# How physics simulations work: simulating particles in electric and magnetic fields

Nikolaj Roager Christensen

Student Colloquium in Physics and Astronomy, Aarhus University

March 2021

# Numerical integration of ordinary differential equations: motion of charged particles in electromagnetic fields

#### Introduction

Theory and physical background 10 min

Euler's Method and the 4th order Runge-Kutta Method 15 min Euler's Method Higher order Runge-Kutta methods Demonstration, particles in a solenoid

Embedded algorithms, and adaptive step size 15 min Dormand Prince 5 (4) method Demonstration: magnetic dipole

#### Conclusion

If time permits, the cyclotron

► Numerical simulations are important.

- ► Numerical simulations are important.
- ► Testing setups, non-analytical systems.

- ► Numerical simulations are important.
- ► Testing setups, non-analytical systems.
- Demonstration, charged particles in electric and magnetic fields.
- ► Analytically known and not.

- ▶ Numerical simulations are important.
- ► Testing setups, non-analytical systems.
- Demonstration, charged particles in electric and magnetic fields.
- Analytically known and not.
- ► Simulations are not experiments!

► Some repetition from Electrodynamics

- ► Some repetition from Electrodynamics
- ► The Lorentz force+N2:

$$ec{F} = q(ec{v} imes ec{B} + ec{E}).$$
 $ec{mr}(t) = q(ec{r}(t) imes ec{B}(ec{r},t) + ec{E}(ec{r},t)).$ 

- ► Some repetition from Electrodynamics
- ► The Lorentz force+N2:

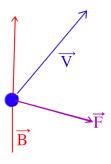
$$ec{F} = q(ec{v} imes ec{B} + ec{E}).$$
 $ec{mr}(t) = q(\dot{ec{r}}(t) imes ec{B}(ec{r},t) + ec{E}(ec{r},t)).$ 

▶ Only 1 particle! so pre-programmed depending on the setup.

- ► Some repetition from Electrodynamics
- ► The Lorentz force+N2:

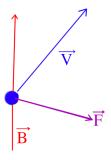
$$ec{F} = q(ec{v} imes ec{B} + ec{E}).$$
 $ec{mr}(t) = q(\dot{\vec{r}}(t) imes ec{B}(ec{r},t) + ec{E}(ec{r},t)).$ 

- ▶ Only 1 particle! so pre-programmed depending on the setup.
- ► Could use potentials  $\phi(\vec{r},t)$   $\vec{A}(\vec{r},t)$  and Hamiltonian, or Lagrangian.
- ▶ Other systems would have other differential equations.



► Magnetic forces do no work:

$$dW_{\vec{B}} = \vec{F}_B \cdot d\vec{r} \propto (\vec{v} \times \vec{B}) \cdot \vec{v} = 0.$$

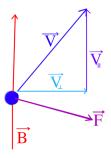


► Magnetic forces do no work:

$$dW_{\vec{B}} = \vec{F}_B \cdot d\vec{r} \propto (\vec{v} \times \vec{B}) \cdot \vec{v} = 0.$$

ightharpoonup  $(\vec{v} = \vec{v}_{\perp} + \vec{v}_{\parallel})$ :

$$|\vec{F}_B| = |q(\vec{v} \times \vec{B})| = |qv_{\perp}B|.$$



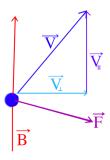
► Magnetic forces do no work:

$$dW_{\vec{B}} = \vec{F}_B \cdot d\vec{r} \propto (\vec{v} \times \vec{B}) \cdot \vec{v} = 0.$$

ightharpoonup  $(\vec{v} = \vec{v}_{\perp} + \vec{v}_{\parallel})$ :

$$|\vec{F}_B| = |q(\vec{v} \times \vec{B})| = |qv_{\perp}B|.$$

► Same as Centripetal force: Cyclotron motion



► Magnetic forces do no work:

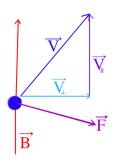
$$dW_{\vec{B}} = \vec{F}_B \cdot d\vec{r} \propto (\vec{v} \times \vec{B}) \cdot \vec{v} = 0.$$

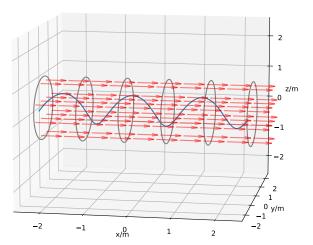
$$ightharpoonup$$
  $(\vec{v} = \vec{v}_{\perp} + \vec{v}_{\parallel})$ :

$$|\vec{F}_B| = |q(\vec{v} \times \vec{B})| = |qv_{\perp}B|.$$

- ► Same as Centripetal force: Cyclotron motion
- Cyclotron radius and frequency:

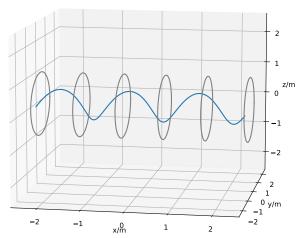
$$R = \frac{v_{\perp}m}{|a|B}$$
  $\omega_c = \frac{|q|B}{m}$ .





