

# 1 Dirac–Frenkel variational principal (DFVP)

Time dependant Shrödinger equation:

$$i\hbar \frac{\partial |\Phi_{ex}\rangle}{\partial t} = \hat{H} |\Phi_{ex}\rangle \quad (1)$$

We will consider the following mean value:

$$W = \frac{\langle \Phi | \hat{H} - i\hbar \frac{\partial}{\partial t} | \Phi \rangle}{\langle \Phi | \Phi \rangle}$$

which, if calculated on exact solutions  $|\Phi_{ex}\rangle$  of (1), equals to zero, and its variation:  $\delta W = 0$

If  $|\Phi_{ex}\rangle$  is an exact solution of (1), then the norm conservation condition is satisfied:

$$\frac{\partial \langle \Phi_{ex} | \Phi_{ex} \rangle}{\partial t} = 0$$

But we do not know beforehand about norm of arbitrary function  $|\Phi\rangle$ . We can consider mean value of  $i\hbar \frac{\partial}{\partial t}$ :

$$\langle \omega \rangle = \frac{\langle \Phi | i\hbar \frac{\partial}{\partial t} | \Phi \rangle}{\langle \Phi | \Phi \rangle}$$

Let us calculate difference between  $\langle \omega \rangle$  and its complex conjugate:

$$\langle \omega \rangle - \langle \omega \rangle^* = \frac{i\hbar (\langle \Phi | \frac{\partial \Phi}{\partial t} \rangle + \langle \frac{\partial \Phi}{\partial t} | \Phi \rangle)}{\langle \Phi | \Phi \rangle}$$

$$\frac{i}{\hbar} (\langle \omega \rangle^* - \langle \omega \rangle) = \frac{\partial}{\partial t} \ln \langle \Phi | \Phi \rangle$$

Now we will solve this equation to obtain time dependance of norm:

$$N(t) = N(0)e^P, \text{ where } P = \frac{i}{\hbar} \int_0^t (\langle \omega \rangle^* - \langle \omega \rangle) dt'$$

We can see, that norm conservation condition is satisfied, when  $P = 0$ ,  $\langle \omega \rangle \in \mathbb{R}$  and hence  $W \in \mathbb{R}$ . But if it is an approximate solution we can not guarantee conservation of norm!

But let us assume, that we can construct function  $|\Phi'\rangle$ , that differs from  $|\Phi\rangle$  by angular multiplier:

$$|\Phi'\rangle = |\Phi\rangle \cdot e^Q, \text{ where } Q = \frac{i}{\hbar} \int_0^t \alpha(t') dt', \alpha(t') \in \mathbb{C}$$

We need to find parameter  $\alpha(t)$ , so that norm  $\langle \Phi' | \Phi' \rangle$  is conserved. We will consider a mean value once again:

$$\langle \omega' \rangle = \frac{\langle \Phi' | i\hbar \frac{\partial}{\partial t} | \Phi' \rangle}{\langle \Phi' | \Phi' \rangle} = \frac{\langle \Phi | e^{-Q} i\hbar \cdot \frac{i}{\hbar} \alpha(t) e^Q + i\hbar \frac{\partial}{\partial t} | \Phi \rangle}{\langle \Phi | \Phi \rangle} = \langle \omega \rangle - \alpha$$

If  $\langle \omega' \rangle \in \mathbb{R}$ , then:

$$0 = \langle \omega' \rangle - \langle \omega' \rangle^* = \langle \omega \rangle - \langle \omega \rangle^* - 2i\text{Im}(\alpha)$$

$$i\text{Im}(\alpha) = \frac{1}{2} (\langle \omega \rangle - \langle \omega \rangle^*)$$

$$\alpha = \text{Re}(\alpha) + i\text{Im}(\alpha) = \text{Re}(\alpha) + \frac{1}{2} (\langle \omega \rangle - \langle \omega \rangle^*)$$

$$\langle \omega' \rangle = \langle \omega \rangle - \alpha = \langle \omega \rangle - \frac{1}{2} (\langle \omega \rangle - \langle \omega \rangle^*) - \text{Re}(\alpha) = \frac{1}{2} (\langle \omega \rangle + \langle \omega \rangle^*) - \text{Re}(\alpha)$$

$$|\Phi'\rangle = |\Phi\rangle \cdot e^Q = |\Phi\rangle \cdot e^R \cdot e^{-0.5P}, \text{ where } R = \frac{i}{\hbar} \int_0^t \text{Re}(\alpha(t')) dt'$$

