

BJT Internal Capacitances &

①

HIGH FREQUENCY MODEL

The physical flow of carriers inside a BJT get stored at p-n junctions and this in some way affect their instantaneous response to change in applied field. Such behaviour can be modelled by representing that by an equivalent capacitance (charge storage phenomenon) that would affect high frequency behaviour of BJT as an amplifier.

BASE CHARGING or Diffusion Capacitance C_{de}

When BJT is in active or saturation region, minority carrier charges are stored in base region. The charge Q_n in terms of collector current i_c

$$Q_n = \frac{W^2 i_c}{2 D_n} = \tau_F \cdot i_c$$

where τ_F = device constant = $\frac{W^2}{2 D_n}$

W = effective width of Base

D_n = electron diffusivity in Base

τ_F = FORWARD BASE TRANSIT TIME

of the order of 10 picoseconds to 100 picoseconds.

Since i_c is related to V_{BE} exponentially Q_n also depends upon V_{BE} nonlinearly. (2)
 This mechanism can be represented by a SMALL SIGNAL DIFFUSION CAPACITANCE C_{de}

$$C_{de} = \frac{dQ_n}{dV_{BE}} = \tau_F \frac{di_c}{dV_{BE}} = \tau_F \cdot g_m$$

$$= \tau_F \cdot \frac{I_C}{V_T}$$

Thus C_{de} increases with I_C or Bias current.
 ————— x ————

BASE EMITTER JN CAPACITANCE

B-E jn. depletion layer capacitance C_{je} is given by

$$C_{je} = \frac{C_{je0}}{\left(1 - \frac{V_{BE}}{V_{oe}}\right)^m}$$

← value of C_{je} at 0V

V_{oe} = EBJ built-in Voltage $\approx 0.9V$

m = grading coeff. ≈ 0.5

V_{BE} = Forward junction DC voltage or bias voltage.

Approx. $C_{je} \approx 2 C_{je0}$

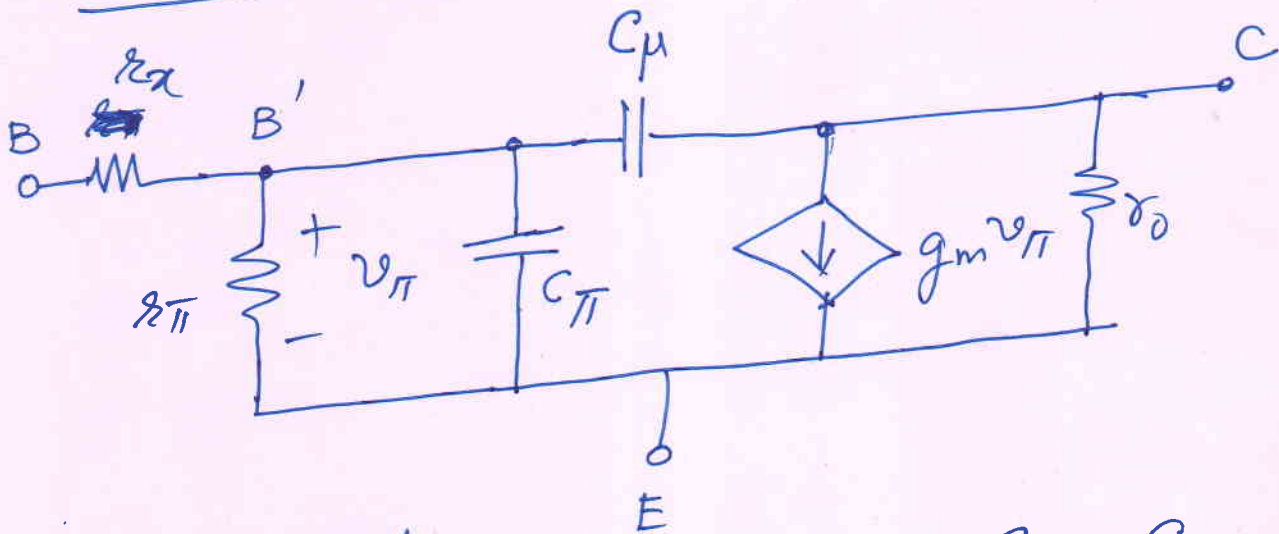
Collector Base Jn. Capacitance C_{μ} (3)

