

# Spectroscopy



Suzanne Ramsay ([sramsay@eso.org](mailto:sramsay@eso.org))  
David Dunlap Instrumentation School,  
July 2019

# Outline of the talk

## ■ Who am I?

- A career in pictures

## ■ Spectroscopy

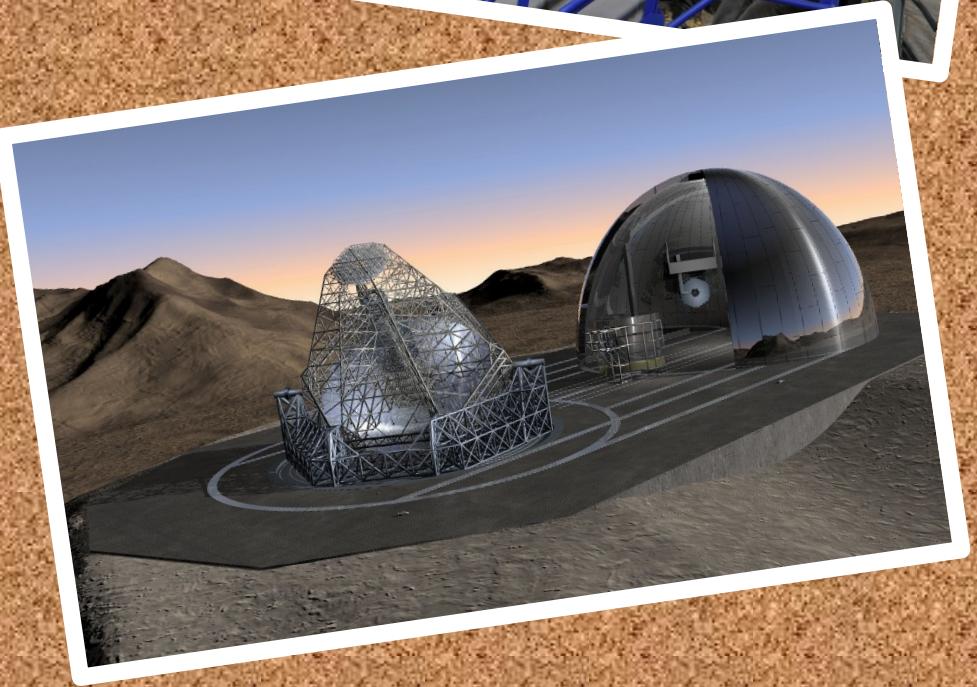
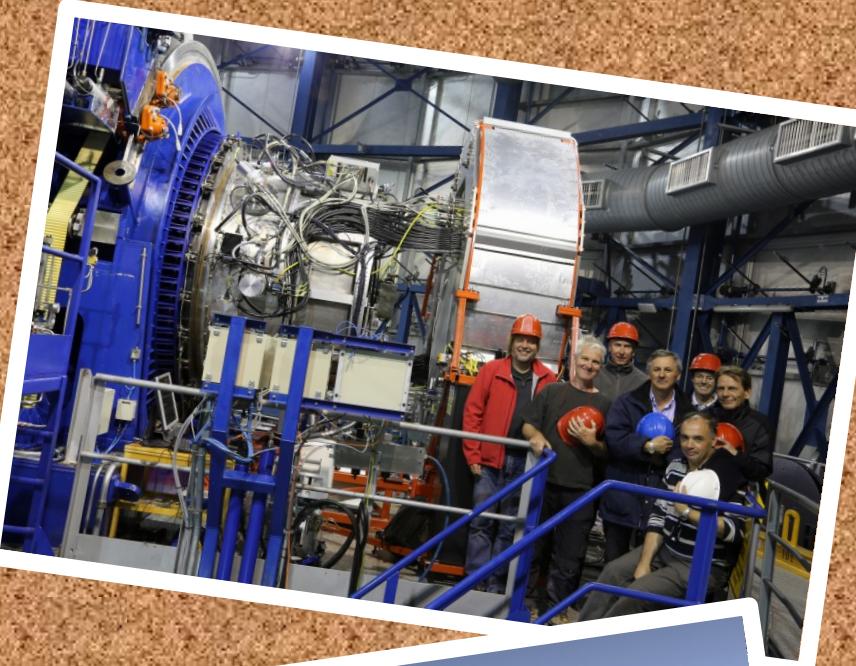
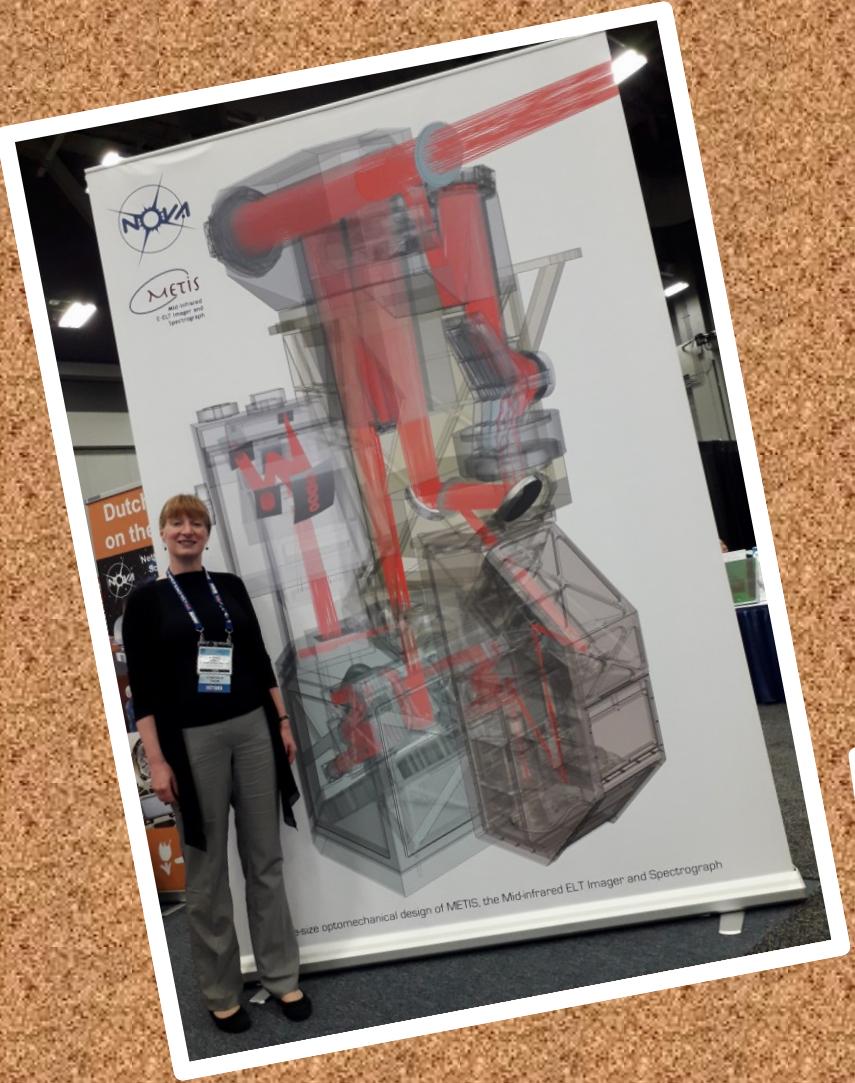
- Refresh: Spectrograph basics
- Dispersion technologies in use today
- Variations on the simple spectrograph
  - Multi-object; Integral field; Cross-dispersed; Ultra-stable
- Back of the envelope spectrograph design

## ■ Real life example(s): spectrographs for an Extremely Large Telescope



# A career in pictures



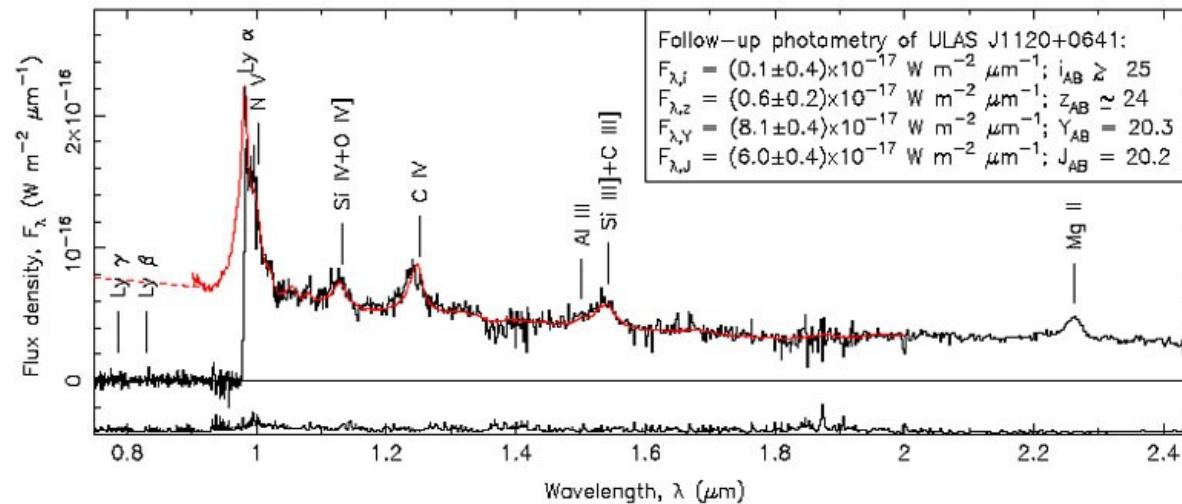




# Spectroscopy

# Why spectroscopy?

- Imaging informs us about what is out there and provides basic information
  - Imaging in multiple filters provides e.g. basic information about stellar temperatures or photometric redshifts
- To take the next step and start to understand the detailed astrophysics, we need to take spectra



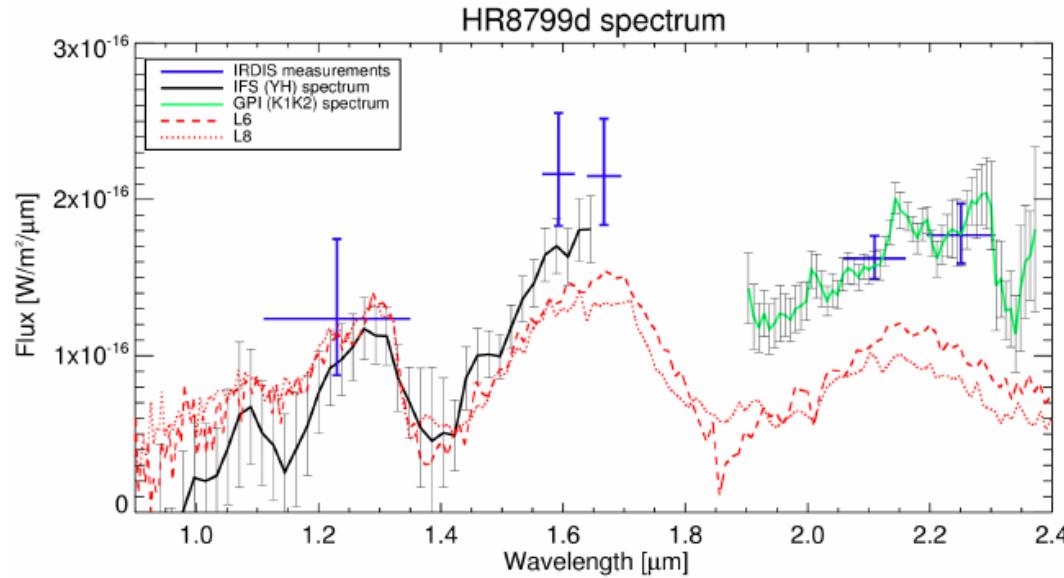
Z= 7.085 quasar; Mortlock et al. 2011

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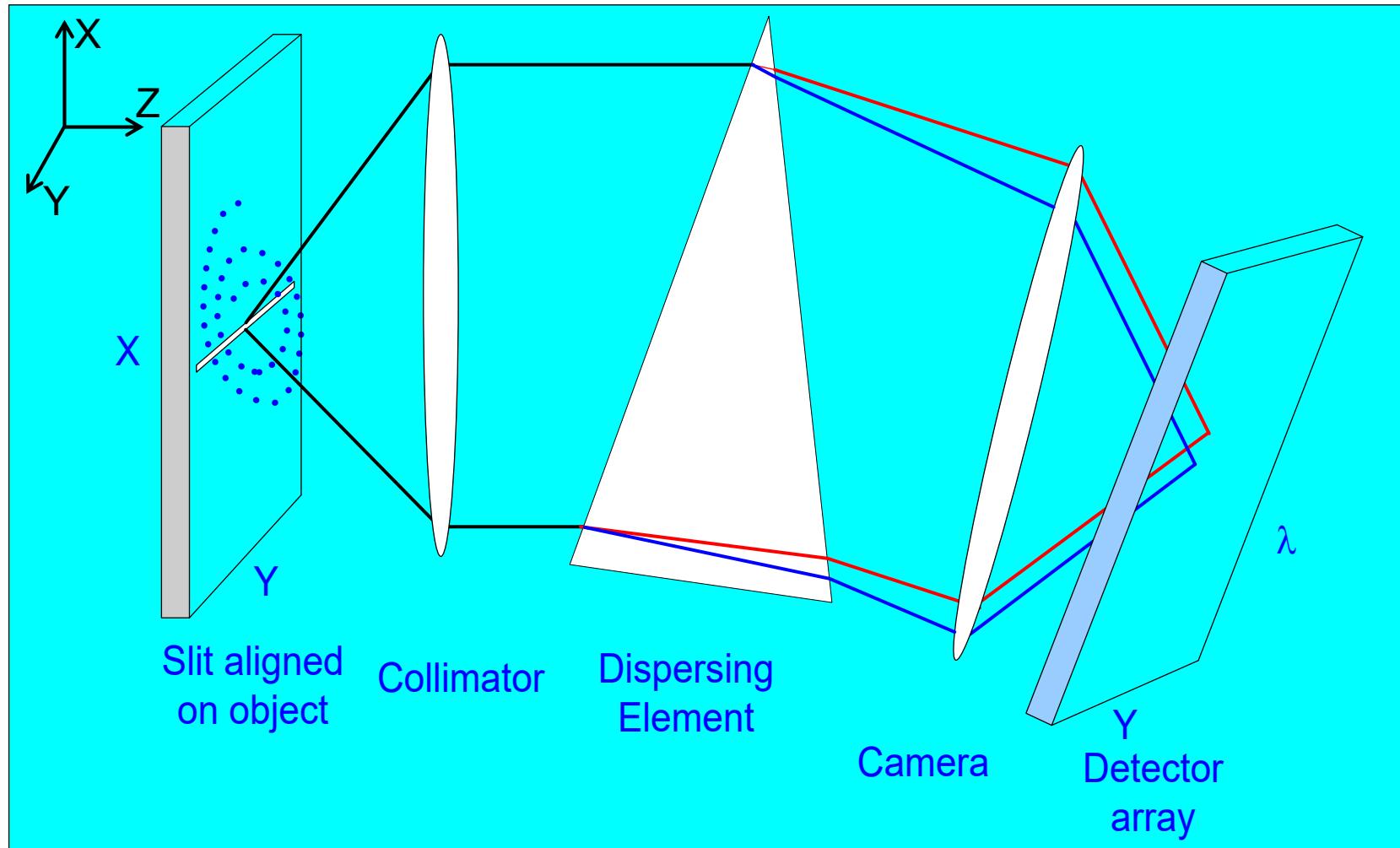
# Why spectroscopy?

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  - Imaging in multiple filters provides e.g. basic information about stellar temperatures or photometric redshifts
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Exoplanet spectrum compared with models of brown dwarfs, Zurlo et al. 2015

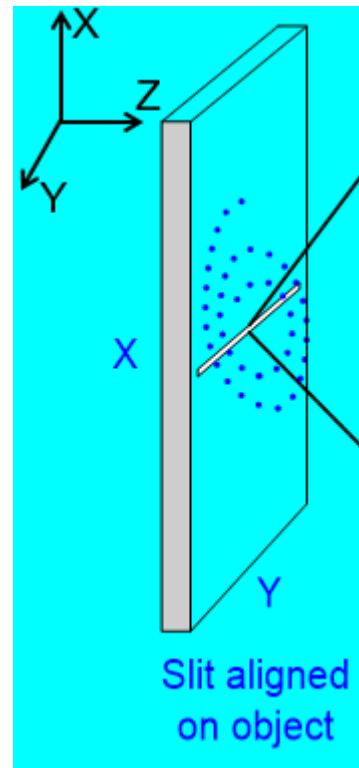
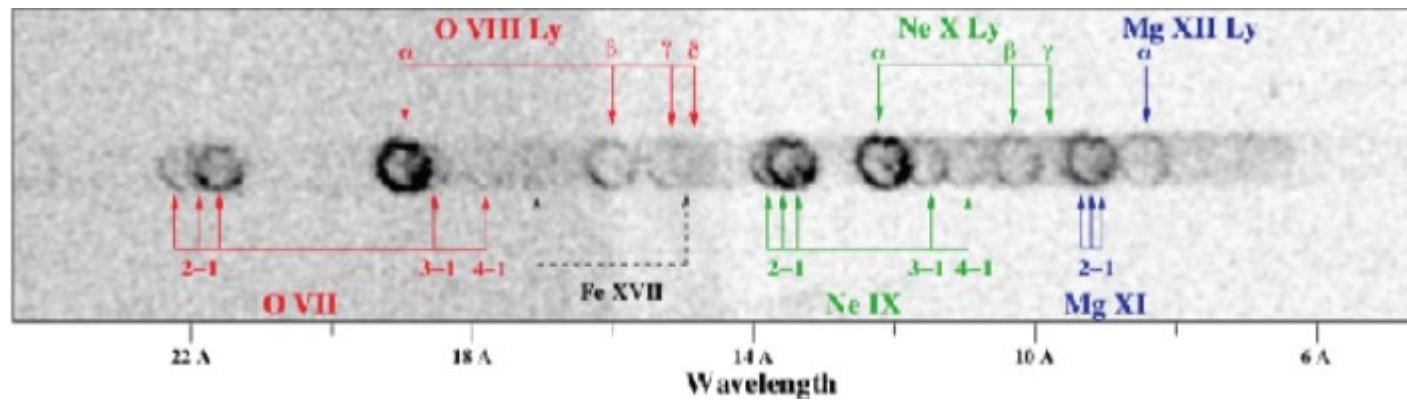
# Basic Spectrometer layout



entrance aperture/slit + collimator + dispersing component  
(if possible at a pupil)

# Entrance slit

- The entrance slit is the ‘field stop’ for the spectrograph
  - It restricts the angles at which light enters the spectrograph
  - Controls stray-light, avoids overlap of the spectra on the detector
- ‘Slitless’ spectroscopy works well for isolated sources and low background radiation (Hubble Space Telescope, Chandra)

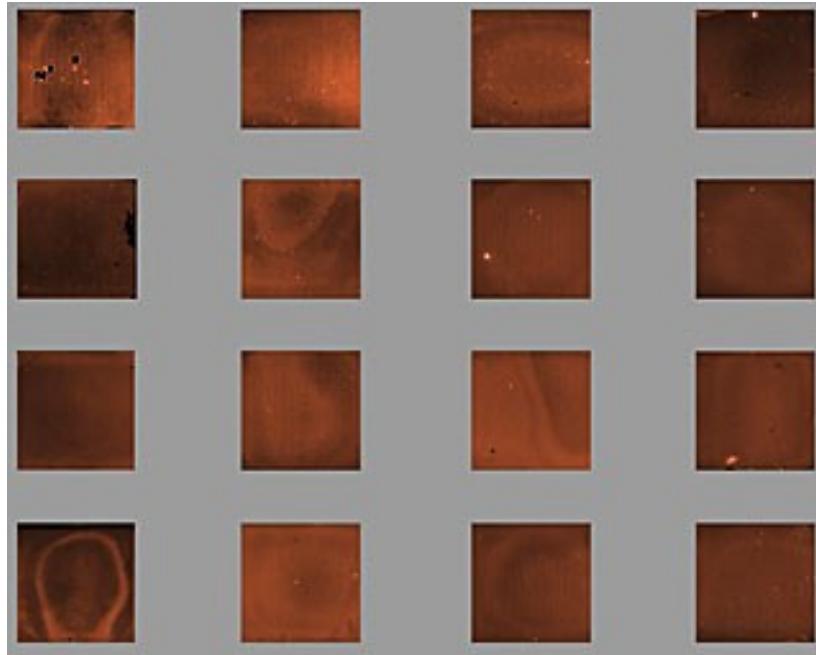


Flanagan et al.  
2007

- At ground based observatories, the background radiation is much brighter than the astronomical targets