As  $N(t) = N(0) \cdot e^P$ , we have:

$$|\Phi'\rangle = |\Phi\rangle \cdot e^R \cdot \left( \frac{N(0)}{N(t)} \right)^{1/2}$$

As we've discussed previously, for an exact solution of (1)  $|\Phi_{ex}\rangle$  mean value  $W$  equals to zero. Let us consider mean values  $W'$ , calculated on function  $|\Phi'\rangle$ :

$$W' = \langle H \rangle - \langle \omega' \rangle = \langle H \rangle - \frac{1}{2} (\langle \omega \rangle + \langle \omega \rangle^*) + \text{Re}(\alpha)$$

Now we need to understand, what  $\alpha$  should be to make  $W'$  equal to zero:

$$\text{Re}(\alpha) = -\langle H \rangle + \frac{1}{2} (\langle \omega \rangle + \langle \omega \rangle^*)$$

$$\alpha = -\langle H \rangle + \frac{1}{2} (\langle \omega \rangle + \langle \omega \rangle^*) + \frac{1}{2} (\langle \omega \rangle - \langle \omega \rangle^*) = \langle \omega \rangle - \langle H \rangle = -W$$

But if  $\alpha = -W$ , we will obtain:

$$|\Phi'\rangle = |\Phi\rangle \cdot e^{-\frac{i}{\hbar} \int_0^t W dt}$$

$$\langle \Phi' | \Phi' \rangle_t = \langle \Phi | \Phi \rangle_t \cdot e^{-\frac{i}{\hbar} \int_0^t (W - W^*) dt} = \langle \Phi | \Phi \rangle_0 \cdot e^P \cdot e^{-P} = \langle \Phi | \Phi \rangle_0$$

So we have built functions  $|\Phi'\rangle$ , that have conserved norm and lead to zero  $W'$ .

Let us consider for simplicity function  $|\Psi\rangle$  to be from the family of functions  $|\Phi'\rangle$ . This function has conserved norm, and mean value  $W$ , calculated on this

function, is real. As the norm of  $|\Psi\rangle$  doesn't change with time, we can write the following equation:

$$\delta\langle\Psi|\Psi\rangle = \langle\delta\Psi|\Psi\rangle + \langle\Psi|\delta\Psi\rangle = 0$$

We will consider only variations  $|\delta\Psi\rangle$ , that are orthogonal to  $|\Psi\rangle$ . Then:

$$\langle\delta\Psi|\Psi\rangle = 0, \quad \langle\Psi|\delta\Psi\rangle = 0 \quad (2)$$

Now we can write down variation  $\delta W$ . For simplicity, we shall denote  $\langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle$  as  $A$  and  $\langle\Psi|\Psi\rangle$  as  $B$ :

$$W = \frac{A}{B}, \quad \delta W = \frac{B \cdot \delta A - A \cdot \delta B}{B^2} = \frac{\delta A - W\delta B}{B}$$

$$\begin{aligned} \delta A - W\delta B &= \langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle + \langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\delta\Psi\rangle - W\langle\delta\Psi|\Psi\rangle - W\langle\Psi|\delta\Psi\rangle = \\ &= \langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle + \langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\delta\Psi\rangle - W\delta\langle\Psi|\Psi\rangle \end{aligned}$$

As  $\delta\langle\Psi|\Psi\rangle = 0$ ,  $\delta W = 0$ , we obtain:

$$\langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle = 0 \quad \text{— Dirac–Frenkel variational principal} \quad (3)$$

$$\langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\delta\Psi\rangle = 0$$

We shall consider the second equation:

$$\begin{aligned} \langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\delta\Psi\rangle &= \left\langle \left( \hat{H} - i\hbar\frac{\partial}{\partial t} \right) \Psi \middle| \delta\Psi \right\rangle - i\hbar \left\langle \frac{\partial\Psi}{\partial t} \middle| \delta\Psi \right\rangle - i\hbar \left\langle \Psi \middle| \frac{\partial}{\partial t} \delta\Psi \right\rangle = \\ &= \langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle^* - i\hbar\frac{\partial}{\partial t}\langle\Psi|\delta\Psi\rangle = 0 \end{aligned}$$

Thus, the second equation is a mere consequence of Dirac–Frenkel variational principal (3) and condition (2).

