

Chapter 1

Discontinuous Galerkin model order reduction of geometrically parametrized Stokes flow

Nirav Vasant Shah, Martin Hess and Gianluigi Rozza

Abstract The present work focuses on geometrical parametrization and reduced order modeling of Stokes flow. The importance of Stokes flow, advantages of discontinuous Galerkin method are introduced first. We also discuss the concept of geometric parametrization and its application along with importance of reduced order model technique. The full order model is based on discontinuous Galerkin method interior penalty formulation. The concepts of broken Sobolev spaces, relevant norms, jump and mean operator are introduced. The weak formulation is derived based in suitable space to obtain the full order model. We then introduce the concept of geometric parametrization. The operators are transformed from fixed domain to parameter dependent domain by exploring affine parameter dependence which results in efficient assembly of system matrix. Thereafter, proper orthogonal decomposition is applied to obtain basis for function space for reduced order model. By using Galerkin projection the linear system to be solved is projected onto reduced space. During the process, offline-online decomposition is used to separate computation of expensive parameter independent part and fast parameter independent part. Finally the technique is applied to test problem. The numerical outcomes presented include the experimental error analysis, eigenvalue computation and measurement of online simulation time. [13]

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1.1 Introduction

The subject of mathematical applications in fluid mechanics starts with one of the variants of the Navier-Stokes equations, such as the Stokes equation. Almost all processes of fluid mechanics require considerations related to the Navier-Stokes equations. Navier-Stokes equation is non-linear, characterizing flow fluctuations. However, in case of laminar flow, i.e. when fluctuations are negligible, this linearized form of the Navier-Stokes equation is the Stokes equation.

Discontinuous Galerkin method (DGM) has found traction as numerical method for elliptic problems **pereire reference** as well as hyperbolic problems **Book on compressible flow reference**. This is due to its several advantages over Finite Element Method (FEM) and Finite Volume Method (FVM). In fact, DG method is considered as combination of FEM and FVM. DGM uses polynomial approximation of suitable degree providing higher accuracy as well as allows discontinuity at the interface, by the concept of numerical flux, allowing greater flexibility. This fact makes DGM naturally attractive to problems such as shock capturing due to presence of steep gradients or discontinuities. Additionally, since the Dirichlet conditions are applied as boundary penalty, it avoids necessity to work with subspace of FEM. Several variants of DGM exist based on computational advantages such as sparsity pattern or extension of computational stencil, complexity of numerical implementation etc.

Geometric parametrization has emerged as important application of Parametric Partial Differential Equations (PPDEs) and as alternative to shape optimization. The concept of geometric parametrization allows to transfer operator evaluated on one domain to another domain efficiently. For linear equations, this means exploiting affine parameter dependence as will be shown in later section. Model Order Reduction (MOR) on the other hand allows reducing the size of the system to be solved and working with the smaller system containing only dominant components and discarding the non-dominant components. It is pertinent to mention that identifying "dominant" components is critical to the success of model order reduction strategy. Optimization of engineering components using geometric parametrization combined with MOR for PPDEs has given quite useful results in the fields such as mechanical, naval and aeronautic designs. Also, the faster computations obtained by MOR has helped in many query context, real time computation and quick transfer of computational results to industrial problems.

In the present work, we first introduce Discontinuous Galerkin Interior Penalty Method (DG-IPM). We subsequently introduce notion of parametrization characterizing geometry of the domain under consideration, exploit affine parameter dependence and its application in the context of offline-online decomposition. We then apply Proper Orthogonal Decomposition (POD) for constructing reduced basis space and apply Galerkin projection to project the system of equations on the space constructed by POD. Finally we present a test problem to demonstrate the introduced method and present numerical result.

1.2 Geometric parametrization

Consider domain $\Omega = \Omega(\mu) \in \mathbb{R}^d$ as open bounded domain. The parameter set $\mu \in \mathbb{P}$, where \mathbb{P} is parameter space, completely characterizes the domain. Also, consider a parameter set $\bar{\mu} \in \mathbb{P}$, as the known parameter set and $\Omega(\bar{\mu})$ as the reference domain, whose configuration is completely known. The mapping $F(\cdot, \mu) : \Omega(\bar{\mu}) \rightarrow \Omega(\mu)$ links reference domain and parametrized domain. Divide the domain $\Omega(\mu)$ into n_{su} subdomains such that $\Omega(\mu) = \bigcup_{i=1}^{n_{su}} \Omega_i(\mu)$, $\Omega_i(\mu) \cap \Omega_j(\mu) = \emptyset$, for $i \neq j$. The boundary of $\Omega(\mu)$, $\partial\Omega(\mu)$ is divided into Neumann boundary $\Gamma_N(\mu)$ and Dirichlet boundary $\Gamma_D(\mu)$ i.e. $\partial\Omega(\mu) = \Gamma_N(\mu) \cup \Gamma_D(\mu)$.