# Optical/NIR background emission for ground-based observatories

■ VIRCAM raw data



■ Processed image

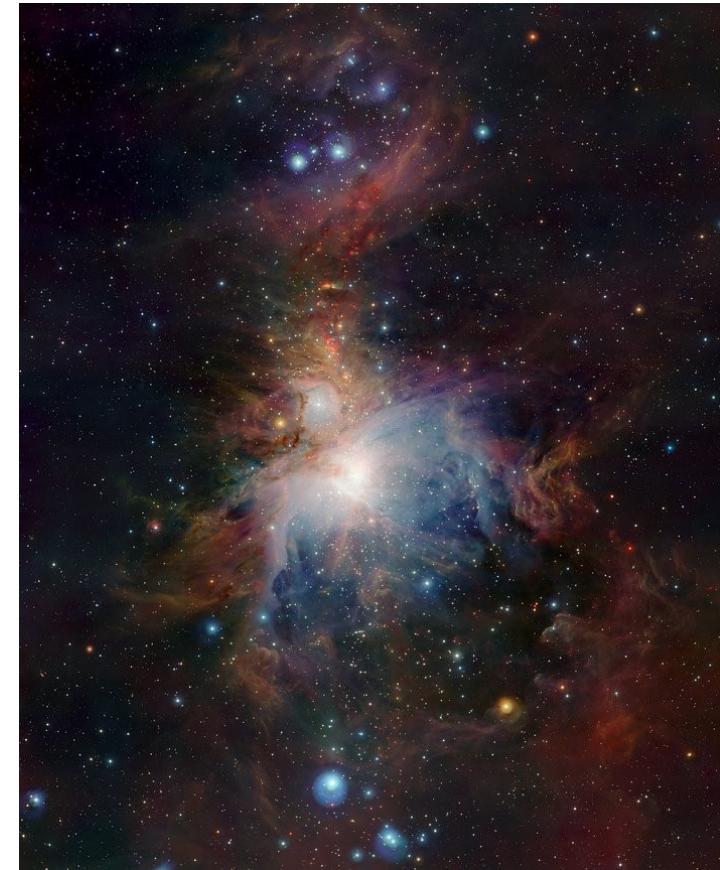
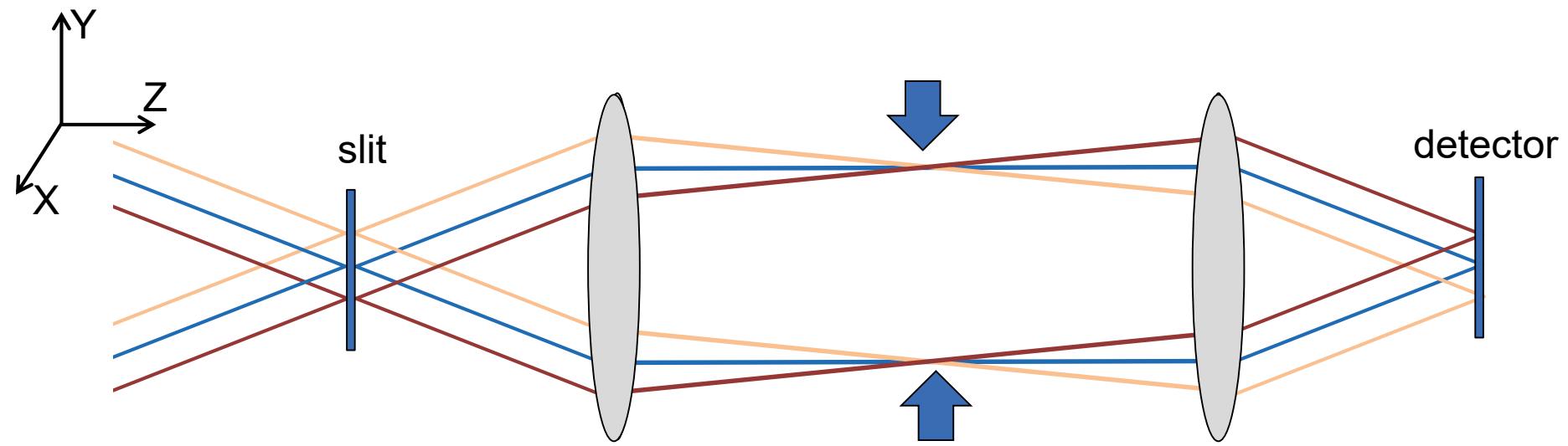


Image of the Orion (Nebula Messier 42)  
Z,J,Ks filters used  
1degree x 1.degree field  
10minutes per filter

Credit: ESO/J.Emerson/VISTA  
Processing: Cambridge Astronomical Survey Unit

# The pupil in any optical system

- The rays from any position along the spectrograph slit fill the pupil

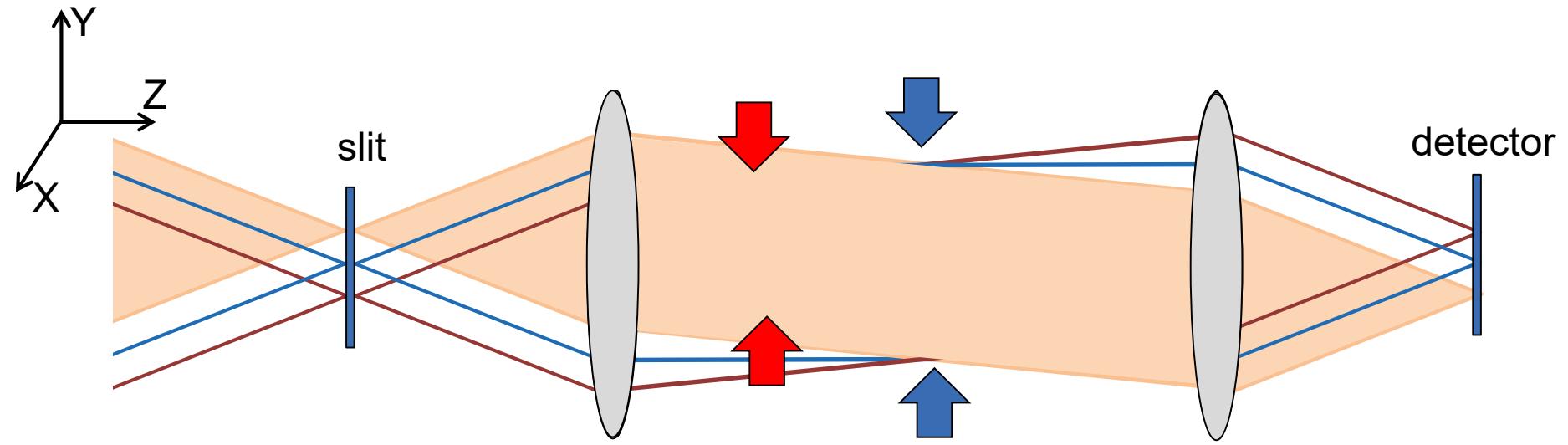


- The pupil is the narrowest point along the optical train

- A pupil or aperture stop placed here restricts the amount of light entering the optical system.
- Gratings or filters are placed as closed as possible to the filter. Any non-uniformities in their coatings or optical quality are experienced by each point along the slit equally>>more uniform results

# Vignetting

- The rays from any position along the spectrograph slit fill the pupil



- The pupil is the narrowest point along the optical train

- A pupil or aperture stop placed here restricts the amount of light entering the optical system.
- Gratings or filters are placed as close as possible to the filter. Any non-uniformities in their coatings or optical quality are experienced by each point along the slit equally>>more uniform results
- The beam is “vignetted” if something blocks the light

# Conservation of étendue

- Etendue (also called throughput) must be conserved to build an efficient instrument
- At any given aperture, étendue is the product of the area of the aperture ( $A$ ) and the solid angle of the rays accepted ( $\Omega$ )
- In particular, for an astronomical instrument:

$$A\Omega_{tel} = A\Omega_{det}$$

*Using a few assumptions and simple geometrical optics, we can calculate some important instrument parameters....but we will get to that later....*

# F-number

## ■ Focal ratio or F-number (F/#)

- The ratio of the diameter of a lens to its focal length
- “Fast” optics: small f-number (f/2, f/3)
- “Slow” optics: large f-number (VLT telescope: f-17)

Larger Aperture  
More Light

f/1.4



f/2



f2.8



f/4



f/5.6

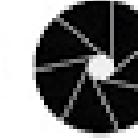


f/8



Smaller Aperture  
Less Light

f/11



f/16

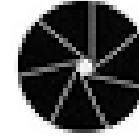
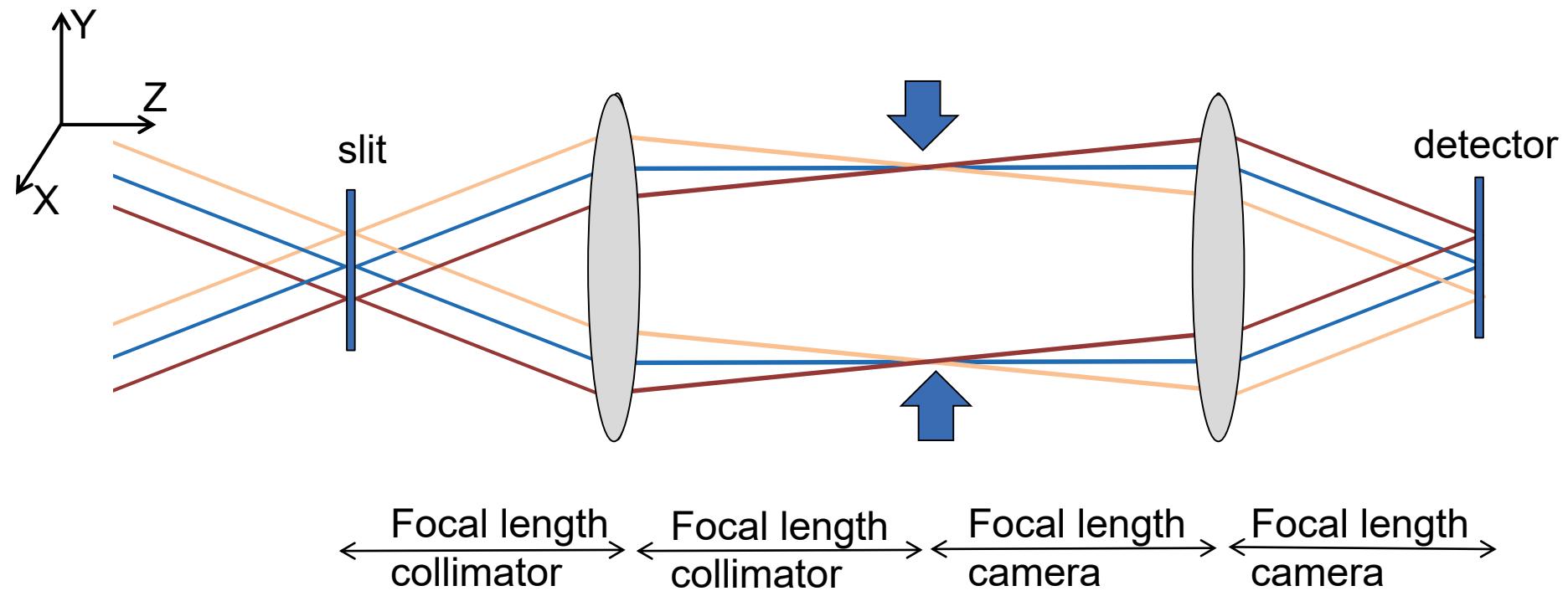


Image Credit: robin@lenscraft.co.uk

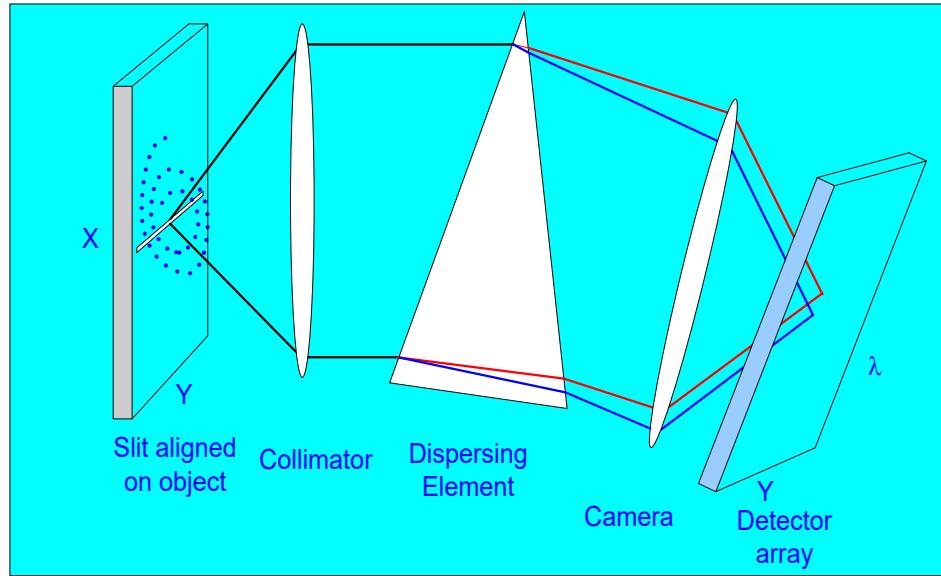
<https://lenscraft.co.uk/landscape-photography-tutorials/tutorial-what-is-an-f-stop/>

# The pupil in any optical system

- The rays from any position along the spectrograph slit fill the pupil



# Dispersing elements



# Prisms

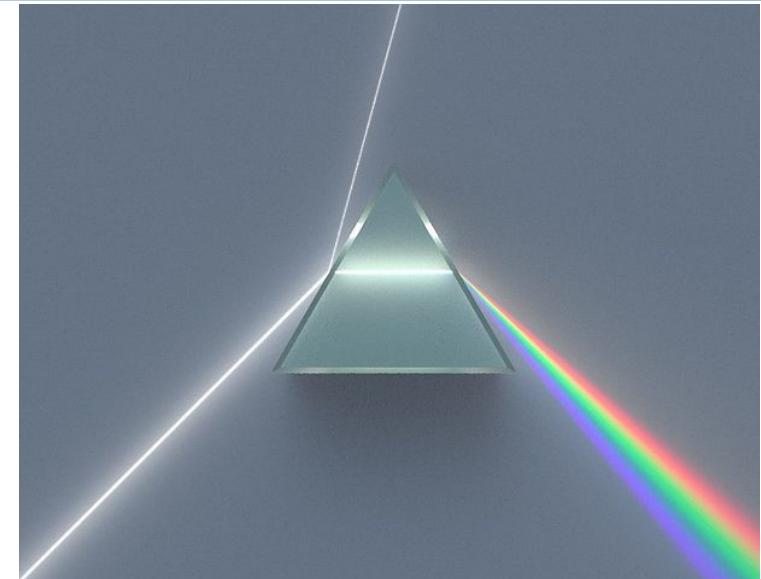
- The starting point for thinking about spectroscopy is often the example of white light dispersed by a prism

- Angular dispersion by a prism

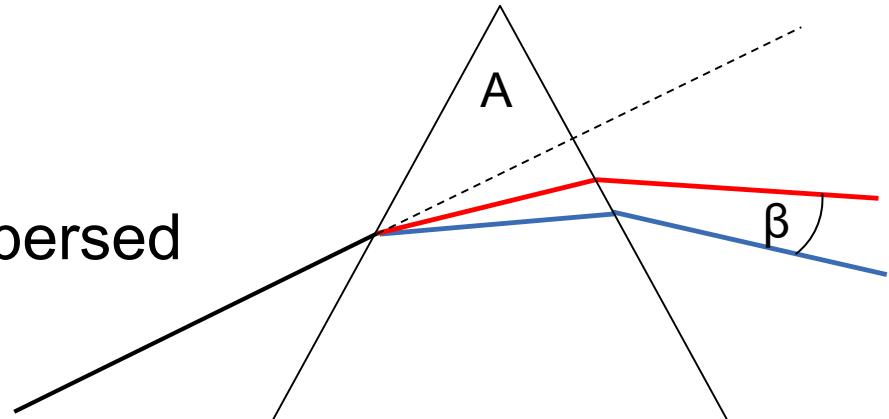
$$\frac{\partial \beta}{\partial \lambda} = A \frac{\partial n}{\partial \lambda}$$

- Hard to get high spectral resolution from a prism

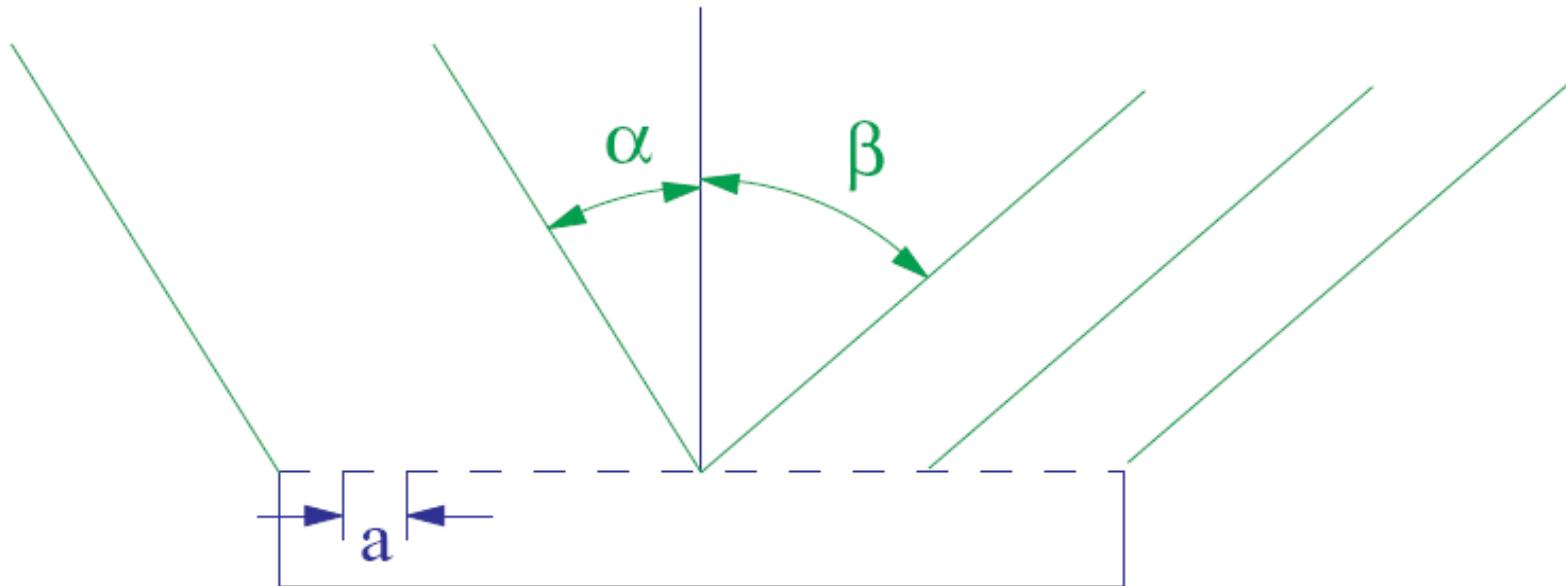
- Applications for cross-dispersed spectroscopy, see later



Dispersive\_Prism\_Illustration\_by\_Spigget.jpg, CC BY-SA 3.0, <https://commons.wikimedia.org/w/index.php?curid=12098156>



# Diffraction by a grating



Grating equation

$$m \frac{\lambda}{a} = \sin \alpha + \sin \beta$$

$m$ =grating order  
 $a$ =groove spacing

$\alpha$ =angle of incidence  
 $\beta$ =angle of diffraction

Angular dispersion

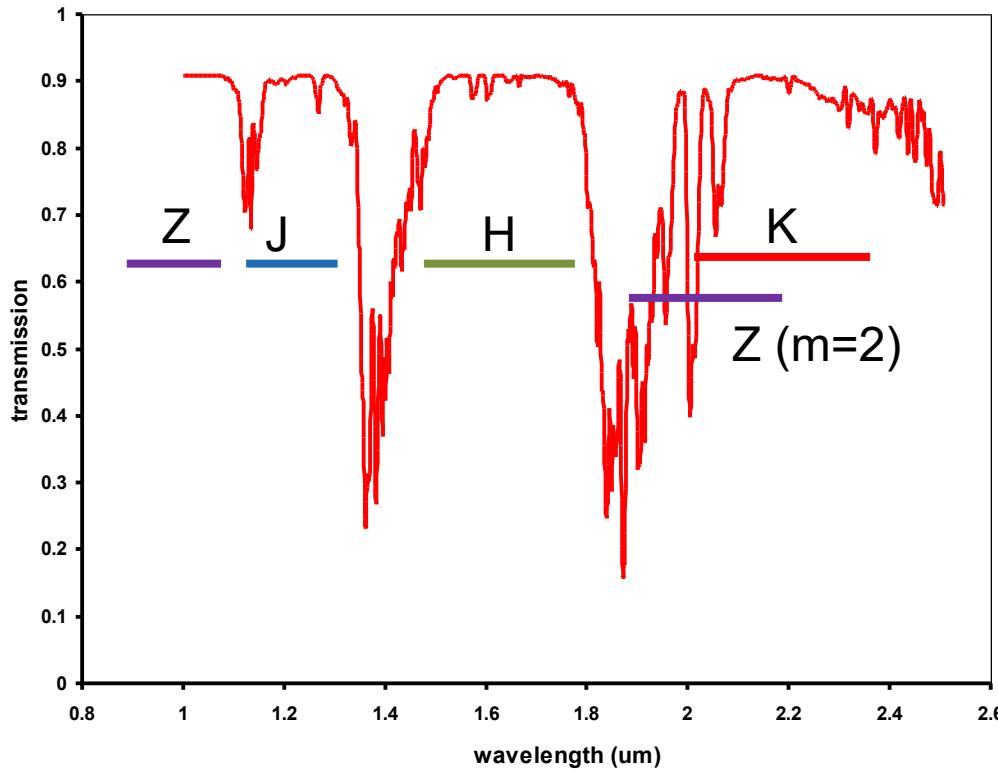
$$\frac{d\beta}{d\lambda} = \frac{m}{a \cos \beta}$$

typical  $a \sim$  few  $\mu\text{m}$   
Few hundred grooves/mm

# Diffraction by a grating

$$m \frac{\lambda}{a} = \sin \alpha + \sin \beta$$

For the same  $\alpha, \beta$  there is more than one combination of  $m, \lambda$   
 >> different wavelengths from different orders will overlap at the detectors



Near-infrared windows/bands:

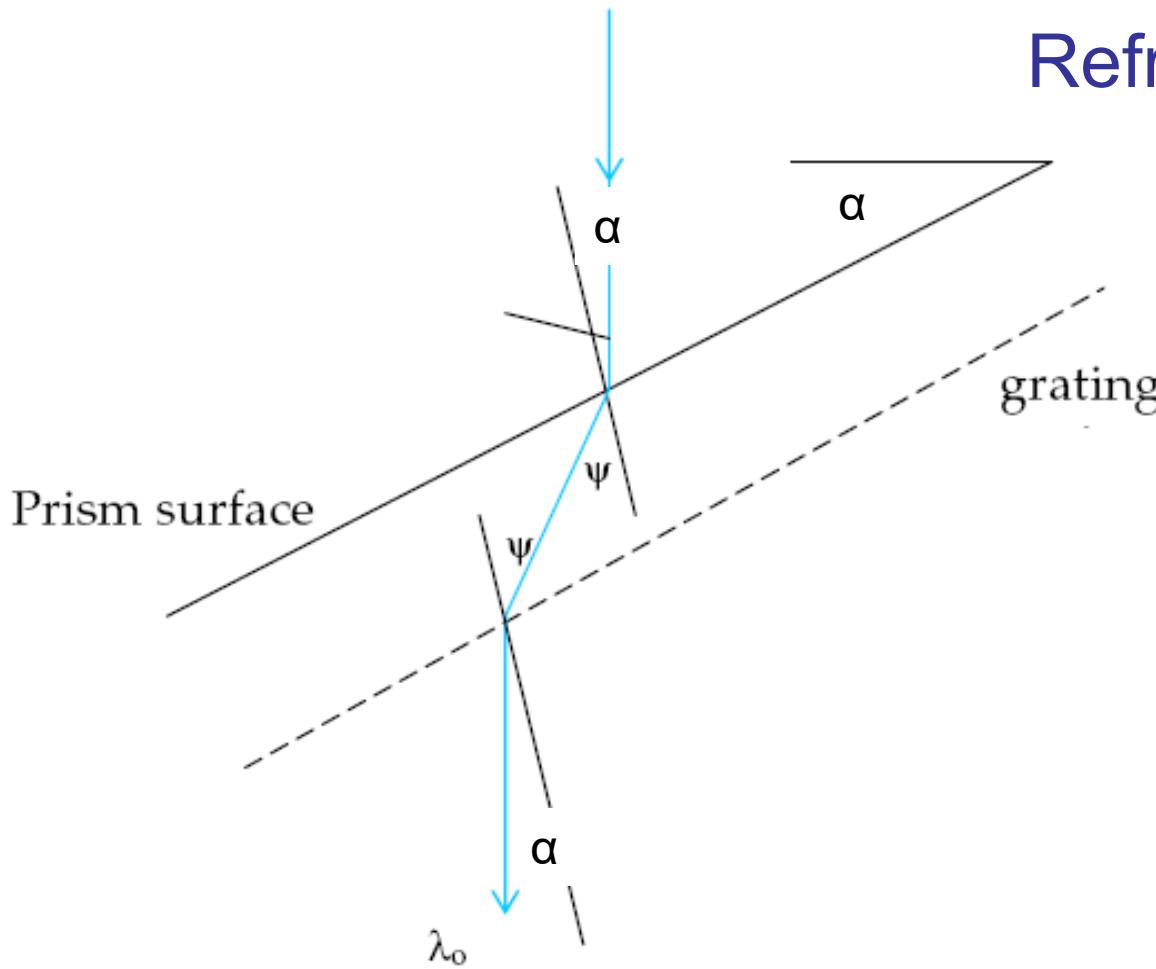
Z – 0.93  $\mu\text{m}$  – 1.015  $\mu\text{m}$  – 1.10  $\mu\text{m}$   
 J – 1.17  $\mu\text{m}$  - 1.25  $\mu\text{m}$  – 1.33  $\mu\text{m}$   
 H – 1.49  $\mu\text{m}$  - 1.64  $\mu\text{m}$  – 1.780  $\mu\text{m}$   
 K – 2.03  $\mu\text{m}$  - 2.2  $\mu\text{m}$  – 2.370  $\mu\text{m}$

Z-band in second order:

Z (m=2) – 1.86  $\mu\text{m}$  – 2.20  $\mu\text{m}$

Solution? Order sorting:  
 Use it: cross-dispersion  
 Loose it: blocking filters

# Grisms



Refraction in the prism:

$$\sin \Psi = n \sin \alpha$$

$n$  from the grating(\*):

$$= d(\sin \alpha - \sin \beta)$$

Prism' equation:

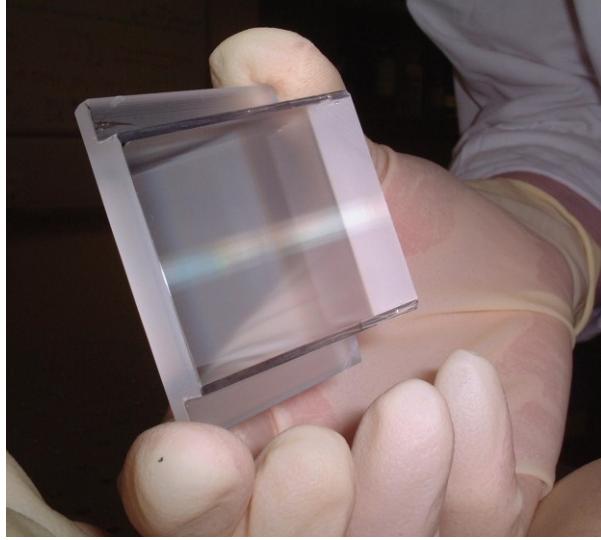
$$= (n - 1)d \sin \alpha$$

NB: dependence on the refractive index of the glass  
(\*) grating equation for a grating used in transmission

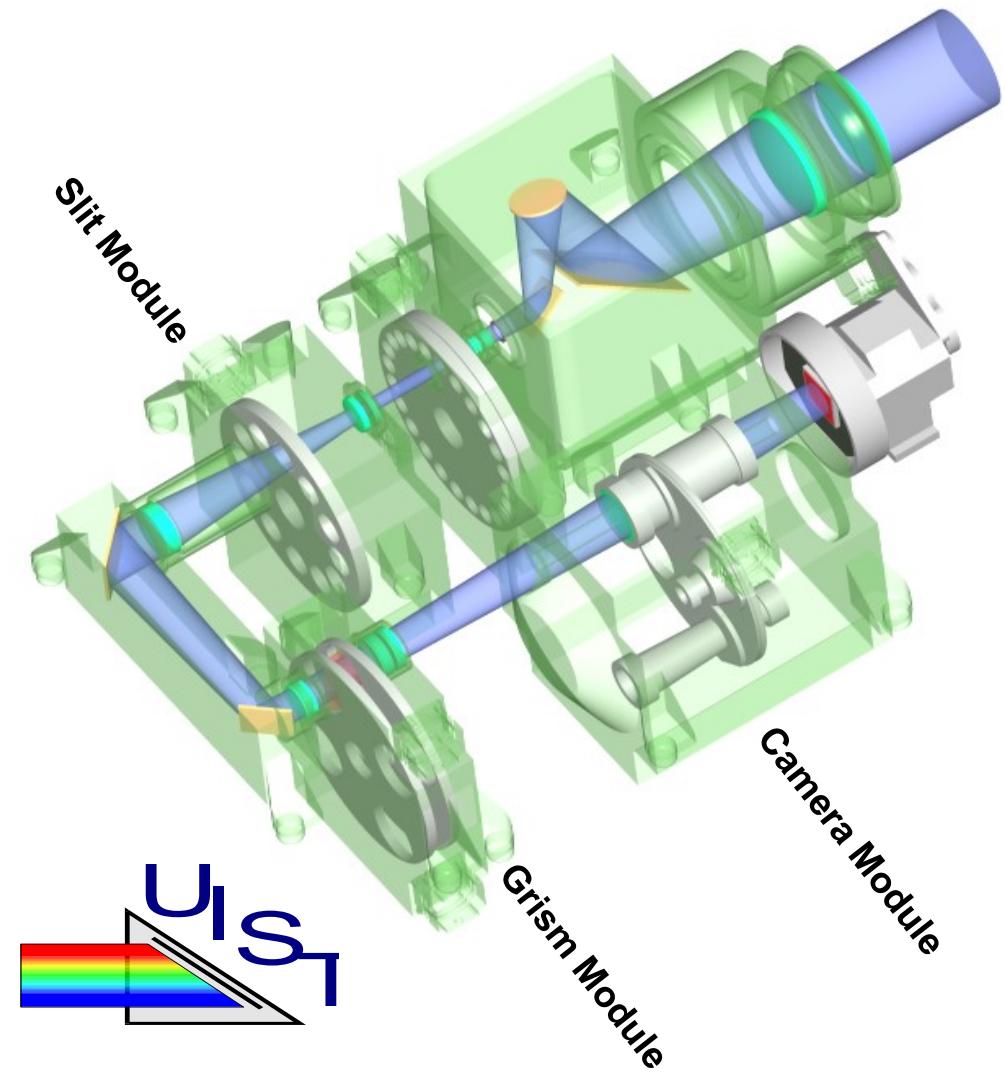
# Grisms

## ■ Advantages

- Compact
- simplified instrument layout
- *Potential* for higher R for a given beam size



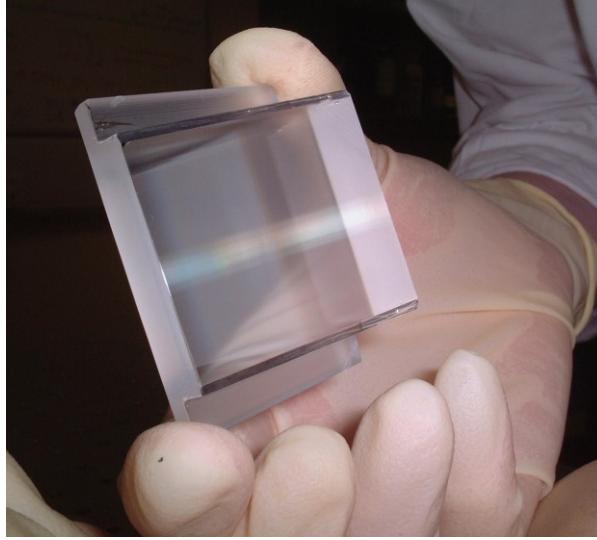
A fused silica grism (from Hyperfine), grating replicated onto the front surface



# Grisms

## ■ Advantages

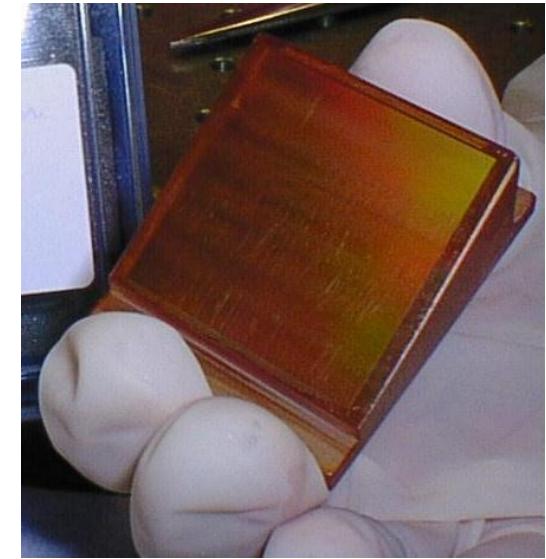
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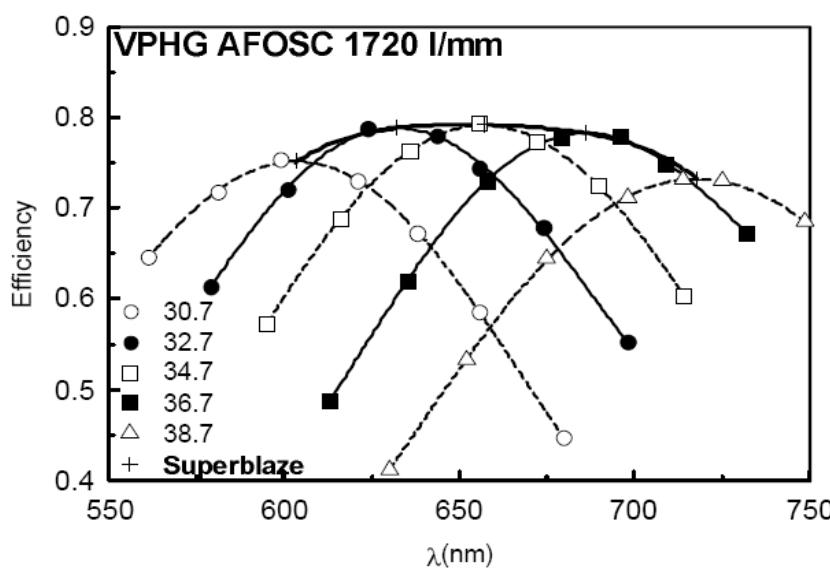
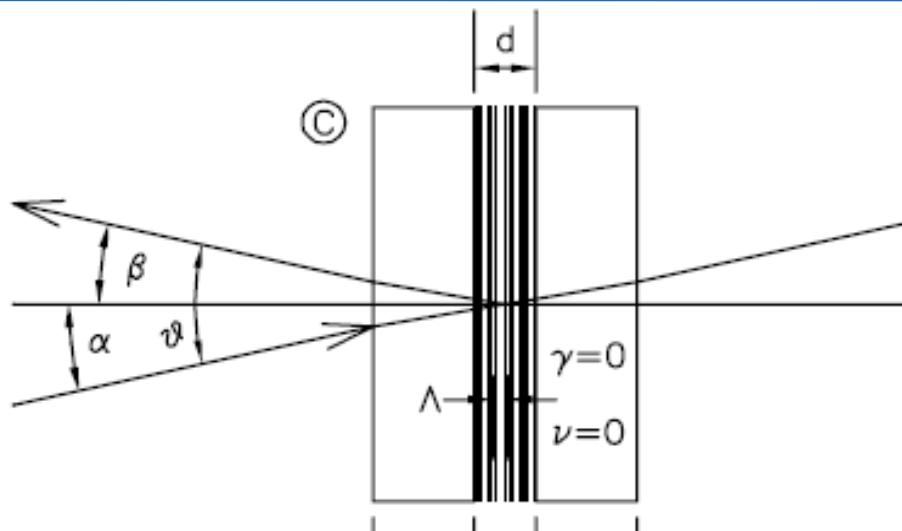
## ■ Disadvantages:

- difficult materials
  - Fused silica,  $n=1.45$  (NIR)
  - (HST)  $\text{CaF}_2$ ,  $n=1.45$  (NIR)
  - KRS5,  $n=2.4$  (NIR)
- Relatively low transmission (60%)

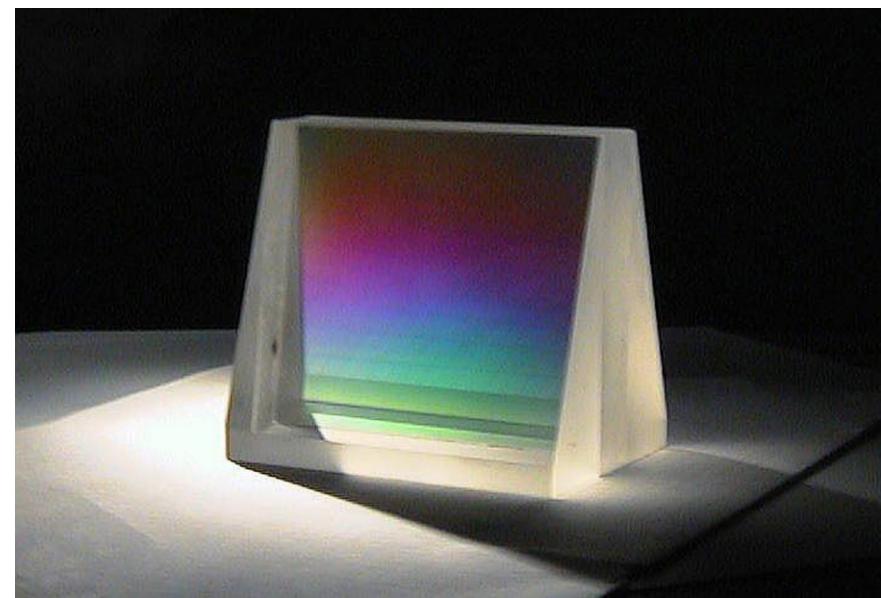


A grism made from KRS5, mechanically ruled grating (from Zeiss, Jena)

# Volume Phase Holographic gratings



$$m\nu\lambda = \sin(\alpha) - \sin(\beta)$$



Use for astronomy proposed about 20 years ago (Barden et al 1998); now standard for modern opt-IR spectrographs.

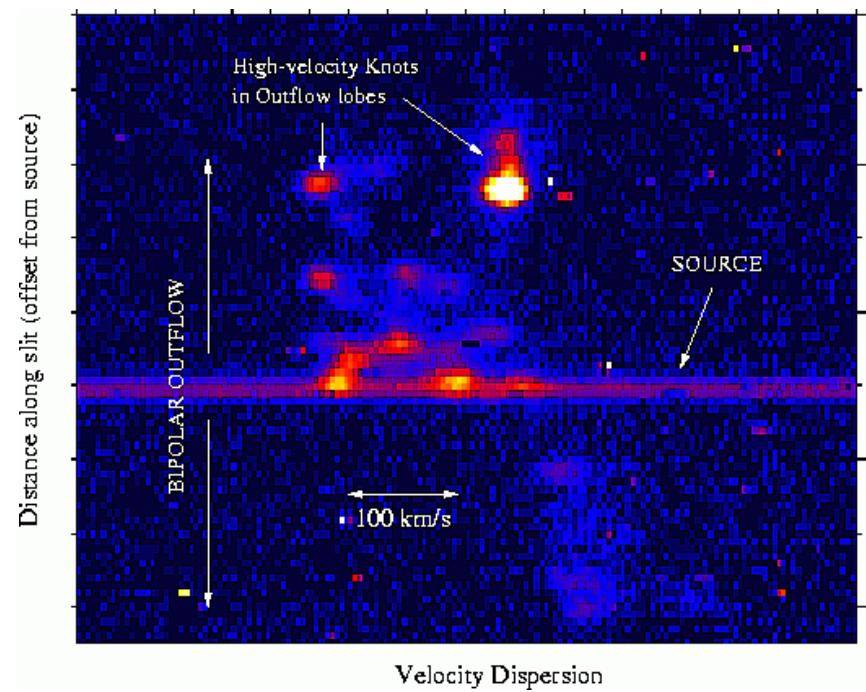
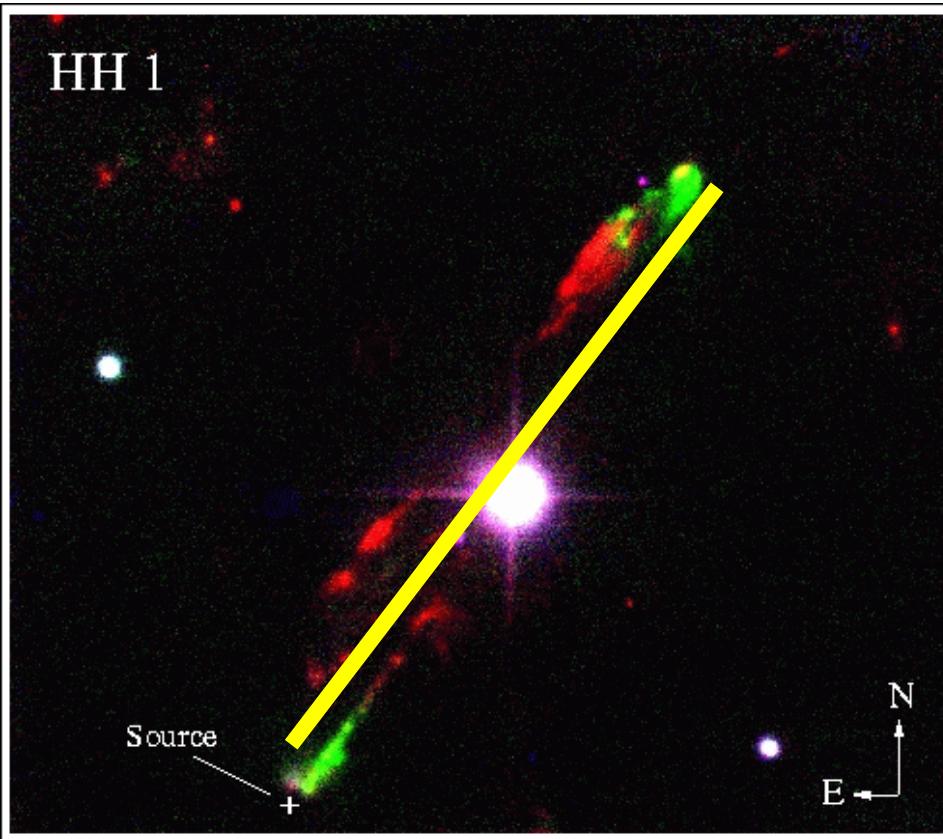
Very high efficiency.