Previously we have discussed the case of orthogonal variation  $|\delta\Psi\rangle$ . Arbitraty variations  $|\delta\Psi\rangle$  can be rewritten as sum of  $|\Psi\rangle$  and  $|\delta_\perp\Psi\rangle$ :

$$|\delta\Psi\rangle = c_{||}|\Psi\rangle + c_\perp|\delta_\perp\Psi\rangle$$

Variation of  $W$  will have the following look:

$$\begin{aligned} \delta W &= \langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\Psi\rangle + \langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t}|\delta\Psi\rangle - W\langle\delta\Psi|\Psi\rangle - W\langle\Psi|\delta\Psi\rangle = \\ &= \langle\delta\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t} - W|\Psi\rangle + \langle\Psi|\hat{H} - i\hbar\frac{\partial}{\partial t} - W|\delta\Psi\rangle = \end{aligned}$$

$$= 2\text{Re}(c_{||})\langle\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}-W|\Psi\rangle - c_{\perp}^*\langle\delta_{\perp}\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}-W|\Psi\rangle - c_{\perp}\langle\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}-W|\delta_{\perp}\Psi\rangle$$

The first term equals to zero, because:

$$W = \frac{\langle\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}|\Psi\rangle}{\langle\Psi|\Psi\rangle}, \quad \langle\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}-W|\Psi\rangle = 0$$

The last two terms are equal to zero due to Dirac–Frenkel variational principle.

We have discussed arbitrary variations of wave function  $|\Psi\rangle$ , that don't affect parameters of hamiltonian  $\hat{H}$ . To write DFVP in the most genral form, we need to consider variations  $|\delta\Psi\rangle$  of the following form:

$$|\delta\Psi\rangle = |\frac{\partial\Psi}{\partial\varepsilon}\rangle\delta\varepsilon, \quad \hat{H} = \hat{H}(\varepsilon)$$

To preserve the form of DFVP for that kind of variation, we need to introduce one more condition, that should be met by approximate wave function  $|\Psi\rangle$ . Let us consider the following equation:

$$\langle\Phi_{ex}(t)|\frac{\partial\hat{H}}{\partial\varepsilon}|\Phi_{ex}(t)\rangle = i\hbar\langle\Phi_{ex}(t)|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle$$

where  $|\Phi_{ex}(t)\rangle$  — exact solution of time dependant Shrödinger equation. That equation is a statement of time dependant Hellman–Feynman theorem (*tdHFT*). To prove it, we will consider the following matrix element:

$$\begin{aligned} \frac{\partial}{\partial\varepsilon}\langle\Phi_{ex}(t)|\hat{H}|\Phi_{ex}(t)\rangle &= \langle\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}|\hat{H}|\Phi_{ex}(t)\rangle + \langle\Phi_{ex}(t)|\frac{\partial\hat{H}}{\partial\varepsilon}|\Phi_{ex}(t)\rangle + \langle\Phi_{ex}(t)|\hat{H}|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle = \\ &= i\hbar\langle\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}|\frac{\partial\Phi_{ex}(t)}{\partial t}\rangle - i\hbar\langle\frac{\partial\Phi_{ex}(t)}{\partial t}|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle + \langle\Phi_{ex}(t)|\frac{\partial\hat{H}}{\partial\varepsilon}|\Phi_{ex}(t)\rangle \\ \langle\Phi_{ex}(t)|\frac{\partial\hat{H}}{\partial\varepsilon}|\Phi_{ex}(t)\rangle &= i\hbar\frac{\partial}{\partial\varepsilon}\langle\Phi_{ex}(t)|\frac{\partial\Phi_{ex}(t)}{\partial t}\rangle + i\hbar\langle\frac{\partial\Phi_{ex}(t)}{\partial t}|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle - i\hbar\langle\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}|\frac{\partial\Phi_{ex}(t)}{\partial t}\rangle = \\ &= i\hbar\langle\Phi_{ex}(t)|\frac{\partial^2\Phi_{ex}(t)}{\partial\varepsilon\partial t}\rangle + i\hbar\langle\frac{\partial\Phi_{ex}(t)}{\partial t}|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle = \\ &= i\hbar\frac{\partial}{\partial t}\langle\Phi_{ex}(t)|\frac{\partial\Phi_{ex}(t)}{\partial\varepsilon}\rangle \end{aligned}$$

Now we need to consider  $\delta W$  in terms of new type of variations:

$$\delta W = \frac{\delta A - W\delta B}{B}$$

As we remember,  $\delta B = 0$ . Thus, to ensure  $\delta W = 0$ , we need  $\delta A = 0$ :

$$\delta\langle\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}|\Psi\rangle = \langle\delta\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}|\Psi\rangle + \langle\delta\Psi|\hat{H}-i\hbar\frac{\partial}{\partial t}|\Psi\rangle^* - i\hbar\frac{\partial}{\partial t}\langle\Psi|\frac{\partial\Psi}{\partial\varepsilon}\rangle\delta\varepsilon + \langle\Psi|\frac{\partial\hat{H}}{\partial\varepsilon}|\Psi\rangle\delta\varepsilon$$

We need to make our function  $|\Psi\rangle$  to behave in such a way, that *tdHFT* will be valid. Then we will obtain DFVP.

