# How physics simulations work: simulating particles in electric and magnetic fields

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# How physics simulations work: simulating particles in electric and magnetic fields

Introduction

Theory and physical background

Euler's Method and the 4th order Runge-Kutta Method Euler's Method Higher order Runge-Kutta methods Demonstration, particles in a solenoid

Embedded algorithms, and adaptive step size Dormand Prince 5 (4) method Demonstration: magnetic dipole

Conclusion

► Numerical simulations are important.

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- Demonstration, charged particles in electric and magnetic fields.
- Analytically known and not.
- ► Simulations are not experiments!

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- ► The Lorentz force+N2:

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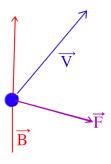
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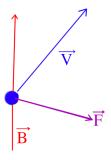
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- ▶ Only 1 particle! so pre-programmed depending on the setup.
- ► Could use potentials  $\phi(\vec{r},t)$   $\vec{A}(\vec{r},t)$  and Hamiltonian, or Lagrangian.
- ▶ Other systems would have other differential equations.



► Magnetic forces do no work:

$$dW_{\vec{B}} = \vec{F}_B \cdot d\vec{r} \propto (\vec{v} \times \vec{B}) \cdot \vec{v} = 0.$$

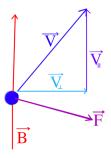


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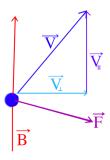
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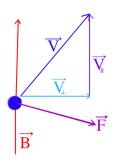
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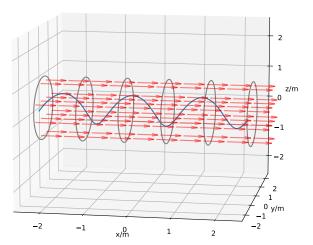
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- ► Same as Centripetal force: Cyclotron motion
- Cyclotron radius and frequency:

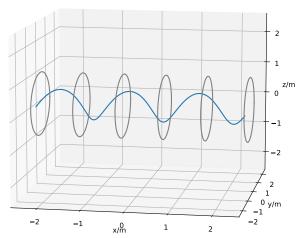
$$R = \frac{v_{\perp}m}{|a|B}$$
  $\omega_c = \frac{|q|B}{m}$ .





#### "Cyclotron motion"

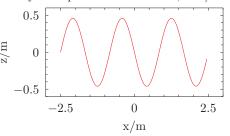
Solenoid with N=1000 turns per m, I=5 A, r=1 m,  $|\vec{B}|\approx 6$  mT. Proton with  $E_{kin}=1$  MeV/c<sup>2</sup> ( $|v|\approx 3.195\times 10^5$  m/s)

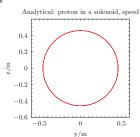


"Cyclotron motion"

$$R pprox 0.5 \, \mathrm{m} \, \mathrm{sin}( heta)$$
  $T = \frac{2\pi}{\omega_c} pprox 10 \, \mu \mathrm{s}$ 

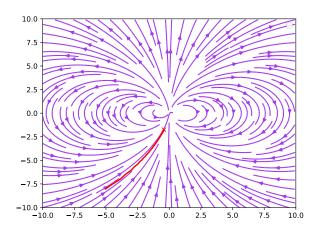
Analytical: proton in a solenoid, side/front-view





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$$R pprox 0.5 \, \mathrm{m} \, \mathrm{sin}( heta) \quad T = rac{2\pi}{\omega_c} pprox 10 \, \mathrm{\mu s}$$



(Actually from my simulation)

- ► Sources: Zeigler et al. Theory of Modeling and Simulation (Third edition) chapter 3
- ► We have:

$$\ddot{\vec{r}}(t) = \frac{q}{m}(\dot{\vec{r}}(t) \times \vec{B}(\vec{r},t) + \vec{E}(\vec{r},t)).$$

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► Here:

$$\mathbf{X} = \begin{pmatrix} \vec{r} \\ \dot{\vec{r}} \end{pmatrix} \quad f_{ode}(\mathbf{X},t) = \begin{pmatrix} \dot{\vec{r}} \\ \frac{q}{m} (\dot{\vec{r}} \times \vec{B}(\vec{r},t) + \vec{E}(\vec{r},t)) \end{pmatrix}.$$