# Variations on the classical spectrograph

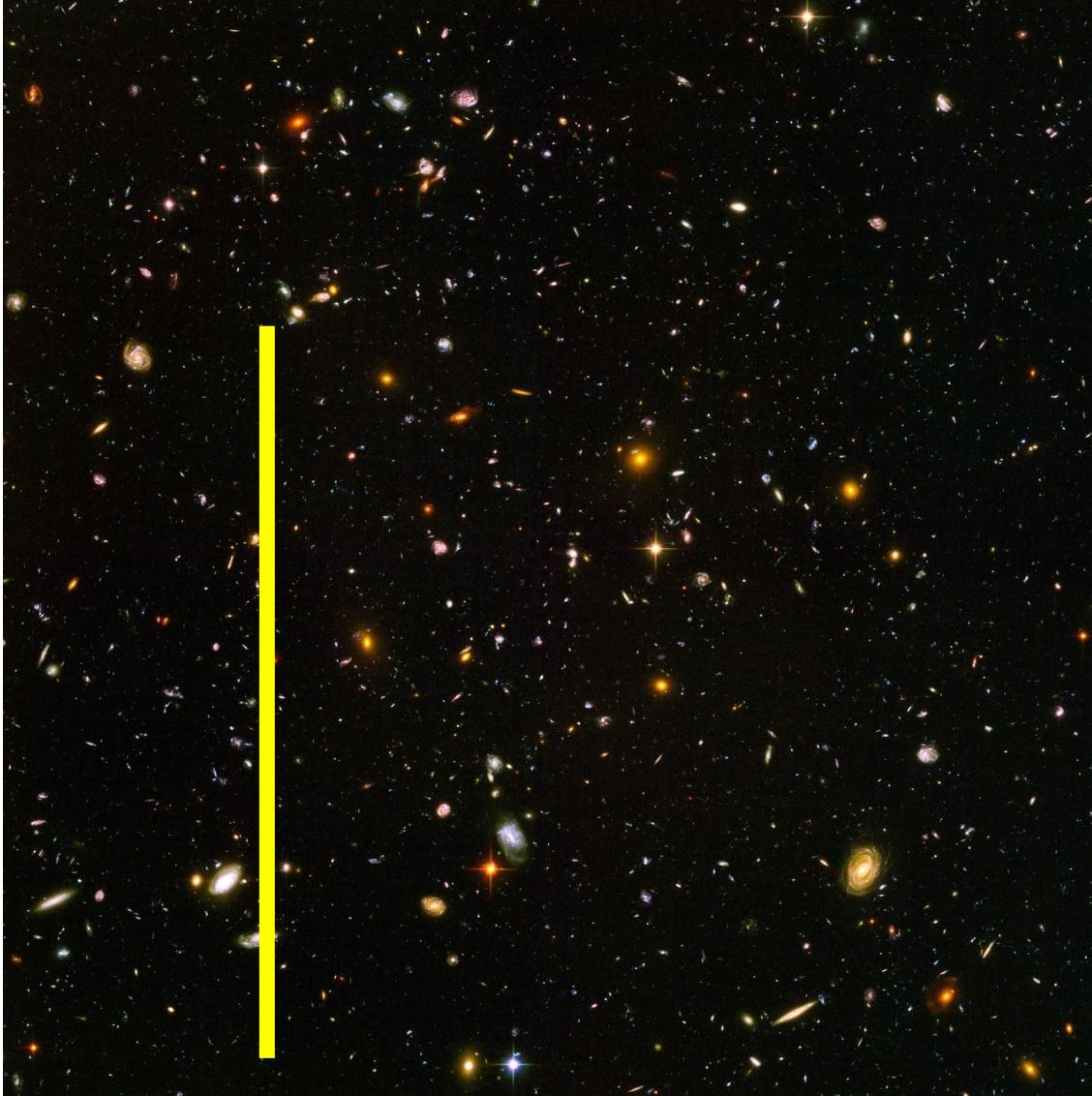
# The long slit: strengths/weaknesses

- With typical plate scales of 0.2arcsec/pixel and 2048 x 2048 pixel detectors, modern spectrometers have slit lengths of 120arcmin
- Some very beautiful astronomical objects will fill this slit, but many not!



Data courtesy UK Infrared Telescope, Tom Ray and Chris Davis

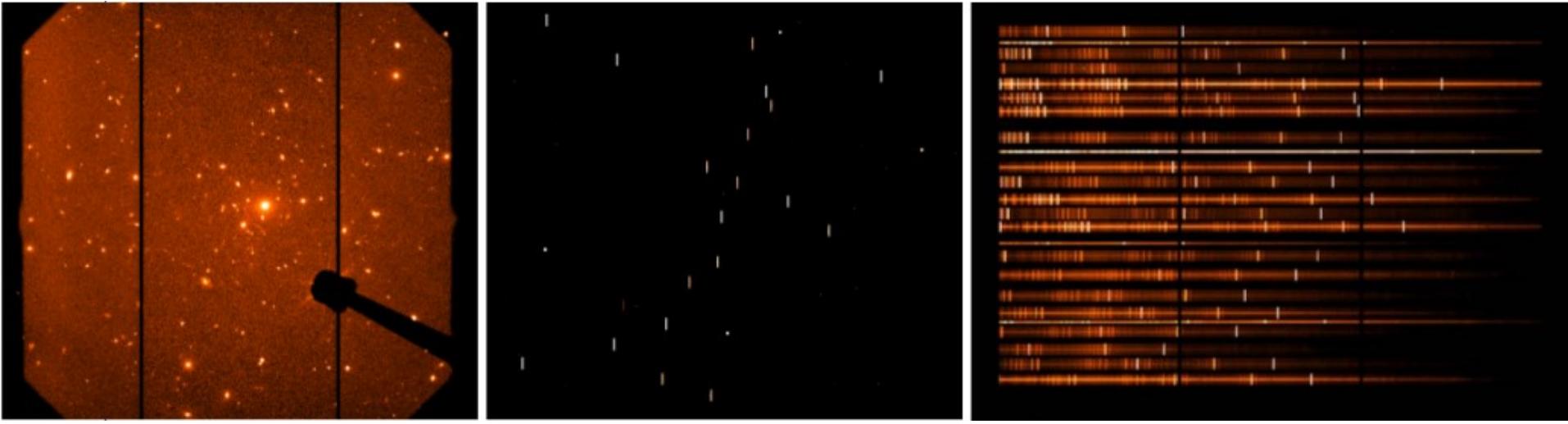
# A more typical set of objects



- This is the Hubble Ultra-Deep field, filled with almost 10,000 galaxies at all ages (all redshifts)
- The image is 3.1arcmin x 3.1 arcmin
- Multi-object spectroscopy offers a more efficient solution for observing a field like this
- (but not 10s of thousands of objects....100s of objects)

# Multi-object spectroscopy (1)

## ■ Multi-object spectroscopy with slit masks



Pre-imaging may be required to map or identify objects

A slit mask is made with slitlets for each target and inserted into the instrument focal plane

Spectra of the multiple objects are obtained with excellent usage of the array and some restrictions on the object selection

*Images are from the Gemini multi-object spectrograph*

David Dunlap School, 2019, Spectroscopy

# Multi-object spectroscopy (2)

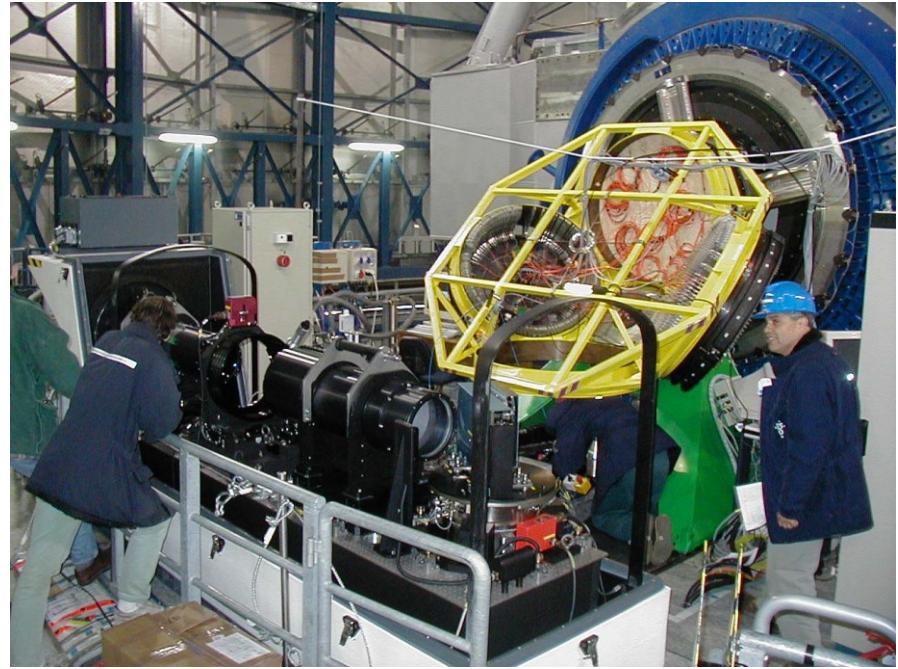
## ■ VIMOS at ESO's 8-m Very Large Telescope



Slit (mask) spectrographs are sensitive to faint objects. VIMOS was designed with galaxy surveys in mind. e.g. VIPERS ([vipers.inaf.it](http://vipers.inaf.it))

# Multi-object spectroscopy (3)

## ■ Multi-object spectroscopy with fibres: FLAMES



The OzPoz fibre coupled to the GIRAFFE or UVES spectrographs are ideally suited to surveys of stars e.g. the GAIA-ESO survey to  $10^5$  sources ([www.gaia-eso.eu](http://www.gaia-eso.eu))

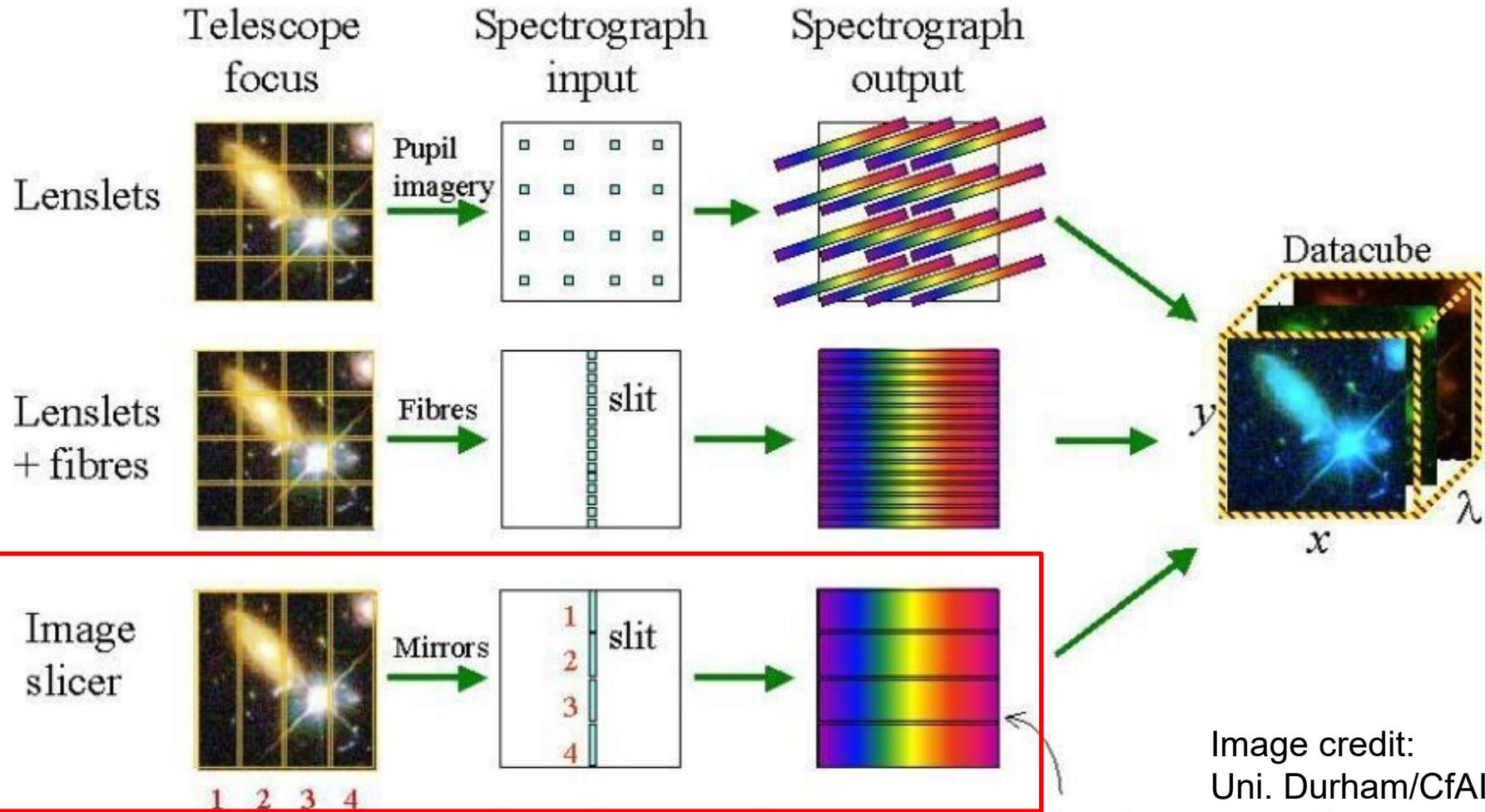
# Back to our galaxy field



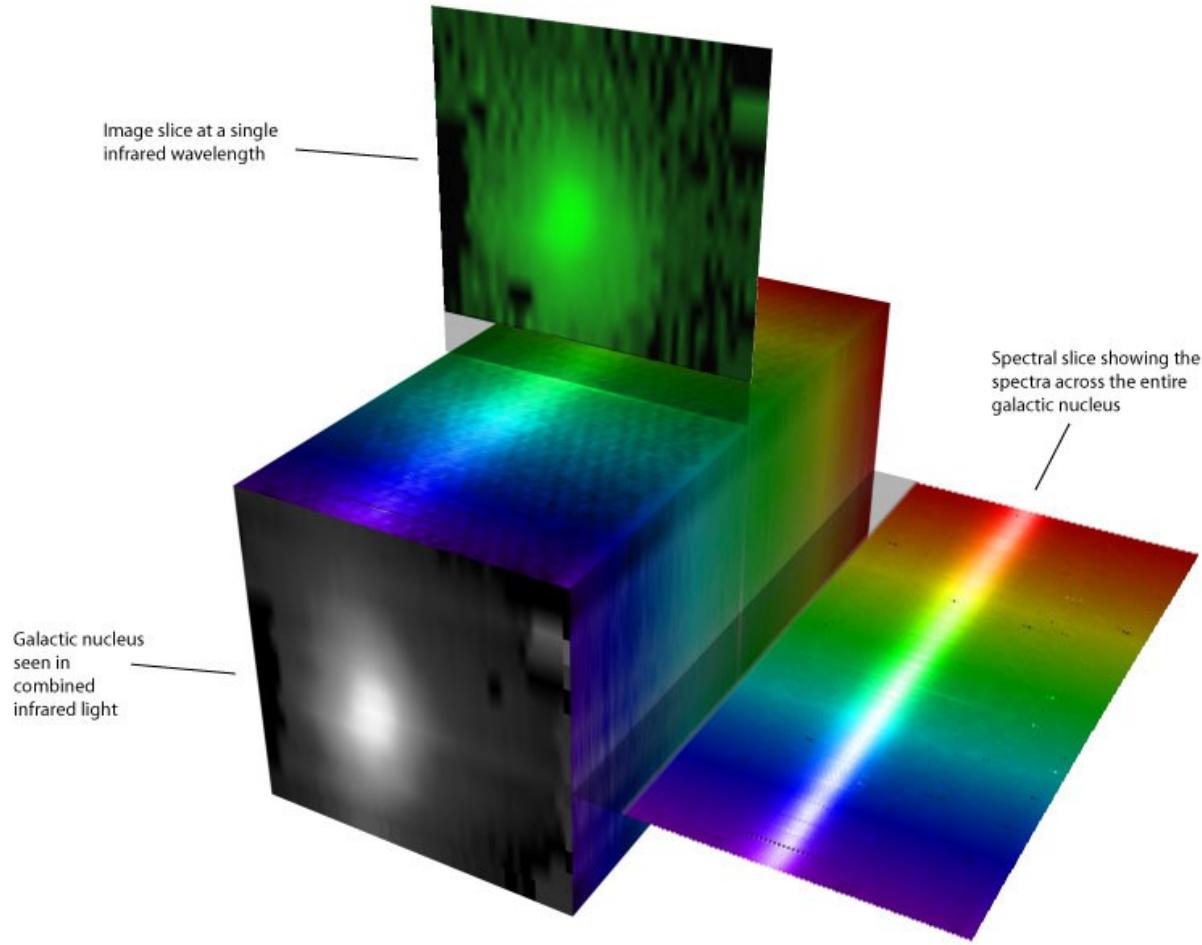
- MOS spectrographs can survey wide areas taking integrated light spectra of hundreds of objects
- But astronomers are also interested in the details of the assembly and evolution of these galaxies
- For this, spatially resolved observations of these small objects is desired
- >> integral field or 3D spectroscopy is ideal for this....

# Integral Field Spectroscopy

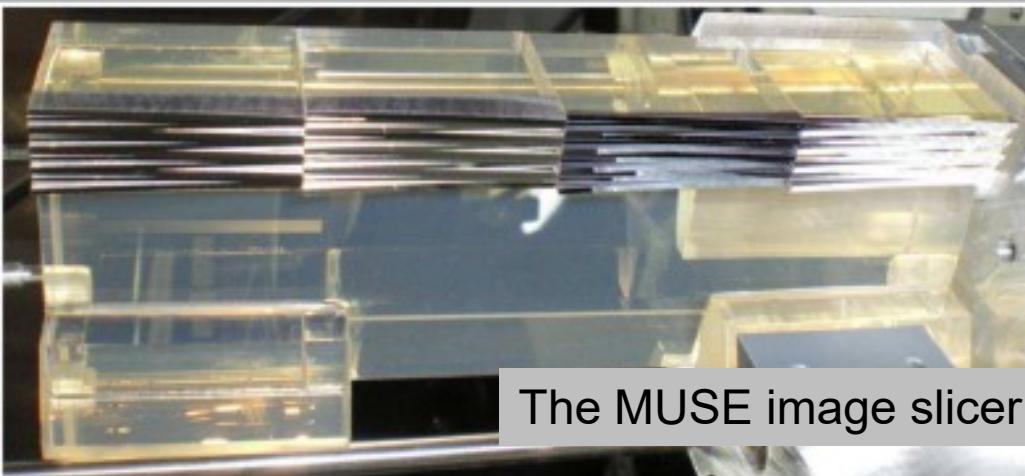
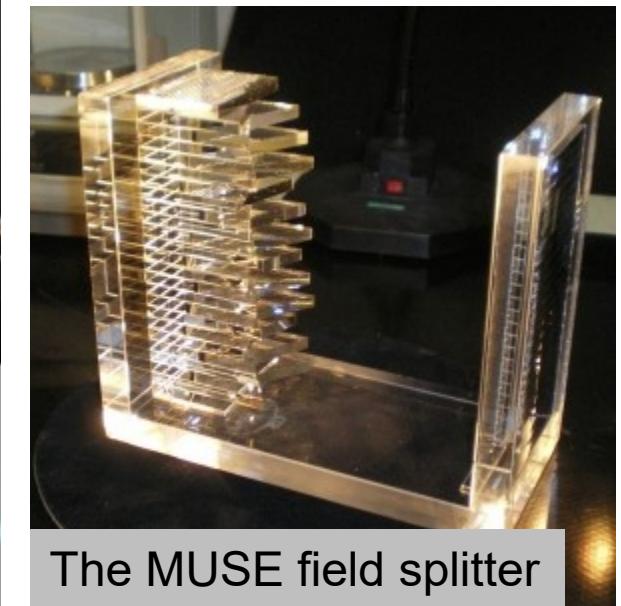
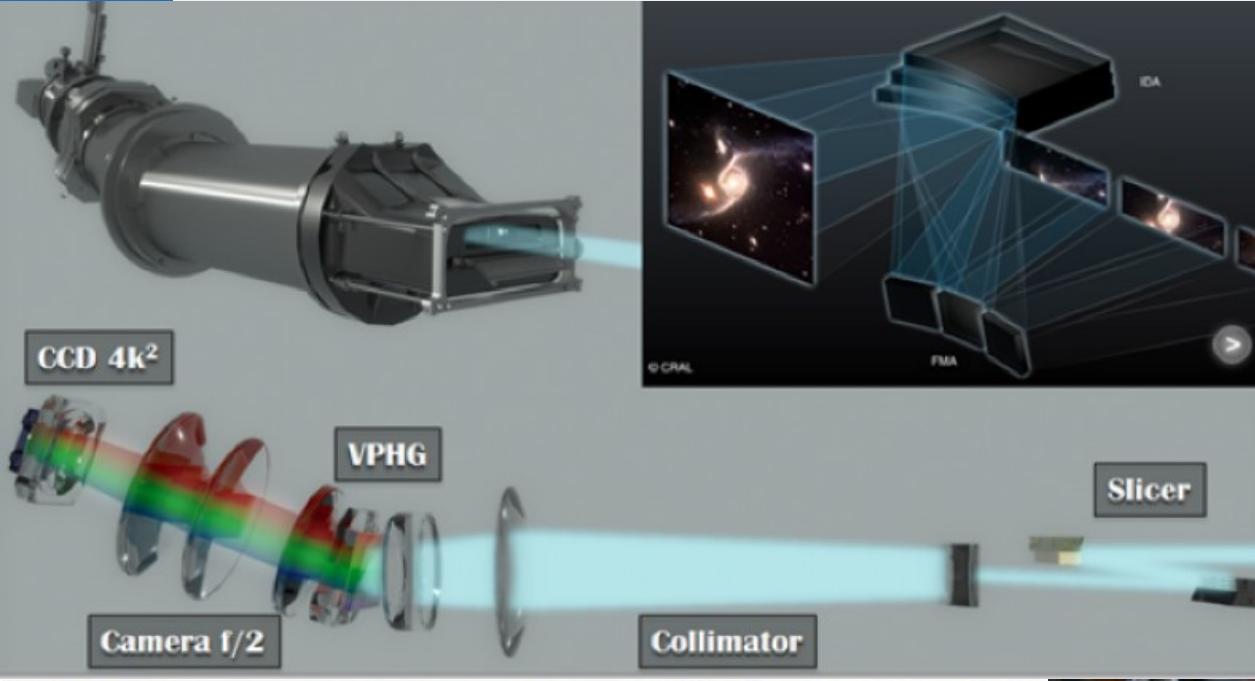
## ■ Integral field or 3D spectroscopy



# Data cube from the UIST IFU



# MUSE at the VLT



# MUSE at the VLT



The Pillars of Creation as observed by HST  
[http://hubblesite.org/image/3471/news\\_release/2015-01](http://hubblesite.org/image/3471/news_release/2015-01)



The Pillars of Creation as observed by MUSE. Three colour images extracted from the datacube.