## 2 Equations of motions in DFVP formalism

Let us assume, that function  $|\Psi\rangle$  can be spanned over linear combination of basis functions  $\{|\phi_k(\vec{\lambda})\rangle\}_{k=1}^N$ :

$$|\Psi\rangle = \sum_{k=1}^N C_k(t) |\phi_k(\vec{\lambda})\rangle, \dim\{\vec{\lambda}\} = M$$

Then we can consider a variation of  $|\Psi\rangle$ :

$$|\delta\Psi\rangle = \sum_{k=1}^N \left( \delta C_k |\phi_k\rangle + C_k \sum_{j=1}^M \left| \frac{\partial \phi_k}{\partial \lambda_{kj}} \right\rangle \delta \lambda_{kj} \right)$$

Thus, using Dirac–Frenkel variational principle, we will obtain:

$$\delta C_m^* \sum_{k=1}^N \langle \phi_m | \hat{H} - i\hbar \frac{\partial}{\partial t} | C_k \phi_k \rangle = 0$$

$$\delta \lambda_{mj}^* \sum_{k=0}^N C_m^* \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \right| \hat{H} - i\hbar \frac{\partial}{\partial t} | C_k \phi_k \rangle = 0$$

As variations are independant and arbitrary, we will get two sets of equations:

$$\sum_{k=1}^N \langle \phi_m | \hat{H} - i\hbar \frac{\partial}{\partial t} | C_k \phi_k \rangle = 0$$

$$\sum_{k=0}^N C_m^* \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \right| \hat{H} - i\hbar \frac{\partial}{\partial t} | C_k \phi_k \rangle = 0$$

Let us consider the first equation:

$$\sum_{k=1}^N C_k \langle \phi_m | \hat{H} | \phi_k \rangle - i\hbar \langle \phi_m | \phi_k \rangle \dot{C}_k - i\hbar \sum_{l=1}^M \langle \phi_m | \frac{\partial \phi_k}{\partial \lambda_{kl}} \rangle \dot{\lambda}_{kl} = 0$$

$$i\hbar \sum_{k=1}^N \mathbb{S}_{mk} \dot{C}_k = \sum_{k=1}^N \left( \mathbb{H}_{mk} - i\hbar \sum_{l=1}^M \langle \phi_m | \frac{\partial \phi_k}{\partial \lambda_{kl}} \rangle \dot{\lambda}_{kl} \right) C_k$$

$$i\hbar \mathbb{S} \dot{\vec{C}} = (\mathbb{H} - i\hbar \tau) \vec{C}$$

$$\dot{\vec{C}} = -\frac{i}{\hbar} \mathbb{S}^{-1} (\mathbb{H} - i\hbar \tau) \vec{C}$$

$$\dot{C}_k = -\frac{i}{\hbar} \sum_{n,r=1}^N \mathbb{S}_{kn}^{-1} (\mathbb{H}_{nr} - i\hbar\tau_{nr}) C_r$$

And the second one:

$$\begin{aligned} & \sum_{k=1}^N \left( C_m^* C_k \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \hat{H} \middle| \phi_k \right\rangle - i\hbar \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \phi_k \right\rangle C_m^* \dot{C}_k - i\hbar \sum_{l=1}^M C_m^* C_k \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \frac{\partial \phi_k}{\partial \lambda_{kl}} \right\rangle \dot{\lambda}_{kl} \right) = 0 \\ & \sum_{k=1}^N \left( C_m^* C_k \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \hat{H} \middle| \phi_k \right\rangle - \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \phi_k \right\rangle \sum_{n,r=1}^N \mathbb{S}_{kn}^{-1} \mathbb{H}_{nr} C_m^* C_r \right) = \\ & = \sum_{k=1}^N \left( -i\hbar \sum_{n,r=1}^N \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \phi_k \right\rangle \mathbb{S}_{kn}^{-1} \sum_{l=1}^M \left\langle \phi_n \middle| \frac{\partial \phi_r}{\partial \lambda_{rl}} \right\rangle \dot{\lambda}_{rl} C_m^* C_r + i\hbar \sum_{l=1}^M C_m^* C_k \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \frac{\partial \phi_k}{\partial \lambda_{kl}} \right\rangle \dot{\lambda}_{kl} \right) \end{aligned}$$