#### The ODE to solve

```
auto ODE = [...](const state_type Data, state_type &
   dDatadt, const double t){
    //Extract position and velocity from data
    vec pos = vec(Data[0],Data[1],Data[2]);
    vec V = vec(Data[3],Data[4],Data[5]);
    //Lorentz+Newtons 2nd law
    vec F = Charge*(Fields.get Efield(pos,t)+
        cross(V,Fields.get_Bfield(pos,t)));
    vec dVdt = F*Inv mass;
    //Save derivative of data
    dDatadt[0]=V.x;dDatadt[1]=V.y;dDatadt[2]=V.z;
    dDatadt[3]=dVdt.x;Datadt[4]=dVdt.y;dDatadt[5]=
   dVdt.z;
    };
```

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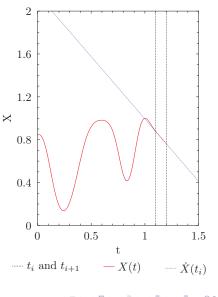
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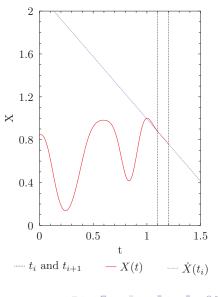
• We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .

► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



- We know only  $\mathbf{X}(t_0)$  and  $t_0$  and  $f_{ode}$ .
- Can we find  $\mathbf{X}(t_0 + h)$  for h > 0?.

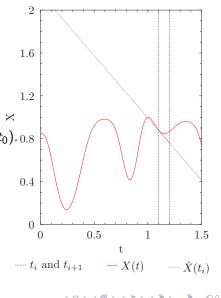
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- ightharpoonup Can we find  $\mathbf{X}(t_0+h)$  for h > 0?.
- ► Euler's Method:

$$\mathbf{X}(t_0+h) = \mathbf{X}(t_0) + hf_{ode}(\mathbf{X}(t_0), t_0)._{0.8}$$

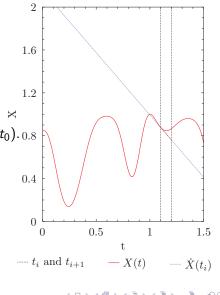
- ► Multiple steps  $t_0, t_1 = t_0 + h, \ldots t_i, \ldots$
- ► Bernard P. Zeigler et al. Theory of Modeling and Simulation (Third edition), chapter 3



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- Multiple steps  $t_0, t_1 = t_0 + h, \ldots t_i, \ldots$
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- Argument suggests we need  $f_{ode}$ ,  $f_{ode}$ , ..., we don't!

#### The Runge Kutta steppers

► In general:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t_{i+1}} \dot{\mathbf{X}}(t) dt.$$

► L. Zheng, X. Zhang, Modeling and Analysis of Modern Fluid Problems, 2017, chapter 8:

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- ▶ More generally, (Explicit and single step), Runge-Kutta family:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \int_{t_i}^{t'} f_{ode}(\mathbf{X}(t), t) dt + \ldots \int_{t^{(m)}}^{t_{i+1}} f_{ode}(\mathbf{X}(t), t) dt.$$



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$$= \sum_{j=1}^{m} h_j f_{ode}(\mathbf{X}(\tau_j), \tau_j).$$

- ▶ Guess  $\tau_1 = t_i$  use  $f_{ode}(\mathbf{X}(t_i), t_i)$  to approximate  $\mathbf{X}(\tau_2)$  to approximate  $\mathbf{X}(\tau_3)$  etc.
- ► L. Zheng, X. Zhang, Modeling and Analysis of Modern Fluid Problems, 2017, chapter 8:

More commonly written:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j \mathbf{K}_j$$

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► Estimate **K**<sub>j</sub>:

$$\mathbf{K}_{1} = f_{ode}(\mathbf{X}(t_{i}), t_{i})$$
 $\mathbf{K}_{2} = f_{ode}(\mathbf{X}(t_{i}) + ha_{21}\mathbf{K}_{1}, t_{i} + c_{2}h)$ 
 $\mathbf{K}_{3} = f_{ode}(\mathbf{X}(t_{i}) + ha_{31}\mathbf{K}_{1} + ha_{32}\mathbf{K}_{2}, t_{i} + c_{3}h)$ 
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- ▶ Taylor series analogy: Local error  $h^{p+1}$ , global  $h^p$ .
- ► Martha L. Abell, James P. Braselton, Differential Equations with Mathematica (Fourth Edition), 2016:

# The General explicit Runge Kutta method

► Expressed in Butcher tableu:

# Higher order methods

▶ 2nd order (Heun's method):

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \frac{h}{2}(\mathbf{k}_1 + \mathbf{k}_2)$$

$$\mathbf{k}_1 = f_{ode}(\mathbf{X}(t_i), t_i)$$

$$\mathbf{k}_2 = f_{ode}(\mathbf{X}(t_i) + h\mathbf{k}_1, t_i + h).$$

# Higher order methods

➤ 4th order, often simply called the Runge Kutta method:

$$\mathbf{X}(t_{i+1}) - \mathbf{X}(t_i) = \frac{h}{6}(\mathbf{k}_1 + 2\mathbf{k}_2 + 2\mathbf{k}_3 + \mathbf{k}_4)$$

$$\mathbf{k}_1 = f_{ode}(\mathbf{X}(t_i), t_i)$$

$$\mathbf{k}_2 = f_{ode}(\mathbf{X}(t_i) + \frac{h}{2}\mathbf{k}_1, t_i + \frac{h}{2})$$

$$\mathbf{k}_3 = f_{ode}(\mathbf{X}(t_i) + \frac{h}{2}\mathbf{k}_2, t_i + \frac{h}{2})$$
Wikipedia-user HilberTraum,
$$\mathbf{k}_4 = f_{ode}(\mathbf{X}(t_i) + h\mathbf{k}_3, t_i + h). \text{ published under creative commins: CC BY-SA 4.0}$$

### **Euler Implementations**

```
state_type Data = Data0;
state type dDatadt;
size t time res = T/timestep;
for (size t i = 1; i < time res; ++i)</pre>
{
    double t=i*dt;
    ODE(Data,dDatadt,t);
    //Euler time evolution
    Data+=timestep*dDatadt;
    save_step( Data , i*timestep );
};
```

# RK4 Implementations (1/2)

```
state_type Data = Data0;
state_type temp=Data0;
state_type K1,K2,K3,K4;
size_t time_res = T/timestep;
for (size t i = 1; i < time res; ++i)</pre>
{
    double t=i*timestep;
    //substep 1
    ODE(Data, K1,t);
    //substep 2
    temp=Data+timestep*K1/2;
    ODE(temp,K2,t+timestep/2);
```

# RK4 Implementations (2/2)

```
//substep 3
temp=Data+timestep*K2/2;
ODE(temp,K3,t+timestep/2);
//substep 4
temp=Data+timestep*K3;
ODE(temp,K4,t+timestep);
//Read data
Data+=timestep*(K1+2.0*K2+2.0*K3+K4)/6.0;
save_step( Data , i*timestep );}
```

▶ Test, same proton in a solenoid use  $\theta = 60^{\circ}$  reference, had:

$$R pprox 0.5\,\mathrm{m}\,\mathrm{sin}( heta) pprox 0.45\,\mathrm{m}$$
  $T = rac{2\pi}{\omega_c} pprox 10\,\mathrm{\mu s}.$ 

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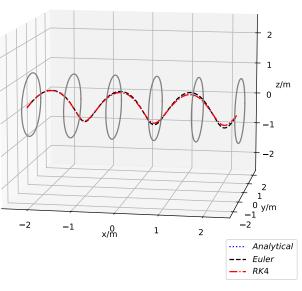
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- ► Compare Analytic, Euler, Runge-Kutta 4.
- ► Consider  $\theta = 60^{\circ}$  at different h.
- Check error on  $|\vec{v}|$ ,  $R = \sqrt{y^2 + z^2}$ .

