



---

# Capacitance of the metal-oxide-semiconductor (MOS) structure

Dr. Sanatan Chattopadhyay

Department of Electronic Science

University of Calcutta



# Si based MOS is the key

## As a material

Si is abundant in nature



High quality native oxide ( $\text{SiO}_2$ )



Appropriate mechanical strength

## Market

Microelectronics market



80% is dominated by CMOS



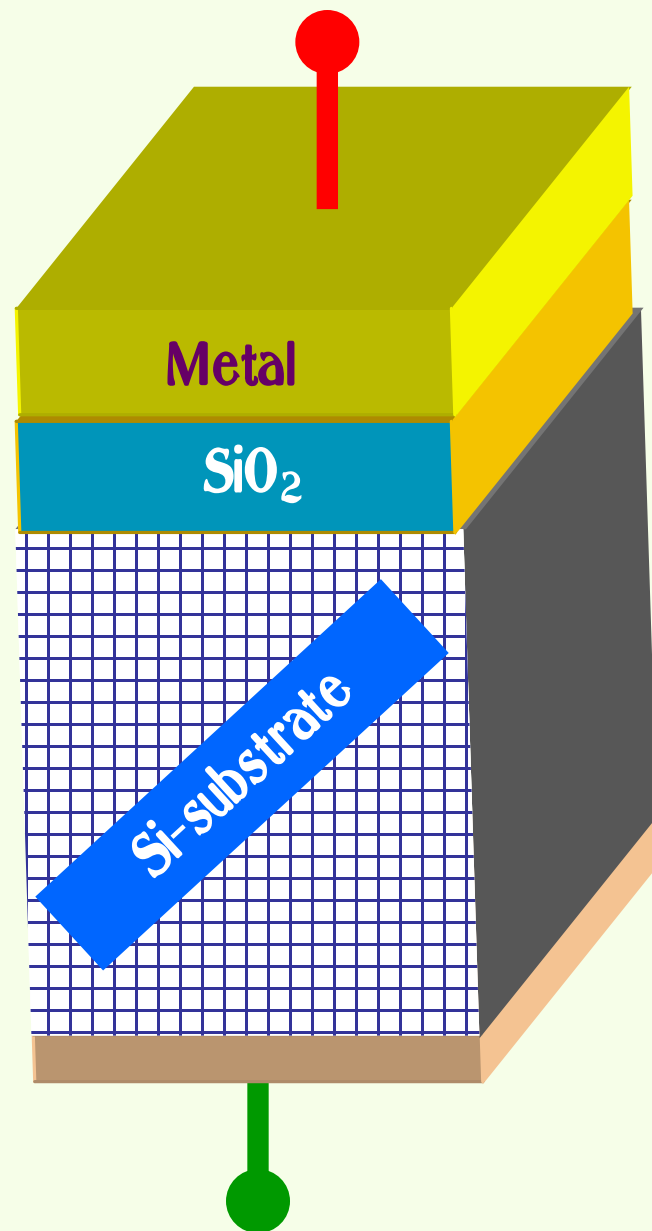
97% is covered by Si

## Applications:

Memory and Storages; Sensors; Diagnostics; Almost in all types of electronic circuits; Filters; Etc. ....



# Metal-Oxide-Semiconductor (MOS) structure



**MOS structure**

## Schematic view:

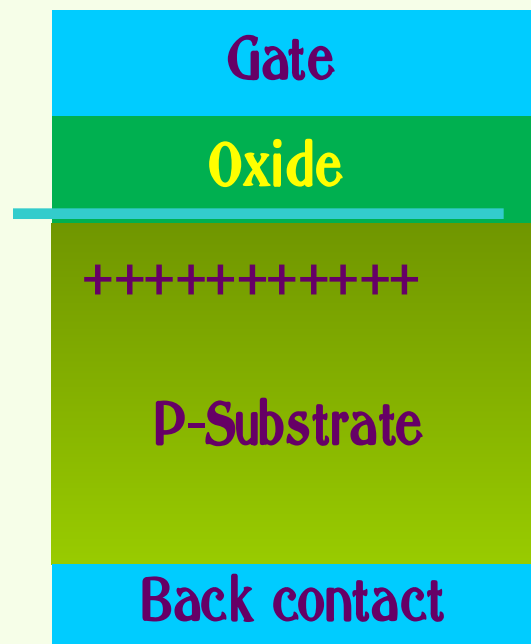
- A Si substrate (p- or n-type)
- $\text{SiO}_2$  is grown on it (@800 - 900°C)
- Ohmic contacts are taken from the top and bottom
- Sometimes, A Poly-Si layer is grown on  $\text{SiO}_2$ , on which gate ohmic contact is taken.



# Distribution of charges: different biases

Holes are attracted.

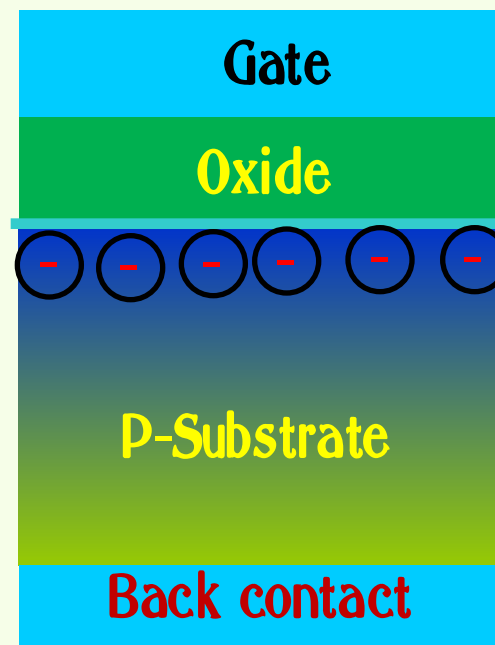
$$V_G < 0$$



Accumulation

Holes are repelled.

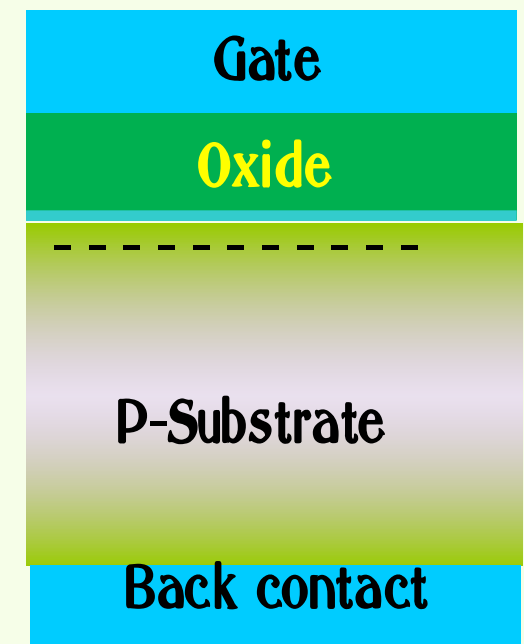
$$V_G > 0$$



Depletion

Electrons are attracted.