#### "Cyclotron motion"

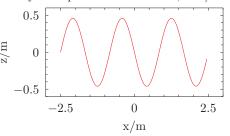
Solenoid with N=1000 turns per m, I=5 A, r=1 m,  $|\vec{B}|\approx 6$  mT. Proton with  $E_{kin}=1$  MeV/c<sup>2</sup> ( $|v|\approx 3.195\times 10^5$  m/s)

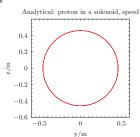


"Cyclotron motion"

$$R pprox 0.5 \, \mathrm{m} \, \mathrm{sin}( heta)$$
  $T = \frac{2\pi}{\omega_c} pprox 10 \, \mu \mathrm{s}$ 

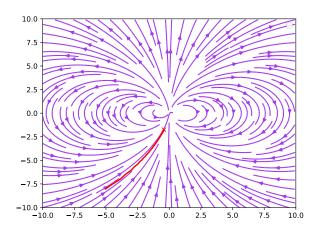
Analytical: proton in a solenoid, side/front-view





"Cyclotron motion"

$$R pprox 0.5 \, \mathrm{m} \, \mathrm{sin}( heta) \quad T = rac{2\pi}{\omega_c} pprox 10 \, \mathrm{\mu s}$$



(Actually from my simulation)

### Ordinary differential equation's.

- ► Sources: Zeigler et al. Theory of Modeling and Simulation (Third edition) chapter 3
- ► We have:

$$\ddot{\vec{r}}(t) = \frac{q}{m}(\dot{\vec{r}}(t) \times \vec{B}(\vec{r},t) + \vec{E}(\vec{r},t)).$$

### Ordinary differential equation's.

- ➤ Sources: Zeigler et al. Theory of Modeling and Simulation (Third edition) chapter 3
- ► We have:

$$\vec{r}(t) = \frac{q}{m}(\vec{r}(t) \times \vec{B}(\vec{r},t) + \vec{E}(\vec{r},t)).$$

► Algorithms exists for ODEs:

$$\dot{\mathbf{X}} = f_{ode}(\mathbf{X}(t), t).$$

### Ordinary differential equation's.

- ► Sources: Zeigler et al. Theory of Modeling and Simulation (Third edition) chapter 3
- ► We have:

$$\ddot{\vec{r}}(t) = \frac{q}{m}(\dot{\vec{r}}(t) \times \vec{B}(\vec{r},t) + \vec{E}(\vec{r},t)).$$

► Algorithms exists for ODEs:

$$\dot{\mathbf{X}} = f_{ode}(\mathbf{X}(t), t).$$

► Here:

$$\mathbf{X} = \left( egin{array}{c} ec{r} \ \dot{ec{r}} \ \dot{ec{r}} \end{array} 
ight) \quad f_{ode}(\mathbf{X},t) = \left( egin{array}{c} \dot{ec{r}} \ \dfrac{\dot{ec{r}}}{m} (\dot{ec{r}} imes ec{B}(ec{r},t) + ec{E}(ec{r},t)) \end{array} 
ight).$$

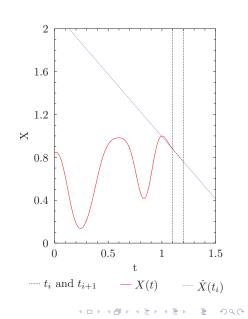


#### The ODE to solve

```
auto ODE = [...](const state_type Data, state_type &
   dDatadt, const double t){
    //Extract position and velocity from data
    vec pos = vec(Data[0],Data[1],Data[2]);
    vec velocity = vec(Data[3],Data[4],Data[5]);
    //Lorentz+Newtons 2nd law
    vec F = Charge*(Fields.get_Efield(pos,t)+
        cross(velocity,Fields.get Bfield(pos,t)));
    vec dVdt = F*Inv mass;
    //Save derivative of data
    dDatadt[0]=velocity.x;
```

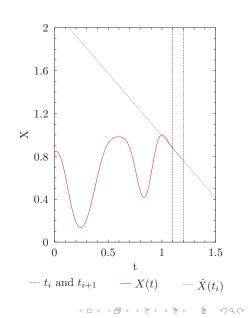
• We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .

► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



- We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .
- Can we find  $\mathbf{X}(t_0 + h)$  for h > 0?.

