Solving Radial Schrödinger equation numerically

1 Radial Schrödinger equation

$$\frac{d^{2}\Psi\left(r\right)}{dr^{2}}+\left[\frac{2\mu E}{\hbar^{2}}-\frac{l\left(l+1\right)}{r^{2}}-U\left(r\right)\right]\Psi\left(r\right)=0\tag{1}$$

where,

$$U\left(r\right) = \frac{2\mu}{\hbar^2} V\left(r\right)$$

2 Dimensionless form

put, $r=xa_0$ where, $a_0=\frac{4\pi\epsilon_0\hbar^2}{2\mu e^2}\approx 0.592$ Å and x is dimensionless.

$$\frac{d^{2}\Psi\left(x\right)}{a_{0}^{2}dx^{2}}+\left\lceil \frac{2\mu E}{\hbar^{2}}-\frac{l\left(l+1\right)}{a_{0}^{2}x^{2}}-U\left(x\right)\right]\Psi\left(x\right)=0$$

multiplying, a_0 on both sides gives:

$$\begin{split} &\Longrightarrow \frac{d^2\Psi\left(x\right)}{dx^2} + \left[\frac{2\mu a_0^2E}{\hbar^2} - \frac{l\left(l+1\right)}{x^2} - a_0^2U\left(x\right)\right]\Psi\left(x\right) = 0 \\ &\Longrightarrow \frac{d^2\Psi\left(x\right)}{dx^2} + \left[\frac{E}{\frac{\hbar^2}{2\mu a_0^2}} - \frac{l\left(l+1\right)}{x^2} - W\left(x\right)\right]\Psi\left(x\right) = 0 \\ &\Longrightarrow \frac{d^2\Psi\left(x\right)}{dx^2} + \left[\epsilon - \frac{l\left(l+1\right)}{x^2} - W\left(x\right)\right]\Psi\left(x\right) = 0 \quad \left(\epsilon = \frac{E}{E_1}\right) \text{ where, } E_1 = -\frac{\hbar^2}{2\mu a_0^2} \approx -13.6 \text{ eV} \end{split}$$

where the negative energy represents bound states. For Hydrogen atom,

$$W(x) = a_0^2 U\left(r\right) = \frac{2\mu a_0^2}{\hbar^2} V\left(r\right) = \frac{2\mu a_0^2}{\hbar^2} \left(\frac{-e^2}{4\pi\epsilon_0 r}\right) = -\frac{2\mu a_0^2}{\hbar^2} \left(\frac{e^2}{4\pi\epsilon_0 a_0 x}\right) = -a_0 \frac{1}{\frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}} \frac{1}{x} = -\frac{2}{x} \quad \left(a_0 = \frac{4\pi\epsilon_0 \hbar^2}{2\mu e^2}\right) = -a_0 \frac{1}{\frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -a_0 \frac{1}{2\mu c^2} \frac{1}{x} = -\frac{2}{x} \quad \left(\frac{a_0}{2\mu c^2} + \frac{4\pi\epsilon_0 \hbar^2}{2\mu c^2}\right) = -\frac{2\mu c^2}{x} = -\frac{$$

Therefore, the final dimensionless radial Schrödinger equation for Hydrogen atom is:

$$\frac{d^{2}\Psi\left(x\right)}{dx^{2}} + \left[\epsilon - \frac{l\left(l+1\right)}{x^{2}} + \frac{2}{x}\right]\Psi\left(x\right) = 0\tag{2}$$