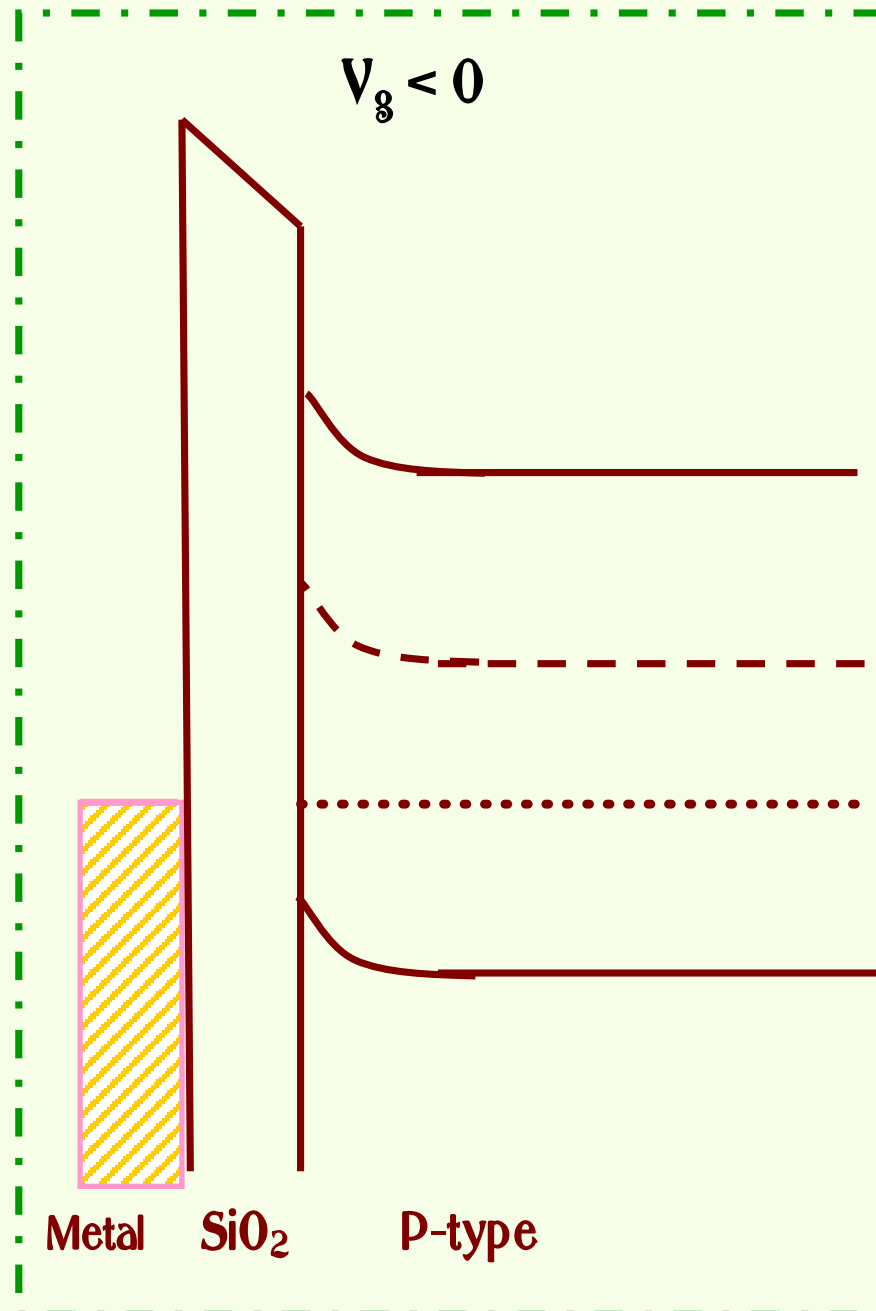
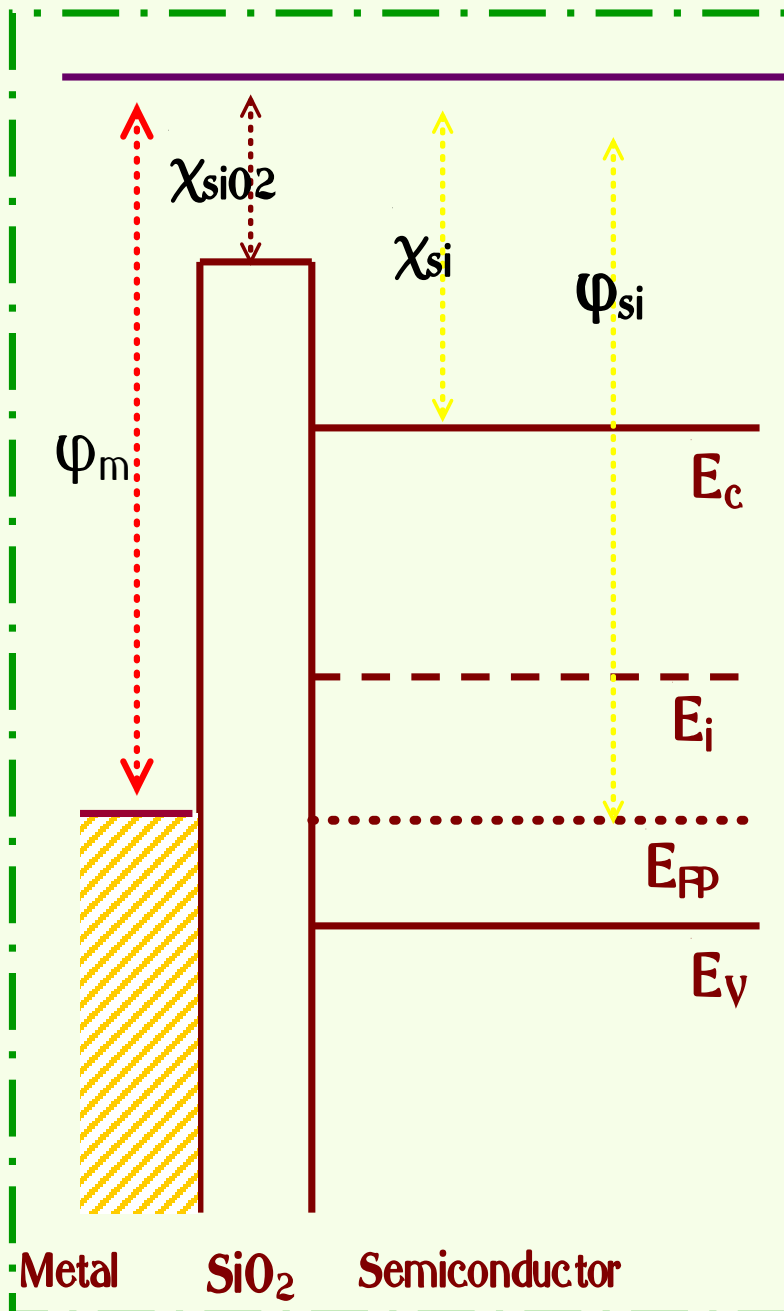
$$V_G \gg 0$$



Inversion



# Band diagram in MOS



When the band bends downward,  $\psi_s > 0$ ; and when it bends upward,  $\psi_s < 0$ .

- For  $\psi_s < 0$ , accumulation of holes (band bends upward) for p-type substrate.
- For  $\psi_s > 0$ , accumulation of electrons for n-type substrate.



# Graphical variation of potential, field, charge

- For charge neutrality of the system, it is required that,

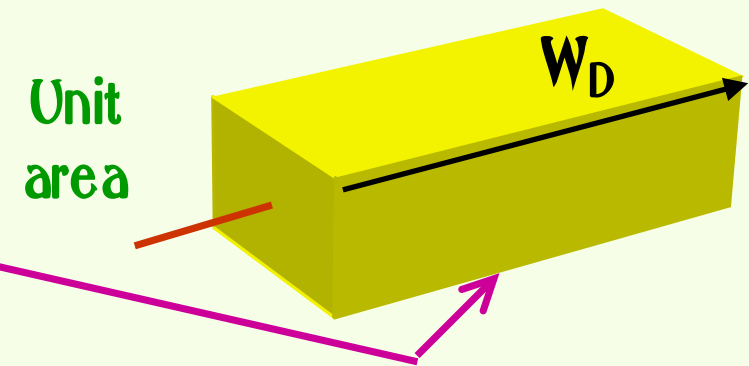
$$Q_M = -(Q_n + \boxed{qN_A W_D}) = -Q_s$$

where

- $Q_M$ : charges/unit area on the metal,
- $Q_n$ : electrons/unit area near the surface (inversion region)
- $qN_A W_D$ : the ionized acceptors/unit area in the space-charge region with depletion width
- $Q_s$ : total charges/unit area in the semiconductor.
- Clearly, in the absence of any work-function difference, the applied voltage will partly appear across the insulator and partly across the semiconductor. Thus,

