Exercise TI-01

January 11, 2023

```
[]: import sisl
import numpy as np
import matplotlib.pyplot as plt
%matplotlib inline
import functools
```

1 Quantum Hall Effect

In this exercise, we will build on **TB_07** and **A_03** to simulate the quantum Hall effect. In **A_03**, we learned how to created a Hall bar device and in **TB_07**, how to deal with magnetic fields.

Here, we will extract the Hall resistance from the transmissions calculated with TBtrans using the Landauer-Büttiker formalism.

1.1 Exercise Overview:

- 1. Create a Hall bar (see A 03)
- 2. Construct Hamiltonians and add magnetic fields (see TB_07)
- 3. Calculate the transmission with TBtrans
- 4. Extract the Hall resistance (R_H) .
- 5. Extract the logitudinal resistance (R_L) .

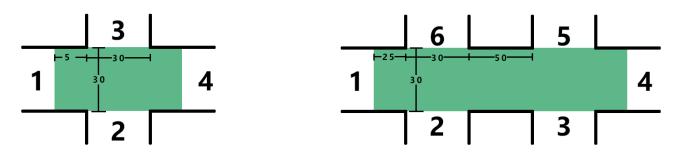
1.2 Exercise

1.2.1 1. Create a Hall bar

In order to be able to observe the quantum Hall effect, the size of the Hall bar needs to be big enough. For a 4(6) lead device reasonable dimensions are: 1. 4-lead device (square lattice): - Width of electrodes (perpendicular to the semi-infinite axis): 30 atoms - Offset of the electrodes 2(3) from the corner of the central part: >= 5 atoms

- 2. 6-lead device (square lattice):
 - Width of electrodes (perpendicular to the semi-infinite axis): 30 atoms
 - Spacing between electrodes on the same side: 50 atoms
 - Offset of the electrodes 2(3,5,6) from the corner of the central part: > 25 atoms
- 3. 6-lead Graphene Hall bar:
 - Width of electrodes (perpendicular to the semi-infinite axis): ~30 Å
 - Spacing between electrodes on the same side: ~ 50 Å
 - Offset of the electrodes 2(3,5,6) from the corner of the central part: ~ 15 Å

4-lead device 6-lead device



```
[]: # Create a Hall bar
```

1.2.2 2. Construct Hamiltonian and add magnetic fields

The required field strengths may vary depending on the size of the Hall bar. We should start with a corse grid, and create a finer grid once we have identified the correct range. A good starting point might by B = 1 / np.arange(10,31).

```
def peierls(self, ia, atoms, atoms_xyz=None, B=None):
    idx = self.geometry.close(ia, R=[0.1, 1.01], atoms=atoms,
    atoms_xyz=atoms_xyz)
    # Onsite
    self[ia, idx[0]] = 4

# Hopping
    if B == 0:
        self[ia, idx[1]] = -1
    else:
        xyz = self.geometry.xyz[ia]
        dxyz = self.geometry[idx[1]]
        self[ia, idx[1]] = - np.exp(-0.5j * B * (dxyz[:, 0] - xyz[0])*(dxyz[: 4,1] + xyz[1]))
```

```
[]: H0 = sisl.Hamiltonian(geom, dtype=np.float64)
H0.construct(functools.partial(peierls, B=0.))

HB = sisl.Hamiltonian(geom, dtype=np.complex128)

rec_phis = np.arange(10,31)
for rec_phi in rec_phis:
    HB.construct(functools.partial(peierls, B=1/rec_phi))
    dH = ...
```

1.2.3 3. Calculate the transmission with TBtrans

The folder of this exercise contains the skeleton of an input file for a 4-lead (RUN-4.fdf) and 6-lead device (RUN-6.fdf), as well as a script to run TBtrans for all values of the magnetic field (run.sh).

Depending on the size of the Hall bar, this step might require a considerable amount of time.

1.2.4 4. Extract the Hall resistance (R_H)

In Hall effect we measure the build up of a potential difference between the measurement electrodes as response to an electric current. The Hall resistance (R_H) in a 4 lead Hall bar like the one shown above is given by

$$R_H = \left. \frac{V_2 - V_3}{I_1} \right|_{I_2 = I_3 = 0}.$$

In TBtrans the chemical potentials of all electrodes can be specified and the currents are calculated as a response to the biases. Rather than trying many combinations of chemical potentials to find the Hall resistance, we use transmission curves to calculate R_H .

To start, we express the lead currents I_i in terms of applied biases V_i and the transmissions T_{ij} between leads i and j

$$I_i = \sum_j G_{ij}(V_i - V_j)$$
 where $G_{ij} = \frac{2e^2}{h}T_{ij}$.