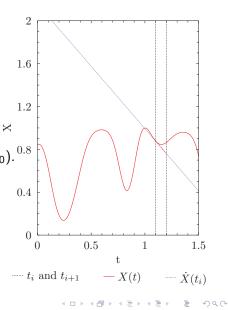
► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



- We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .
- Can we find  $\mathbf{X}(t_0 + h)$  for h > 0?.
- (Explicit Forward) Euler's Method:

$$\mathbf{X}(t_0+h) = \mathbf{X}(t_0) + hf_{ode}(\mathbf{X}(t_0), t_0).$$
<sup>0.8</sup>

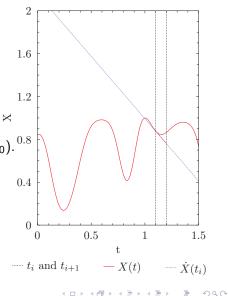
- Multiple steps  $t_0, t_1 = t_0 + h, \dots t_i, \dots$
- ► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



- We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .
- Can we find  $\mathbf{X}(t_0 + h)$  for h > 0?.
- (Explicit Forward) Euler's Method:

$$\mathbf{X}(t_0+h) = \mathbf{X}(t_0) + hf_{ode}(\mathbf{X}(t_0), t_0).$$
<sup>0.8</sup>

- Multiple steps  $t_0, t_1 = t_0 + h, \dots t_i, \dots$
- ► Can this be justified?
- ► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



► Multiple justifications for why.

- ► Multiple justifications for why.
- ► First 2 terms in Taylor series Zeigler et al.:

$$\mathbf{X}(t_0 + h) = \mathbf{X}(t_0) + hf_{ode}(\mathbf{X}(t_0), t_0) + h^2 \dots + \dots$$

- ► Multiple justifications for why.
- ► First 2 terms in Taylor series Zeigler et al.:

$$\mathbf{X}(t_i+h)=\mathbf{X}(t_i)+hf_{ode}(\mathbf{X}(t_i),t_i)+h^2\ldots+\ldots$$

▶ Local error  $h^2$ . Global error:  $h^1$ .

- ► Multiple justifications for why.
- ► First 2 terms in Taylor series Zeigler et al.:

$$\mathbf{X}(t_i+h)=\mathbf{X}(t_i)+hf_{ode}(\mathbf{X}(t_i),t_i)+h^2\ldots+\ldots$$

- ▶ Local error  $h^2$ . Global error:  $h^1$ .
- ightharpoonup Argument suggests we need  $f_{ode}, f_{ode}, \dots$

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t_{i+1}} \dot{\mathbf{X}}(t_i) dt.$$

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t_{i+1}} f_{ode}(\mathbf{X}(t_i), t_i) dt.$$

▶ Mean Value theorem for integrals  $t_i \le \tau \le t_{i+1}$ .

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t_{i+1}} f_{ode}(\mathbf{X}(t_i), t_i) dt = h f_{ode}(\mathbf{X}(\tau), \tau).$$

▶ Mean Value theorem for integrals  $t_i \le \tau \le t_{i+1}$ .

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = hf_{ode}(\mathbf{X}(\tau), \tau).$$

- ▶ Mean Value theorem for integrals  $t_i \le \tau \le t_{i+1}$ .
- Guess  $\tau = t_i$ . Good guess, if  $f_{ode} \approx 0$ .

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = hf_{ode}(\mathbf{X}(\tau), \tau).$$

- ▶ Mean Value theorem for integrals  $t_i \le \tau \le t_{i+1}$ .
- ▶ Guess  $\tau = t_i$ . Good guess, if  $f_{ode} \approx 0$ .
- ▶ More generally, (*Explicit* and *single step*), Runge-Kutta family:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t'} f_{ode}(\mathbf{X}(t_i), t_i) dt + \ldots \int_{t^{(m)}}^{t_{i+1}} f_{ode}(\mathbf{X}(t_i), t_i) dt.$$



# The Runge Kutta steppers

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = hf_{ode}(\mathbf{X}(\tau), \tau).$$

- ▶ Mean Value theorem for integrals  $t_i \le \tau \le t_{i+1}$ .
- Guess  $\tau = t_i$ . Good guess, if  $\dot{f}_{ode} \approx 0$ .
- ▶ More generally, (*Explicit* and *single step*), Runge-Kutta family:

$$egin{aligned} \mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) &= \int_{t_i}^{t'} f_{ode}(\mathbf{X}(t_i), t_i) dt + \ldots \int_{t^{(m)}}^{t_{i+1}} f_{ode}(\mathbf{X}(t_i), t_i) dt, \ &= \sum_{j=1}^m h_j f_{ode}(\mathbf{X}( au_j), au_j). \end{aligned}$$

- ▶ Use  $f_{ode}(\mathbf{X}(t_i), t_i)$  to approximate  $\mathbf{X}(\tau_1)$ .
- ► L. Zheng, X. Zhang, Modeling and Analysis of Modern Fluid Problems, 2017, chapter 8:

More commonly written:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j \mathbf{K}_j$$

► Martha L. Abell, James P. Braselton, Differential Equations with Mathematica (Fourth Edition), 2016:

More commonly written:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{i=1}^m b_j \mathbf{K}_j$$

► Estimate **K**<sub>j</sub>:

$$\mathbf{K}_{1} = f_{ode}(\mathbf{X}(t_{i}), t_{i})$$
 $\mathbf{K}_{2} = f_{ode}(\mathbf{X}(t_{i}) + ha_{21}\mathbf{K}_{1}, t_{i} + c_{2}h)$ 
 $\mathbf{K}_{3} = f_{ode}(\mathbf{X}(t_{i}) + ha_{31}\mathbf{K}_{1} + ha_{32}\mathbf{K}_{2}, t_{i} + c_{3}h)$ 
 $\vdots$ 