$$V = V_{ox} + \psi_s$$

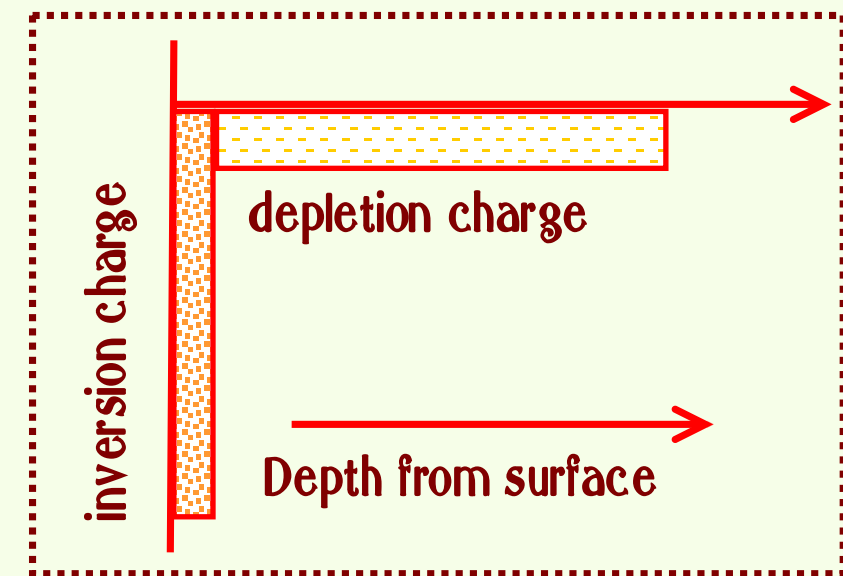
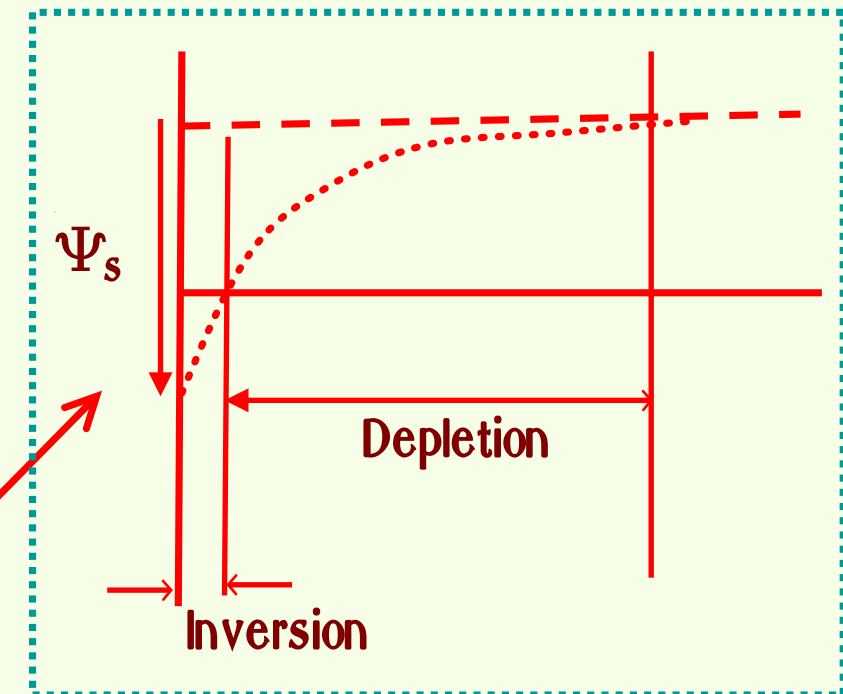
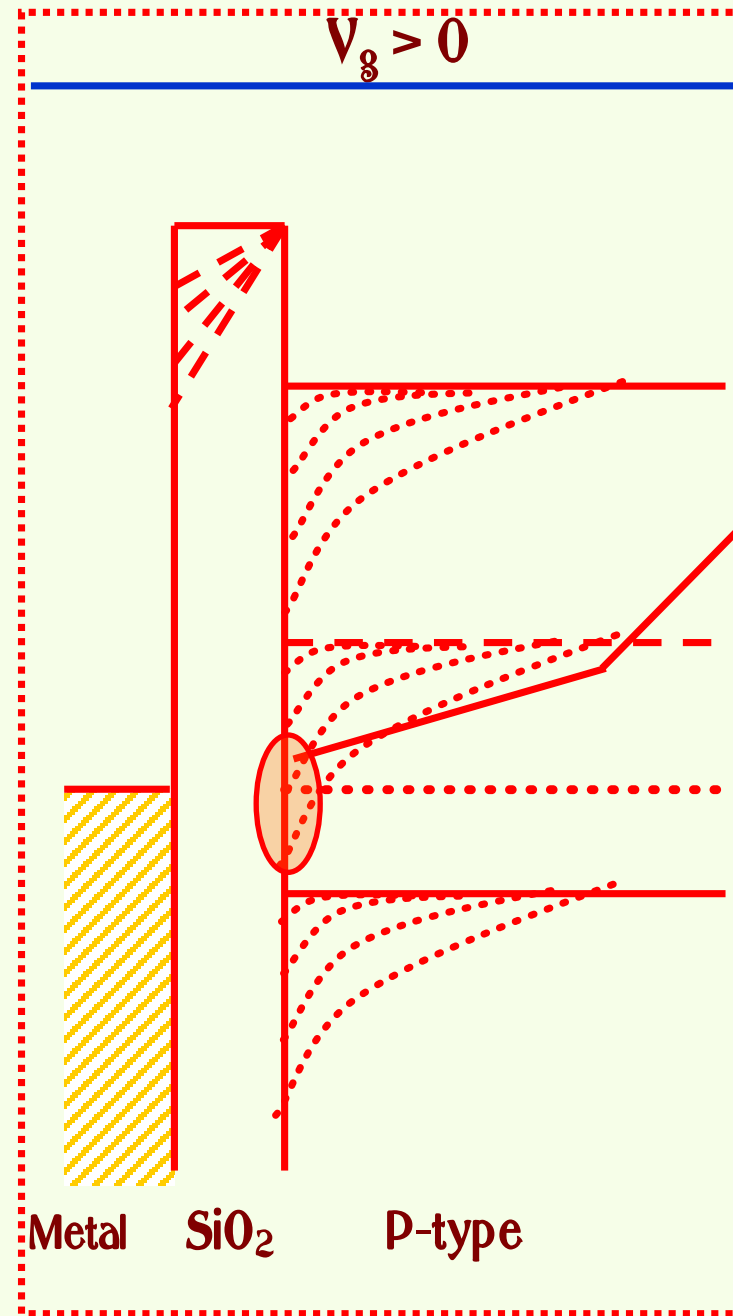
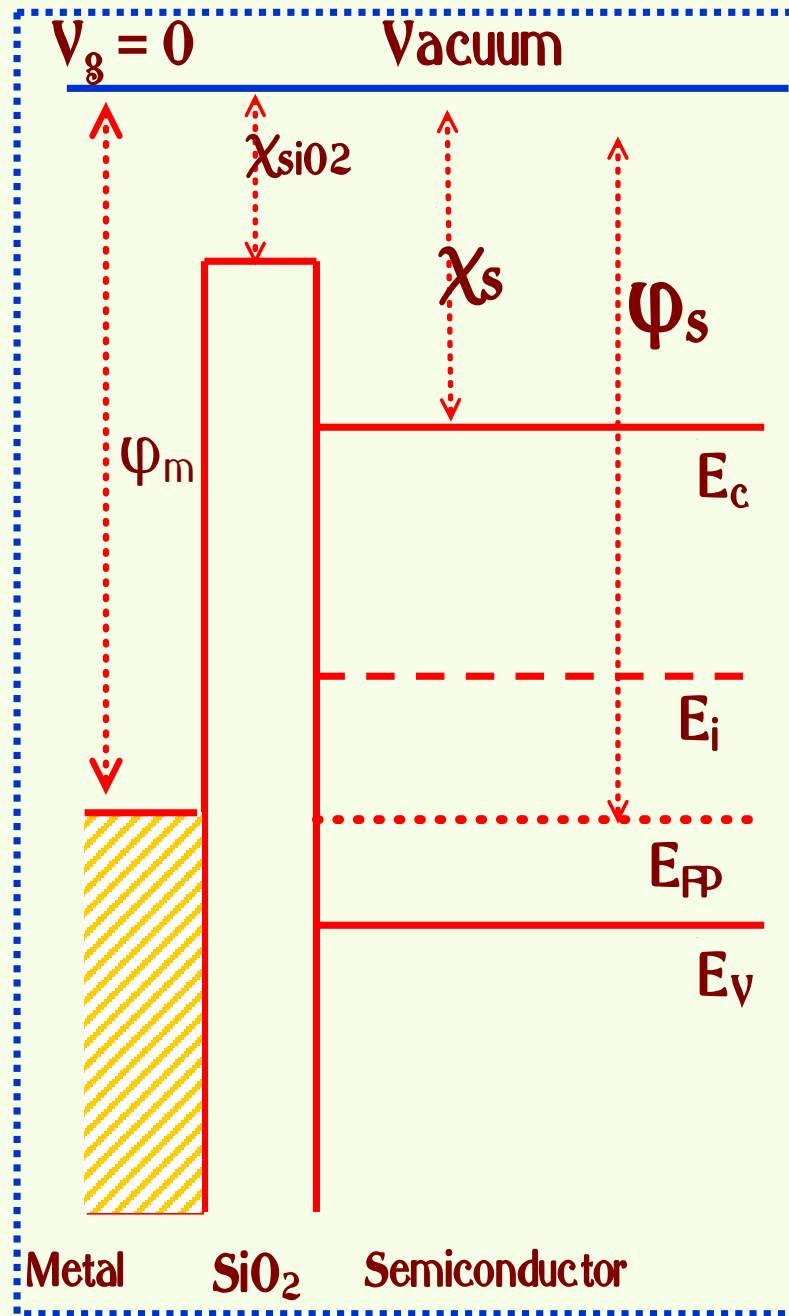
where  $V_i$  is the potential across the insulator and is given by







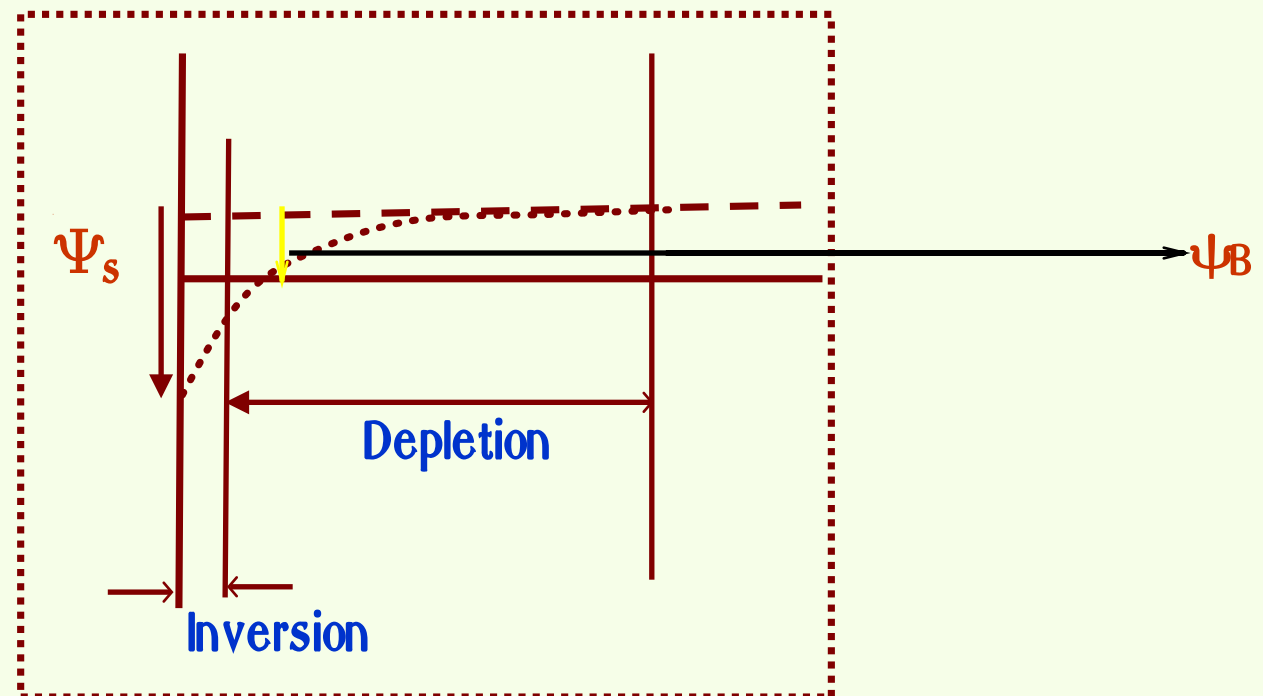
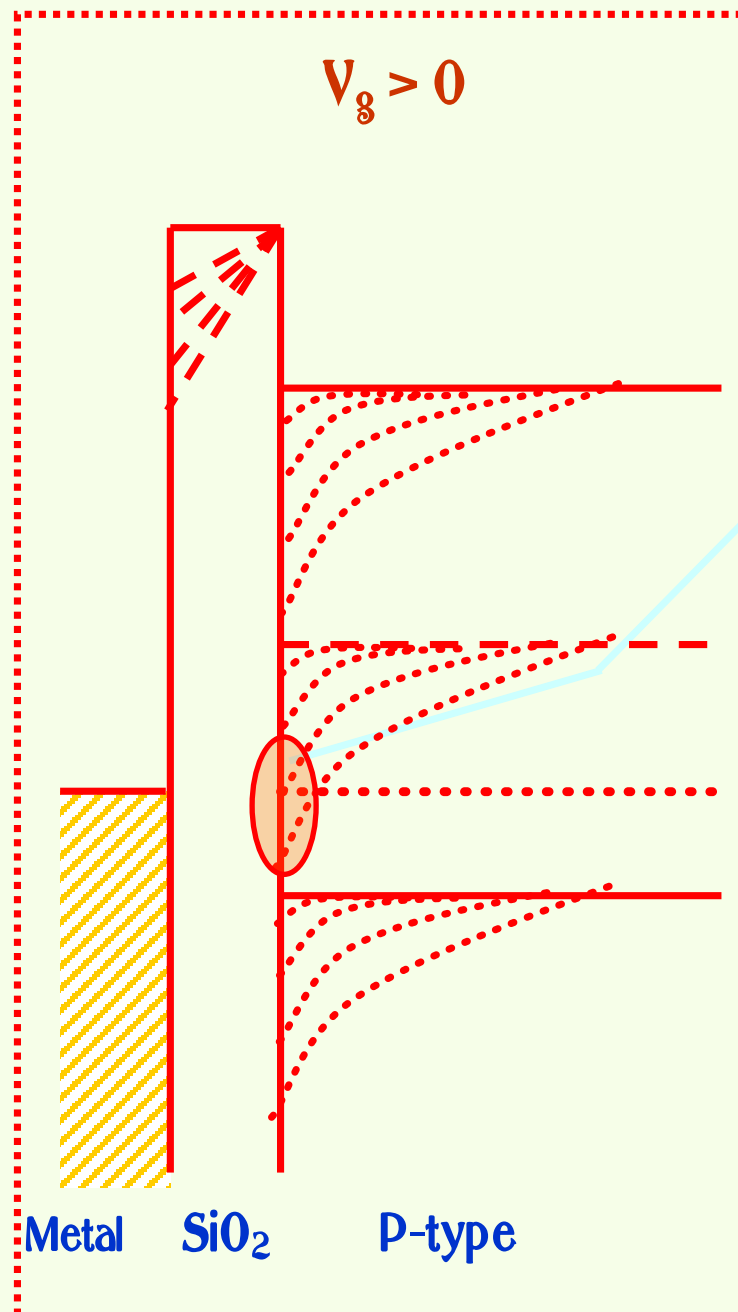
# How does an inversion layer form?