We rewrite this relation as

$$\mathbf{I} = \mathcal{G}\mathbf{V} \quad \text{, where} \quad \mathcal{G}_{ii} = \sum_{i \neq j} G_{ij} \quad \text{and} \quad \mathcal{G}_{ij} = -G_{ij}.$$

Since the currents only depend on bias differences, we can set one of them to zero without loss of generality (here $V_4=0$). Further, Kirchhoff's current law allow us to eliminate one of current (here $I_4=-I_1-I_2-I_3$). This leaves us with an invertible 3×3 matrix equation.

Using the inverse **R** of \mathcal{G} , we can express V_2 and V_3 in terms of the lead currents I_i and calculate the Hall conductance:

$$\mathbf{V} = \mathcal{G}^{-1}\mathbf{I} = \mathbf{R}\mathbf{I} \quad \Rightarrow \quad V_i = R_{i1}I_1 + R_{i2}I_2 + R_{i3}I_3,$$

and finally, we find the Hall resistance:

$$R_{H} = \left. \frac{R_{21}I_{1} + R_{22}I_{2} + R_{23}I_{3} - (R_{31}I_{1} + R_{32}I_{2} + R_{33}I_{3})}{I_{1}} \right|_{I_{2} = I_{3} = 0} \tag{1}$$

$$=R_{21}-R_{31} \tag{2}$$

The derivation for the 6-lead device is analogous and yields:

$$R_H = R_{21} - R_{61}$$

If everything is set up correctly, the quantization of the Hall resistance should be visible.

$$R_H = \frac{h}{2ne^2}$$
 for square lattice Hall bar $n \in \mathbb{N}$ (3)

$$R_H = \frac{h}{2(2n-1)e^2} \qquad \text{for graphene Hall bar} \qquad n \in \mathbb{N}$$
 (4)

1.2.5 5. Extract longitudinal resistance R_L

The longitudinal resistance can be extracted using the same approach used for the Hall resistance. With a 6-lead Hall bar we can replace the V_3 with V_6 and get it immediately.

$$R_L = R_{21} - R_{31}$$

With a 4-lead Hall bar we need to create a new device with electrodes 2 and 3 on the same side of the Hall bar.

If the energy mesh in TBTrans and the mesh for magnetic field strength are fine engough, spikes in R_L should be observable at each step of R_H .

```
[]: # Create short-hand function to open files
gs = sisl.get_sile
# No magnetic field
tbt0 = gs('M_0/siesta.TBT.nc')
# All magnetic fields in increasing order
tbts = [gs('M_{})/siesta.TBT.nc'.format(rec_phi)) for rec_phi in rec_phis]
```

```
[]: def calc_G(tbtsile, n):
         # Construct G
         G = np.zeros((tbtsile.nE, n, n))
         for i in range(n):
             for j in range(n):
                 if i == j: continue
                 Gij = tbtsile.transmission(i,j)
                 G[:,i,j] = -Gij
                 G[:,i,i] += Gij
         return G
     # Number of electrodes used in the Hall bar
     nEl = 6
     # Calculate Conductance matrix G
     G = np.asarray([calc_G(tbt, nEl) for tbt in tbts])
     # Remove one row and column from the matrix
     referenceElectrode = 3
     G = np.delete(np.delete(G, referenceElectrode, axis=2), referenceElectrode,
      ⇒axis=3)
```

```
# Invert G to get resistance matrix R

# Hint: If the matrix inversion might fail for some values of rec_phi, because the magentic field is too strong.

R = np.linalg.inv(G)

if nEl == 6:
    # For 6-lead device
    RH = R[:,:,1,0] - R[:,:,4,0]  # R_21 - R_61
    RL = R[:,:,1,0] - R[:,:,2,0]  # R_21 - R_31

elif nEl == 4:
    # For 4-lead device
    RH = R[:,:,1,0] - R[:,:,2,0]  # R_21 - R_31
```

```
[]: # Plot the Hall resistance
     # - as a function of energy for a fixed magnetic field strength
     # - as a function of magnetic field strength for a fixed energy
     \# E idx = \dots
     # phi_i dx = \dots
     fig, axs = plt.subplots(1,2, sharey=True, figsize=(12,4))
     axs[0].set_title('$\phi = \{:.5f\}\$'.format(1/rec_phis[phi_idx]))
     axs[0].set_ylabel(r'$R \left[h/2e^2\right]$')
     axs[0].set_xlabel('$E/t$')
     axs[0].axvline(x=E[E_idx], ls='--', c='k')
     # axs[0].plot(...)
     axs[1].set_title(f'E = {E[E_idx]:.5f}')
     axs[1].axvline(x=1/rec_phis[phi_idx], ls='--', c='k')
     axs[1].set_xlabel('$\phi/\phi_0$')
     # axs[1].plot(...)
     axs[1].legend(loc='center right', bbox_to_anchor=(1.5,0.5), ncol=1)
     plt.show()
```