In the case of affine transformation, F is of the form,

$$x = F(\hat{x}, \mu) = G_F(\mu)\hat{x} + c_F(\mu); \forall x \in \Omega, \hat{x} \in \hat{\Omega}. \quad (1.1)$$

The inverse map T is expressed in the form,

$$\hat{x} = T(x, \mu) = G_T(\mu)x + c_T(\mu); \forall x \in \Omega, \hat{x} \in \hat{\Omega}. \quad (1.2)$$

1.3 Discontinuous Galerkin formulation

Each subdomain is divided into N_{el} number of triangular elements τ_k such that $\Omega = \bigcup_{k=1}^{N_{el}} \tau_k$. The triangulation \mathcal{T} is the set of all triangular elements i.e. $\mathcal{T} = \{\tau_k\}_{k=1}^{N_{el}}$. The internal boundary $\Gamma = \{\partial\tau_k\}_{k=1}^{N_{el}} \setminus \partial\Omega$. We represent \vec{n} as the outward pointing normal to an edge of element.

The Stokes's equation in strong form can be stated as,

$$-\nu\Delta\vec{u} + \nabla p = \vec{f}, \text{ in } \Omega, \quad (1.3)$$

$$\nabla \cdot \vec{u} = 0, \quad (1.4)$$

$$\vec{u} = \vec{u}_D, \text{ on } \Gamma_D, \quad (1.5)$$

$$-p\vec{n} + \nu\vec{n} \cdot \nabla\vec{u} = \vec{t}, \text{ on } \Gamma_N. \quad (1.6)$$

The vector variable velocity \vec{u} and scalar pressure p are the unknowns. ν is the material property known as kinematic viscosity. Vector \vec{f} is external force term or source term. \vec{u}_D is the Dirichlet velocity and \vec{t} is the Neumann value.

Before introducing weak form let us introduce broken Sobolev spaces for variables.

The space for velocity is

$$\mathbb{V} = \{\vec{\phi} \in (L^2(\mathcal{T}))^d \mid \vec{\phi} \in (P^D(\tau_k))^d, \tau_k \in \mathcal{T}\}. \quad (1.7)$$

The space for pressure is

$$\mathbb{Q} = \{\psi \in (L^2(\mathcal{T})) \mid \psi \in (P^{D-1}(\tau_k)), \tau_k \in \mathcal{T}\}. \quad (1.8)$$

Here, $P^D(\tau_k)$ denotes space of polynomials of degree at most $D \geq 2$ over τ_k .

In order to approximate the numerical flux we need the concept of Jump and Average operator. The superscript + refers to quantity in the element itself and the superscript - refers to quantity in the neighboring element (Figure 1.1).

The presence of normal vector \vec{n} in jump and average operator introduced below allows symmetric formulation and also ensures that jump of a vector is vector and jump of a scalar is scalar.

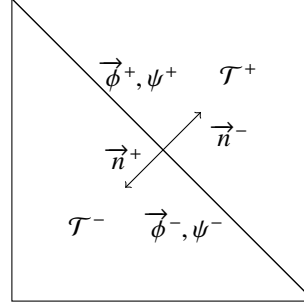


Fig. 1.1: Element self (+) and neighbouring element (-)

- For vector quantity $\vec{\phi}$:
 - Jump operator: $\left[\vec{\phi} \cdot \vec{n}\right] = \vec{\phi}^+ \cdot \vec{n}^+ + \vec{\phi}^- \cdot \vec{n}^-$ on Γ , $\left\{\vec{\phi}\right\} = \frac{\vec{\phi}^+ + \vec{\phi}^-}{2}$ on $\partial\Omega$.
 - Average operator: $\left\{\vec{\phi}\right\} = \frac{\vec{\phi}^+ + \vec{\phi}^-}{2}$ on Γ , $\left[\vec{\phi}\right] = \vec{\phi}$ on $\partial\Omega$.
- For scalar quantity ψ :
 - Jump operator: $[\psi] = \psi^+ + \psi^-$ on Γ , $[\psi] = \psi$ on $\partial\Omega$.
 - Average operator: $\{\psi\} = \frac{\psi^+ + \psi^-}{2}$ on Γ , $\{\psi\} = \psi$ on $\partial\Omega$.