3 Numerov's algorithm

Discrete Taylor series expansion of Ψ_{i+1} about x_i

$$\Psi_{i+1} = \Psi_i + (x_{i+1} - x_i) \Psi' \Big|_{x_i} + \frac{(x_{i+1} - x_i)^2}{2!} \Psi'' \Big|_{x_i} + \frac{(x_{i+1} - x_i)^3}{3!} \Psi''' \Big|_{x_i} + \cdots$$

$$= \Psi_i + h \Psi'_i + \frac{h^2}{2!} \Psi''_i + \frac{h^3}{3!} \Psi'''_i + \frac{h^4}{4!} \Psi''''_i + \cdots$$
(3)

where,

$$\Psi'\Big|_{x_i} = \frac{d\Psi(r)}{dr}\Big|_{x_i} = \Psi'_i \quad \text{and } h = (x_{i+1} - x_i)$$

$$\tag{4}$$

Discrete Taylor series expansion of Ψ_{i-1} about x_i

$$\Psi_{i-1} = \Psi_i + (x_{i-1} - x_i) \Psi' \bigg|_{x_i} + \frac{(x_{i-1} - x_i)^2}{2!} \Psi'' \bigg|_{x_i} + \frac{(x_{i-1} - x_i)^3}{3!} \Psi''' \bigg|_{x_i} + \cdots$$

$$= \Psi_i - h \Psi'_i + \frac{h^2}{2!} \Psi''_i - \frac{h^3}{3!} \Psi'''_i + \frac{h^4}{4!} \Psi''''_i + \cdots$$
(5)

Adding Eq. (3) and Eq. (5) we get:

$$\Psi_{i+1} + \Psi_{i-1} = 2\Psi_i + h^2 \Psi_i'' + \frac{h^4}{12} \Psi_i'''' + \cdots$$

$$\implies \Psi_{i+1} - 2\Psi_i + \Psi_{i-1} = h^2 \Psi_i'' + \frac{h^4}{12} \Psi_i'''' + \cdots$$

$$\implies \Psi_i'' = \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} - \frac{h^2}{12} \Psi_i''''$$
(7)

ignoring the 4th order term we get:

$$\Psi_i'' = \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} \tag{8}$$

How to find the 4th order partial derivative term, Ψ'''' ? Write the LHS of Eq. (2) as:

$$\frac{d^2\Psi(x)}{dx^2} + \left[\epsilon - \frac{l(l+1)}{x^2} + \frac{2}{x}\right]\Psi(x) = s(x)$$

$$\Rightarrow \frac{d^2\Psi(x)}{dx^2} + k^2(x)\Psi(x) = s(x) \quad \left(k^2(x) = \epsilon - \frac{l(l+1)}{x^2} + \frac{2}{x}\right)$$

$$\Rightarrow \frac{d^2\Psi(x)}{dx^2} = s(x) - k^2(x)\Psi(x)$$
(9)

$$\Longrightarrow \frac{d^{4}\Psi\left(x\right)}{dx^{4}} = \frac{d^{2}}{dx^{2}}\left[s(x) - k^{2}\left(x\right)\Psi\left(x\right)\right]$$

Simplify it as

$$\frac{d^{4}\Psi(x)}{dr^{4}} = \frac{d^{2}}{dx^{2}}t(x) \quad \left[t(x) = s(x) - k^{2}(x)\Psi(x)\right]$$

$$\frac{d^{4}\Psi(x)}{dr^{4}} = \frac{t_{i+1} - 2t_{i} + t_{i-1}}{h^{2}} \quad \text{(Using Eq. (8))}$$
(10)

Deriving the update equation: Putting, Eq. (10) in Eq. (7) we get:

$$\Psi_i'' = \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} - \frac{1}{12} \left(t_{i+1} - 2t_i + t_{i-1} \right) \tag{11}$$

Putting the above equation in the Eq. (9) we get:

$$\frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} - \frac{1}{12} \left(t_{i+1} - 2t_i + t_{i-1} \right) = s_i - k_i^2 \Psi_i$$