# Band diagram at inversion condition



$\psi_s = 0$	Flat-band condition
$\psi_B > \psi_s > 0$	Depletion of holes (bands bending downward)
$\psi_s = \psi_B$	Fermi-level at mid-gap, $E_F = E_i(0)$ , $n_p(0) = p_p(0) = n_i$
$2\psi_B > \psi_s > \psi_B$	Weak inversion (electron enhancement)
$\psi_s > 2\psi_B$	Strong inversion [ $n_p(0) = p_{p0} +$ or $N_A$ ]

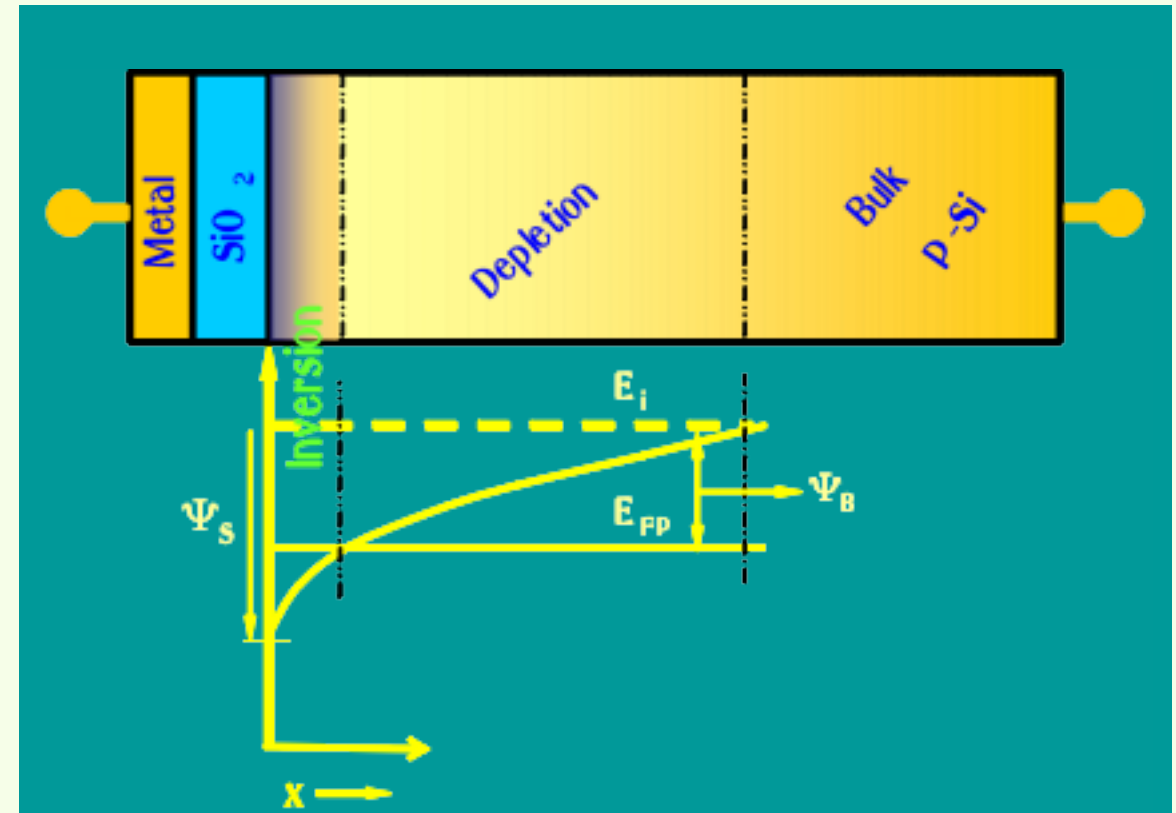


# Derivation: space charge

Redistribution of carriers due to the application of bias:

$$n_p(x) = n_{p0} e^{\frac{q\psi(x)}{kT}} = n_{p0} e^{\beta\psi(x)}$$

$$p_p(x) = n_{p0} e^{-\frac{q\psi(x)}{kT}} = n_{p0} e^{-\beta\psi(x)}$$



Where  $\psi_s$  is '+' when the band is bent downward,  $n_{p0}$  and  $p_{p0}$  are the equilibrium densities of electrons and holes, respectively, in the bulk of the semiconductor, and

$$\beta = q/kT$$

At the surface the densities are:

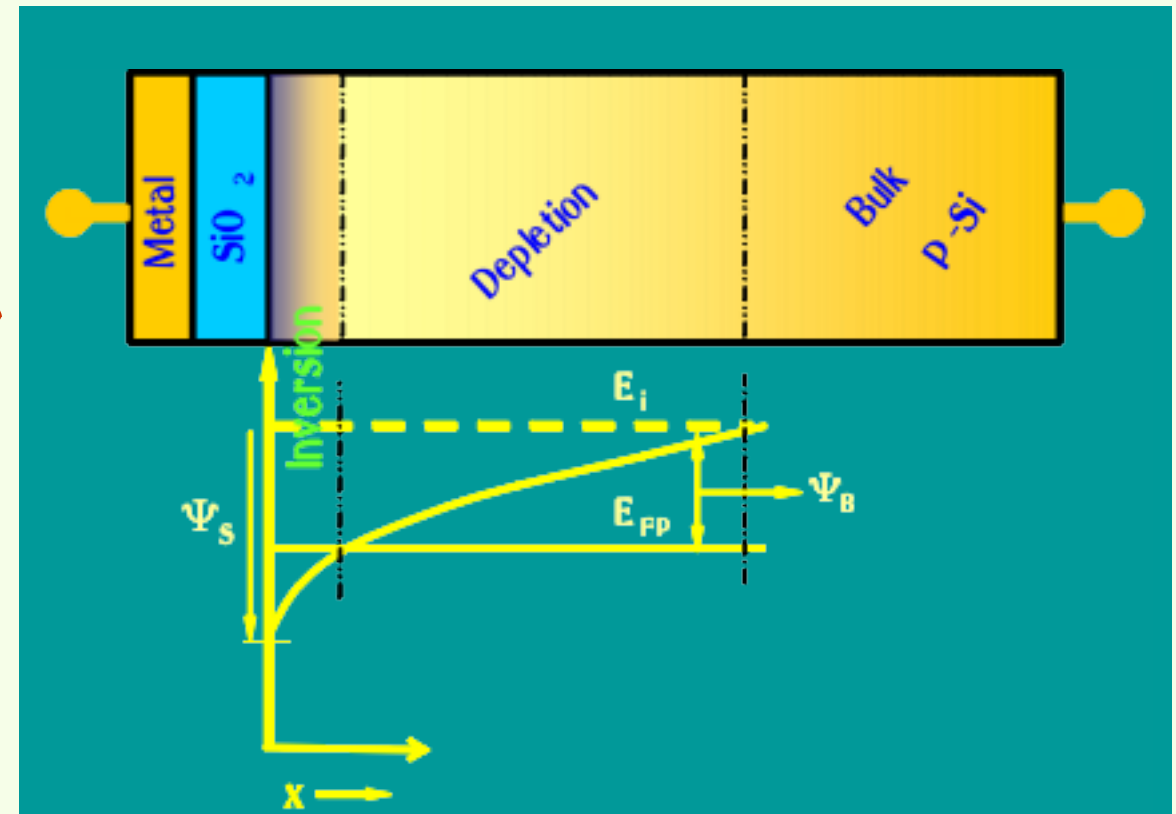
$$n_p(0) = n_{p0} e^{\beta\psi_s}$$

$$p_p(0) = p_{p0} e^{-\beta\psi_s}$$



# Derivation: space charge

From the previous discussions and with the help of the above equations, the following of surface potential can be distinguished:



$\psi_s < 0$ : Accumulation of holes (band bends downward)

$\psi_s = 0$ : Flat-band voltage

$\psi_{BP} > \psi_s > 0$ : Depletion of holes (band bends downward)

$\psi_s = \psi_{BP}$ : Fermi-level at midgap,  $E_F = E_i(0)$ ;  $n_p(0) = p_p(0) = n_i$ .