► Martha L. Abell, James P. Braselton, Differential Equations with Mathematica (Fourth Edition), 2016:

More commonly written:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{i=1}^m b_j \mathbf{K}_j$$

► Estimate **K**<sub>j</sub>:

$$egin{aligned} \mathbf{K}_1 &= f_{ode}(\mathbf{X}(t_i), t_i) \ \mathbf{K}_2 &= f_{ode}(\mathbf{X}(t_i) + h a_{21} \mathbf{K}_1, t_i + c_2 h) \ \mathbf{K}_3 &= f_{ode}(\mathbf{X}(t_i) + h a_{31} \mathbf{K}_1 + h a_{32} \mathbf{K}_2, t_i + c_3 h) \ &\vdots \end{aligned}$$

- ► Want exact for p'th order polynomial.
- ► Martha L. Abell, James P. Braselton, Differential Equations with Mathematica (Fourth Edition), 2016:

More commonly written:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j \mathbf{K}_j$$

► Estimate **K**<sub>j</sub>:

$$\mathbf{K}_{1} = f_{ode}(\mathbf{X}(t_{i}), t_{i})$$
 $\mathbf{K}_{2} = f_{ode}(\mathbf{X}(t_{i}) + ha_{21}\mathbf{K}_{1}, t_{i} + c_{2}h)$ 
 $\mathbf{K}_{3} = f_{ode}(\mathbf{X}(t_{i}) + ha_{31}\mathbf{K}_{1} + ha_{32}\mathbf{K}_{2}, t_{i} + c_{3}h)$ 
 $\vdots$ 

- ► Want exact for p'th order polynomial.
- ▶ Taylor series analogy: error  $h^p$ .
- ► Martha L. Abell, James P. Braselton, Differential Equations with Mathematica (Fourth Edition), 2016:

### Higher order methods

▶ 2nd order (Heun's method):

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \frac{h}{2}(\mathbf{k}_1 + \mathbf{k}_2)$$

$$\mathbf{k}_1 = f_{ode}(\mathbf{X}(t_i), t_i)$$

$$\mathbf{k}_2 = f_{ode}(\mathbf{X}(t_i) + h\mathbf{k}_1, t_i + h)$$

# Higher order methods

➤ 4th order, often simply called the Runge Kutta method:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \frac{h}{6}(\mathbf{k}_1 + 2\mathbf{k}_2 + 2\mathbf{k}_3 + \mathbf{k}_4)$$

$$\mathbf{k}_1 = f_{ode}(\mathbf{X}(t_i), t_i)$$

$$\mathbf{k}_2 = f_{ode}(\mathbf{X}(t_i) + \frac{h}{2}\mathbf{k}_1, t_i + \frac{h}{2})$$

$$\mathbf{k}_3 = f_{ode}(\mathbf{X}(t_i) + \frac{h}{2}\mathbf{k}_2, t_i + \frac{h}{2})$$
Wikipedia-user HilberTraum,
$$\mathbf{k}_4 = f_{ode}(\mathbf{X}(t_i) + h\mathbf{k}_3, t_i + h)$$
 published under creative commins: CC BY-SA 4.0

#### The General explicit Runge Kutta method

► General explicit, single step, fixed size, Runge Kutta method

# The General explicit Runge Kutta method

- ► General explicit, single step, fixed size, Runge Kutta method
- ► Expressed in Butcher tableu:

#### **Euler Implementations**

```
state_type Data = Data0;
state type dDatadt;
size t time res = T/timestep;
for (size t i = 1; i < time res; ++i)</pre>
{
    double t=i*dt;
    ODE(Data,dDatadt,t);
    //Euler time evolution
    Data+=timestep*dDatadt;
    save_step( Data , i*timestep );
};
```

### RK4 Implementations (1/2)

```
state_type Data = Data0;
state_type temp=Data0;
state_type K1,K2,K3,K4;
size_t time_res = T/timestep;
for (size t i = 1; i < time res; ++i)</pre>
{
    double t=i*timestep;
    //substep 1
    ODE(Data, K1,t);
    //substep 2
    temp=Data+timestep*K1/2;
    ODE(temp,K2,t+timestep/2);
```

## RK4 Implementations (2/2)

```
//substep 3
temp=Data+timestep*K2/2;
ODE(temp,K3,t+timestep/2);
//substep 4
temp=Data+timestep*K3;
ODE(temp,K4,t+timestep);
//Read data
Data+=timestep*(K1+2.0*K2+2.0*K3+K4)/6.0;
save_step( Data , i*timestep );
```

## "Correct" way

```
#include <boost/array.hpp>
#include <boost/numeric/odeint.hpp>
using namespace boost::numeric::odeint;
typedef boost::array< double, 6 > state_type;
size_t steps = integrate_const(
   runge kutta4< state type >(),
   ODE. //Lorentz-force
   Data0 ,//{pos0,v0}
   0.0 . //t0=0
   T , //max time
   timestep ,//length of each step
    save_step //User defined save data function
);
```