The weak form of Stokes equation is as follow,

$$a_{IP}(\vec{u}, \vec{\phi}) + b(\vec{\phi}, p) + \left(\{p\}, [\vec{n} \cdot \vec{\phi}]\right)_{\Gamma \cup \Gamma_D} = l_{IP}(\vec{\phi}), \quad (1.9)$$

$$a_{IP}(\vec{u}, \vec{\phi}) = \left(\nabla \vec{u}, \nabla \vec{\phi}\right) + C_{11} \left([\vec{u}], [\vec{\phi}]\right)_{\Gamma \cup \Gamma_D} - \nu \left(\{\nabla \vec{u}\}, [\vec{n} \otimes \vec{\phi}]\right)_{\Gamma \cup \Gamma_D} - \nu \left([\vec{n} \otimes \vec{u}], \{\nabla \vec{\phi}\}\right)_{\Gamma \cup \Gamma_D}, \quad (1.10)$$

$$b(\phi, \psi) = - \int_{\mathcal{T}} \psi \nabla \cdot \vec{\phi}, \quad (1.11)$$

$$l_{IP}(\vec{\phi}) = \left(\vec{f}, \vec{\phi}\right) + \left(\vec{t}, \vec{\phi}\right)_{\Gamma_N} + C_{11} \left(\vec{u}_D, \vec{\phi}\right)_{\Gamma_D} - \left(\vec{n} \otimes \vec{u}_D, \nu \nabla \vec{\phi}\right)_{\Gamma_D}. \quad (1.12)$$

The penalty paramter $C_{11} > 0$ in $a_{IP}(\vec{u}, \vec{\phi})$ is an empirical constant to be kept large enough to maintain coercivity of bilinear form.

The weak for of continuity equation is as follow,

$$b(\vec{u}, \psi) + (\psi, [\vec{n} \cdot \vec{u}])_{\Gamma \cup \Gamma_D} = (\psi, \vec{n} \cdot \vec{u}_D)_{\Gamma_D} . \quad (1.13)$$

In matrix form system of equations can be written as,

$$\begin{pmatrix} A & B \\ B^T & 0 \end{pmatrix} \begin{pmatrix} U \\ P \end{pmatrix} = \begin{pmatrix} F_1 \\ F_2 \end{pmatrix} . \quad (1.14)$$

Stiffness matrix Solution vector Right hand side (Known)

Here, $A_{ij} = a_{IP}(\vec{\phi}_i, \vec{\phi}_j)$, $B_{ij} = b(\vec{\phi}_i, \psi_j) + (\{\psi_j\}, [n \cdot \phi_i])_{\Gamma \cup \Gamma_D}$, $F_1 = l_{IP}(\vec{\phi}_i)$ and $F_2 = (\psi_j, \vec{n} \cdot \vec{u}_D)_{\Gamma_D}$.

1.4 Affine expansion

We evaluate and solve the Stokes equation weak formulation on reference domain $\Omega_{\bar{\mu}}$. Given a parameter set $\mu \neq \bar{\mu}$ we need to evaluate the linear systems of equation (1.14) on new domain $\Omega(\mu)$. To accomplish this we use affine expansion using linear nature of equation and diving $\Omega(\bar{\mu})$ into triangular subdomains $\Omega_i(\bar{\mu})$, $i = \{1, 2, \dots, n_{su}\}$ as explained earlier in the section geometric parametrization [Section 1.2]. The transformation of operators in affine is standard in literature **ADD affine expansion literature** and only two relevant section as specific to DG formulation will be mentioned here as below.

In order to transfer the terms containing jump and average operator following approach is used in present analysis.

$$\left(\{\nabla \vec{\phi}\}, [\vec{n} \otimes \vec{\phi}] \right) = \left(\nabla \vec{\phi}^+, \vec{n}^+ \otimes \vec{\phi}^+ \right) + \left(\nabla \vec{\phi}^+, \vec{n}^- \otimes \vec{\phi}^- \right) + \left(\nabla \vec{\phi}^-, \vec{n}^+ \otimes \vec{\phi}^+ \right) + \left(\nabla \vec{\phi}^-, \vec{n}^- \otimes \vec{\phi}^- \right) . \quad (1.15)$$

Each term on right hand side of equation (1.15) can now be transformed using affine map.

The coercivity term $C_{11} \left([\vec{\phi}, \vec{u}] \right)_{\Gamma \cup \Gamma_D}$ is not transformed but used as evaluated on reference domain $\Omega(\bar{\mu})$. The affine transformation is given by,

$$C_{11} \left([\vec{\phi}(\hat{x}), \vec{u}(\hat{x})] \right)_{\Gamma(\mu) \cup \Gamma_D(\mu)} = C_{11} \alpha \left([\vec{\phi}(\mathbf{F}(\hat{x})), \vec{u}(\mathbf{F}(\hat{x}))] \right)_{\Gamma(\bar{\mu}) \cup \Gamma_D(\bar{\mu})} , \quad (1.16)$$

$$\alpha = \frac{meas(\Gamma(\mu) \cup \Gamma_D(\mu))}{meas(\Gamma(\bar{\mu}) \cup \Gamma_D(\bar{\mu}))} , \quad \hat{x} \in \Omega(\bar{\mu}), \quad x \in \Omega(\mu) .$$

Since, C_{11} is empirical coefficient replacing $C_{11}\alpha$ with C_{11} will not change the formulation as long as coercivity of a_{IP} over parameter space \mathbb{P} is maintained. In the present analysis, C_{11} is not calculated exactly but only order of magnitude required for C_{11} is estimated and C_{11} is kept one magnitude larger as safeguard against round-off errors. As long as the domain $\Omega(\mu)$ is not deformed much compared to its reference configuration $\Omega(\bar{\mu})$, C_{11} and $C_{11}\alpha$ will have the same order of magnitude. However, large deformations are anyhow not favorable as it will lead to bad mesh quality and in turn, will lead to poor DGM-approximation.