$2\psi_{BP} > \psi_s > \psi_{BP}$ : Weak inversion

$\psi_s > 2\psi_{BP}$ : Strong inversion



# Derivation: space charge

The potential  $\psi(x)$  as a function of distance can be obtained by using the one-dimensional Poisson's equation:

$$\frac{d^2\psi(x)}{dx^2} = -\frac{\rho(x)}{\epsilon_s}$$

Where  $\rho(x)$  is the total space-charge density given by:

$$\rho(x) = q(N_D^+ - N_A^- + p_p - n_p)$$

$N_D^+$  : ionized donors density

$N_A^-$  : ionized acceptor density

Now, in the bulk of the semiconductor, far from the surface, charge neutrality must exist.

Therefore at:  $\psi_p(\infty) = 0$ , we have  $\rho(x) = 0$  and



# Derivation: space charge

$$N_D^+ - N_A^- + p_{po} - n_{po} = 0; \text{ or, } N_D^+ - N_A^- = n_{po} - p_{po}$$

Now,  $n_p(0) = n_{po} e^{\beta\psi_s}$  and,  $p_p(0) = p_{po} e^{\beta\psi(x)}$

$$\therefore \frac{d^2\psi(x)}{dx^2} = -\frac{q}{\epsilon_s} (n_{po} - p_{po} + p_p - n_p)$$

$$\therefore \frac{d^2\psi(x)}{dx^2} = \frac{q}{\epsilon_s} (p_{po} - p_p + n_p - n_{po})$$

$$\therefore \frac{d^2\psi(x)}{dx^2} = \frac{q}{\epsilon_s} (p_{po} - p_{po} e^{-\beta\psi_s} + n_{po} e^{\beta\psi_s} - n_{po})$$

$$\therefore \frac{d^2\psi(x)}{dx^2} = \frac{q}{\epsilon_s} \left[ (p_{po} (1 - e^{-\beta\psi_s}) + n_{po} (e^{\beta\psi_s} - 1)) \right] = \frac{qp_{po}}{\epsilon_s} \left[ \frac{n_{po}}{p_{po}} (e^{\beta\psi_s} - 1) + (1 - e^{-\beta\psi_s}) \right]$$



# Derivation: space charge

$$\text{or, } \frac{d}{d\psi} \left( \frac{d\psi}{dx} \right) \left( \frac{d\psi}{dx} \right) = \frac{qp_{po}}{\epsilon_s} \left[ \frac{n_{po}}{p_{po}} (e^{\beta\psi_s} - 1) + (1 - e^{-\beta\psi_s}) \right]$$

$$\text{or, } \int_0^{\frac{d\psi}{dx}} \left( \frac{d\psi}{dx} \right) d \left( \frac{d\psi}{dx} \right) = \frac{qp_{po}}{\epsilon_s} \int_0^{\psi} \left[ \frac{n_{po}}{p_{po}} (e^{\beta\psi_s} - 1) + (1 - e^{-\beta\psi_s}) \right] d\psi$$

$$\text{or, } \frac{1}{2} \left( \frac{d\psi}{dx} \right)^2 = \frac{qp_{po}}{\epsilon_s} \left[ \frac{n_{po}}{p_{po}} \left( \frac{1}{\beta} e^{\beta\psi_s} - \psi \right) + \left( \psi - \frac{1}{\beta} e^{-\beta\psi_s} \right) \right]_0^{\psi}$$

$$\text{or, } \left( \frac{d\psi}{dx} \right)^2 = \frac{2qp_{po}}{\beta\epsilon_s} \left[ \frac{n_{po}}{p_{po}} (e^{\beta\psi_s} - \beta\psi - 1) + (e^{-\beta\psi_s} + \beta\psi - 1) \right]$$

$$\text{or, } \frac{d\psi}{dx} = \sqrt{\frac{2p_{po}kT}{\epsilon_s}} \cdot \sqrt{\frac{n_{po}}{p_{po}} (e^{\beta\psi_s} - \beta\psi - 1) + (e^{-\beta\psi_s} + \beta\psi - 1)}$$



# Derivation: space charge

Therefore, the surface field of the MOS structure,

$$E_s = -\frac{d\psi}{dx}$$

From the surface field, we can deduce the total space charge per unit area by applying Gauss's law:

$$Q_s = -E_s \epsilon_s = \epsilon_s \cdot \sqrt{\frac{2 p_{po} k T}{\epsilon_s}} \cdot \sqrt{\frac{n_{po}}{p_{po}} \left( e^{\beta \psi_s} - \beta \psi_s - 1 \right) + \left( e^{-\beta \psi_s} + \beta \psi_s - 1 \right)}$$

$$\text{or, } Q_s = \sqrt{2 \epsilon_s p_{po} k T} \cdot \sqrt{\frac{n_{po}}{p_{po}} \left( e^{\beta \psi_s} - \beta \psi_s - 1 \right) + \left( e^{-\beta \psi_s} + \beta \psi_s - 1 \right)}$$

$$\text{or, } Q_s = \sqrt{2 \epsilon_s p_{po} k T} \cdot F \left( \beta \psi_s, \frac{n_{po}}{p_{po}} \right)$$

$$\text{where, } F \left( \beta \psi_s, \frac{n_{po}}{p_{po}} \right) = \sqrt{\frac{n_{po}}{p_{po}} \left( e^{\beta \psi_s} - \beta \psi_s - 1 \right) + \left( e^{-\beta \psi_s} + \beta \psi_s - 1 \right)}$$





# Derivation: space charge

A typical variation of space charge density  $Q_s$  will be as follows:

i).  $\psi$  is negative:  $Q_s$  is positive  $\rightarrow$  corresponds to accumulation and function  $F$  is dominated by the first term:

$$\Rightarrow Q_s \propto e^{\frac{q|\psi_s|}{2kT}}$$

ii).  $\psi = 0$ :  $Q_s = 0 \rightarrow$  corresponds to flatband condition.

iii). For  $2\psi_B > \psi_s > 0$ :  $Q_s$  is negative and we get depletion and weak inversion condition. The function  $F$  is now dominated by the second term.

$$\Rightarrow Q_s \propto \sqrt{\psi_s}$$

iv). For  $\psi_s > 2\psi_B$ , we will have strong inversion condition with  $F$  dominated by the fourth term.

