter objects. The distance in terms of an object's absolute (inherent) magnitude  $M$  and apparent (measured from Earth) magnitude  $m$  is given by

$$5 \cdot \log_{10} d = m - M + 5 \quad (7)$$

The point is that because Type Ia supernovae all start from white dwarfs, we think they should all have the same absolute magnitudes  $M$ . Therefore, if we see them far away, we can measure distances to the far away galaxies where they occur. This is exactly what we've been doing in our project: measuring the distances and redshifts to hundreds of these supernovae. They fit nicely on a curve, which we can also model using the Friedmann equation. Specifically for the case of a flat universe with a cosmological constant (we call this Lambda CDM cosmology):

$$d_L(z; H_0, \Omega_M) = \frac{1+z}{H_0} \int_0^z [(1+z')^3 \Omega_M + (1 - \Omega_M)]^{-1/2} dz' \quad (8)$$

Astropy does this all under the hood (thank you astropy), so we don't need to worry about deriving this stuff. The equation also changes slightly for the case of evolving dark energy.

## 4 DESI Results

What about these DESI results? What does DESI even do, and why should we care? DESI makes a 3D map of the universe by measuring the positions of a bunch of galaxies across the sky. By measuring how these galaxies are clustered, we can measure something called the Baryonic Acoustic Oscillation scale, which is like a little bump in the shape of the correlation function of these galaxies. (Too much jargon - what this really means is that if you look at some fixed comoving distance from a given galaxy, you expect to see more galaxies there than usual because of some cool early universe physics that became the seeds for structure in the universe). DESI measures this scale very precisely at different redshifts, and gets an independent measure of distance versus redshift, and tries to fit cosmologies to this data, just like we did with the supernova. They find that the best-fitting model may actually be one of evolving dark energy, with  $w_0 > -1$  and  $w_a < 0$ . DESI's constraints/conclusions are much stronger when they combine their results with other experiments/observations.

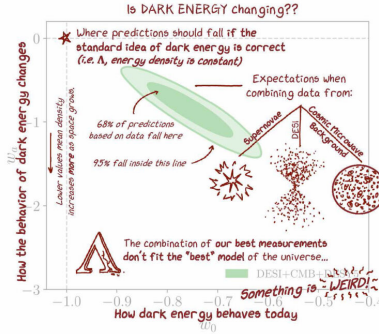


Figure 2: A nice image of the DESI results that better describes what they mean. Take a look at [this article](#) for more

In fact, the results are strongest when they combine with a supernova dataset, which measures the distances at lower redshifts than DESI.

On the other hand, a different supernova data set weakens the DESI results, which is very interesting. The difference comes down to if Type Ia supernova are really as standard as we think they are. This is a very active and contentious area of research, which requires a lot of precise astrophysics to better understand Type Ia supernovae, and also better calibration of the cosmic distance ladder.

The other thing about supernova (which is a very dramatic field in cosmology) is something called the Hubble tension. One of the guys who won the Nobel prize for the first discovery has continued to make measurements of the Hubble constant with supernovae. What his group has consistently been finding is that  $H_0 \approx 72$ . A bunch of other experiments/probes of cosmology have (including DESI) have consistently found  $H_0 \approx 68$ . This doesn't seem like a big deal, but it really is, as our error bars are getting smaller and smaller, and we have no real reason to believe these two things are different. This problem has been referred to as the Hubble tension, and has inspired a lot of cosmology research (including DESI) over the past few decades.

Hopefully this project seems a little cooler and a little easier to understand now! Please ask me more questions if you have any, and I hope you both had fun doing this project!