1.5 Reduced basis method

1.5.1 Snapshot proper orthogonal decomposition

We present now snapshot proper orthogonal decomposition method. Here, “snapshot” means solution calculated by discontinuous Galerkin method. We calculate solution based on $\mu_n, n \in \{1, \dots, n_s\}$ i.e. n_s snapshots are generated. We also introduce inner product matrices $\mathbf{M}_v \in \mathbb{R}^{\vec{u}_{ndofs} \times \vec{u}_{ndofs}}$ and $\mathbf{M}_p \in \mathbb{R}^{p_{ndofs} \times p_{ndofs}}$, where \vec{u}_{ndofs} and p_{ndofs} are number of velocity and pressure degrees of freedom respectively.

$$\mathbf{M}_v = \int_{\Omega} \vec{\phi}_i \cdot \vec{\phi}_j + \sum_{k=1}^{N_{el}} \int_{\tau_k} \nabla \vec{\phi}_i : \nabla \vec{\phi}_j, \quad i, j = 1, \dots, \vec{u}_{ndofs},$$

$$\mathbf{M}_p = \int_{\Omega} \psi_i \psi_j, \quad i, j = 1, \dots, p_{ndofs}.$$

We also introduce matrices storing velocity snapshots \mathbf{S}_v and storing pressure snapshots \mathbf{S}_p . We discuss the method only for velocity snapshots. The method is similar for pressure snapshots. We note the size of matrices, useful for matrix operations presented hereafter.

$$\mathbf{S}_v \in \mathbb{R}^{\vec{u}_{ndofs} \times n_s}, \quad \mathbf{S}_p \in \mathbb{R}^{p_{ndofs} \times n_s},$$

$$\mathbf{M}_v \in \mathbb{R}^{\vec{u}_{ndofs} \times \vec{u}_{ndofs}}, \quad \mathbf{M}_p \in \mathbb{R}^{p_{ndofs} \times p_{ndofs}}.$$

1.5.2 Spectral decomposition of snapshots

We denote the dimension of reduced basis as N and assert that $N \ll n_s$. We now perform the spectral decomposition of $\mathbf{S}_v^T \mathbf{M}_v \mathbf{S}_v$,

$$\mathbf{S}_v^T \mathbf{M}_v \mathbf{S}_v = \mathbf{V} \mathbf{\Theta} \mathbf{V}^T. \quad (1.17)$$

The columns of V are eigenvectors and Θ has eigenvalues θ_i , $1 \leq i \leq n_s$ such that,

$$\Theta_{ij} = \theta_i \delta_{ij} . \quad (1.18)$$

We also note that $\theta_i > 0$ and $\theta_1 \geq \theta_2 \geq \dots \geq \theta_{n_s}$ i.e. the eigenvalues are in sorted order. We form the reduced basis by linear combination of the snapshot vector,

$$\mathbf{B}_v = \mathbf{S}_v \mathbf{A} , \mathbf{A} \in \mathbb{R}^{n_s \times N} . \quad (1.19)$$

Here, \mathbf{B}_v is defined such that, if $\vec{\phi} \in \mathbb{R}^{n \times d}$, $\mathbf{B}_v \in \mathbb{R}^{n \times N}$, the reduced basis for velocity $\vec{\phi}_N \in \mathbb{R}^{d \times N}$ is formed by,

$$\vec{\phi}_N = \vec{\phi}^T \mathbf{B}_v . \quad (1.20)$$

Considering orthonormality of reduced basis $\vec{\phi}_N$ with respect to inner product \mathbf{M}_v ,

$$\langle \vec{\phi}_N, \vec{\phi}_N \rangle_{\mathbf{M}_v} = \mathbf{B}_v^T \mathbf{M}_v \mathbf{B}_v = \mathbf{I} . \quad (1.21)$$

Considering above orthonormality, we express matrix \mathbf{A} as,

$$\mathbf{A} = \mathbf{V} \Theta^{-\frac{1}{2}} \mathbf{R} , \mathbf{R} \in \mathbb{R}^{n_s \times N} , \mathbf{R}^T \mathbf{R} = \mathbf{I} . \quad (1.22)$$

where, \mathbf{I} is identity matrix of suitable size.