\therefore CB is reversed biased, its junction has a capacitance like effect modelled as

$$C_{\mu} = \frac{C_{\mu 0}}{\left(1 + \frac{V_{CB}}{V_{oc}}\right)^m} \rightarrow C_{\mu} \text{ at } 0V$$

$V_{oc} \rightarrow$ CBJ built-in Voltage
 $\approx 0.75V$
 $m \approx 0.2 \text{ to } 0.3V$

HIGH-FREQUENCY HYBRID- π MODEL.



Newly Added:

C_{π} = emitter base capacitance = $C_{de} + C_{je}$

C_{μ} = collector-base capacitance

r_x = small value to model resistance of silicon material \approx few tens of ohms.

At low frequencies $r_x + r_{\pi} \approx r_{\pi}$ therefore it was never considered or brought in discussion.
 At high freq. C_{π} will bypass r_{π} & r_x will matter.

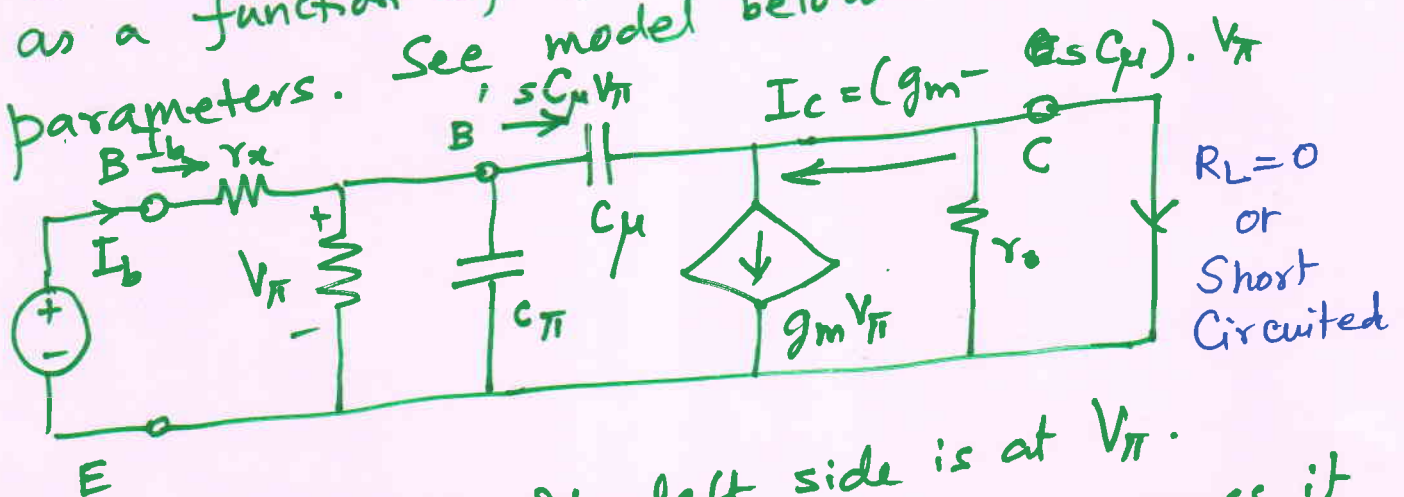
Typically C_{π} = few tens of Picofarads
 C_{μ} = fraction to few Picofarads
 r_x = tens of ohms.

4

Voltages and currents described now onwards will be function of frequency hence we will use uppercase letters with lowercase subscripts like V_{π} or I_c .

Expression of Upper Cutoff Frequency

Derive an expression for short circuit current gain β (also called as h_{fe}) or CE S.C.C. Gain as a function of frequency, in terms of model parameters. See model below with $R_L = 0$. S.C.