Mcleod et al (2015)

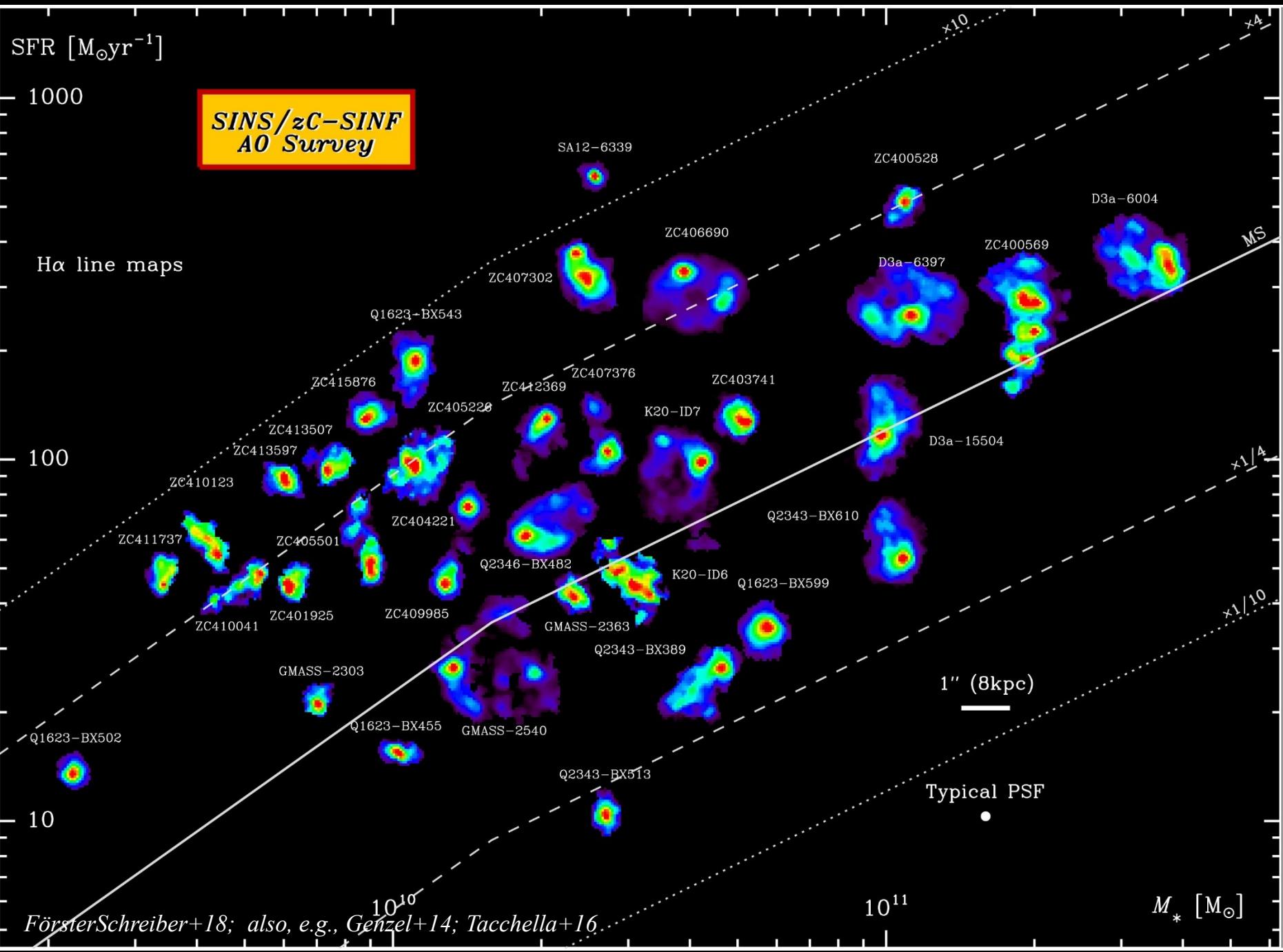
# The MUSE ultra-deep field

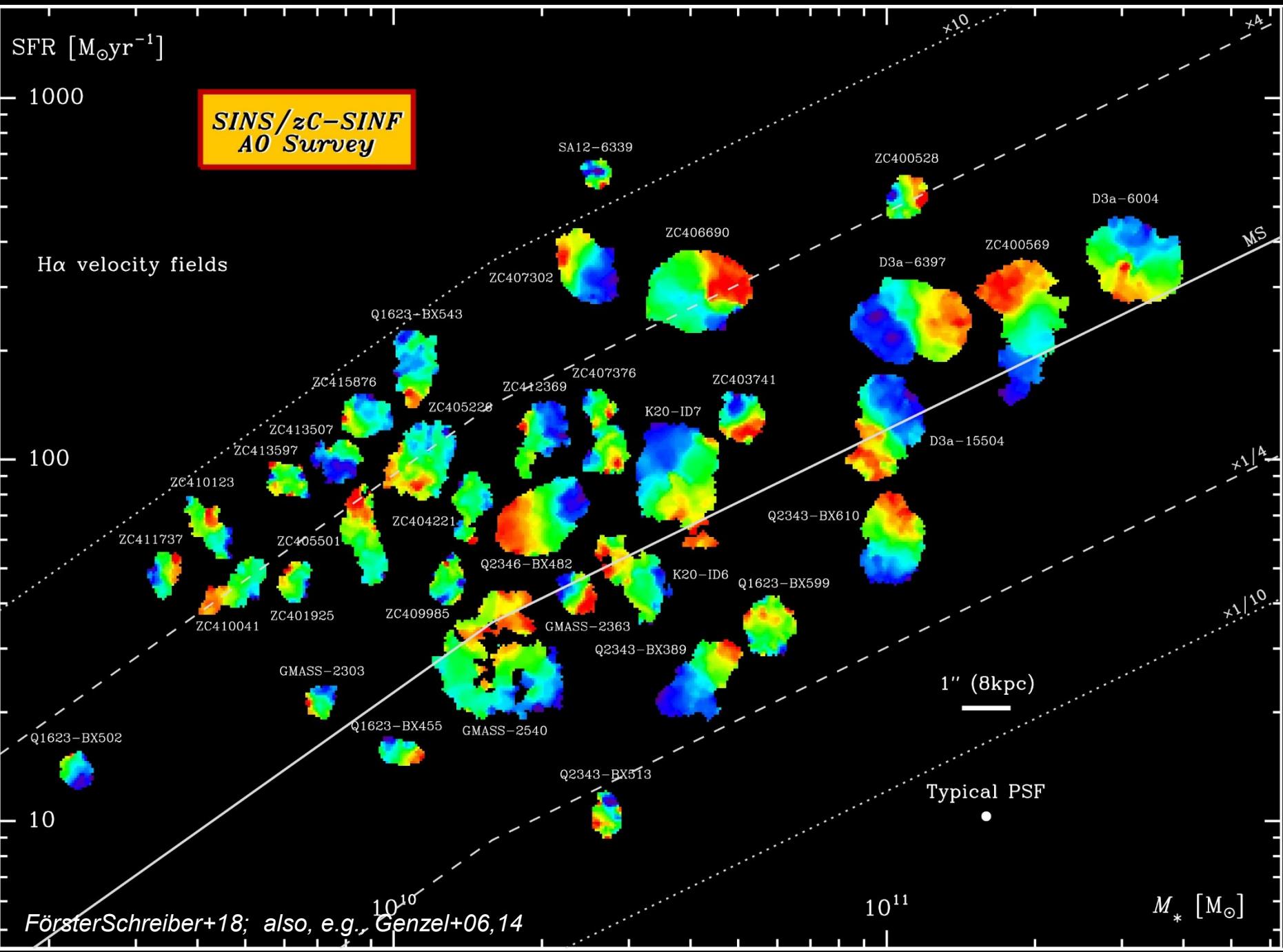


There is a spectrum for every pixel in this image!

10h exposure; 9 x the MUSE 1arcmin field of view.

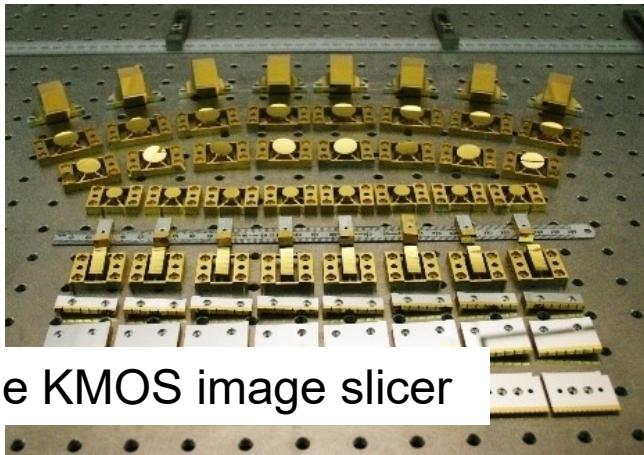
1600 galaxies of which 72 not detected by HST (Lyman-alpha emitters).





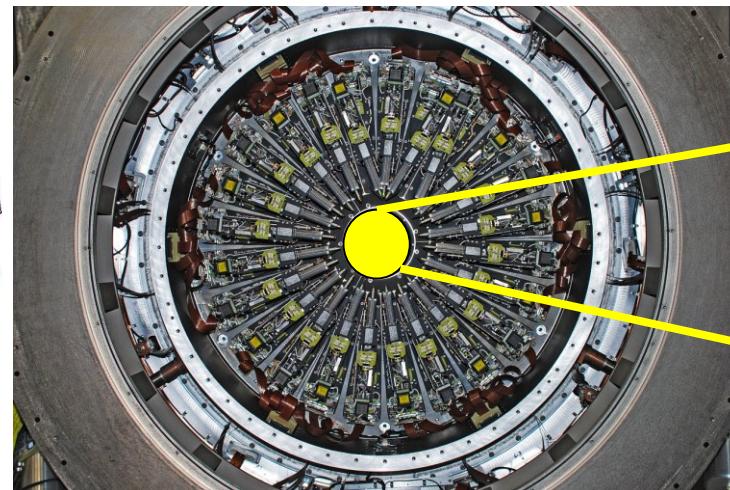
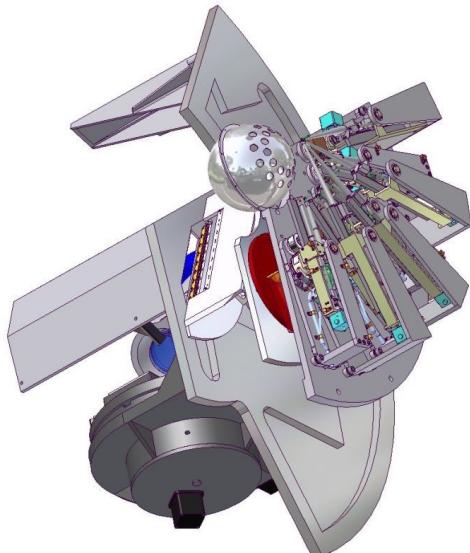
# Multi-IFU spectrographs

## ■ Multi-IFU spectroscopy in the NIR with KMOS

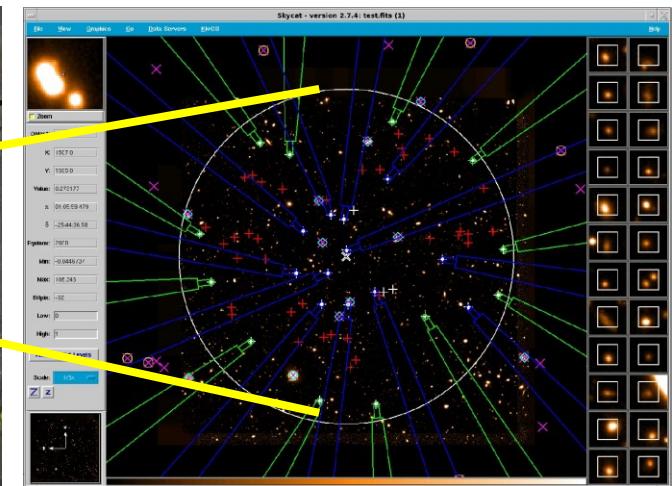


The KMOS image slicer

24 pick-off arms select targets over the VLT 7arcmin field of view then direct the beam to 24 IFUs and 3 spectrographs.



electroscopy



SFR

$\text{SFR}_{\text{MS}, \log(M_*) = 10.5}$

KMOS<sup>3D</sup>

10

1

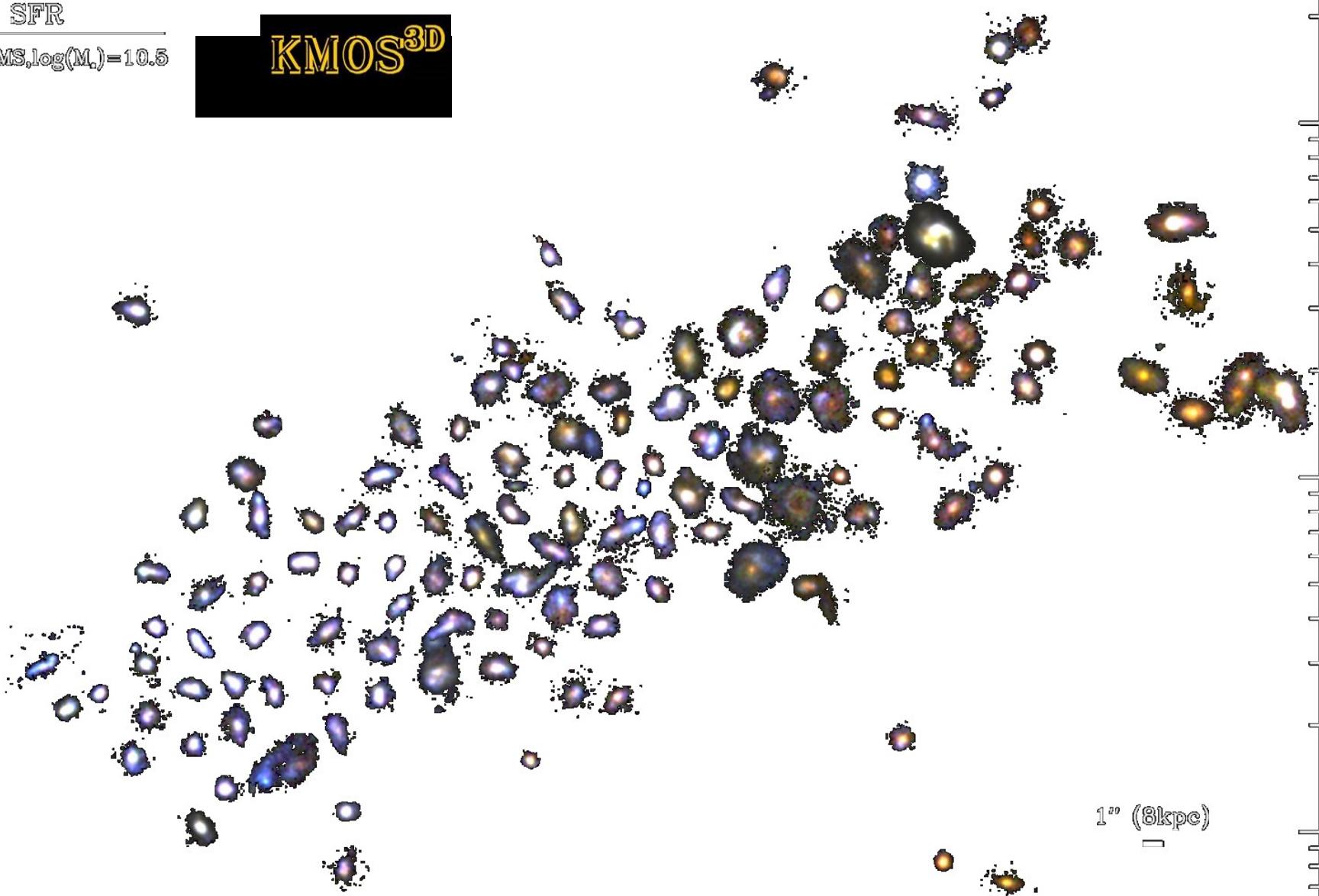
0.1

$10^{10}$

$10^{11}$

$M_* [M_\odot]$

$1''$  (8kpc)



SFR

$\text{SFR}_{\text{MS}, \log(M_*) = 10.5}$

KMOS<sup>3D</sup>

10

1

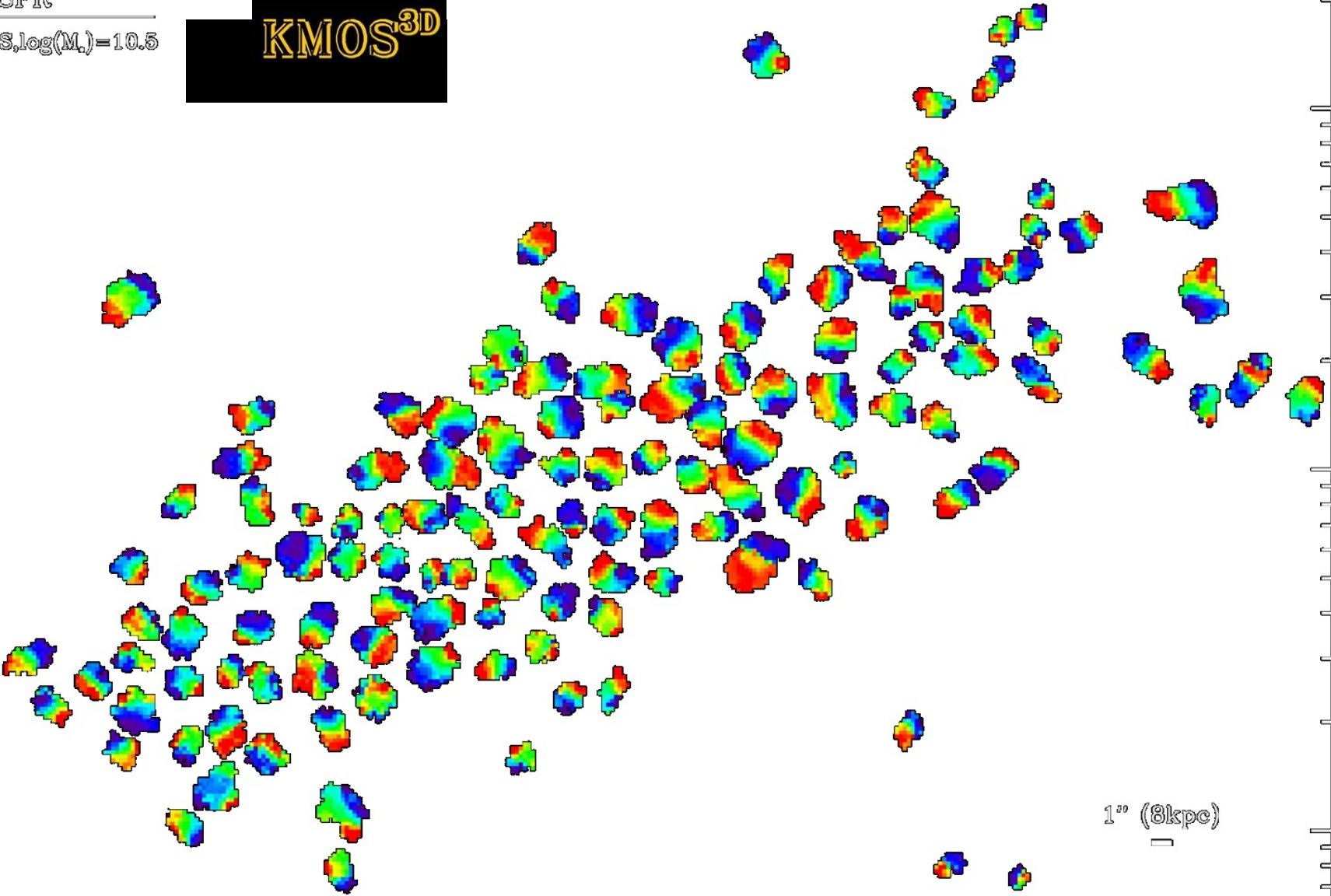
0.1

$10^{10}$

$10^{11}$

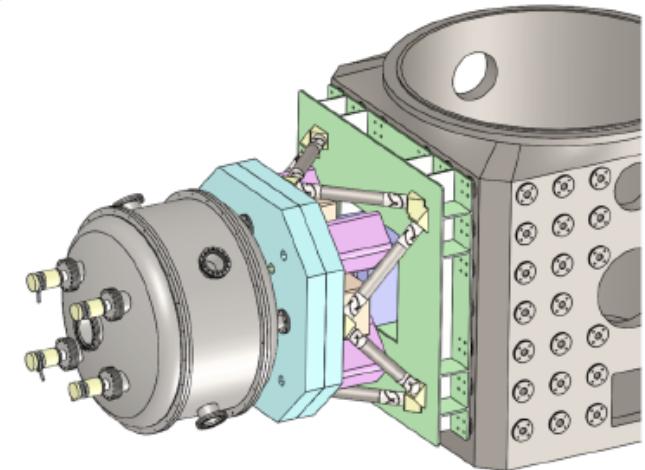
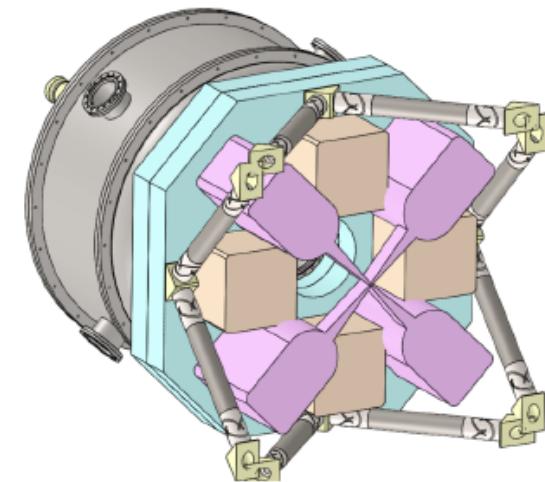
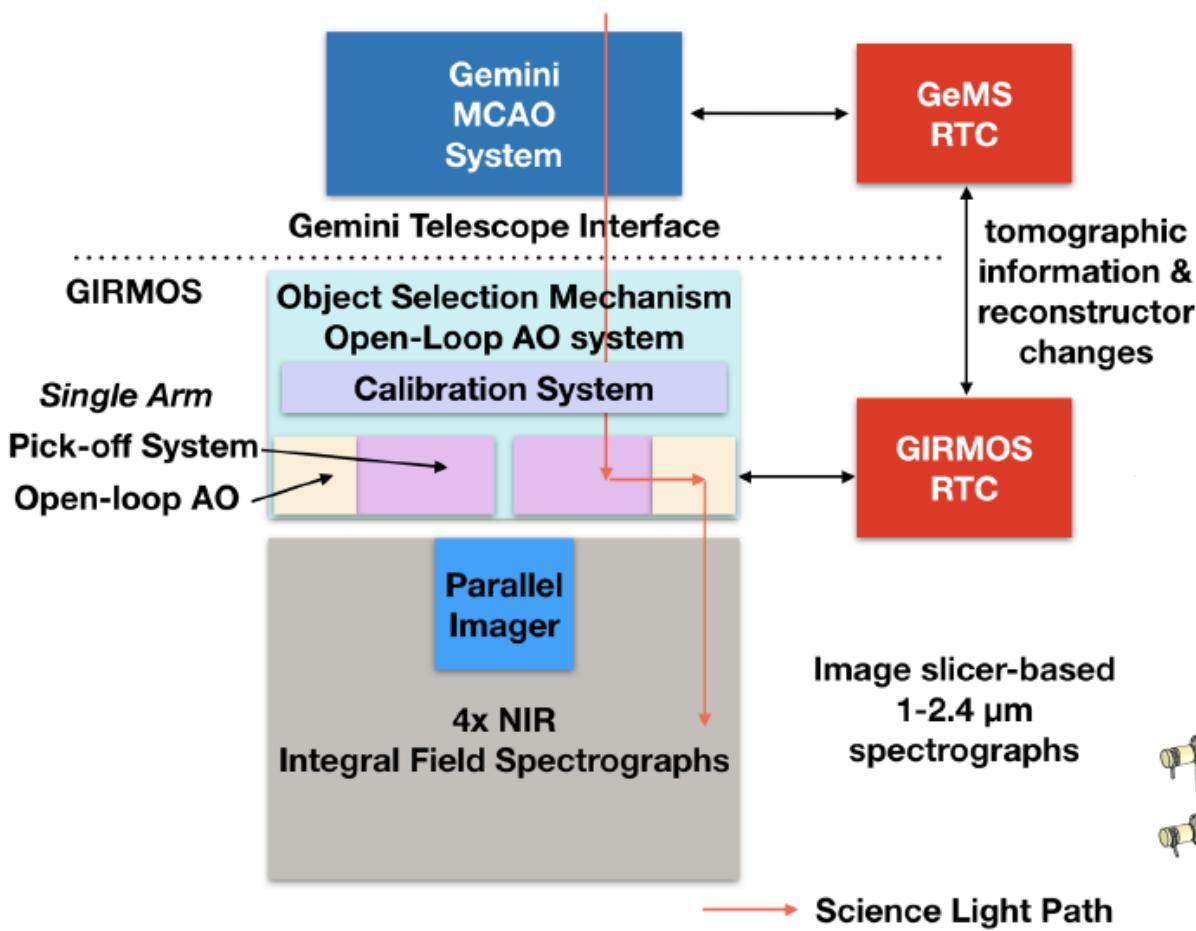
$M_* [M_\odot]$