We set now \mathbf{R} as,

$$\mathbf{R} = [\mathbf{I}_{N \times N}; \mathbf{0}_{(n_s \times N)}] \text{ and accordingly } \mathbf{B}_v = \mathbf{S}_v \mathbf{V} \Theta^{-\frac{1}{2}} \mathbf{R} . \quad (1.23)$$

1.5.3 Galerkin reduced basis formulation

We now present the reduced bilinear form as,

$$a(u_N, \phi_N; \mu) + b(p_N, \phi_N; \mu) = f_1(\phi_N, \mu), \quad (1.24)$$

$$b(u_N, \psi_N; \mu) = f_2(\psi_N, \mu). \quad (1.25)$$

In discrete form, we form reduced equation as,

$$\begin{pmatrix} \mathbf{B}_v^T \mathbf{A}(\mu) \mathbf{B}_v & \mathbf{B}_v^T \mathbf{B}(\mu) \mathbf{B}_v \\ \mathbf{B}_p^T \mathbf{B}(\mu)^T \mathbf{B}_v & \mathbf{0} \end{pmatrix} \begin{pmatrix} U_N \\ P_N \end{pmatrix} = \begin{pmatrix} \mathbf{B}_v^T F_1(\mu) \\ \mathbf{B}_p^T F_2(\mu) \end{pmatrix}, \quad (1.26)$$

$\tilde{\mathbf{K}} \qquad \qquad \qquad \zeta \qquad \qquad \qquad \tilde{\mathbf{F}}$

and accordingly we solve following variational form for reduced degrees of freedom ζ ,

$$\tilde{\mathbf{K}} \zeta = \tilde{\mathbf{F}} , \quad (1.27)$$

and calculate reduced solutions \vec{u}_N and p_N as,

$$\vec{u}_N = \mathbf{B}_v U_N, \quad p_N = \mathbf{B}_p P_N \quad (1.28)$$

1.6 Offline-online procedure

The offline-online procedure is used to separate computationally intensive parameter independent offline procedure and faster parameter dependent online procedure **cite CRBM**. During the offline phase n_s snapshots are computed and reduced basis spaces \mathbf{B}_v and \mathbf{B}_p are created. The offline procedure is outlined in section 1.5.1 and section 1.5.2. During the online phase the systems of equations are projected on reduced space using Galerkin projection, the smaller systems of equation obtained by Galerkin projection is solved and the reduced basis solution is computed as outlined in section 1.5.3. The parameter dependent matrices in equation (1.26) are evaluated by using affine decomposition.

1.7 Numerical example

We perform the POD-Galerkin method as mentioned in section 1.5.1 - section 1.5.3. The boundary $x = 0$ is Dirichlet boundary with inflow velocity at point $(0, y)$ as $u = (y(1 - y), 0)$. The boundary $x = 1$ is a Neumann boundary with zero Neumann value i.e. $t = (0, 0)$. Other boundaries are Dirichlet boundary with no slip condition. The source term is $f = (0, 0)$.

The geometric parameters were coordinates of tip of the obstacle on bottom edge.

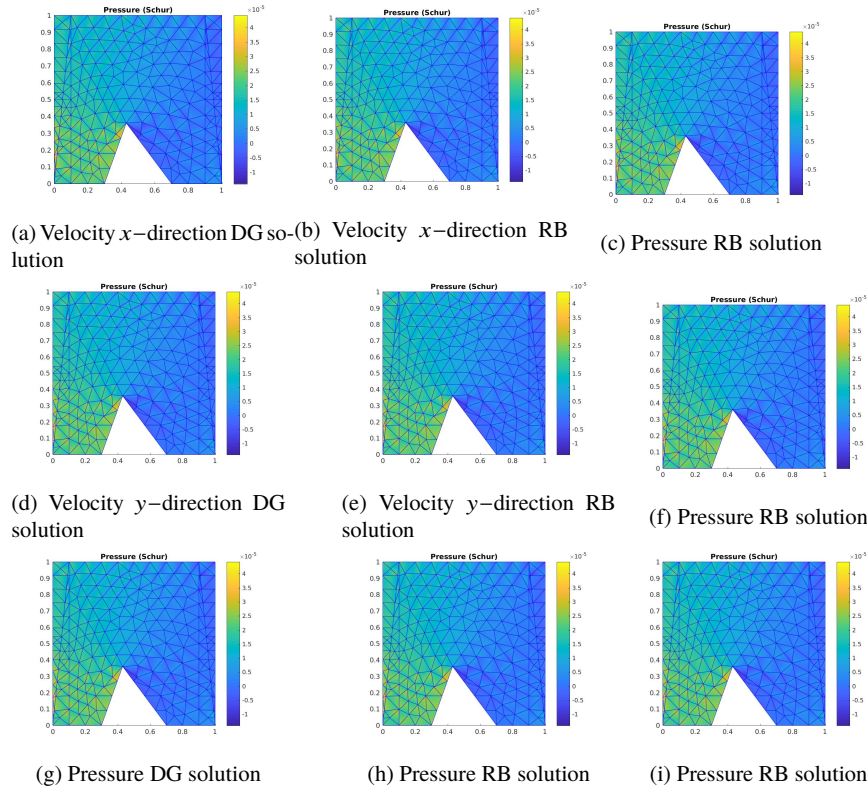
$$(x, y) = (\mu_1, \mu_2), \quad (x_{ref}, y_{ref}) = (0.5, 0.3),$$

with (x_{ref}, y_{ref}) as the reference parameters.