See capacitor C_{μ} . Its left side is at V_{π} .
 Its right side is at GND \therefore voltage across it is V_{π} . \therefore Current through $C_{\mu} = \frac{V_{\pi}}{X_{C_{\mu}}} = \frac{V_{\pi}}{1/C_{\mu}s}$

$= C_{\mu}s \cdot V_{\pi}$
 $\therefore r_o$ is bypassed by a short circuit, it does not matter at all in calculations

Now look at currents in Node C :

(5)

$$\text{coll. } I_c = g_m V_{\pi} - s C_{\mu} V_{\pi}.$$

$$\text{current} = (g_m - s C_{\mu}) V_{\pi}$$

V_{π} is voltage across r_{π} due to I_b .

$$V_{\pi} = I_b (r_{\pi} \parallel C_{\pi} \parallel C_{\mu})$$

\therefore voltage across C_{μ} is also V_{π} .

$$V_{\pi} = \frac{I_b}{\frac{1}{r_{\pi}} + s C_{\pi} + s C_{\mu}}$$

Calculate h_{fe}

$$h_{fe} = \frac{I_c}{I_b} = \frac{(g_m - s C_{\mu}) \cdot V_{\pi}}{(\frac{1}{r_{\pi}} + s C_{\pi} + s C_{\mu}) \cdot V_{\pi}}$$

Assuming that at high frequencies $g_m \gg \omega C_{\mu}$ & we neglect $s C_{\mu}$ in numerator.

$$h_{fe} = \frac{g_m \cdot r_{\pi}}{1 + s(C_{\pi} + C_{\mu}) r_{\pi}}$$
$$= \frac{\beta_0}{1 + s(C_{\pi} + C_{\mu}) r_{\pi}}$$

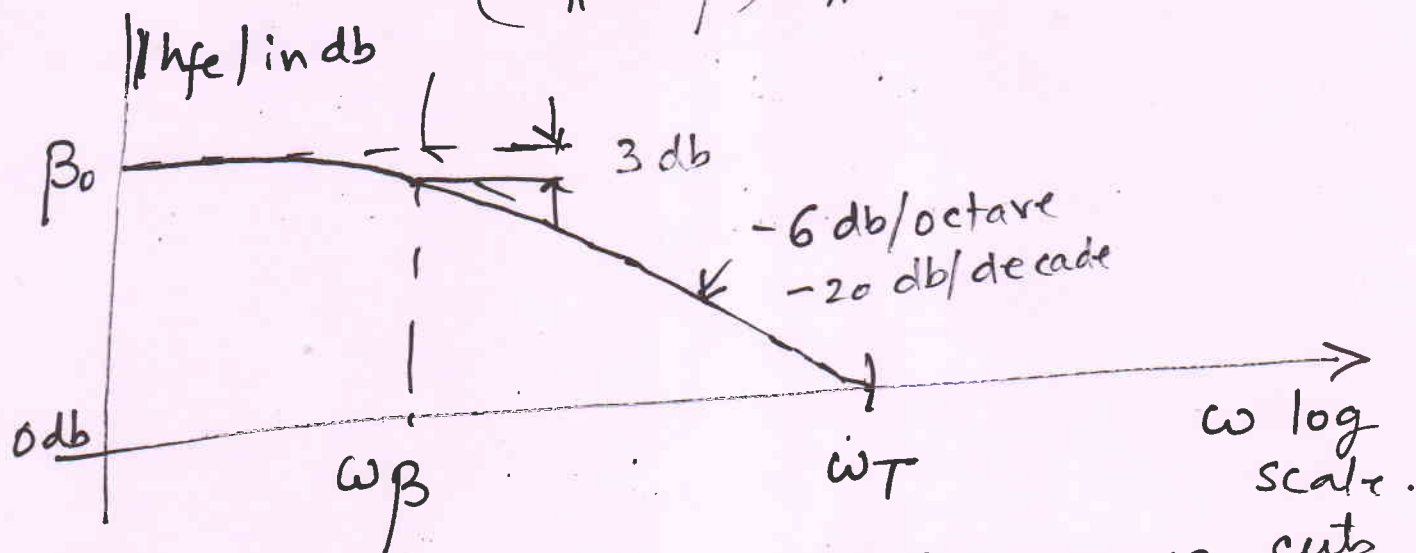
where $\beta_0 =$ low frequency β .
 $= g_m \cdot r_{\pi}$