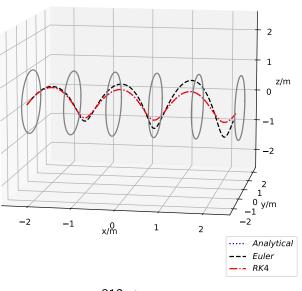
# At a glance, 3D view

$$h = t_{i+1} - t_i = 0.01 \, \mu s$$



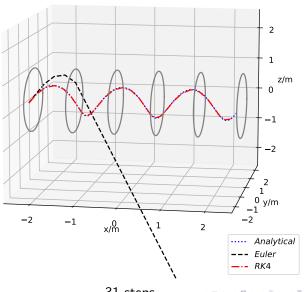
# At a glance, 3D view

$$h = t_{i+1} - t_i = 0.1 \,\mu s$$

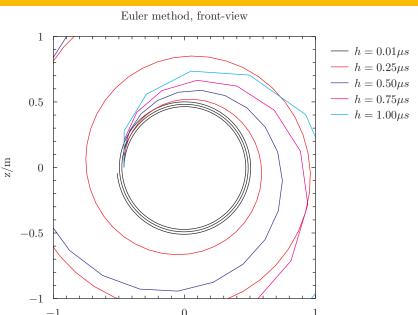


# At a glance, 3D view

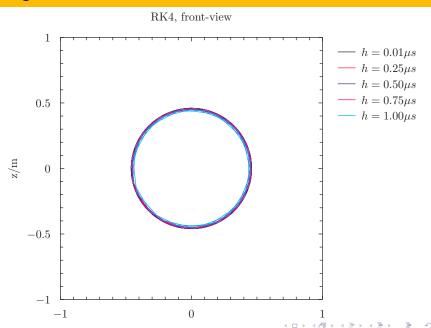
$$h = t_{i+1} - t_i = 1.0 \,\mu s$$



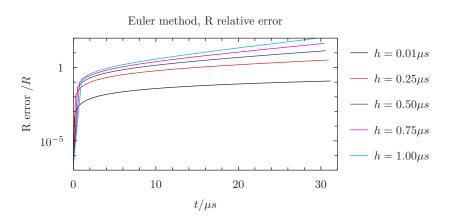
## At a glance, front view, no border



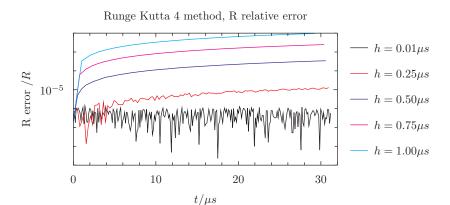
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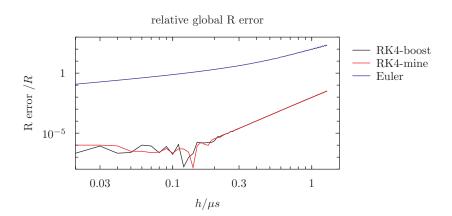
#### Constant radius?



### Constant radius?

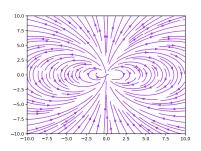


### Error as function of h

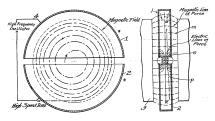


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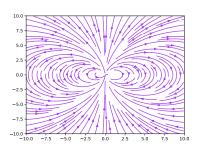


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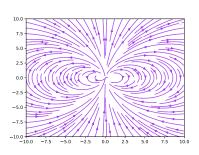
Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

- ► Usually don't know error.
- ► Hard to pick *h* , and may change:
- ► Inhomogeneous fields (here, a true dipole)
- ► Time dependent fields (here the cyclotron, bad example)
- ► Let the computer pick *h* for error.



► True dipole:

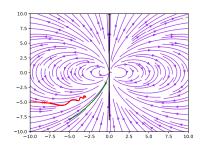
$$\vec{B} = \frac{\mu_0}{4\pi r^3} \left[ 3\hat{\vec{r}} (\vec{m} \cdot \hat{\vec{r}} - \vec{m}) \right].$$



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► Weak field, little change, long steps.

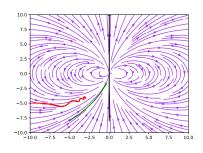


Arbitrarily set  $\frac{\mu_0}{4\pi}|\vec{m}|=0.155\,\mathrm{T/m^3}~.$  Protons with speed around  $10\,000\,\mathrm{m/s}$ 

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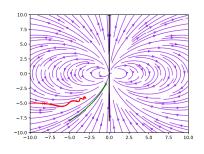


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- strong field, large change, short steps.
- ► No analytical solution.