The training set was generated by random generation of 100 parameters between the interval $[0.4, 0.6] \times [0.4, 0.6]$. The test set contained 10 parameters between the interval $[0.4, 0.6] \times [0.4, 0.6]$.

ADD : experimental results: Eigenvalue drop, error vs basis function, RB,DG Solution and error at one parameter value

1. Livelihood and survival mobility are oftentimes coutcomes of uneven socioeconomic development.
 - a. Livelihood and survival mobility are oftentimes coutcomes of uneven socioeconomic development.
 - b. Livelihood and survival mobility are oftentimes coutcomes of uneven socioeconomic development.

Fig. 1.2: DG and RB solution $[\mu_x \mu_y] = [0.43 \ 0.36]$

2. Livelihood and survival mobility are oftentimes outcomes of uneven socioeconomic development.

Subparagraph Heading

In order to avoid simply listing headings of different levels we recommend to let every heading be followed by at least a short passage of text. Use the \LaTeX automatism for all your cross-references and citations as has already been described in Sect. ??, see also Fig. 1.7.

For unnumbered list we recommend to use the `itemize` environment – it will automatically be rendered in line with the preferred layout.

- Livelihood and survival mobility are oftentimes outcomes of uneven socioeconomic development, cf. Table 1.1.
 - Livelihood and survival mobility are oftentimes outcomes of uneven socioeconomic development.

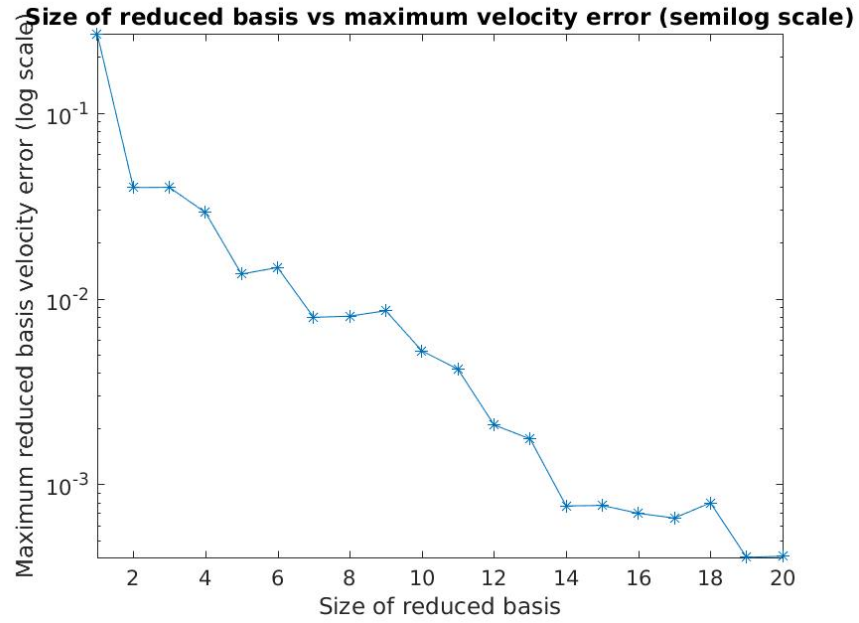


Fig. 1.3: Size of reduced basis for velocity vs Online simulation time (semilog scale)

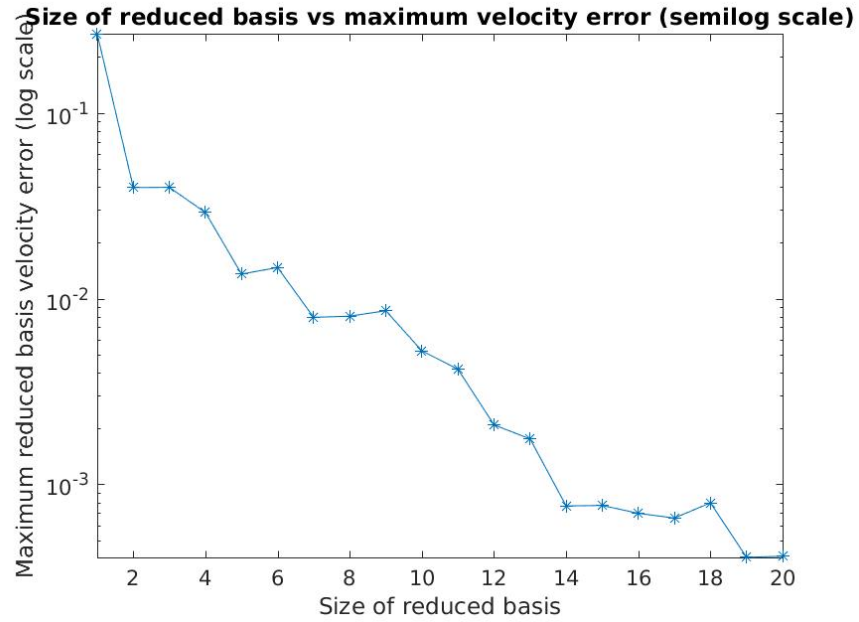


Fig. 1.4: Size of reduced basis for velocity vs Online simulation time (semilog scale)