▶ Test, same proton in a solenoid use  $\theta = 60^{\circ}$  reference, had:

$$R pprox 0.5 \, \mathrm{m} \, \mathrm{sin}( heta) pprox 0.45 \, \mathrm{m}$$
  $T = rac{2\pi}{\omega_c} pprox 10 \, \mathrm{\mu s}$ 

▶ Test, same proton in a solenoid use  $\theta = 60^{\circ}$  reference, had:

$$R pprox 0.5 \,\mathrm{m} \,\mathrm{sin}( heta) pprox 0.45 \,\mathrm{m}$$
  $T = \frac{2\pi}{\omega_c} pprox 10 \,\mathrm{\mu s}$ 

► Compare Analytic, Euler, Runge-Kutta 4.



▶ Test, same proton in a solenoid use  $\theta = 60^{\circ}$  reference, had:

$$R \approx 0.5 \,\mathrm{m} \,\mathrm{sin}(\theta) \approx 0.45 \,\mathrm{m}$$
  $T = \frac{2\pi}{\omega_c} \approx 10 \,\mathrm{\mu s}$ 

- ► Compare Analytic, Euler, Runge-Kutta 4.
- ► Consider  $\theta = 60^{\circ}$  at different h.

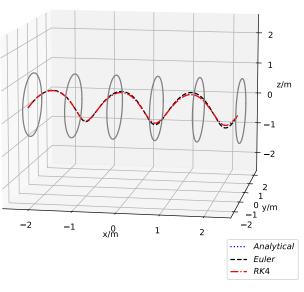
▶ Test, same proton in a solenoid use  $\theta = 60^{\circ}$  reference, had:

$$R \approx 0.5 \, \mathrm{m} \, \mathrm{sin}(\theta) \approx 0.45 \, \mathrm{m}$$
  $T = \frac{2\pi}{\omega_c} \approx 10 \, \mathrm{\mu s}$ 

- ► Compare Analytic, Euler, Runge-Kutta 4.
- ► Consider  $\theta = 60^{\circ}$  at different h.
- Check error on  $|\vec{v}|$ ,  $R = \sqrt{y^2 + z^2}$  and x(t).

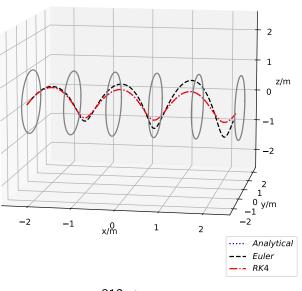
# At a glance, 3D view

$$h = t_{i+1} - t_i = 0.01 \, \mu s$$



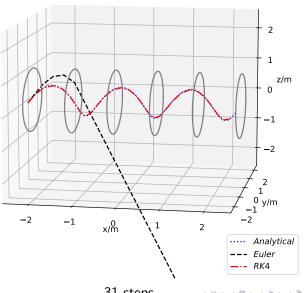
#### At a glance, 3D view

$$h = t_{i+1} - t_i = 0.1 \,\mu s$$

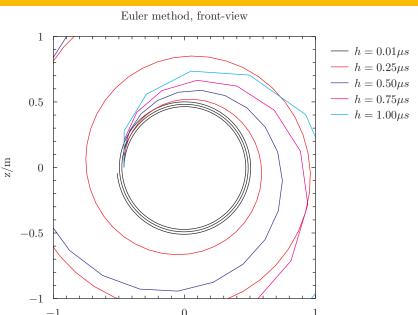


#### At a glance, 3D view

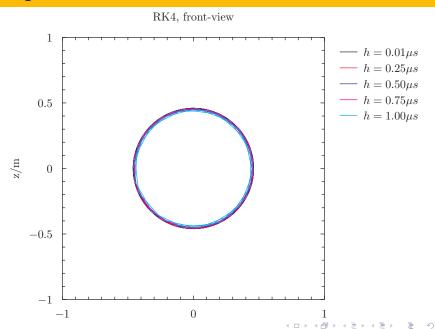
$$h = t_{i+1} - t_i = 1.0 \,\mu s$$



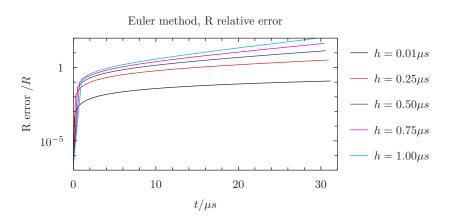
#### At a glance, front view, no border



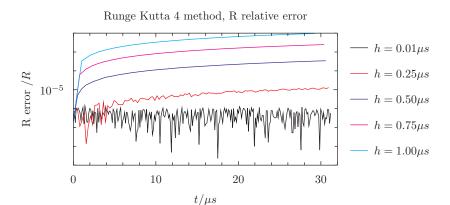
#### At a glance, front view, no border



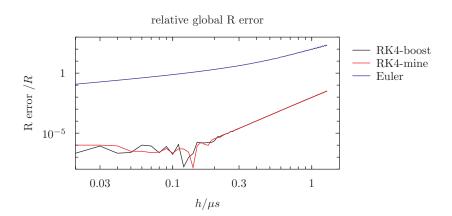
#### Constant radius?