Now, we can see this expression as an STC function at a 3-db or corner freq. (6)

$$\omega = \omega_{\beta} \text{ where}$$

$$\omega_{\beta} = \frac{1}{(C_{\pi} + C_{\mu}) r_{\pi}}$$

See Bode Plot



The frequency at which this curve cuts 0-dB line is called frequency where $|h_{fe}| = 1$ or UNITY GAIN BANDWIDTH ω_T such that $\omega_T = (\beta_0 \cdot \omega_{\beta})$

$$\text{Thus } \omega_T = \frac{\beta_0}{(C_{\pi} + C_{\mu}) r_{\pi}} = \frac{g_m}{(C_{\pi} + C_{\mu})}$$

$$\therefore f_T = \frac{\omega_T}{2\pi} = \frac{g_m}{2\pi(C_{\pi} + C_{\mu})}$$

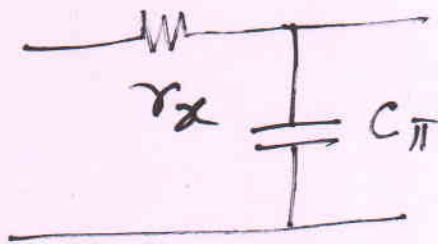
f_T is usually specified in BJT datasheets. It is in MHz to GHz range.

From f_T , we can get C_π & C_μ . (7)

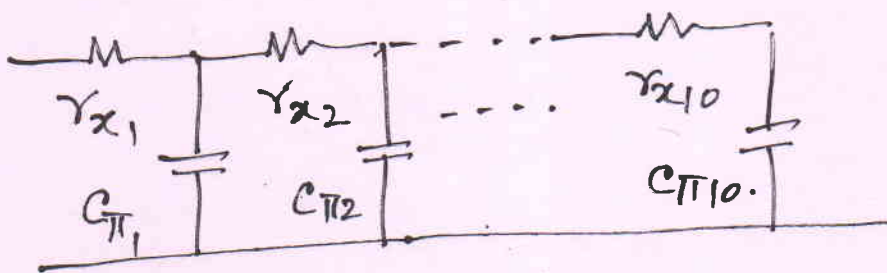
We can measure practically C_μ between B & C and get a measure of C_π .

This STC model we discussed, works well upto $0.2 f_T$. At higher frequencies, we must use distributed R-C model for better estimates:

Lumped model.



distributed model. for say 10 taps...



Numerical on C_μ , C_π and f_T .

3.149 Given npn BJT $I_C = 0.5 \text{ mA}$
 $V_{CB} = 2 \text{ V}$; $\beta_0 = 100$; $V_A = 50 \text{ V}$
 $\tau_p = 30 \text{ ps}$; $C_{je0} = 20 \text{ femto farad}$
 $C_{\mu 0} = \cancel{20}^{30} \text{ femto F}$; $\underline{\hspace{1cm}} = 20 \text{ fF} = .02 \text{ pF}$
 $V_{oc} = 0.75 \text{ V}$, $m_{CBT} = 0.5$; $r_x = 100 \Omega$

Draw hybrid- π model & find f_T .

$r_x = 100 \Omega$; $g_m = I_C / V_T = 0.5 \text{ mA} / 25 \text{ mV} = 20 \text{ mS}$

$r_\pi = \beta_0 / g_m = 100 / 20 \text{ mS} = 5 \text{ k}\Omega$

$r_o = V_A / I_C = 50 \text{ V} / 0.5 \text{ mA} = 100 \text{ k}\Omega$

$C_\mu = C_{\mu 0} / \left(1 + \frac{V_{CB}}{V_{oc}}\right)^m = \frac{\cancel{20}^{30} \text{ fF}}{\left(1 + \frac{2 \text{ V}}{0.75 \text{ V}}\right)^{0.5}} = 15.66 \text{ fF}$