$$\Rightarrow Q_s \propto e^{\frac{q|\psi_s|}{2kT}}$$

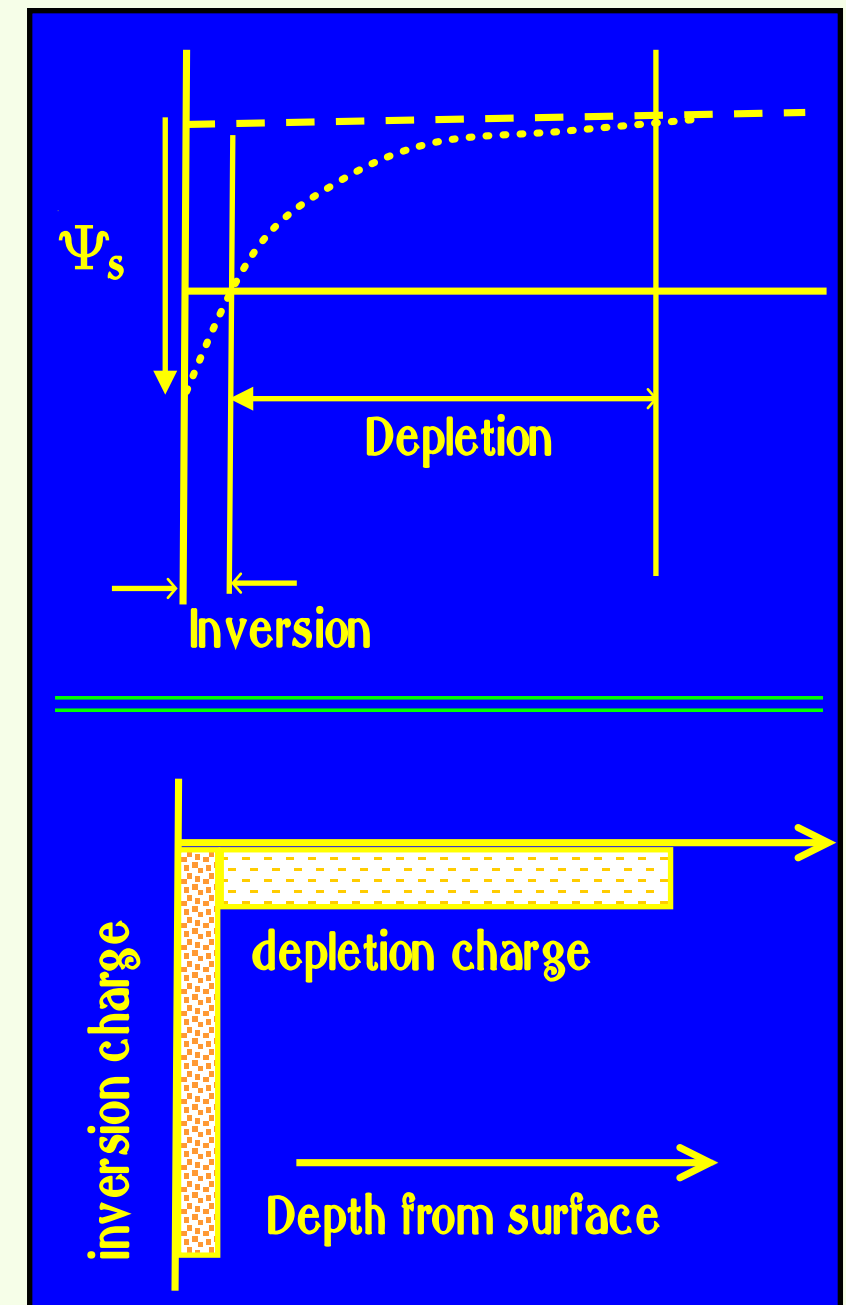
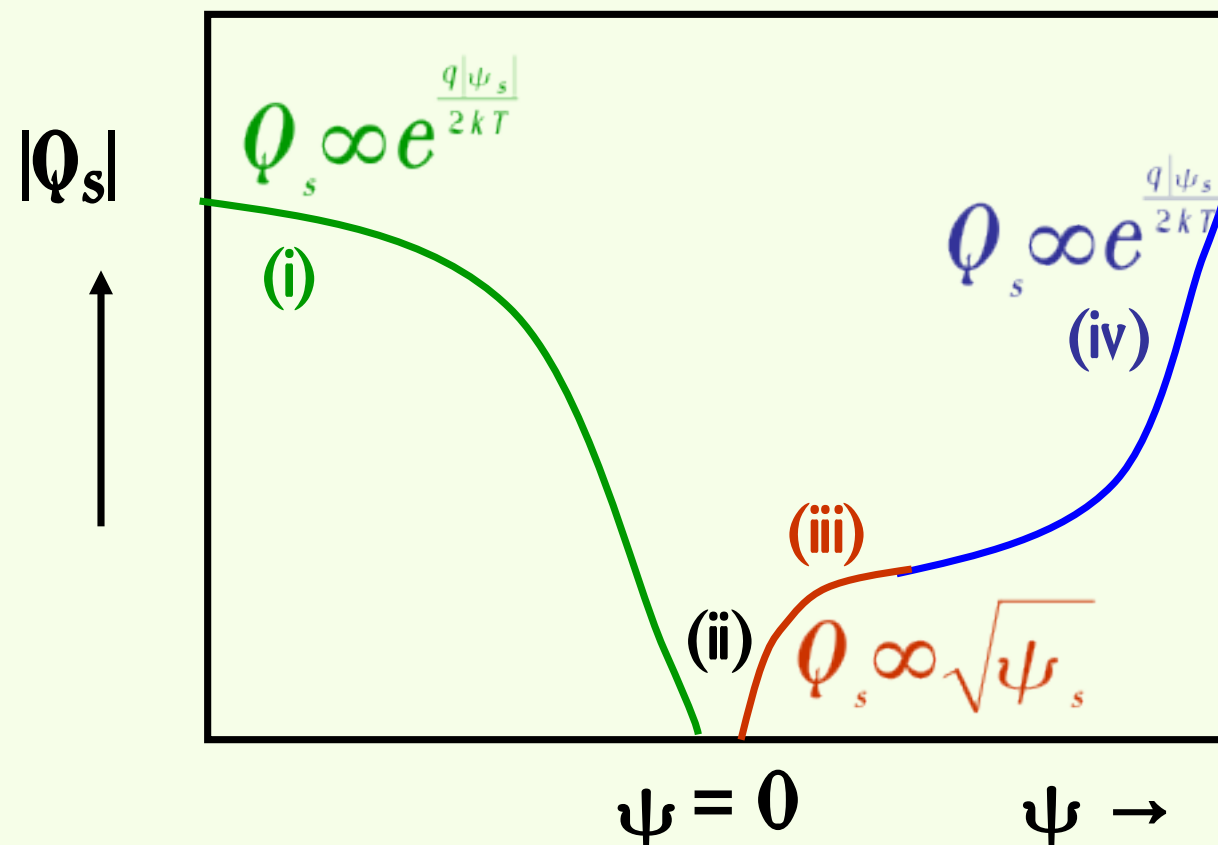




# Derivation: space charge

- Also note that this strong inversion begins at a surface potential,

$$\psi_s(\text{strong inversion}) \approx 2\psi_p = \frac{2kT}{q} \ln\left(\frac{N_A}{n_i}\right)$$

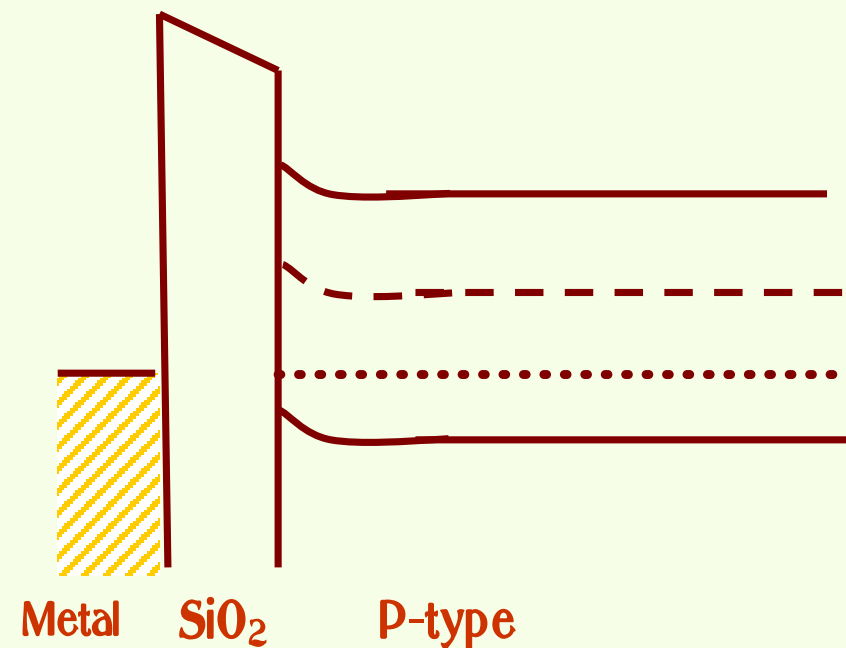
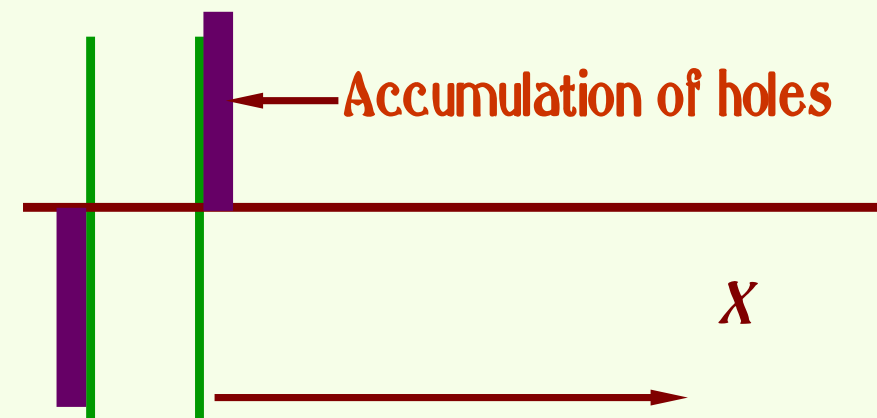
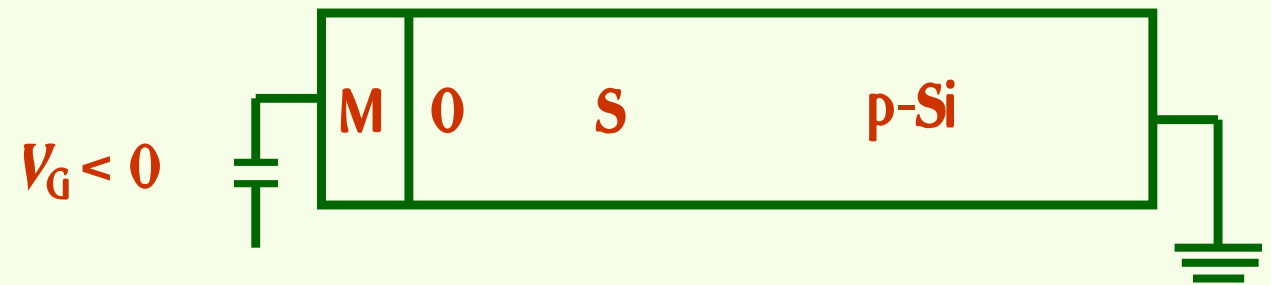
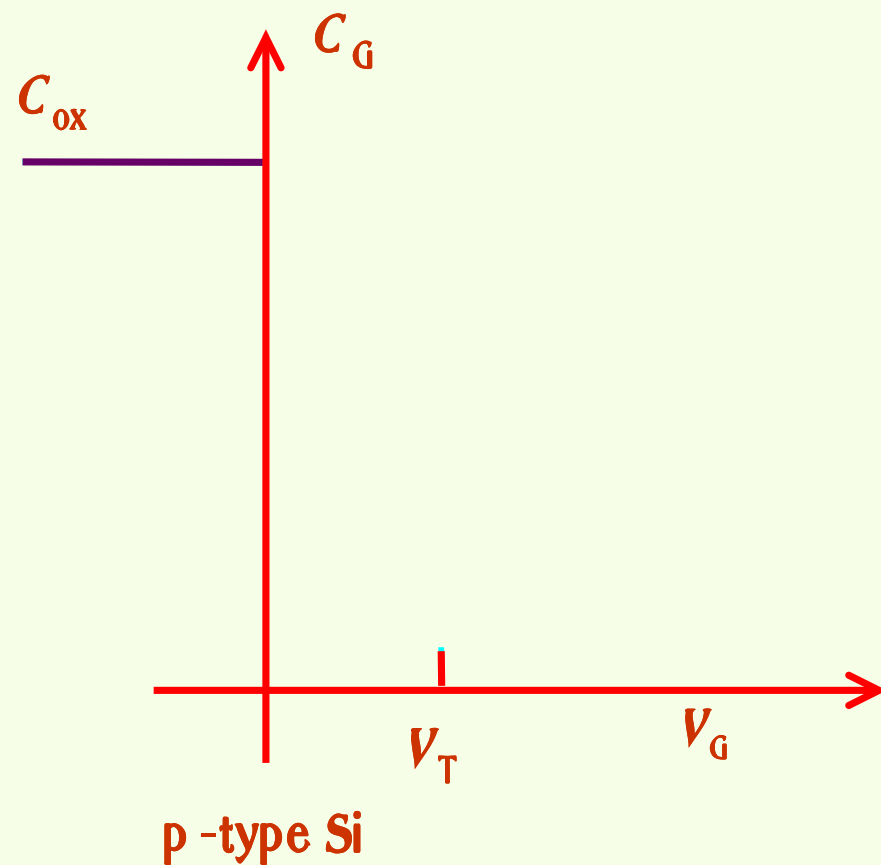