#### Constant radius?



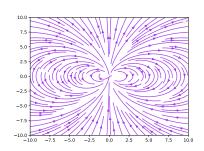
#### Error as function of h



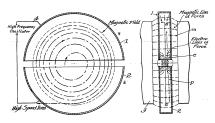
► *h* must be small "enough"

- ► h must be small "enough"
- ► Hard to pick, and may change:

- ► h must be small "enough"
- ► Hard to pick, and may change:
- ► Inhomogeneous fields (here, a true dipole)

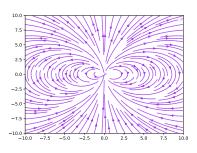


- ► h must be small "enough"
- ► Hard to pick, and may change:
- ► Inhomogeneous fields (here, a true dipole)
- Time dependent fields (here the cyclotron, bad example)



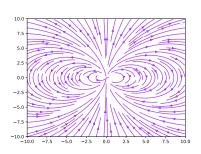
Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

- ► h must be small "enough"
- ► Hard to pick, and may change:
- ► Inhomogeneous fields (here, a true dipole)
- ► Time dependent fields (here the cyclotron, bad example)
- ► Let the computer pick *h*.



► True dipole:

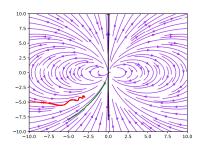
$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3\hat{\vec{r}} (\vec{m} \cdot \hat{\vec{r}} - \vec{m}) \right].$$



► True dipole:

$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3\hat{\vec{r}} (\vec{m} \cdot \hat{\vec{r}} - \vec{m}) \right].$$

Weak field, little change, long steps.

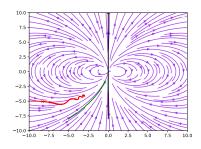


Arbitrarily set  $\frac{\mu_0}{4\pi}|\vec{m}|=0.155\,\mathrm{T/m^3}$  . Protons with speed around  $10\,000\,\mathrm{m/s}$ 

► True dipole:

$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3\hat{\vec{r}} (\vec{m} \cdot \hat{\vec{r}} - \vec{m}) \right].$$

- Weak field, little change, long steps.
- strong field, large change, short steps.

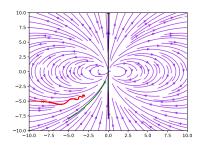


Arbitrarily set  $\frac{\mu_0}{4\pi}|\vec{m}|=0.155\,\mathrm{T/m^3}$  . Protons with speed around  $10\,000\,\mathrm{m/s}$ 

► True dipole:

$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3\hat{\vec{r}} (\vec{m} \cdot \hat{\vec{r}} - \vec{m}) \right].$$

- Weak field, little change, long steps.
- strong field, large change, short steps.
- No analytical solution (afaik.)



Arbitrarily set  $\frac{\mu_0}{4\pi}|\vec{m}|=0.155\,\mathrm{T/m^3}$  . Protons with speed around  $10\,000\,\mathrm{m/s}$ 

### Runge-Kutta Adaptive step size, how

► Reduce *h* until the "error" is small enough

## Runge-Kutta Adaptive step size, how

- ► Reduce *h* until the "error" is small enough
- ▶ Use 2 different order methods,  $\mathbf{x}^{(p+1)}(t_{i+1})$ ,  $\mathbf{x}^{(p)}(t_{i+1})$ .

## Runge-Kutta Adaptive step size, how

- ► Reduce *h* until the "error" is small enough
- ▶ Use 2 different order methods,  $\mathbf{x}^{(p+1)}(t_{i+1})$ ,  $\mathbf{x}^{(p)}(t_{i+1})$ .
- ▶ Approximate "Error" as  $\mathbf{E} = |\mathbf{x}^{(p+1)}(t_{i+1}) \mathbf{x}^{(p)}(t_{i+1})|$ .

### Runge-Kutta Adaptive step size, how

- ► Reduce *h* until the "error" is small enough
- ▶ Use 2 different order methods,  $\mathbf{x}^{(p+1)}(t_{i+1})$ ,  $\mathbf{x}^{(p)}(t_{i+1})$ .
- ▶ Approximate "Error" as  $\mathbf{E} = |\mathbf{x}^{(p+1)}(t_{i+1}) \mathbf{x}^{(p)}(t_{i+1})|$ .
- ► Adjust step size to keep the error(s) small (Implementations differ!).

Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and AppliMathematics

(1)

ode45 in Matlab, scipy.solve\_ivp in Python, RungeKutta\_dopri5 in boost::odeint.

$$\mathbf{X}^{(5)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(5)} \mathbf{k}_j$$
 $\mathbf{X}^{(4)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(4)} \mathbf{k}_j.$ 

$$\mathbf{k}_i = f_{ode}(\mathbf{X}(t_i) + h \sum_{j=1}^{i-1} a_{ki} \mathbf{k}_i, t_i + c_3 h).$$

ode45 in Matlab, scipy.solve\_ivp in Python, RungeKutta\_dopri5 in boost::odeint.

$$\mathbf{X}^{(5)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(5)} \mathbf{k}_j$$
  $\mathbf{X}^{(4)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(4)} \mathbf{k}_j.$ 

$$\mathbf{k}_i = f_{ode}(\mathbf{X}(t_i) + h \sum_{i=1}^{i-1} a_{ki} \mathbf{k}_i, t_i + c_3 h).$$

▶ 7 k<sub>i</sub>'s (actually 6 by clever re-usage).

Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and Applied Mathematics

# Butcher Tableu of Dormand Prince (5) 4

Ci	a <sub>ij</sub>							
$\begin{array}{c} 0 \\ 1 \\ 153 \\ \hline 14 \\ 1580 \\ 9 \\ 1 \\ 1 \end{array}$	$ \begin{array}{r} \frac{1}{5} \\ \frac{3}{40} \\ 40 \\ 45 \\ 19372 \\ \hline 6561 \\ 9017 \\ 3168 \\ 35 \\ 384 \end{array} $	$\begin{array}{c} \frac{9}{40} \\ -\frac{56}{15} \\ -\frac{25360}{2187} \\ -\frac{355}{33} \\ 0 \end{array}$	$\begin{array}{r} -\frac{32}{9} \\ \underline{64448} \\ \underline{6561} \\ \underline{46732} \\ \underline{5247} \\ \underline{500} \\ \underline{1113} \end{array}$	$-\frac{212}{729} \\ \frac{49}{176} \\ \frac{125}{192}$	$- \frac{5103}{18656} \\ - \frac{2187}{6784}$	11 84		
$/b^{(5)} / b^{(4)}$	$   \begin{array}{r}     \frac{35}{384} \\     \underline{5179} \\     \overline{57600}   \end{array} $	0	$\begin{array}{r} 500 \\ \hline 1113 \\ 7571 \\ \hline 16695 \end{array}$	125 192	$-\frac{2187}{6784} \\ \frac{393}{640}$	$-\frac{\frac{11}{84}}{\frac{92097}{339200}}$	$0 \\ \frac{187}{2100}$	1 40

► Step size correction (Dormand Prince):

$$h_{new} = 0.9 h_{old} \left[ \frac{\delta}{||\mathbf{E}||} \right]^{\frac{1}{p+1}}$$

► Step size correction (Dormand Prince):

$$h_{new} = 0.9 h_{old} \left[ \frac{\delta}{||\mathbf{E}||} \right]^{\frac{1}{p+1}}$$

► Step size correction (Me):

$$h_{new} = \min_{j} \left( 0.9 h_{old} \left[ \frac{\delta_{j}}{\mathrm{E}_{j}} \right]^{\frac{1}{p}} \right), \quad \mathrm{E}_{j} > \delta_{j} : \mathrm{reject}$$

$$\delta_{j} = \min(\delta_{abs}, |\mathbf{X}_{j}(t_{i})|\delta_{rel}) \sqrt{\frac{h_{old}}{T}}.$$

► Step size correction (Dormand Prince):

$$h_{new} = 0.9 h_{old} \left[ \frac{\delta}{||\mathbf{E}||} \right]^{\frac{1}{p+1}}$$

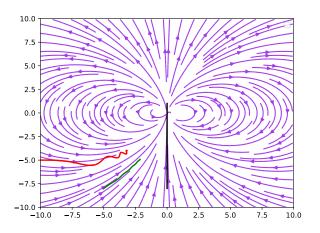
► Step size correction (Me):

$$h_{new} = \min_{j} \left( 0.9 h_{old} \left[ \frac{\delta_{j}}{\mathrm{E}_{j}} \right]^{\frac{1}{p}} \right), \quad \mathrm{E}_{j} > \delta_{j} : \mathrm{reject}$$

$$\delta_{j} = \min(\delta_{abs}, |\mathbf{X}_{j}(t_{i})|\delta_{rel}) \sqrt{\frac{h_{old}}{T}}.$$

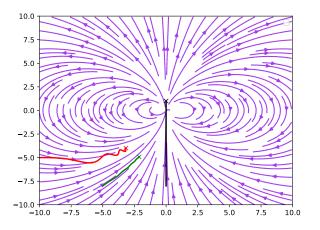
➤ Scale by step size  $\frac{h_{old}}{T}$ , "Fail safe"  $\sqrt{\dots}$ Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and Applied Mathematics

#### Does it work?

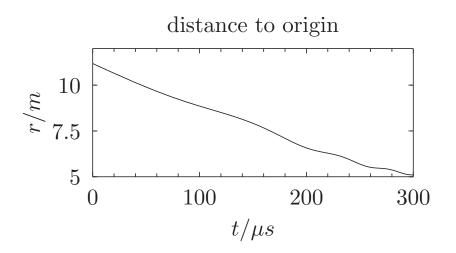


My version relative and absolute error  $10^{-6}$ . 94 steps (+ 14 rejected).