$$\rho_{mk} = C_m^* C_k$$

$$\mathbb{H}_{ml}^{(j0)} = \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \hat{H} \middle| \phi_k \right\rangle$$

$$\mathbb{S}_{mk}^{(j0)} = \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \phi_k \right\rangle$$

$$\mathbb{S}_{nr}^{(0l)} = \left\langle \phi_n \middle| \frac{\partial \phi_r}{\partial \lambda_{rl}} \right\rangle$$

$$\mathbb{S}_{mk}^{(jl)} = \left\langle \frac{\partial \phi_m}{\partial \lambda_{mj}} \middle| \frac{\partial \phi_k}{\partial \lambda_{kl}} \right\rangle$$

$$\begin{aligned} & \sum_{k=1}^N \rho_{mk} \mathbb{H}_{mk}^{(j0)} - \sum_{r=1}^N \rho_{mr} \left( \sum_{k,n=1}^N \mathbb{S}_{mk}^{(j0)} \mathbb{S}_{kn}^{-1} \mathbb{H}_{nr} \right) = \\ & = i\hbar \sum_{l=1}^M \left( \sum_{k=1}^N \rho_{mk} \mathbb{S}_{mk}^{(jl)} \dot{\lambda}_{kl} - \sum_{r=1}^N \rho_{mr} \left( \sum_{k,n=1}^N \mathbb{S}_{mk}^{(j0)} \mathbb{S}_{kn}^{-1} \mathbb{S}_{nr}^{(0l)} \right) \dot{\lambda}_{rl} \right) \\ & \sum_{k=1}^N \rho_{mk} \left( \mathbb{H}_{mk}^{(j0)} - (\mathbb{S}^{(j0)} \mathbb{S}^{-1} \mathbb{H})_{mk} \right) = i\hbar \sum_{l=1}^M \sum_{k=1}^N \rho_{mk} \left( \mathbb{S}_{mk}^{(jl)} - (\mathbb{S}^{(j0)} \mathbb{S}^{-1} \mathbb{S}^{(0l)})_{mk} \right) \dot{\lambda}_{kl} \end{aligned}$$

$$Y_m^j = \sum_{k=1}^N \rho_{mk} \left( \mathbb{H}_{mk}^{(j0)} - (\mathbb{S}^{(j0)} \mathbb{S}^{-1} \mathbb{H})_{mk} \right)$$

$$X_{mk}^{jl} = \rho_{mk} \left( \mathbb{S}_{mk}^{(jl)} - (\mathbb{S}^{(j0)} \mathbb{S}^{-1} \mathbb{S}^{(0l)})_{mk} \right)$$

$$i\hbar \dot{\Lambda} = \mathbb{X}^{-1} \mathbb{Y}, \text{ where } \Lambda \text{ is matrix}$$

Thus, we have two sets of equations:

$$\dot{\vec{C}} = -\frac{i}{\hbar}\mathbb{S}^{-1}(\mathbb{H} - i\hbar\tau)\vec{C}$$

$$\dot{\Lambda} = -\frac{i}{\hbar}\mathbb{X}^{-1}\mathbb{Y}$$

### 3 Gaussian wave packets

For basis functions  $|\phi_k\rangle$  we can choose Gaussian wave packets (GWP):

$$|g_k(\vec{q}_k, \vec{p}_k)\rangle = \left(\frac{\pi}{\omega}\right)^{1/4} e^{\sum_{\alpha=\{x,y,z\}} -\frac{1}{2}\omega(r_\alpha - q_{k\alpha})^2 + ip_{k\alpha}(r_\alpha - q_{k\alpha})}$$

We can transform parameters  $(\vec{q}_k, \vec{p}_k)$  to  $(\vec{\xi}_k, \vec{\eta}_k)$ :

$$\xi_{k\alpha} = \omega q_{k\alpha} + ip_{k\alpha}$$

$$\eta_{k\alpha} = \frac{1}{4} \left( \ln \left[ \frac{\omega}{\pi} \right] - 2\omega q_{k\alpha}^2 \right) - iq_{k\alpha} p_{k\alpha}$$

Thus we will obtain another representation of GWP:

$$|g_k(\vec{\xi}_k, \vec{\eta}_k)\rangle = e^{\sum_{\alpha=\{x,y,z\}} -\frac{1}{2}\omega r_\alpha^2 + \xi_{k\alpha} r_\alpha + \eta_{k\alpha}}$$