Arbitrarily set  $\frac{\mu_0}{4\pi}|\vec{m}|=0.155\,\mathrm{T/m^3}$  . Protons with speed around  $10\,000\,\mathrm{m/s}$ 

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- ► Adjust step size to keep the error(s) small (Implementations differ!).

Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and AppliMathematics

(1)

ode45 in Matlab, scipy.solve\_ivp in Python, RungeKutta\_dopri5 in boost::odeint.

$$\mathbf{X}^{(5)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(5)} \mathbf{k}_j$$
 $\mathbf{X}^{(4)}(t_{i+1}) - \mathbf{X}(t_i) = h \sum_{j=1}^m b_j^{(4)} \mathbf{k}_j.$ 

$$\mathbf{k}_j = f_{ode}(\mathbf{X}(t_i) + h \sum_{k}^{j-1} a_{kj} \mathbf{k}_j, t_i + c_3 h).$$

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7 k<sub>j</sub>'s (actually 6 by clever re-usage).
Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and Applied Mathematics

# Butcher Tableu of Dormand Prince (5) 4

Ci	a <sub>ij</sub>							
$\begin{array}{c} 0 \\ 1 \\ 153 \\ \hline 14 \\ 1580 \\ 9 \\ 1 \\ 1 \end{array}$	$ \begin{array}{r} \frac{1}{5} \\ \frac{3}{40} \\ 40 \\ 45 \\ 19372 \\ \hline 6561 \\ 9017 \\ 3168 \\ 35 \\ 384 \end{array} $	$\begin{array}{c} \frac{9}{40} \\ -\frac{56}{15} \\ -\frac{25360}{2187} \\ -\frac{355}{33} \\ 0 \end{array}$	$\begin{array}{r} -\frac{32}{9} \\ \underline{64448} \\ \underline{6561} \\ \underline{46732} \\ \underline{5247} \\ \underline{500} \\ \underline{1113} \end{array}$	$-\frac{212}{729} \\ \frac{49}{176} \\ \frac{125}{192}$	$- \frac{5103}{18656} \\ - \frac{2187}{6784}$	11 84		
$/b^{(5)} / b^{(4)}$	$   \begin{array}{r}     \frac{35}{384} \\     \underline{5179} \\     \overline{57600}   \end{array} $	0	$\begin{array}{r} 500 \\ \hline 1113 \\ 7571 \\ \hline 16695 \end{array}$	125 192	$-\frac{2187}{6784} \\ \frac{393}{640}$	$-\frac{\frac{11}{84}}{\frac{92097}{339200}}$	$0 \\ \frac{187}{2100}$	1 40

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$$h_{new} = 0.9 h_{old} \left[ \frac{\delta}{||\mathbf{E}||} \right]^{\frac{1}{p+1}}$$

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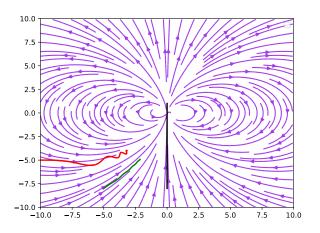
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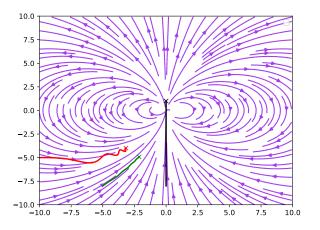
➤ Scale by step size  $\frac{h_{old}}{T}$ , "Fail safe"  $\sqrt{\dots}$ Dormand, J. R.; Prince, P. J. (1980), "A family of embedded Runge-Kutta formulae", Journal of Computational and Applied Mathematics

#### Does it work?

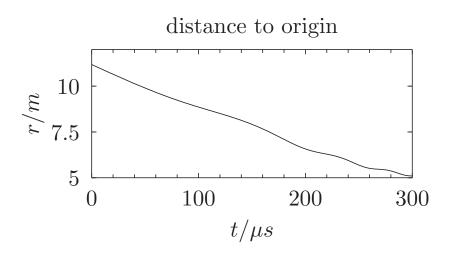


My version relative and absolute error  $10^{-6}$ . 94 steps (+ 14 rejected).