# MOS - capacitor under accumulation

$$V_G < 0; C_G = C_{ox}$$

$$\text{where, } C_{ox} = \frac{\epsilon_{ox} A}{T_{ox}}$$



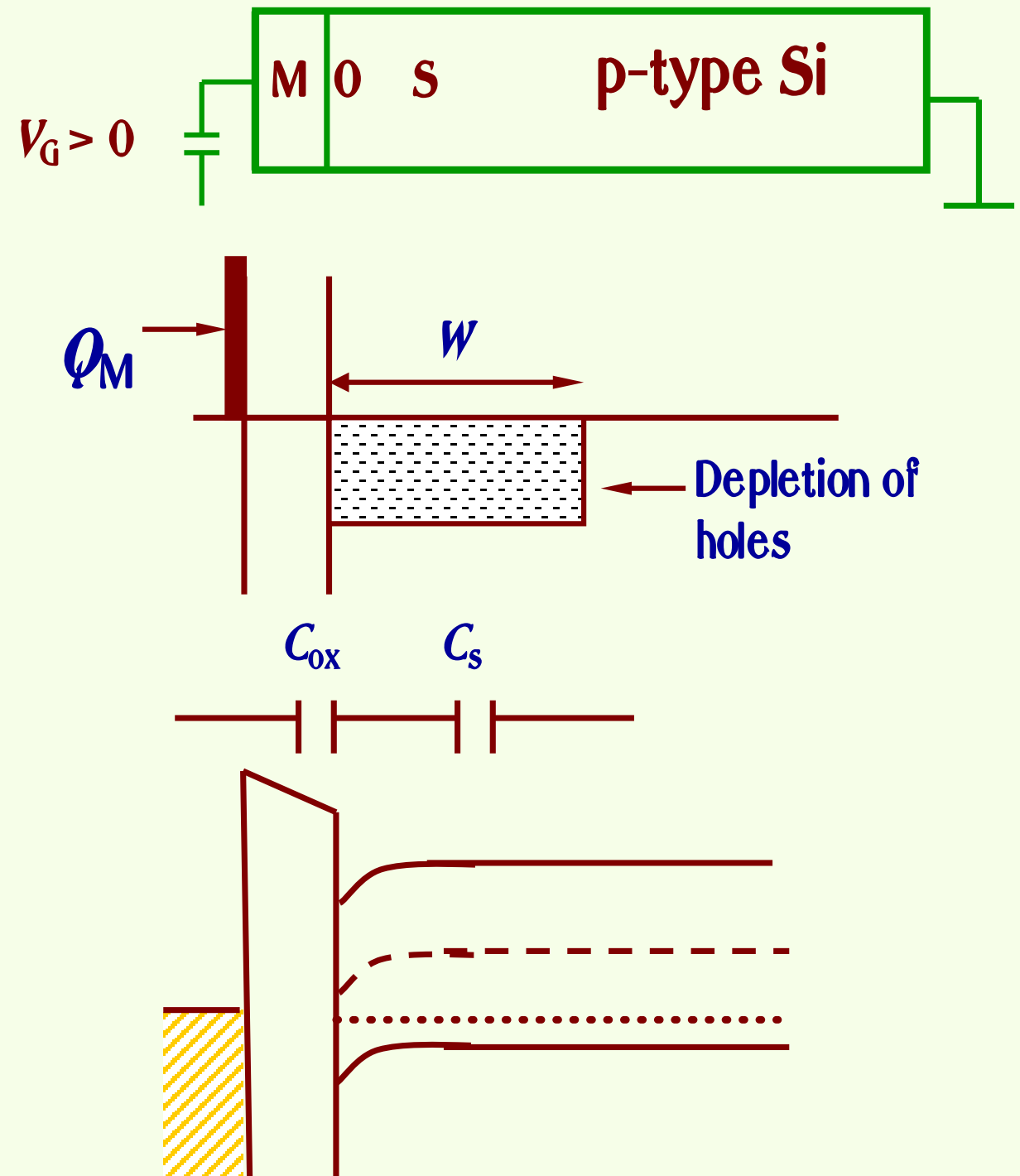
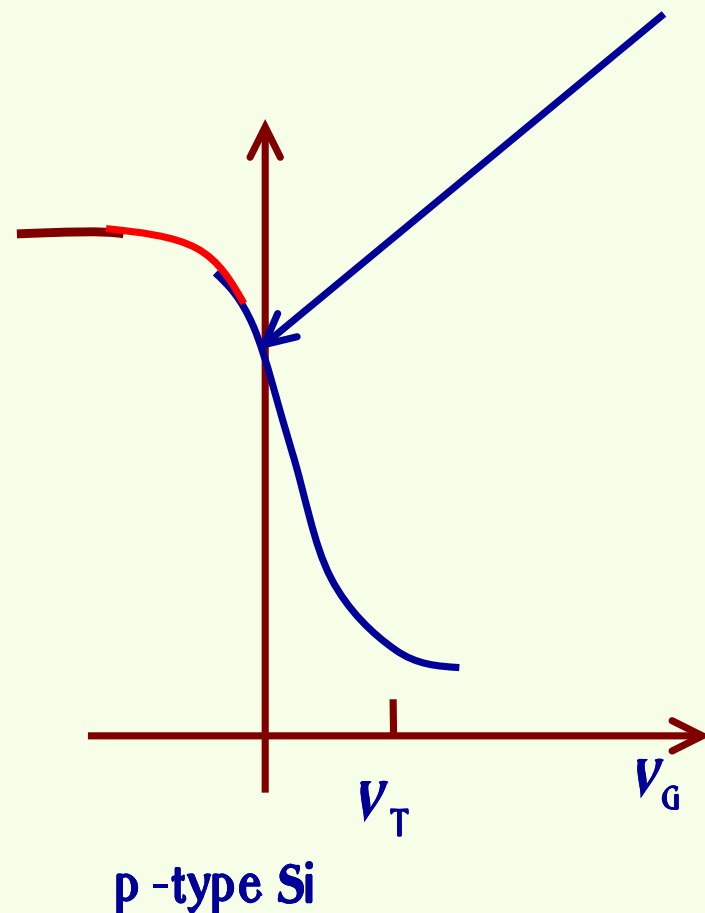