$1''$  (8kpc)



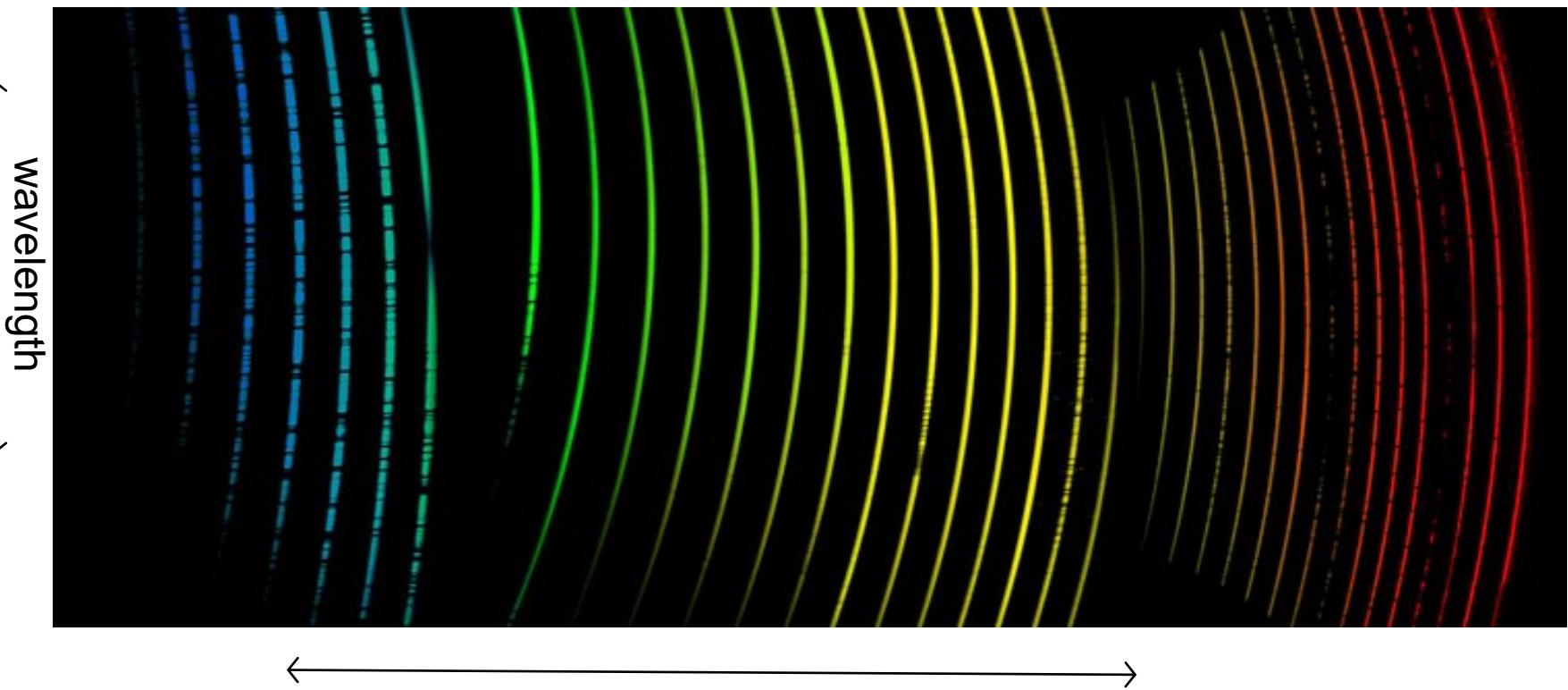
# GIRMOS

■ Sivanandam et al. 2018 (arXiv:1807.03797)



# Cross-dispersed spectrographs

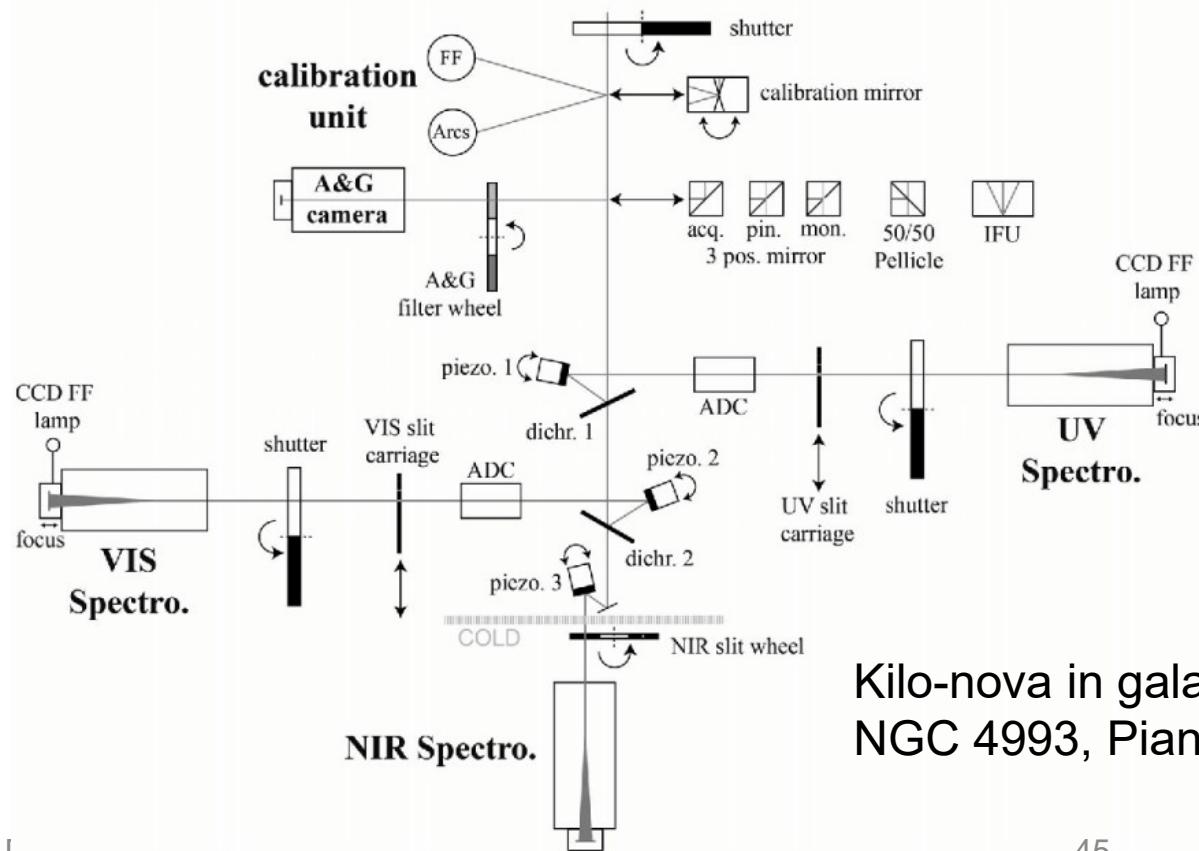
- For the widest wavelength range and small objects: cross-dispersed spectroscopy



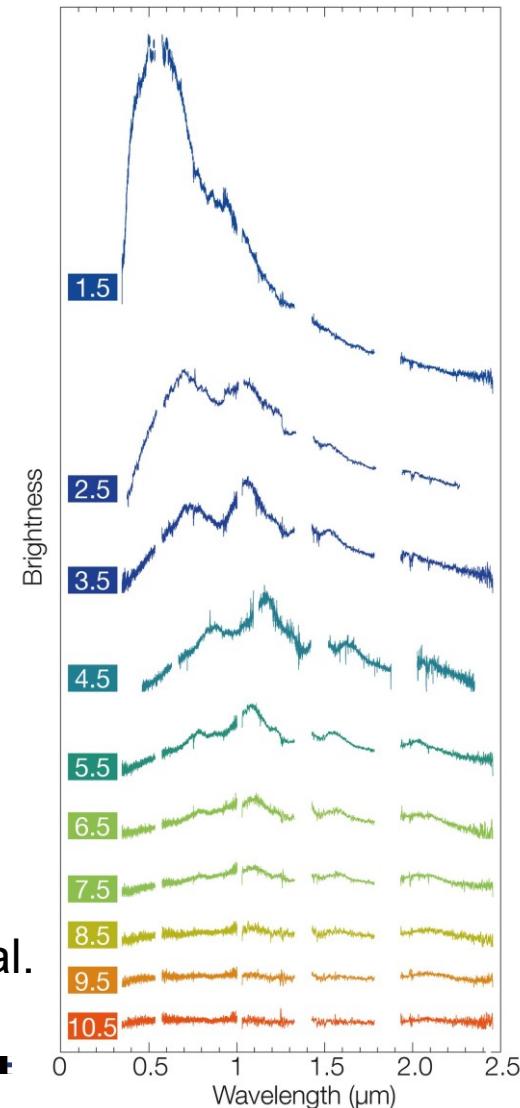
# Cross-dispersed spectrographs

## ■ Cross-dispersed spectroscopy with X-shooter

X-shooter was designed specifically for quick follow up of transient sources, in particular gamma ray bursts and supernovae.

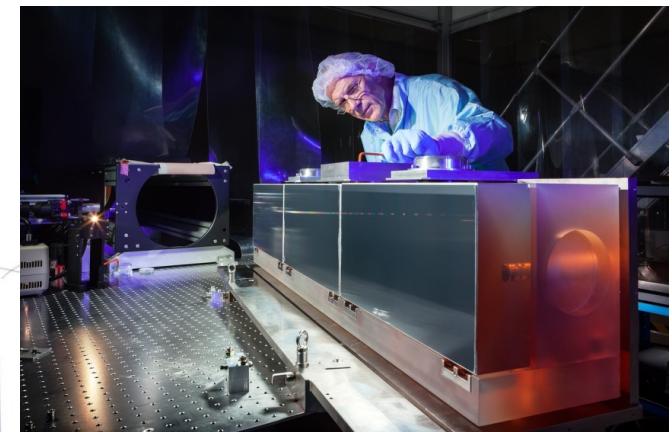
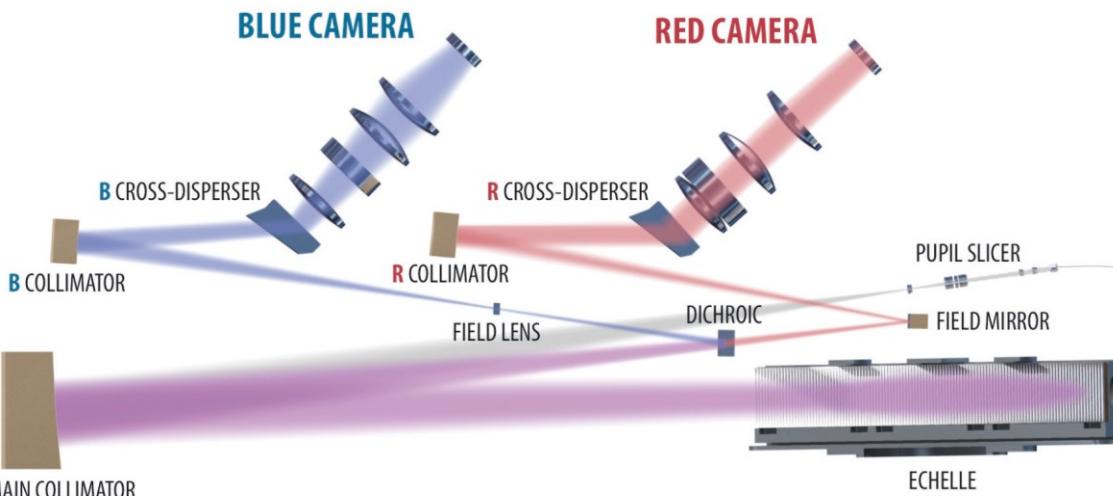
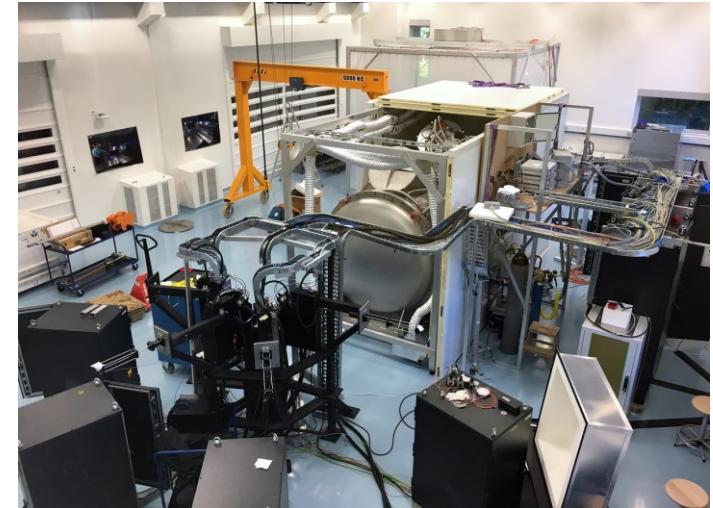
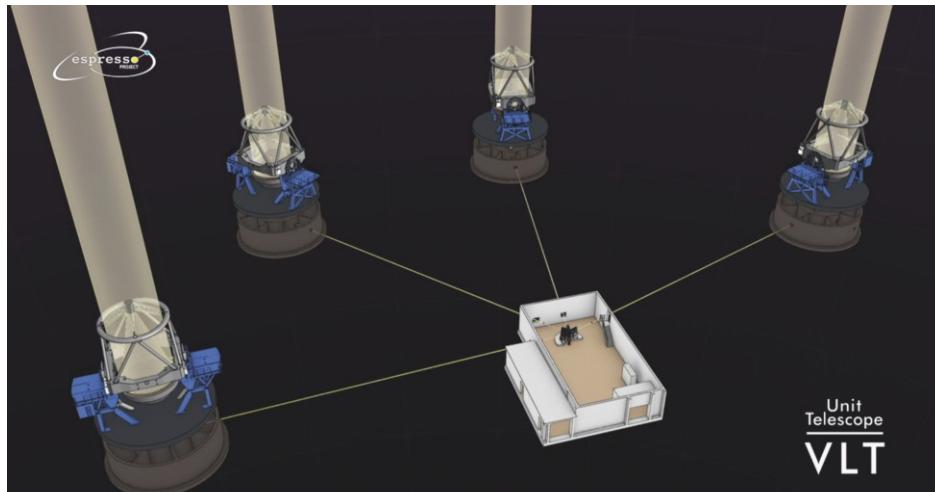


Kilo-nova in galaxy  
NGC 4993, Pian et al.



# Ultra-stable spectrographs

- Typical science case: long term measurement of radial velocities of planets: < m/s precision required



# Worked example

## Spectrograph Design for instrument scientists

# Conservation of étendue

- Etendue (also called throughput) must be conserved to build an efficient instrument
- At any given aperture, étendue is the product of the area of the aperture ( $A$ ) and the solid angle of the rays accepted ( $\Omega$ )
- In particular, for an astronomical instrument

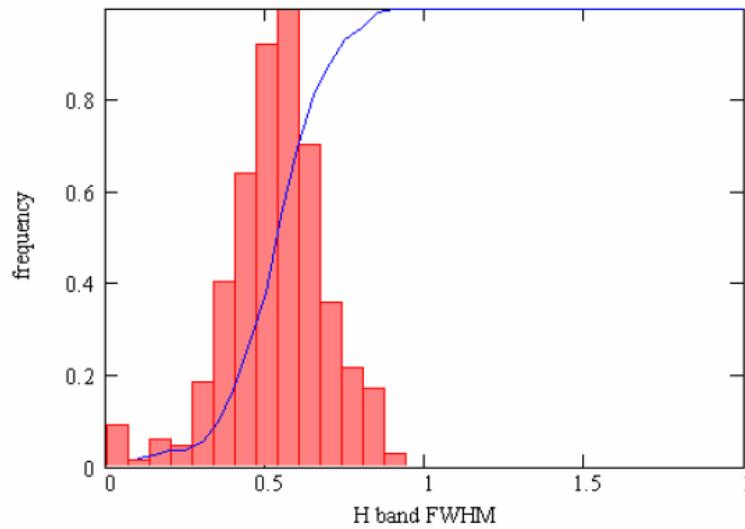
$$A\Omega_{tel} = A\Omega_{det}$$

*Using a few assumptions and simple geometrical optics, we can calculate some important instrument parameters.....*

# Back-of-the-envelope instrument for the ESO VLT (1)

## ■ Scientific goal

- Detection and characterisation of galaxies (in clusters) at redshift  $z=1-2$  (peak period of star formation)
- Velocities resolution  $\sim 10 \text{ km s}^{-1}$  to measure rotation curves
- Spatial/angular resolution set by the seeing limit at the telescope



H median FWHM = 0.53 arcsec  
David Lupton School, 2019, Spectroscopy

## ■ Scientific requirements

- $z=1-2 >>$  most diagnostic lines fall in the  $0.8-2.4 \mu\text{m}$  wavelength range
- Best seeing  $\sim 0.4 \text{ arcsecs}$  in K band ( $2.0-2.4 \mu\text{m}$ )  $>>$  0.2 arcsec per pixel
- Spectral resolving power set by the vel. res'n to  $R (\frac{\lambda}{\delta\lambda}) \cong 3000$

# Back-of-the-envelope instrument (2): camera focal ratio

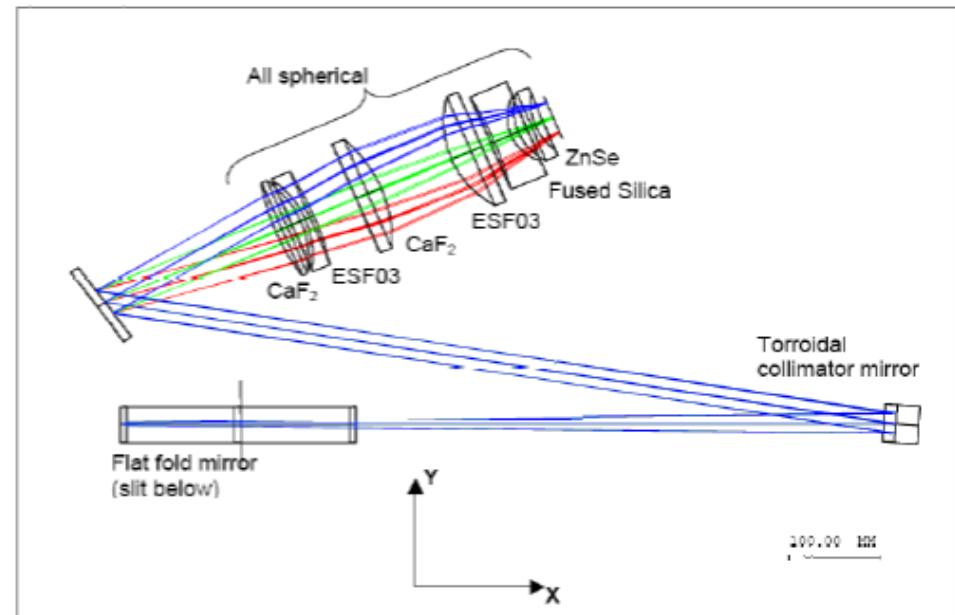
- Conservation of étendue and the detector pixel size determines the camera f-number:

$$D_{tel} \theta_{tel} = d_{pix} \theta_{cam}$$

$$\theta_{cam} = \frac{d_{pix}}{f\_number}$$

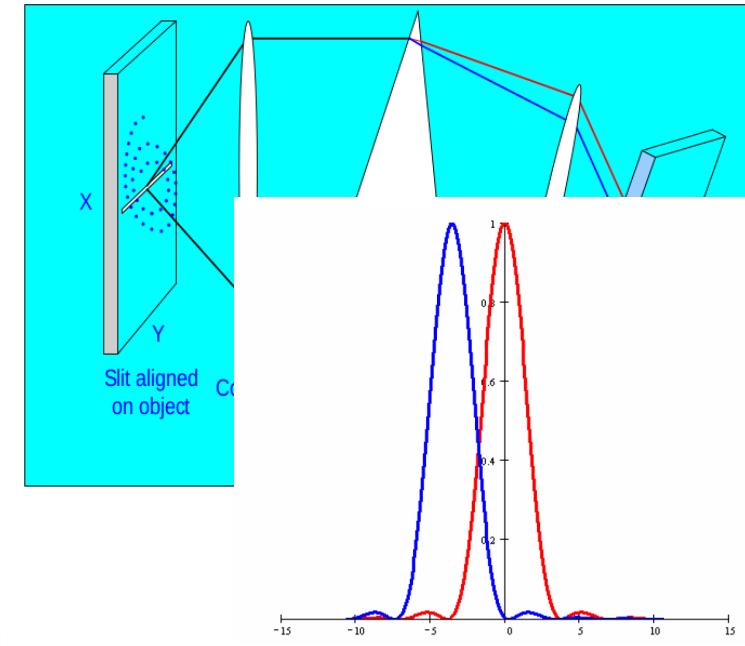
- $\theta_{tel}=0.2\text{arcsecs}$ ,  $D_{tel}=8.2\text{m}$
- $d_{pix}=18\mu\text{m}$
- Camera should be f/2.3
- Such a fast camera is hard but not impossible