#### Does it work?

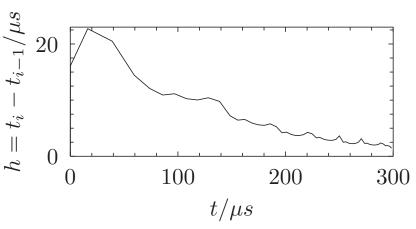


Odeint library relative and absolute error  $10^{-7}$ . 92 steps.

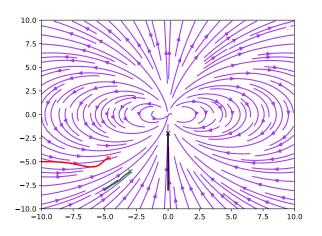


#### Does it work?

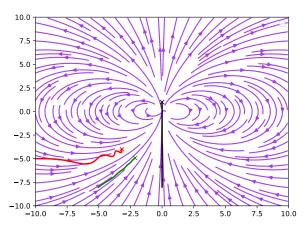
# Adaptive timesteps



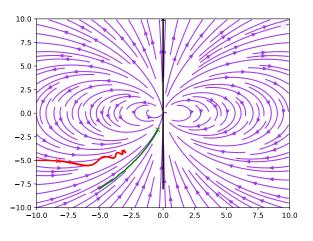
My version, adaptive step size.



 $T = 0.2 \, \text{ms}$ 

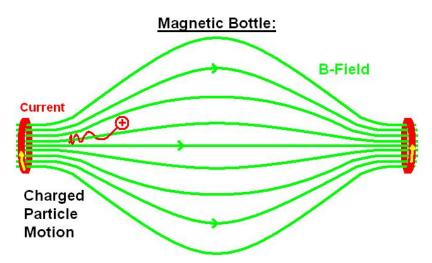


 $T = 0.3 \, \text{ms}$ 

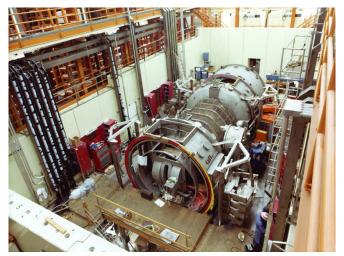


 $T = 0.6 \, \text{ms}$ 

Magnetic "mirror" or "bottle".



Wikipedia user WikiHelper2134, Public Domain.



Tandem Mirror Experiment, The Lawrence Livermore National Laboratory, 1978.

#### Conclusion

► Simulating particles in electric and magnetic fields.

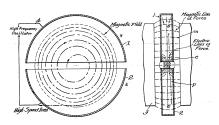
#### Conclusion

- ► Simulating particles in electric and magnetic fields.
- Numerically solving ordinary differential equations
- ► Reasonable agreement with known results.
- ► Can easily be generalized to other systems.

#### Conclusion

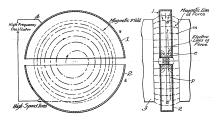
- ► Simulating particles in electric and magnetic fields.
- Numerically solving ordinary differential equations
- ► Reasonable agreement with known results.
- ► Can easily be generalized to other systems.
- Limitation, simulations are not experiments.

► Electric field accelerates, magnetic contains.



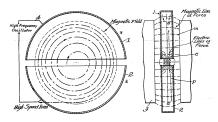
Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

- ► Electric field accelerates, magnetic contains.
- ► Single gab, oscillating field.



Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

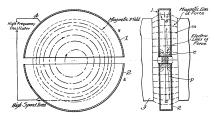
- ► Electric field accelerates, magnetic contains.
- ► Single gab, oscillating field.
- Uses classical Cyclotron frequency



Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

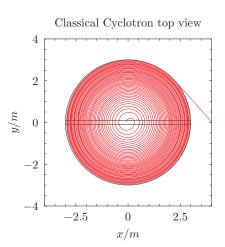
- ► Electric field accelerates, magnetic contains.
- ► Single gab, oscillating field.
- Uses classical Cyclotron frequency
- Analytical final speed, in principle path.

$$\frac{R|q|B}{m} = v_{\perp}$$

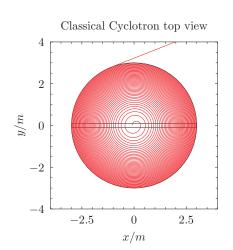


Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

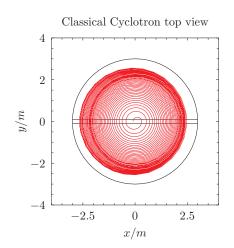
► with fixed step size, looks bad 4999 points



- ▶ with fixed step size, looks bad 4999 points
- my adabtive method, error:  $10^{-6}$
- ► Looks great but 2 million points.



- with fixed step size, looks bad 4999 points
- my adabtive method, error:  $10^{-6}$
- ► Looks great but 2 million points.
- ► odeint library



- ▶ with fixed step size, looks bad 4999 points
- ► my adabtive method, error: 10<sup>-6</sup>
- ► Looks great but 2 million points.
- ► odeint library
- ▶ non-continuous ode are bad.