$C_{je} \approx 2 C_{je0} = 20 \times 20 = 40 \text{ fF}$

$C_{de} = \tau_F \cdot g_m = 30 \text{ ps} \cdot 20 \text{ mS} = 600 \text{ femto F}$

$C_\pi = C_{je} + C_{de} = 40 + 600 = 640 \text{ fF} = 0.640 \text{ pF}$

$f_T = \frac{g_m}{2\pi(C_\pi + C_\mu)} = \frac{20 \text{ mS} \times 10^{-3}}{2 \times 3.14 \left(\frac{640 + 15.66}{66}\right) \times 10^{-15}} \text{ Hz}$

$\boxed{= 4.857 \text{ Giga Hz}}$

Q: 3.150 At 500 MHz signal, $|h_{fe}| = 2.5 @ I_C = 0.2 \text{ mA}$

and at same frequency $|h_{fe}| = 11.6 @ I_C = 1.0 \text{ mA}$ (9)

Given: $C_\mu = 0.05 \text{ pF} \approx 50 \text{ fF}$.

Find f_T at both Collector Currents.

What is r_F and C_{je} value?

————— x —————

Note that

$$|h_{fe}| = \frac{f_T}{\text{freq of measurement}}$$

$$@ I_C = 0.2 \text{ mA} \quad \therefore f_T = |h_{fe}| \cdot f = 2.5 \times 500 \text{ MHz} = 1.25 \text{ GHz}$$

at another I_C value @ $I_C = 1.0 \text{ mA}$

$$f_T = 11.6 \times 500 \text{ MHz} = 5.8 \text{ GHz}$$

Note: Change in bias current from 0.2 mA to 1.0 mA increases f_T from 1.25 to 5.8 GHz.

Thus if we want a higher frequency amplifier, we should keep $I_C =$ higher value.

$$\therefore f_T = g_m / 2\pi (C_\pi + C_\mu)$$

$$\therefore C_\pi = \left(g_m / (2\pi f_T) \right) - C_\mu$$

$$\text{at } I_C = 0.2 \text{ mA} \rightarrow g_m = 0.2 / 25 = 8 \text{ mV}$$

$$C_\pi = \left[8 \text{ mV} / \left(2 \times 3.14 \times 1.25 \text{ GHz} \right) \right] - 50 \text{ fF} = 968 \text{ fF}$$

at $I_C = 1.0 \text{ mA}$ $g_m = 40 \text{ mV}$

(10)

$$C_{\pi} = \left[40 \text{ mV} / (2 \times 3.14 \times 5.8 \text{ GHz}) \right] - 50 \text{ fF}$$

$$= 1047.6 \text{ fF.}$$

Note that C_{π} also increases with I_C .
— x —

Next, $C_{\pi} = C_{je} + \tau_F \cdot g_m$

② $I_C = 0.2 \text{ mA}$

$$968 \text{ fF} = C_{je} + \tau_F \cdot 8 \text{ mV} \quad \text{--- (1)}$$

③ $I_C = 1.0 \text{ mA}$

$$1047.6 \text{ fF} = C_{je} + \tau_F \cdot 40 \text{ mV} \quad \text{--- (2)}$$

Solving (1) and (2) $C_{je} = 950 \text{ fF}$ and
 $\tau_F = 2.47 \text{ ps}$

Q: 3.151 Given: $I_C = 2 \text{ mA}$; $C_{\mu} = 1 \text{ pF}$, $C_{\pi} = 10 \text{ pF}$
 $\beta = 150$. Find f_T and f_{β} .

$$f_T = g_m / (2\pi(C_{\mu} + C_{\pi})) = 80 \text{ mV} / (2 \times 3.14 \times (1 + 10) \text{ pF})$$

$$= \boxed{1.158 \text{ GHz}}$$

$$f_{\beta} = \frac{f_T}{\beta_0} = \frac{1.158 \text{ GHz}}{150} = \boxed{7.725 \text{ MHz.}}$$