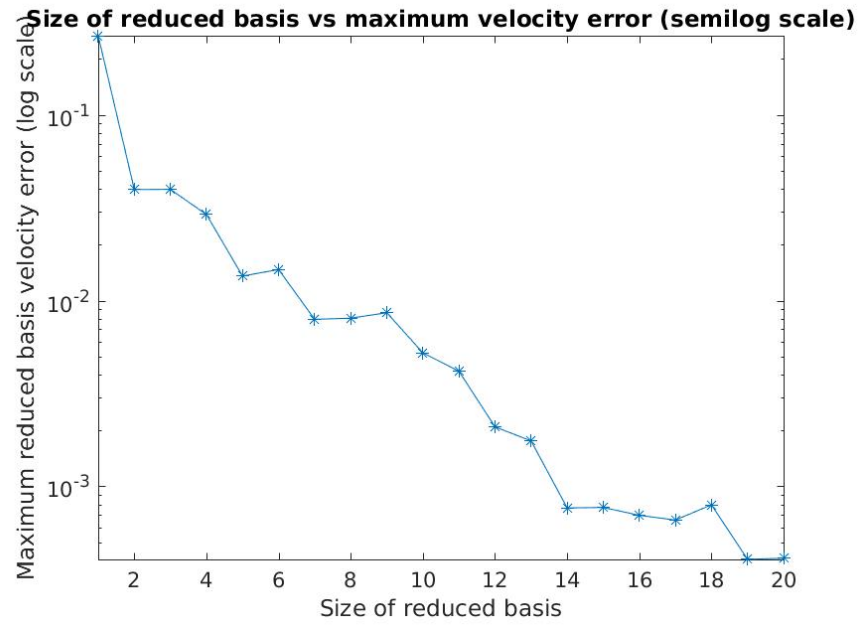


Fig. 1.5: Size of reduced basis for velocity vs Online simulation time (semilog scale)

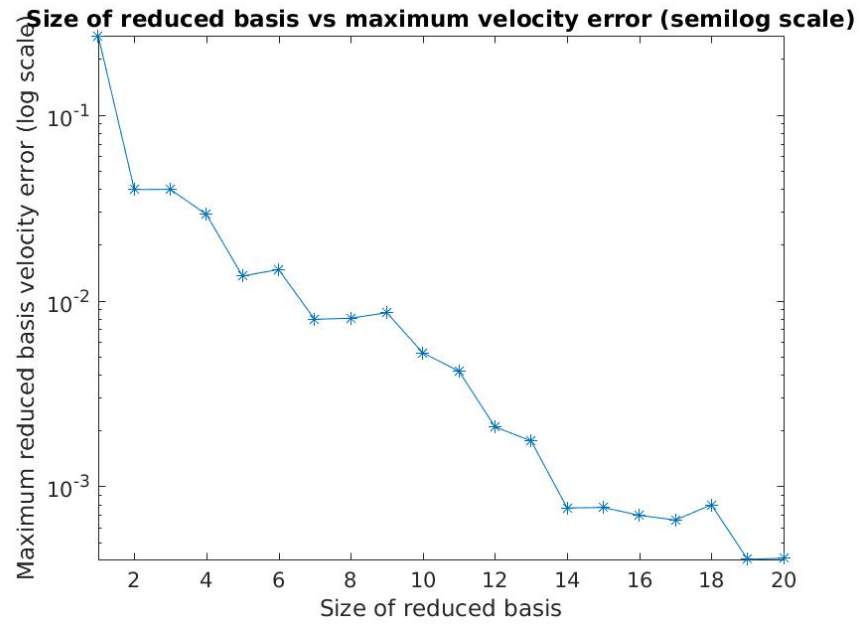


Fig. 1.6: Size of reduced basis for velocity vs Online simulation time (semilog scale)

Fig. 1.7 If the width of the figure is less than 7.8 cm use the `sidecaption` command to flush the caption on the left side of the page. If the figure is positioned at the top of the page, align the sidecaption with the top of the figure – to achieve this you simply need to use the optional argument `[t]` with the `sidecaption` command