Again,

$$t(x) = s(x) - k^{2}(x) \Psi(x) \implies s(x) = t(x) + k^{2}(x) \Psi(x)$$

therefore,

$$\begin{split} \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} &- \frac{1}{12} \left(t_{i+1} - 2t_i + t_{i-1} \right) = t_i + k_i^2 \Psi_i - k_i^2 \Psi_i \\ \Longrightarrow \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} &- \frac{1}{12} \left(t_{i+1} - 2t_i + t_{i-1} \right) = t_i \\ \Longrightarrow \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} &- \frac{1}{12} \left(t_{i+1} + 10t_i + t_{i-1} \right) = 0 \\ \Longrightarrow \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} &- \frac{1}{12} \left(s_{i+1} - k_{i+1}^2 \Psi_{i+1} + 10 \left(s_i - k_i^2 \Psi_i \right) + s_{i-1} - k_{i-1}^2 \Psi_{i-1} \right) = 0 \\ \Longrightarrow \Psi_{i+1} - 2\Psi_i + \Psi_{i-1} &- \frac{h^2}{12} \left(s_{i+1} - k_{i+1}^2 \Psi_{i+1} + 10 \left(s_i - k_i^2 \Psi_i \right) + s_{i-1} - k_{i-1}^2 \Psi_{i-1} \right) = 0 \\ \Longrightarrow \Psi_{i+1} \left(1 + \frac{h^2}{12} k_{i+1}^2 \right) - 2\Psi_i \left(1 - \frac{5h^2}{12} k_i^2 \right) + \Psi_{i-1} \left(1 + \frac{h^2}{12} k_{i-1}^2 \right) - \frac{h^2}{12} \left(s_{i+1} + 10s_i + s_{i-1} \right) = 0 \end{split}$$

For homogeneous case, s(x) = 0, therefore,

$$\frac{h^2}{12}\left(s_{i+1} + 10s_i + s_{i-1}\right) = 0$$

And we are left we the following update equation

$$\Psi_{i+1}\left(1+\frac{h^2}{12}k_{i+1}^2\right)-2\Psi_i\left(1-\frac{5h^2}{12}k_i^2\right)+\Psi_{i-1}\left(1+\frac{h^2}{12}k_{i-1}^2\right)=0$$

where,

$$k_i^2 = \epsilon_n - \frac{l(l+1)}{x_i^2} + \frac{2}{x_i}, \quad \epsilon_n = \frac{E_n}{E_1} = -\frac{1}{n^2} \text{ (where, } E_n = \frac{E_1}{n^2} \ n = 1, 2, 3, \cdots)$$

Put, $k_i^2 = g_i$, and $f_i = \left(1 + \frac{h^2}{12}k_i^2\right) = \left(1 + \frac{h^2}{12}g_i\right)$, such that the above equation can be written as:

$$\begin{split} \Psi_{i+1}f_{i+1} - 2\Psi_i \left(1 - \frac{5h^2}{12}g_i\right) + \Psi_{i-1}f_{i-1} &= 0 \\ \Longrightarrow \Psi_{i+1}f_{i+1} &= -10\Psi_i \left(-\frac{1}{5} + \frac{h^2}{12}g_i\right) - \Psi_{i-1}f_{i-1} \\ \Longrightarrow \Psi_{i+1}f_{i+1} &= -10\Psi_i \left(-1 - \frac{1}{5} + 1 + \frac{h^2}{12}g_i\right) - \Psi_{i-1}f_{i-1} \\ \Longrightarrow \Psi_{i+1}f_{i+1} &= -10\Psi_i \left(-\frac{6}{5} + 1 + \frac{h^2}{12}g_i\right) - \Psi_{i-1}f_{i-1} \\ \Longrightarrow \Psi_{i+1}f_{i+1} &= -10\Psi_i \left(-\frac{6}{5} + f_i\right) - \Psi_{i-1}f_{i-1} \\ \Longrightarrow \Psi_{i+1}f_{i+1} &= \Psi_i \left(12 - 10f_i\right) - \Psi_{i-1}f_{i-1} \\ \Longrightarrow \Psi_{i+1} &= \frac{\Psi_i \left(12 - 10f_i\right) - \Psi_{i-1}f_{i-1}}{f_{i+1}} \end{split}$$

The final update equation is thus:

$$\Psi_{i+1} = \frac{\Psi_i (12 - 10f_i) - \Psi_{i-1} f_{i-1}}{f_{i+1}}$$

where,

$$f_i = \left(1 + \frac{h^2}{12}g_i\right)$$
 and $g_i = \epsilon_n - \frac{l(l+1)}{x_i^2} + \frac{2}{x_i}$, with $\epsilon_n = -\frac{1}{n^2}$ $(n = 1, 2, 3, \dots)$

n being the **Principal quantum number**.