# MOS capacitor under depletion

- Depletion condition:  $V_G > 0$

$$C_{ox} = \frac{\epsilon_{ox} A}{T_{ox}} \quad C_s = \frac{\epsilon_s A}{W}$$

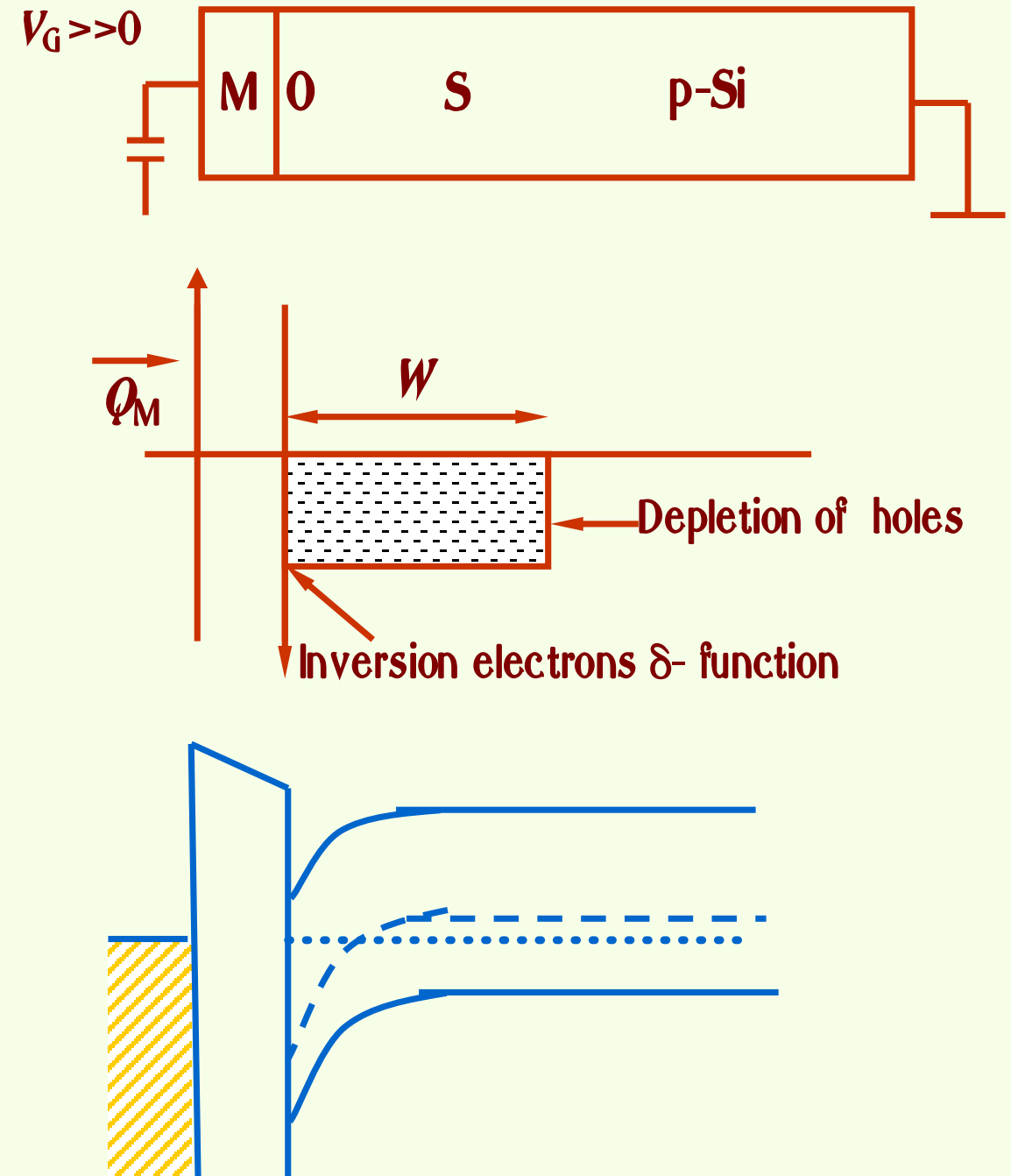
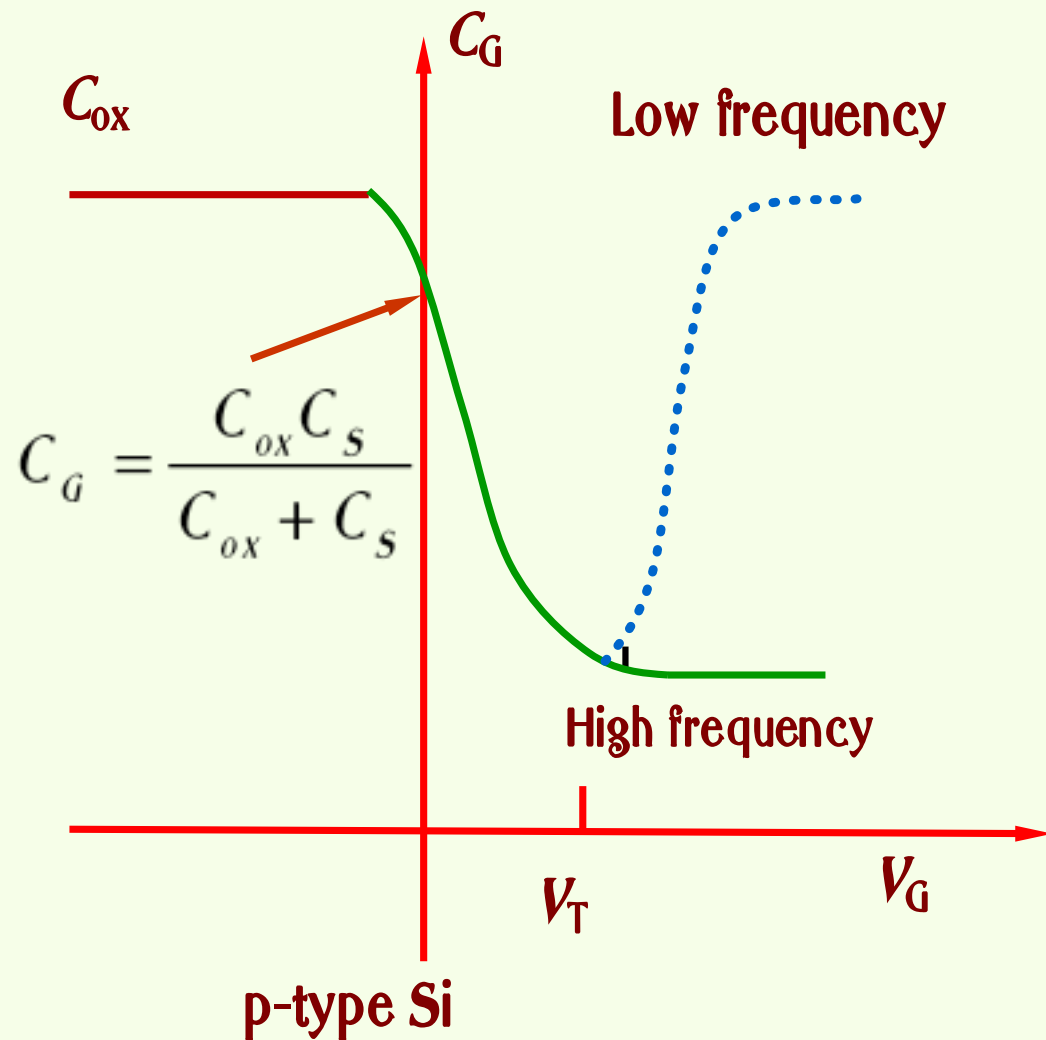
$$\frac{1}{C_G} = \frac{1}{C_{ox}} + \frac{1}{C_s} \Rightarrow C_G = \frac{C_{ox} C_s}{C_{ox} + C_s}$$





# MOS capacitor under inversion

- Inversion condition  $\psi_s = 2\psi_B$





# Graphical variation of potential, field, charge

$$V_{ox} = E_{ox} T_{ox} = \frac{|Q_s| T_{ox}}{\epsilon_{ox}} = \frac{|Q_s|}{C_{ox}}$$

- The total capacitance  $C$  of the system is a series combination of the insulator capacitance

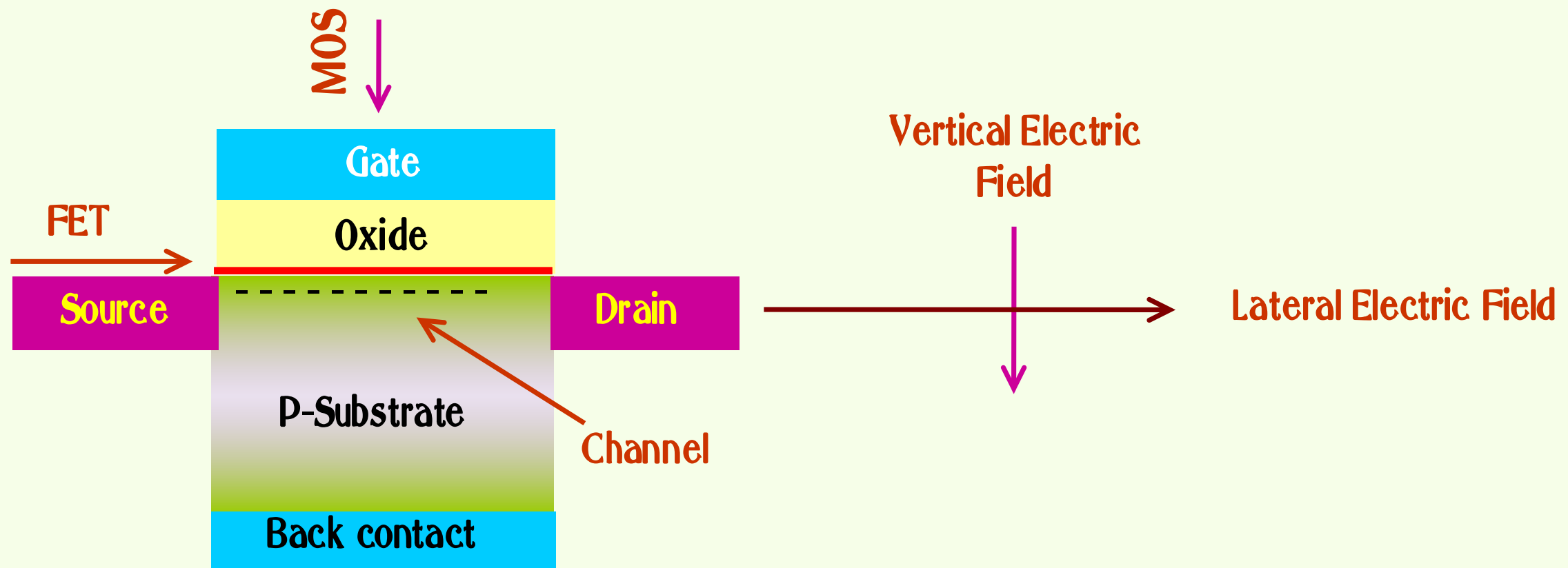
$$C_{ox} = \frac{\epsilon_{ox}}{T_{ox}}$$

and the semiconductor depletion-layer capacitance  $C_D$ :

$$C = \frac{C_{ox} C_s}{C_{ox} + C_s}$$



# MOSFET



Two orthogonal electric fields work together to initiate the operation of a MOSFET. Vertical field applied from the gate creates a channel for the carriers and lateral electric field drags the carriers from source to the drain, leading to generate a current along the channel.