*Optical layout of the KMOS spectrograph*



# Spectral resolving power

- With the f-number of the camera defined and the spectral resolution defined by the astronomers, the pupil size can be calculated
- To resolve two spectral lines at the array  $\Delta\theta = \Delta\lambda \frac{\delta\theta}{\delta\lambda}$
- $\frac{\delta\theta}{\delta\lambda}$  comes from the grating equation  $m\lambda = d \sin(\theta i) + \sin(\theta r) \gg \frac{\delta\theta}{\delta\lambda} = \frac{m}{d \cos(\theta)}$
- Using  $\Delta\theta = \theta_{cam} = \frac{d_{pix}}{focal\ length}$  and the definition of resolving power  $R = \frac{\lambda}{\delta\lambda}$



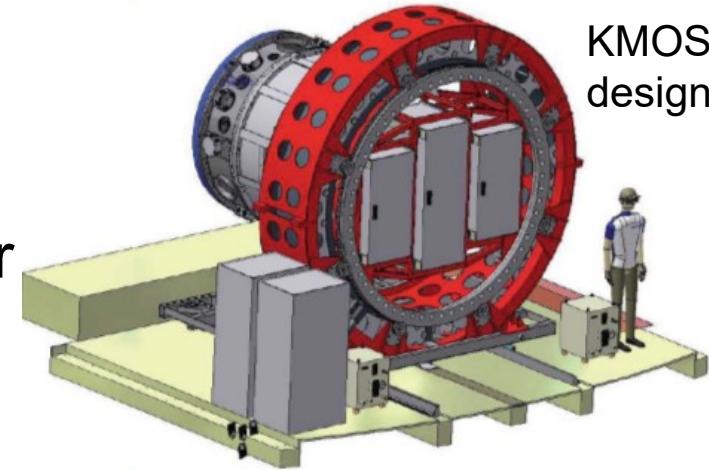
$$R = \frac{2f \tan \theta r}{d_{pix}}$$

*f=focal length*

# Back-of-the-envelope instrument (4): Pupil/beam size

- For KMOS, set  $R \sim 4000$  over 2 pixels to Nyquist sample the spectral lines
- Set  $\theta r = 45$  degrees for a first estimate
- Then  $f \sim 72\text{mm}$  and the beam size for an f/2.3 camera is  $\sim 31\text{mm}$
- In practice the KMOS beam is 33mm diameter
- With these key numbers, a first idea of the size and technical difficulty of the instrument may be estimated but doesn't give the whole story!

$$R = \frac{2f \tan \theta r}{d_{pix}}$$



# A back of the envelope instrument for the ELT

- 0.2arcsecs/pixel, 15 $\mu\text{m}$  pixels, 39-m telescope
- $\sim f/0.4$  on the ELT
- Alternatively, define the camera to be  $f/2$
- The seeing disk is imaged onto  $\sim 150 \mu\text{m}$  at the detector or 10 pixels
- Such oversampling implies larger focal plane>>larger instruments>>increased detectors costs>>power consumption etc....
- The tendency towards larger arrays of smaller pixels on larger telescopes is a real challenge for instrument design!

# At the diffraction limit

## ■ Diffraction limited instruments:

- H-band diffraction limited FWHM = 10mas
- 0.005 arcsecs/18 μm pixel
- ~f/19 on the ELT

■ Aside from the scientific benefits, using adaptive optics can help reduce image sizes to the point where instruments are easier to build

■ For an instrument to operate at the diffraction limit of a telescope set the pixel field of view:

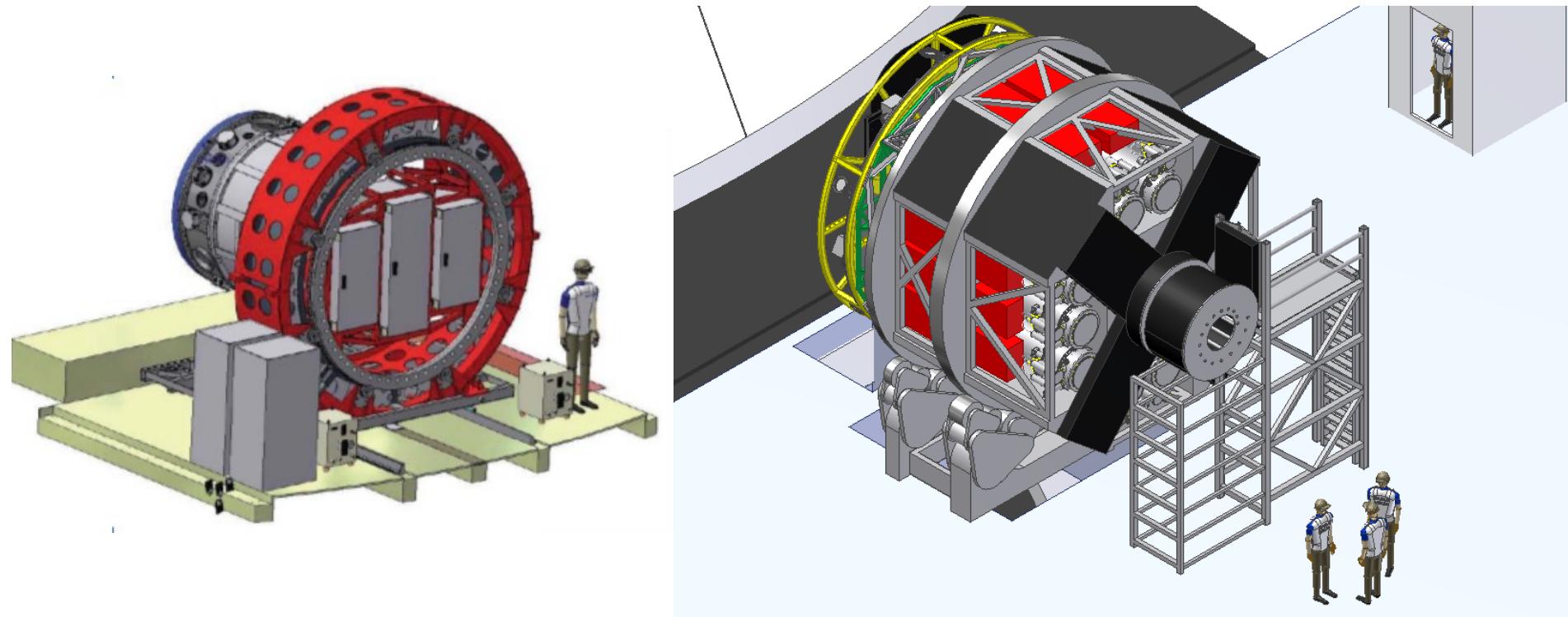
$$\theta_{tel} = \frac{1.22\lambda}{D_{tel}} \quad \text{then} \quad \frac{d_{pix}}{f\_number} \sim \lambda$$

■ A diffraction limited instrument can go on any telescope!

# The challenge of building instruments for an ELT

- Instrument size typically increases with telescope size
  - VLT f-ratio: f/15 plate scale: 0.582mm/arcsec
  - E-ELT f-ratio: f/17.7 plate scale: 3.3mm/arcsec
- Some things get easier
  - Relaxed positioning tolerances
  - KMOS pick-off arms position to 0.2arcsecs = 120μm on VLT; would be 660μm on E-ELT
- Some things get harder....

# KMOS vs EAGLE

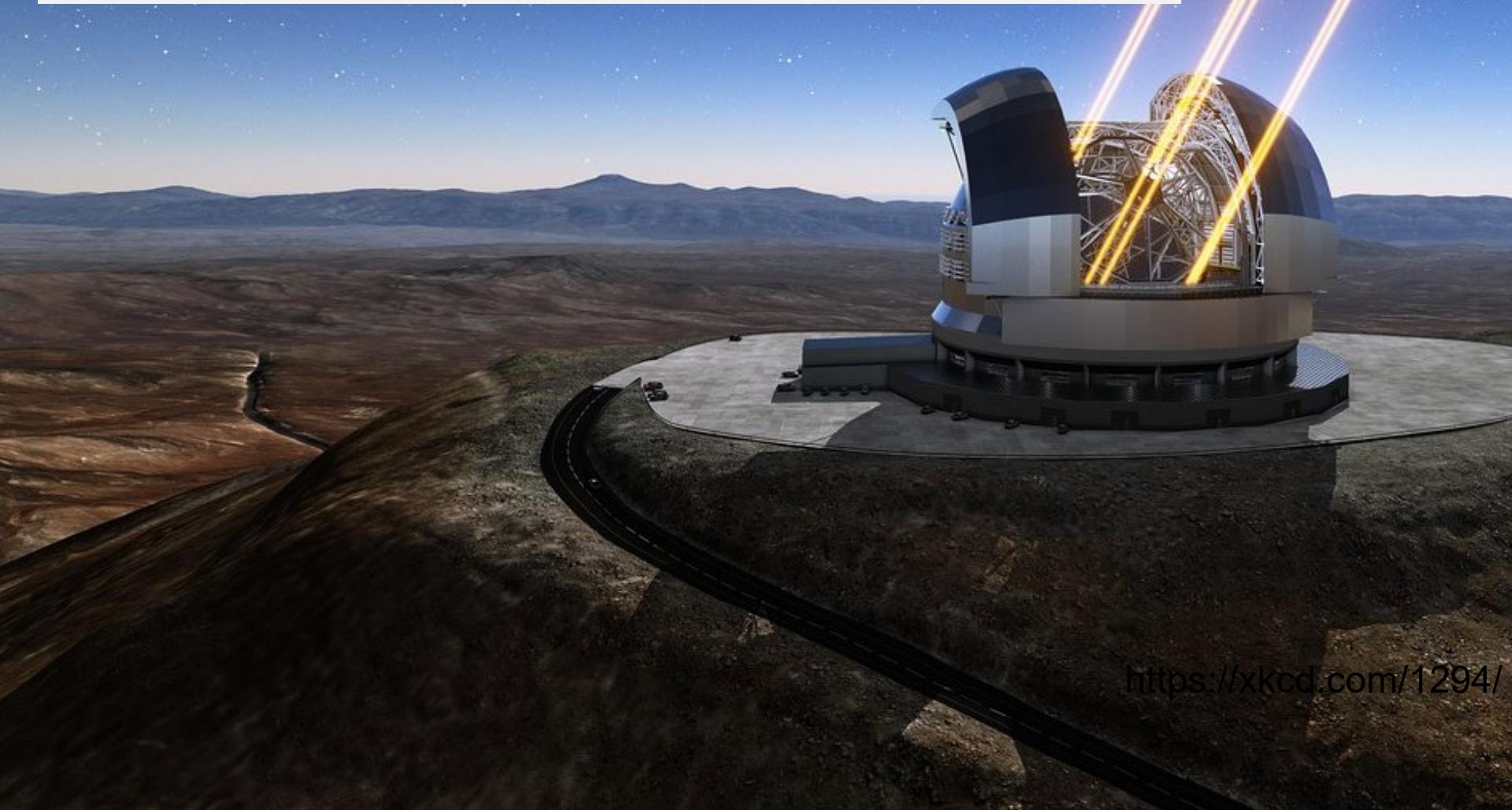


KMOS: 0.2arcsec pixels on 8-m telescope, 7arcmin field of view

EAGLE concept: 40mas pixels on a 39-m telescope, 7arcmin field of view

These instruments have the **SAME AΩ product/étendue** and can select the same numbers of objects (20-24).

# ESO's Extremely Large Telescope and its instruments



<https://xkcd.com/1294/>



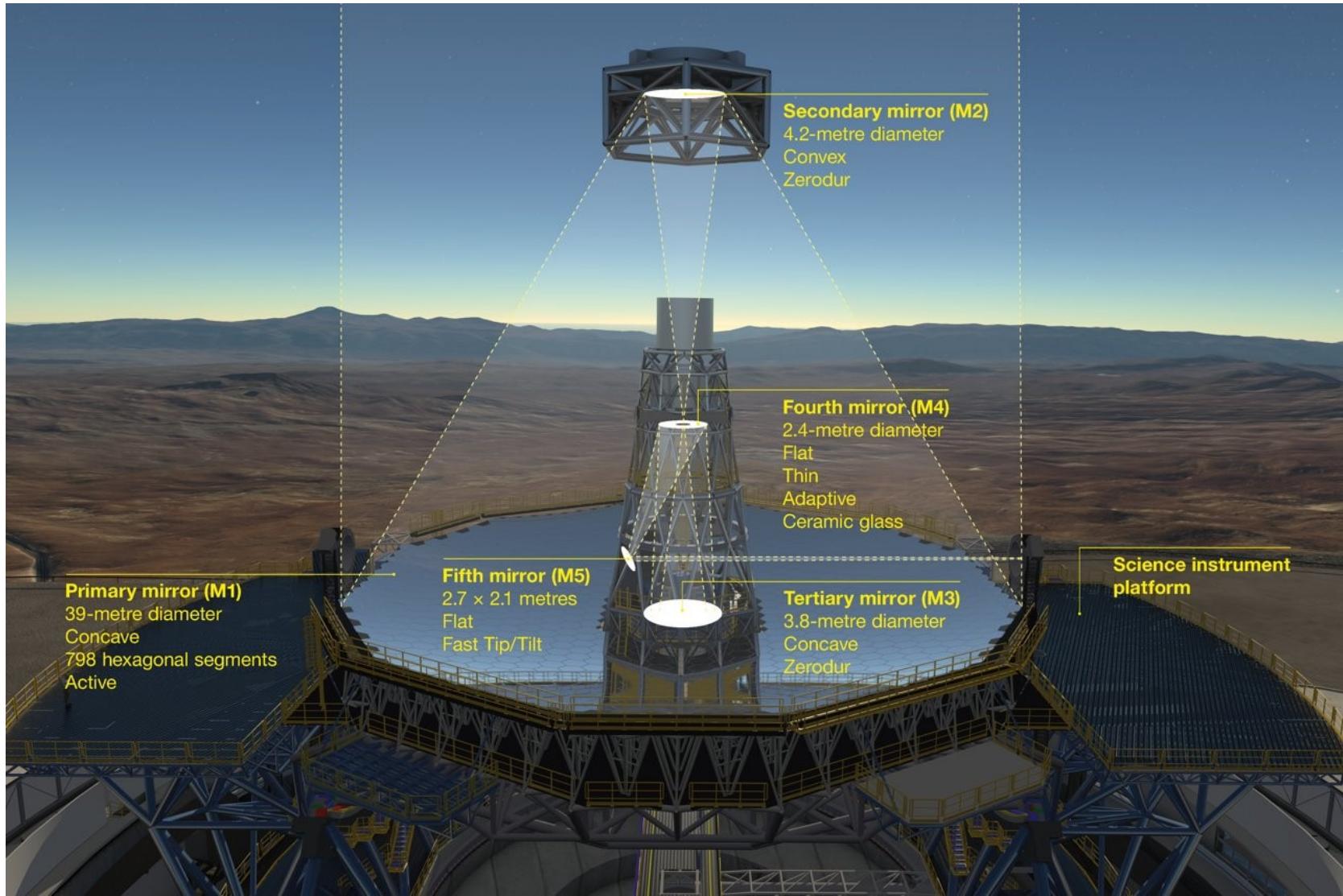
# ESO's Extremely Large Telescope

- THE VERY LARGE TELESCOPE
- THE EXTREMELY LARGE TELESCOPE
- THE OVERWHELMINGLY LARGE TELESCOPE  (CANCELED)
- THE OPPRESSIVELY COLOSSAL TELESCOPE
- THE MIND-NUMBINGLY VAST TELESCOPE
- THE DESPAIR TELESCOPE
- THE CATACLYSMIC TELESCOPE
- THE TELESCOPE OF DEVASTATION
- THE NIGHTMARE SCOPE
- THE INFINITE TELESCOPE
- THE FINAL TELESCOPE