#### Does it work?

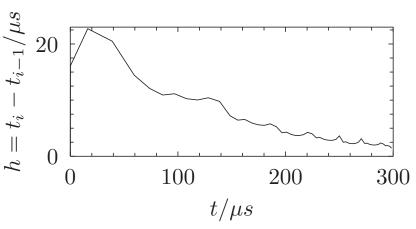


Odeint library relative and absolute error  $10^{-7}$ . 92 steps.

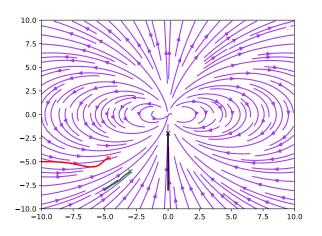


### Does it work?

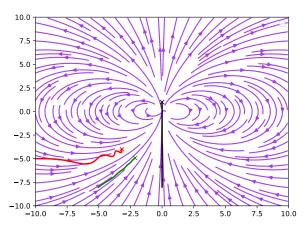
# Adaptive timesteps



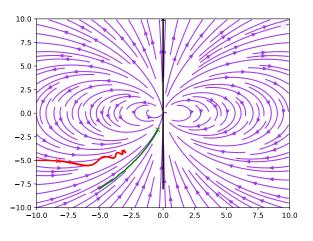
My version, adaptive step size.



 $T = 0.2 \, \text{ms}$ 

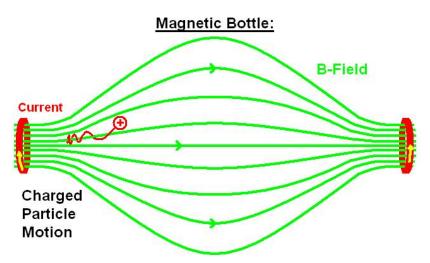


 $T = 0.3 \, \text{ms}$ 

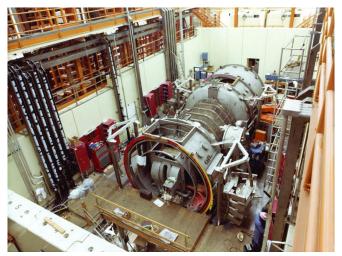


 $T = 0.6 \, \text{ms}$ 

Magnetic "mirror" or "bottle".



Wikipedia user WikiHelper2134, Public Domain.



Tandem Mirror Experiment, The Lawrence Livermore National Laboratory, 1978.

#### Conclusion

► Simulating particles in electric and magnetic fields.

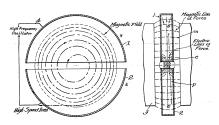
#### Conclusion

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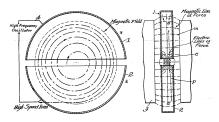
- ► Simulating particles in electric and magnetic fields.
- Numerically solving ordinary differential equations
- ► Reasonable agreement with known results.
- ► Can easily be generalized to other systems.
- Limitation, simulations are not experiments.

► Electric field accelerates, magnetic contains.



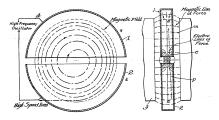
Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

- ► Electric field accelerates, magnetic contains.
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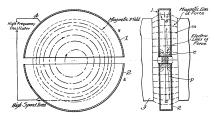
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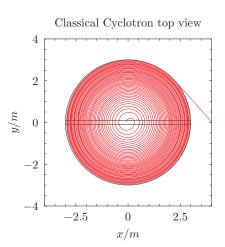
- ► Electric field accelerates, magnetic contains.
- ► Single gab, oscillating field.
- Uses classical Cyclotron frequency
- ► Analytical final speed, in principle path.

$$\frac{R|q|B}{m} = v_{\perp}$$

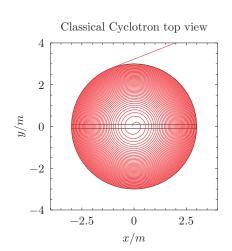


Ernest O. Lawrence, 1934, U.S. Patent 1,948,384; image in Public Domain.

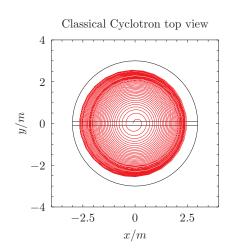
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