4 References

Codes and explanation of the Numerov's method can be found in this github link

How to find the Eigen values numerically?

• Method 1: (Matrix Numerov method) Consider the one dimensional Schrödinger equation written in terms of dimensionless variables:

$$\left[\frac{d^{2}}{dx^{2}} - W(x)\right]\Psi(x) = -\epsilon\Psi(x) \tag{12}$$

where,

$$\epsilon = \frac{E}{E_1}$$
 with, $E_1 = \frac{\hbar^2}{2\mu a_0^2}$ and $W(x) = a_0^2 U(r) = \frac{2\mu a_0^2}{\hbar^2} V(r)$, $(r = xa_0)$

Here, a_0 is some arbitary constant, and not the Bohr's radius.

Eq. (12) is an Eigen value equation. Therefore we need to find the Eigen values of the operator,

$$\left[\frac{d^2}{dx^2} - W(x)\right]$$

Using the result,

$$\frac{d^2}{dx^2}\Psi\left(x\right) \equiv \Psi_i^{\prime\prime} \approx \frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2}$$

Putting the above equation in Eq. (12) we get,

$$\frac{\Psi_{i+1} - 2\Psi_i + \Psi_{i-1}}{h^2} + W_i \Psi_i = -\epsilon \Psi_i \tag{13}$$

We now construct a vector containing all the values of Ψ as:

$$\Psi = \begin{pmatrix} \Psi_1 \\ \Psi_2 \\ \vdots \\ \Psi_{N-1} \\ \Psi_N \end{pmatrix}$$

Such that, Eq. (13) can be written in matrix form as:

$$\frac{1}{h^{2}} \begin{pmatrix}
-2 & 1 & & & \\
1 & -2 & 1 & & \\
& \ddots & \ddots & \ddots & \\
& & 1 & -2 & 1 \\
& & & 1 & -2
\end{pmatrix}
\begin{pmatrix}
\Psi_{1} \\
\Psi_{2} \\
\vdots \\
\Psi_{N-1} \\
\Psi_{N}
\end{pmatrix} + \begin{pmatrix}
W_{1} \\
W_{2} \\
\vdots \\
W_{N-1} \\
W_{N}
\end{pmatrix}
\begin{pmatrix}
\Psi_{1} \\
\Psi_{2} \\
\vdots \\
\Psi_{N-1} \\
\Psi_{N}
\end{pmatrix} = -\epsilon \begin{pmatrix}
\Psi_{1} \\
\Psi_{2} \\
\vdots \\
\Psi_{N-1} \\
\Psi_{N}
\end{pmatrix} \tag{14}$$

The above equation can be written as:

$$\left(\frac{1}{h^2}A + W\right)\Psi = -\epsilon\Psi\tag{15}$$

So all that is left is to find the Eigen values of the matrix, $\left(\frac{1}{h^2}A + W\right)$.