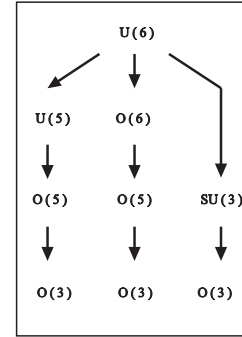


Table 1.1: Please write your table caption here

Classes	Subclass	Length	Action Mechanism
Translation	mRNA ^a	22 (19–25)	Translation repression, mRNA cleavage
Translation	mRNA cleavage	21	mRNA cleavage
Translation	mRNA	21–22	mRNA cleavage
Translation	mRNA	24–26	Histone and DNA Modification

^a Table foot note (with superscript)

- Livelihood and survival mobility are oftentimes coutcomes of uneven socioe-economic development.
- Livelihood and survival mobility are oftentimes coutcomes of uneven socioeco-nomic development.

Run-in Heading Boldface Version Use the \LaTeX automatism for all your cross-references and citations as has already been described in Sect. ??.

Run-in Heading Boldface and Italic Version Use the \LaTeX automatism for all your cross-references and citations as has already been described in Sect. ??.

Run-in Heading Displayed Version

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1.8 Section Heading

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- Type 1 That addresses central themes pertainng to migration, health, and disease. In Sect. ??, Wilson discusses the role of human migration in infectious disease distributions and patterns.
- Type 2 That addresses central themes pertainng to migration, health, and disease. In Sect. ??, Wilson discusses the role of human migration in infectious disease distributions and patterns.

1.8.1 Subsection Heading

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If you want to emphasize complete paragraphs of texts we recommend to use the newly defined class option and environment `svgraybox`. This will produce a 15 percent screened box 'behind' your text.

1.8.1.1 Subsubsection Heading

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Theorem 1 *Theorem text goes here.*

Definition 1 Definition text goes here.

Proof Proof text goes here. □

Paragraph Heading

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Theorem 2 *Theorem text goes here.*

Definition 2 Definition text goes here.

Proof Proof text goes here. □

Trailer Head

If you want to emphasize complete paragraphs of texts in an `Trailer Head` we recommend to use

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\begin{trailer}{Trailer Head}
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\end{trailer}
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? Questions

If you want to emphasize complete paragraphs of texts in an `Questions` we recommend to use

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\end{question}
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> Important

If you want to emphasize complete paragraphs of texts in an **Important** we recommend to use

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\begin{important}{Important}  
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\end{important}
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! Attention

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\begin{warning}{Attention}  
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\end{warning}
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Program Code

If you want to emphasize complete paragraphs of texts in an **Program Code** we recommend to use

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\begin{programcode}{Program Code}  
\begin{verbatim}...\end{verbatim}  
\end{programcode}
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Tips

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\begin{tips}{Tips}  
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Overview

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Background Information

If you want to emphasize complete paragraphs of texts in an **Background Information** we recommend to use

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Legal Text

If you want to emphasize complete paragraphs of texts in an **Legal Text** we recommend to use

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\begin{legalttext}{Legal Text}
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\end{legalttext}
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Acknowledgements If you want to include acknowledgments of assistance and the like at the end of an individual chapter please use the `acknowledgement` environment – it will automatically be rendered in line with the preferred layout.

Appendix

When placed at the end of a chapter or contribution (as opposed to at the end of the book), the numbering of tables, figures, and equations in the appendix section continues on from that in the main text. Hence please *do not* use the `appendix` command when writing an appendix at the end of your chapter or contribution. If there is only one the appendix is designated “Appendix”, or “Appendix 1”, or “Appendix 2”, etc. if there is more than one.

$$a \times b = c \tag{1.29}$$

References

References

References may be *cited* in the text either by number (preferred) or by author/year.¹ If the citation in the text is numbered, the reference list should be arranged in ascending order. If the citation in the text is author/year, the reference list should be *sorted* alphabetically and if there are several works by the same author, the following order should be used:

1. all works by the author alone, ordered chronologically by year of publication
2. all works by the author with a coauthor, ordered alphabetically by coauthor
3. all works by the author with several coauthors, ordered chronologically by year of publication.

The *styling* of references² depends on the subject of your book:

- The *two* recommended styles for references in books on *mathematical, physical, statistical and computer sciences* are depicted in [1, 2, 3, 4, 5] and [6, 7, 8, 9, 10].
- Examples of the most commonly used reference style in books on *Psychology, Social Sciences* are [11, 12, 13, 14, 15].
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- Examples of the basic Springer Nature style used in publications on a wide range of subjects such as *Computer Science, Economics, Engineering, Geosciences, Life Sciences, Medicine, Biomedicine* are [21, 22, 24, 23, 25].

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¹ Make sure that all references from the list are cited in the text. Those not cited should be moved to a separate *Further Reading* section or chapter.

² Always use the standard abbreviation of a journal's name according to the *ISSN List of Title Word Abbreviations*, see <http://www.issn.org/en/node/344>

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