<https://xkcd.com/1294/>



# ELT optics





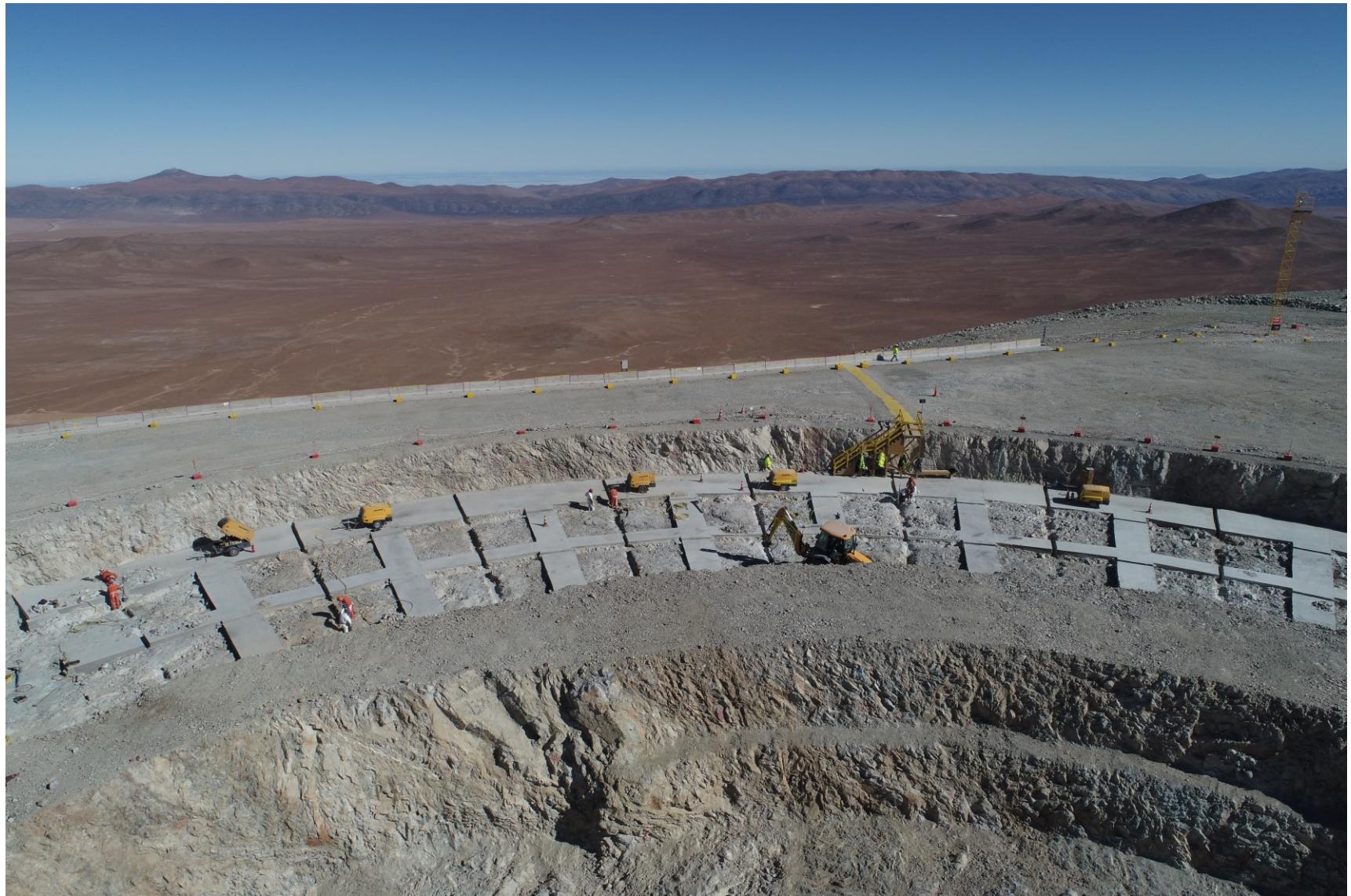
# ESO open house 2011



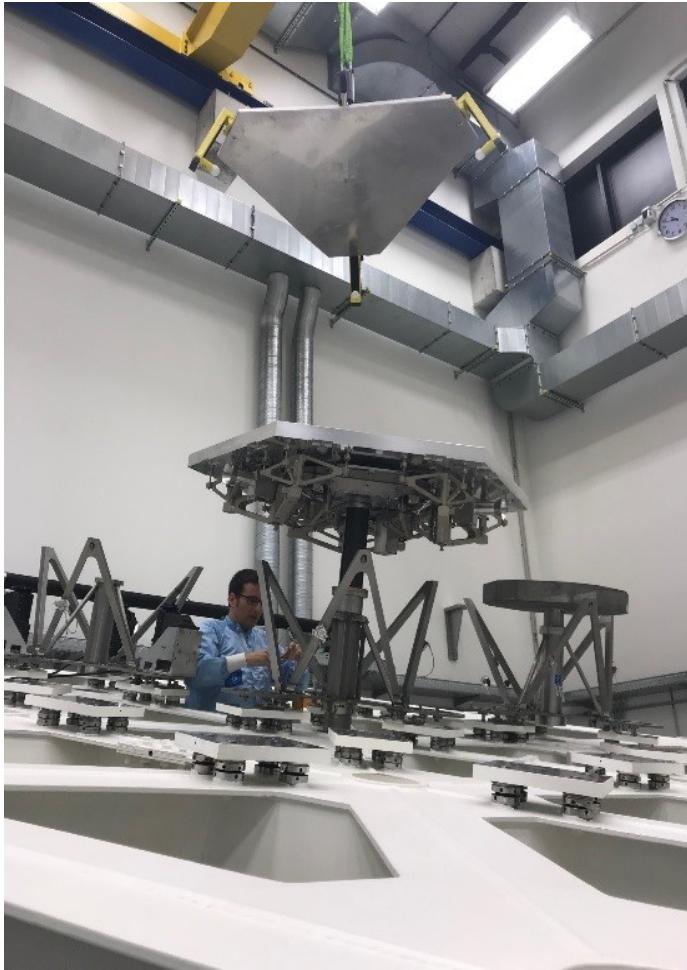
# The ELT Site: Cerro Armazones



# ELT site: construction

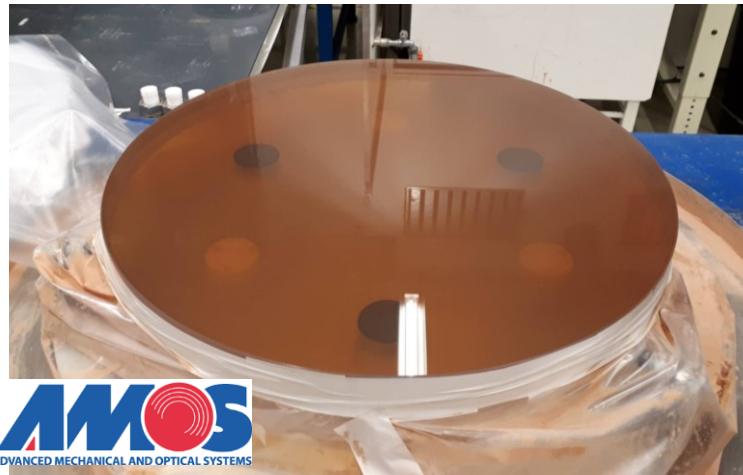
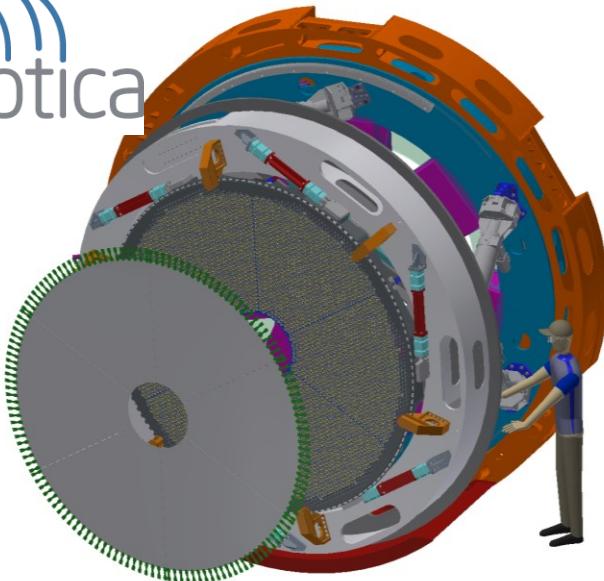


# M1 Mirror Blanks



# M4 deformable mirror

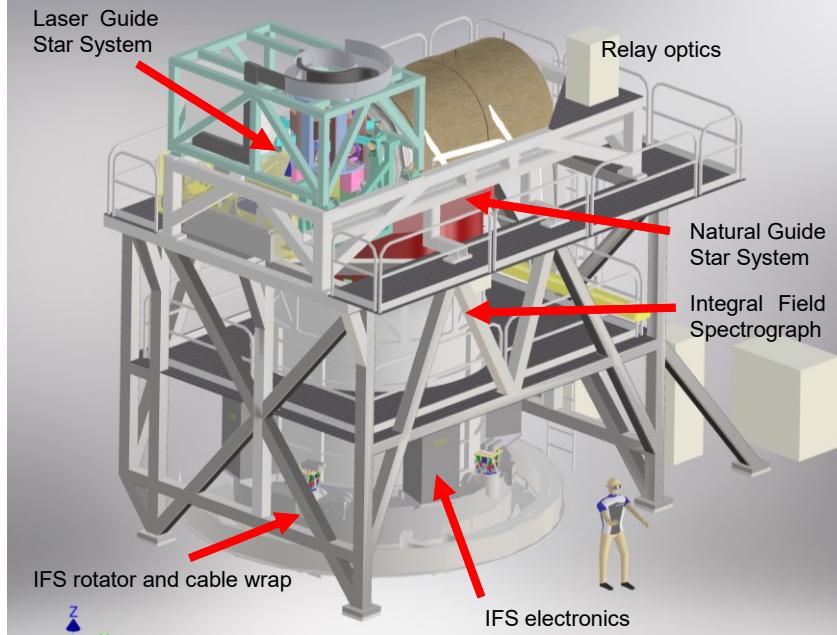
adoptica



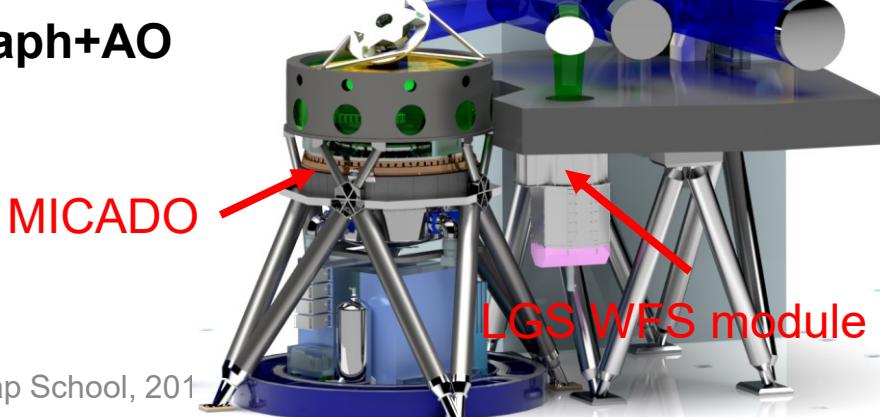
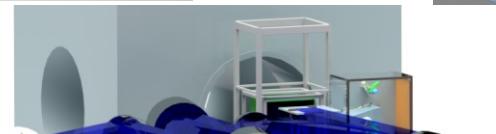
, Q2 2019



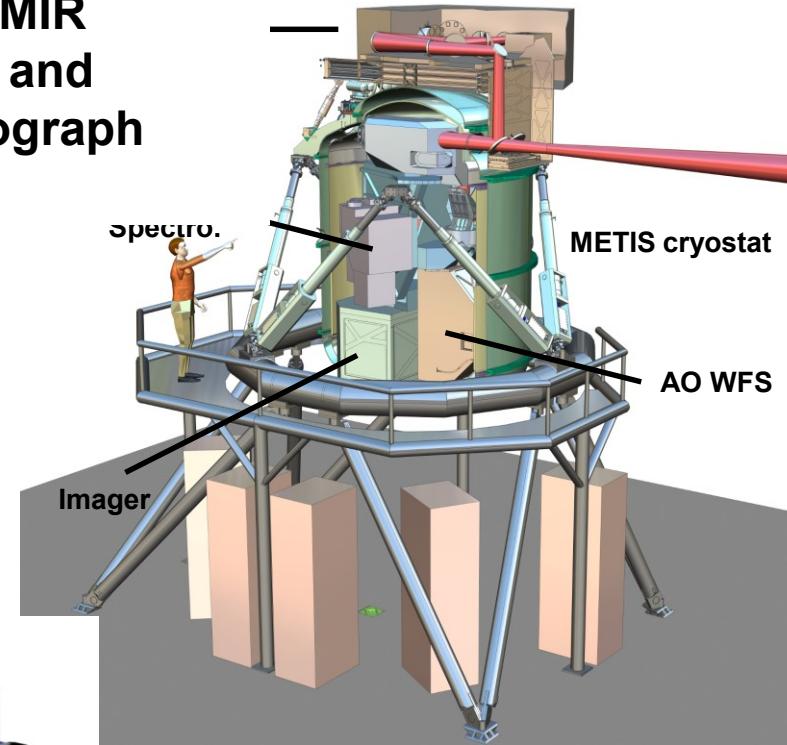
# ELT instruments



**HARMONI** optical  
to NIR 3D  
spectrograph+AO

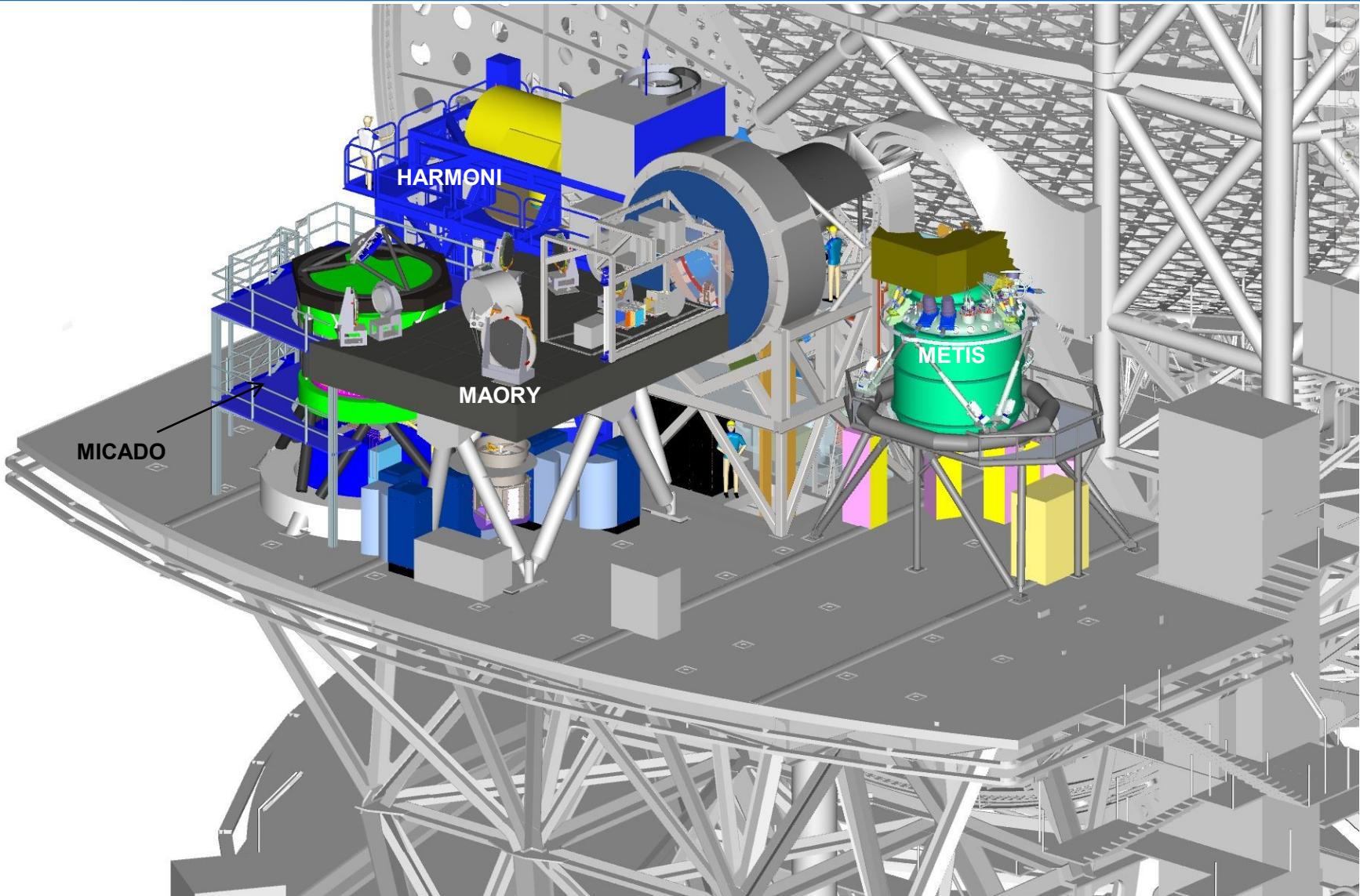


**METIS MIR**  
imager and  
spectrograph  
+AO

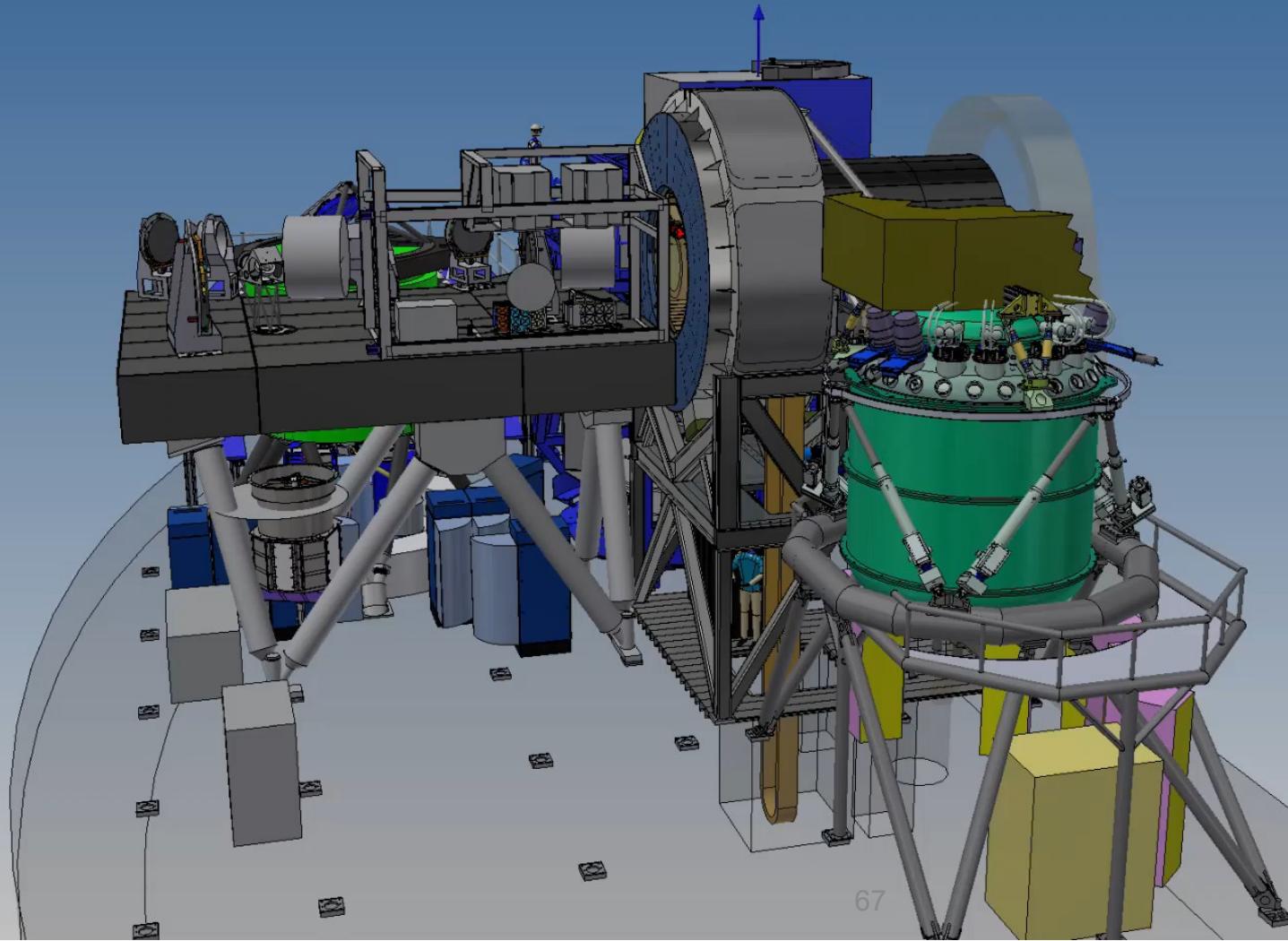


**MICADO** NIR imager and  
slit spectrograph with the  
**MAORY** MCAO module

# Instruments on Nasmyth A



# Instruments on Nasmyth A



- A question for the audience:
- Smart phones drive the development of large format CCDs for their built in cameras.
- Why are these not useful for astronomy? Discuss.

# Links and acknowledgements

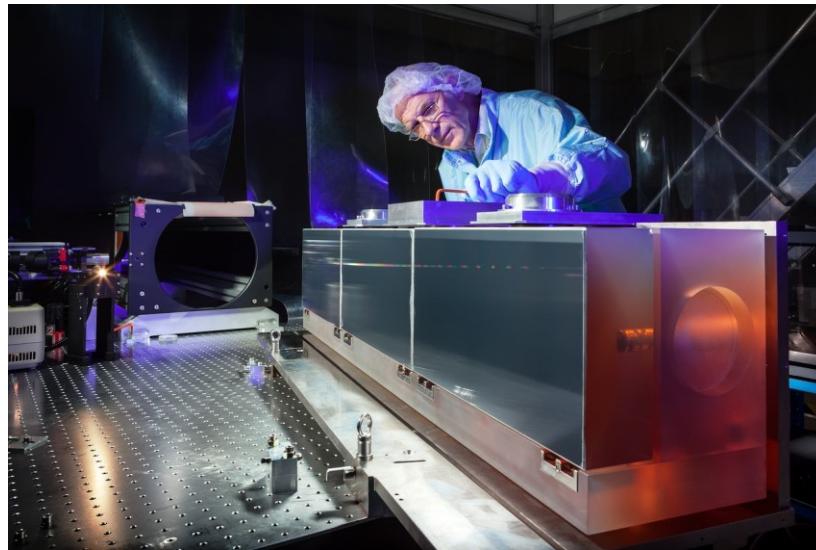
- AO open/closed loop GIF (slides #):  
[www.mpg.mpi-a.de/AO/INTRO/AOWFSintro.html](http://www.mpg.mpi-a.de/AO/INTRO/AOWFSintro.html)
- VISIR detector image from Mills, Beuville & Corrales  
(spie.org/newsroom/3786-evolution-of-long-wavelength-astronomy-sensors)
- KMOS-3d survey: see Wisnioski, E et al (2015),  
ApJ, 799, 209

# What sets the resolving power of a spectrograph?

- Theoretical limit set by the number of rulings
  - Back to school calculation of this from a standard optics book
- Actually for spectrographs with moderate resolving power then the slit width sets the resolving power
  - Later – come to the example of small movements with the slit introducing a fake radial velocity signal

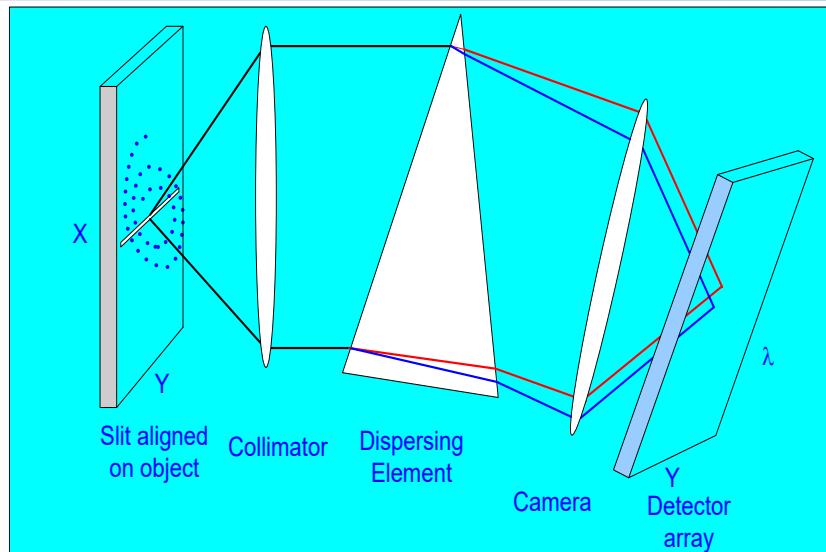
# Large high resolving power gratings

- The grating for the ESPRESSO spectrograph



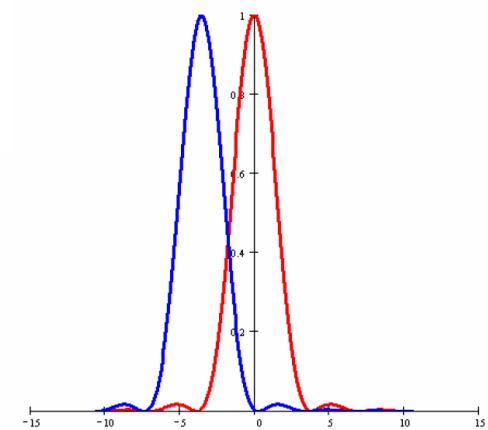
<https://www.eso.org/public/images/potw1537a/>

# Spectral resolving power, R



Angular separation of two wavelengths at the detector

$$\Delta\theta = \Delta\lambda \cdot \frac{d\theta}{d\lambda}$$



From conservation of étendue:  $\theta_p = \frac{D_T \theta_T}{d_p}$

Wavelengths are resolved when these are equal, so:

$$R = \frac{\lambda}{\Delta\lambda} = \frac{d_p}{D_T \theta_T} \cdot \lambda \frac{d\theta}{d\lambda}$$

**NB: max R depends on Pupil size, telescope D and angular size of the object.**