 \mathbf{OR}

$$\begin{split} &\Psi_{i+1}\left(1+\frac{h^2}{12}k_{i+1}^2\right)-2\Psi_i\left(1-\frac{5h^2}{12}k_i^2\right)+\Psi_{i-1}\left(1+\frac{h^2}{12}k_{i-1}^2\right)=0\\ \Longrightarrow &\Psi_{i+1}\left(1+\frac{h^2}{12}\left(\epsilon_n-W_{i+1}\right)\right)-2\Psi_i\left(1-\frac{5h^2}{12}\left(\epsilon_n-W_i\right)\right)+\Psi_{i-1}\left(1+\frac{h^2}{12}\left(\epsilon_n-W_{i-1}\right)\right)=0 \ \left(k_i^2=\epsilon_n-W_i(x_i)\right)\\ \Longrightarrow &\frac{\Psi_{i+1}-2\Psi_i+\Psi_{i-1}}{h^2}-\frac{\Psi_{i+1}W_{i+1}+10\Psi_iW_i+\Psi_{i-1}W_{i-1}}{12}+\epsilon_n\frac{\Psi_{i+1}+10\Psi_i+\Psi_{i-1}}{12}=0\\ \Longrightarrow &-\frac{\Psi_{i+1}-2\Psi_i+\Psi_{i-1}}{h^2}+\frac{\Psi_{i+1}W_{i+1}+10\Psi_iW_i+\Psi_{i-1}W_{i-1}}{12}=\epsilon_n\frac{\Psi_{i+1}+10\Psi_i+\Psi_{i-1}}{12} \end{split}$$

where, for i = 1 the above equation is:

$$-\frac{\Psi_2 - 2\Psi_1 + \Psi_0}{h^2} + \frac{\Psi_2 W_2 + 10\Psi_1 W_1 + \Psi_0 W_0}{12} = \epsilon_n \frac{\Psi_2 + 10\Psi_1 + \Psi_0 W_0}{12}$$

Assume, $\Psi_0 = 0$, this gives:

$$-\frac{\Psi_2 - 2\Psi_1}{h^2} + \frac{\Psi_2 W_2 + 10\Psi_1 W_1}{12} = \epsilon_n \frac{\Psi_2 + 10\Psi_1}{12}$$

for i = 2,

$$-\frac{\Psi_3-2\Psi_2+\Psi_1}{h^2}+\frac{\Psi_3W_3+10\Psi_2W_2+\Psi_1W_1}{12}=\epsilon_n\frac{\Psi_3+10\Psi_2+\Psi_1}{12}$$

for i = 3,

$$-\frac{\Psi_4-2\Psi_3+\Psi_2}{h^2}+\frac{\Psi_4W_4+10\Psi_3W_3+\Psi_2W_2}{12}=\epsilon_n\frac{\Psi_4+10\Psi_3+\Psi_2}{12}$$

Again, for i = 4,

$$-\frac{\Psi_5-2\Psi_4+\Psi_3}{h^2}+\frac{\Psi_5W_5+10\Psi_4W_4+\Psi_3W_3}{12}=\epsilon_n\frac{\Psi_5+10\Psi_4+\Psi_3}{12}$$

If the maximum number of points is 4, then $\Psi_5 = 0$

$$\implies -\frac{-2\Psi_4 + \Psi_3}{h^2} + \frac{10\Psi_4 W_4 + \Psi_3 W_3}{12} = \epsilon_n \frac{10\Psi_4 + \Psi_3}{12}$$

The above equations can be combined together in matrix form as:

$$\begin{split} \Longrightarrow & \left[-\frac{1}{h^2} \begin{pmatrix} -2 & 1 & 0 & 0 \\ 1 & -2 & 1 & 0 \\ 0 & 1 & -2 & 1 \\ 0 & 0 & 1 & -2 \end{pmatrix} + \frac{1}{12} \begin{pmatrix} 10W_1 & W_2 & 0 & 0 \\ W_1 & 10W_2 & W_3 & 0 \\ 0 & W_2 & 10W_3 & W_4 \\ 0 & 0 & W_3 & 10W_4 \end{pmatrix} \right] \begin{pmatrix} \Psi_1 \\ \Psi_2 \\ \Psi_3 \\ \Psi_4 \end{pmatrix} = \frac{\epsilon_n}{12} \begin{pmatrix} 10 & 1 & 0 & 0 \\ 1 & 10 & 1 & 0 \\ 0 & 1 & 10 & 1 \\ 0 & 0 & 1 & 10 \end{pmatrix} \begin{pmatrix} \Psi_1 \\ \Psi_2 \\ \Psi_3 \\ \Psi_4 \end{pmatrix} \\ \Longrightarrow & \left[-\frac{1}{h^2} \left(\mathbb{I}_{-1} - 2\mathbb{I}_0 + \mathbb{I}_1 \right) + \frac{1}{12} \left(\mathbb{I}_{-1} + 10\mathbb{I}_0 + \mathbb{I}_1 \right) W \right] \Psi = \frac{\epsilon_n}{12} \left(\mathbb{I}_{-1} + 10\mathbb{I}_0 + \mathbb{I}_1 \right) \Psi \\ \Longrightarrow & \left[-A + BW \right] \Psi = \epsilon_n B \Psi \quad \left(A = \frac{1}{h^2} \left(\mathbb{I}_{-1} - 2\mathbb{I}_0 + \mathbb{I}_1 \right), \ B = \frac{1}{12} \left(\mathbb{I}_{-1} + 10\mathbb{I}_0 + \mathbb{I}_1 \right) \right) \\ \Longrightarrow & \left[-B^{-1}A + W \right] \Psi = \epsilon_n \Psi \quad \left(W(x) = \frac{2ma_0^2}{\hbar^2} V(x), \ \epsilon_n = \frac{E_n}{E_0}, \ E_0 = \frac{\hbar^2}{2ma_0^2} \right) \end{split}$$

The energy eigen values evaluated using the above equation are going to be dimensionless. The equation below however will give the exact value of the energy eigen value.

$$\left[-\frac{\hbar^2}{2m} B^{-1} A + V \right] \Psi = E \Psi \quad \left(A \equiv \frac{A}{a_0^2} \right)$$

Here, \mathbb{I}_p denotes a matrix who's \mathbf{p}^{th} diagonal is unity

$$I_{-1} = \begin{pmatrix} 0 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \end{pmatrix}, I_{0} = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}, I_{1} = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 1 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 0 & 0 \end{pmatrix}, \text{ and } W = \begin{pmatrix} W_{1} & 0 & 0 & 0 \\ 0 & W_{2} & 0 & 0 \\ 0 & 0 & W_{3} & 0 \\ 0 & 0 & 0 & W_{4} \end{pmatrix}$$

So, we need to find the Eigen values and Eigen vectors of the matrix, $\left[-\frac{\hbar^2}{2m}B^{-1}A + V\right]$.

• Method 2 (shooting method): To find the Energy Eigen values we can also utilize the boundary conditions, such as

$$\Psi(r=0) = 0;$$
 $\Psi(r=\infty) = 0$ (for Hydrogen atom)

Or

$$\Psi(x=\pm\infty)=0$$
 (for one dimensional linear Harmonic oscillator)

The boundary condition at ∞ can be implemented by choosing, some large, $r=r_c$ (say) such that, $\Psi(r_c)=0$.

For Hydrogen atom, if say, E_1 corresponds to a correct energy level, then when we backward integrate the radial wavefunction Numerov's method, we should get $\Psi(0) = 0$. A deviation of the energy from E_1 will result in $\Psi(0) \neq 0$. The basic procedure to search for correct Energy Eigen value is thus as follows: Start with a guess energy, E_1

- 1) The guess energy, E_1 should be smaller than the smallest potential energy. In case of Hydrogen atom, it should be smaller than $-Z^2$.
- 2) With the guess energy integrate the equation and get the value of the wavefunction at r = 0, which we will denote as Ψ_1 . Meanwhile, Set another energy, $E_2 = E_1$.
- 3) Increase the energy E_2 by an amount δE and get a new energy, $E_2 = E_2 + \delta E$.
- -4) Integrate the Scrodinger equation to get the corresponding wavefunction, Ψ_2 and evaluate $\Psi_2(0)$.
- -5) Go back to step 2 if $\Psi_1(0) \times \Psi_2(0) > 0$.
- 6) At this step, we should have the correct energy enclosed in the interval, $[E_1, E_2]$. Use root finding method, e.g. "Newton Raphson method", "Secant method", "Brent's method" etc. to get the correct energy.